

How a cosmologist views the evidence for dark matter.



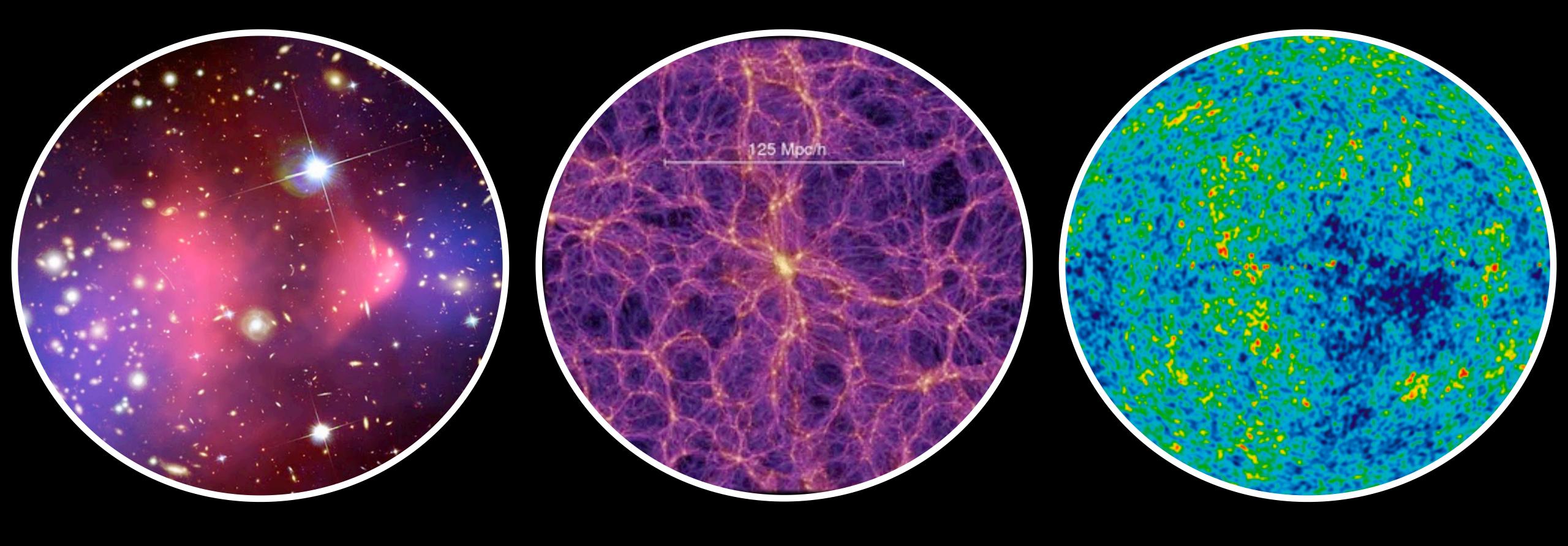
## **Bullet Cluster**

Two galaxy clusters colliding at ~4000 km/s

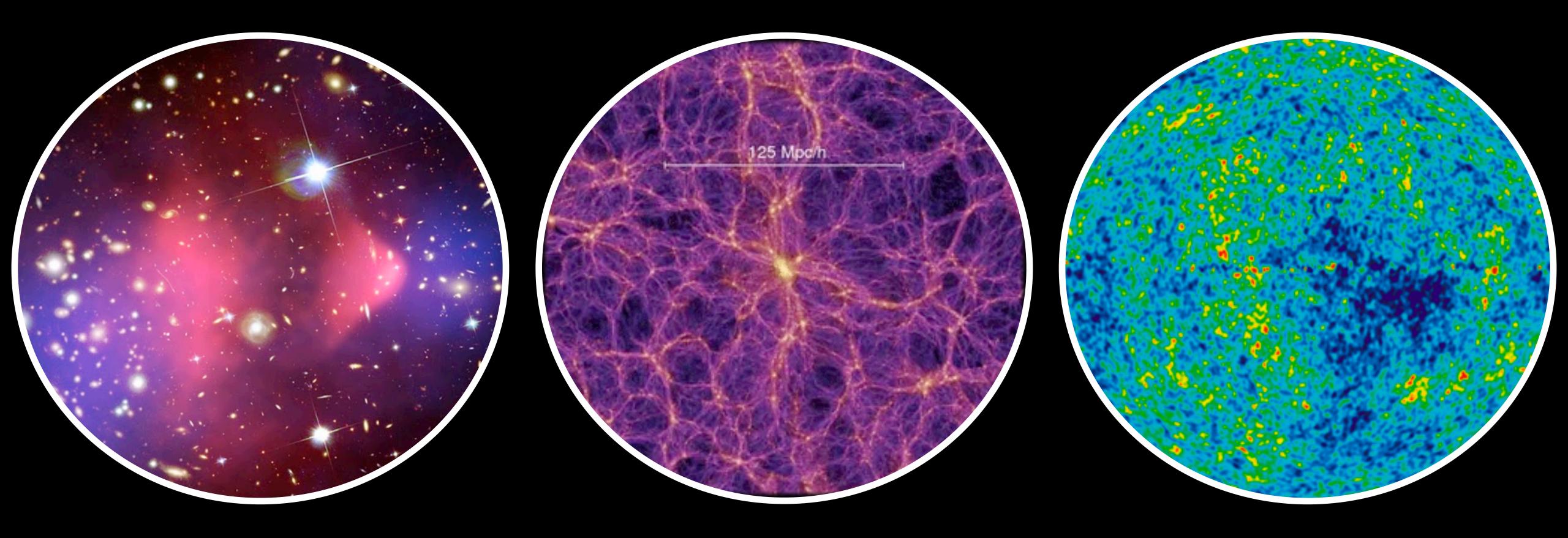
• Hot gas stuck in the middle - stars pass through.

 Total mass distribution traces the stars, which are only ~10% of the baryonic content.

Most mass must be dark!



How a cosmologist views the evidence for dark matter.



## **Early Universe:**

Radiation - Provides pressure that smears out over density.

**Baryonic Matter - Gravitational potential, but responds to radiation.** 

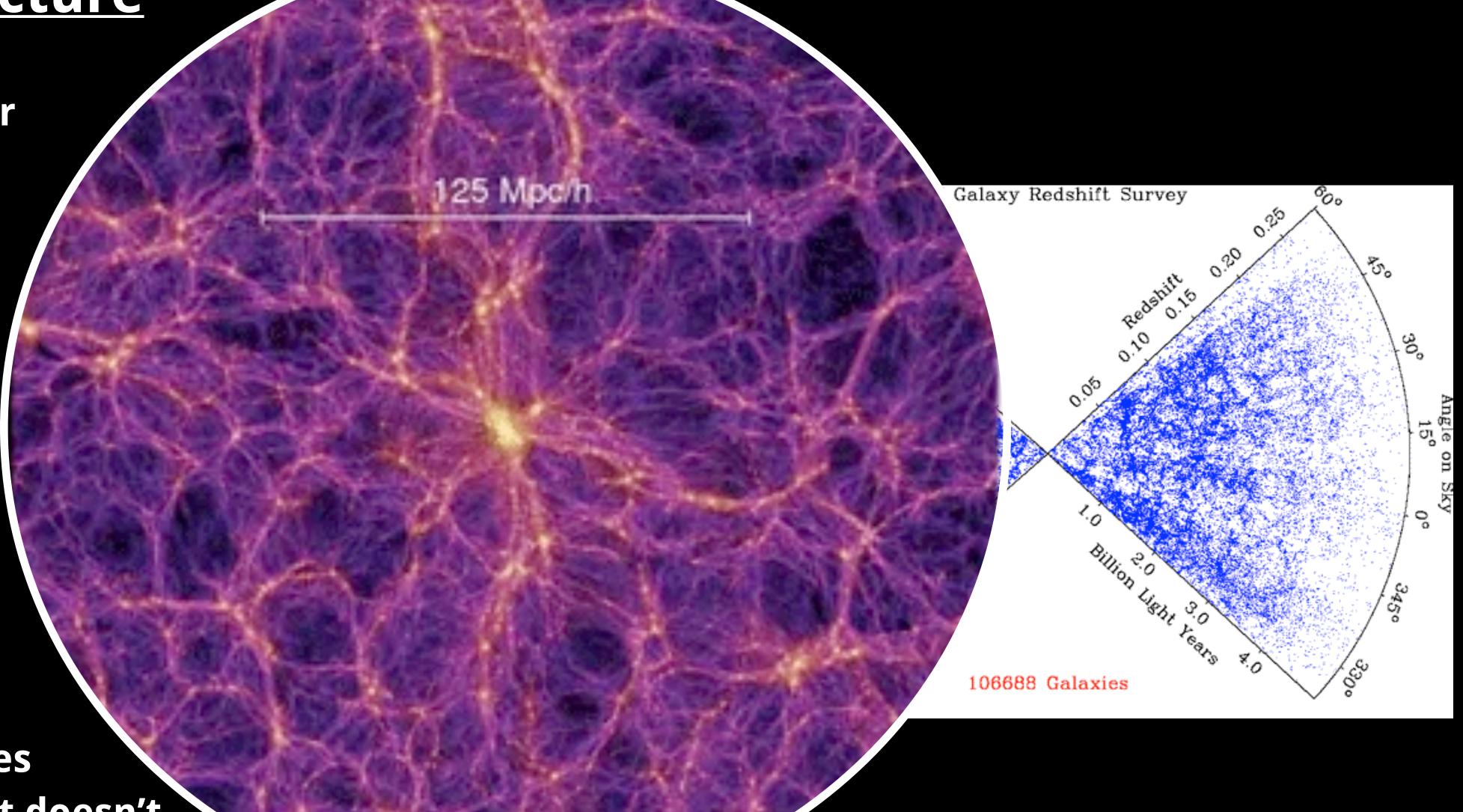
Dark Matter - Gravitational potential, does not respond to radiation.

Large Scale Structure

• When baryonic matter collapse it heats up.

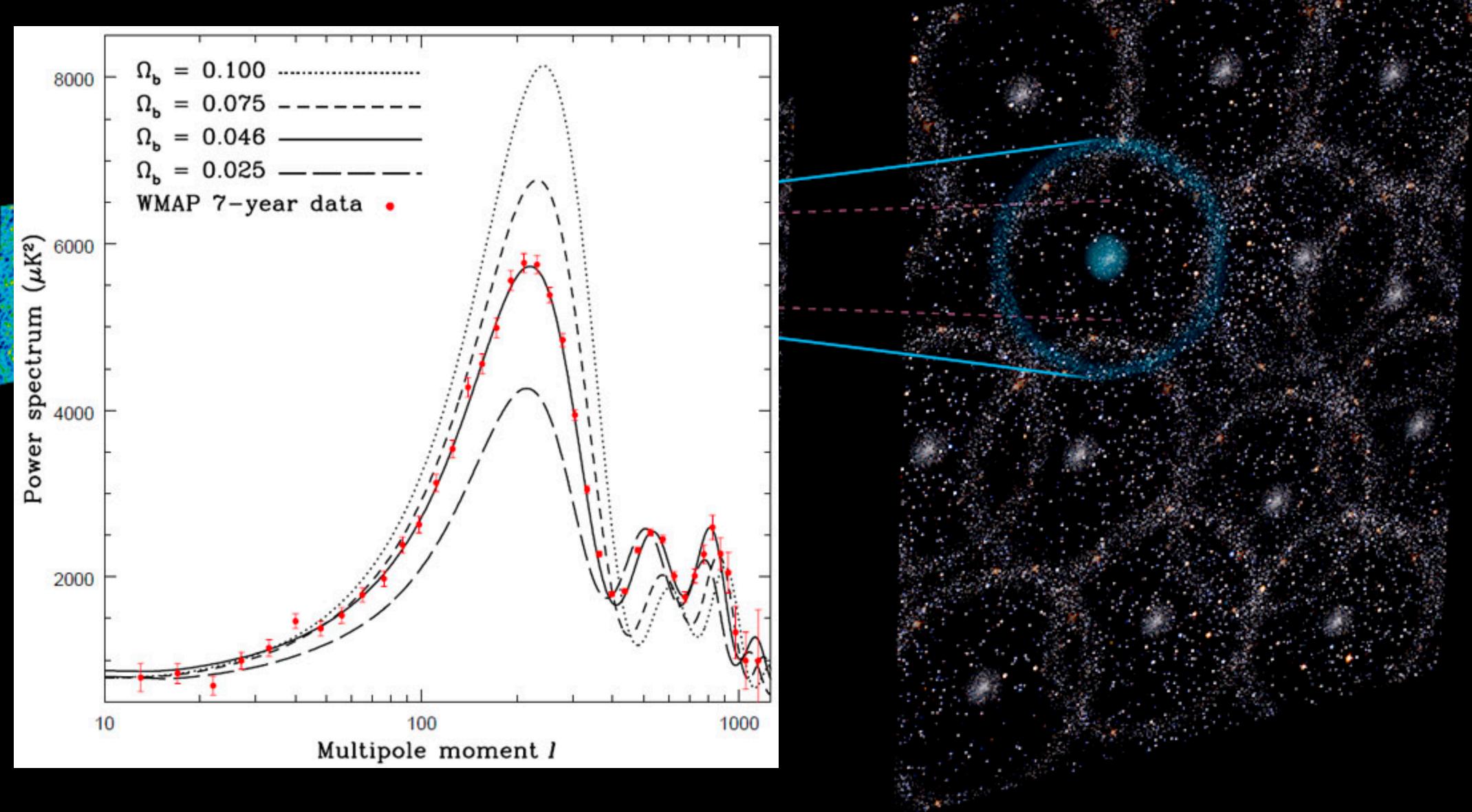
• This produces photons that cause the matter to expand (baryonic acoustic oscillations).

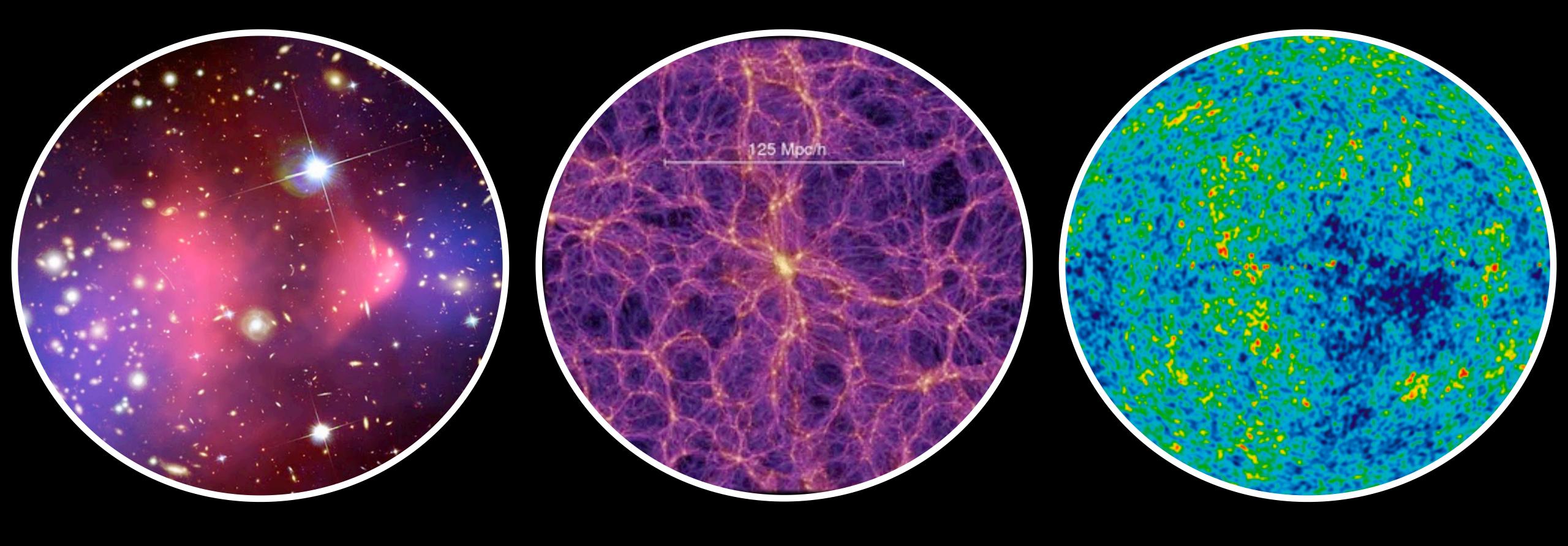
• Distribution of galaxies requires a matter that doesn't interact with photons.



# Large Scale Structure

## Large Scale Structure





How a cosmologist views the evidence for dark matter.

## Cosmic Microwave Background

• Interaction of radiation (photons) with cold matter (baryons) controls angular distribution of structure in the early stages of the universe.

2000

= 0.046 —

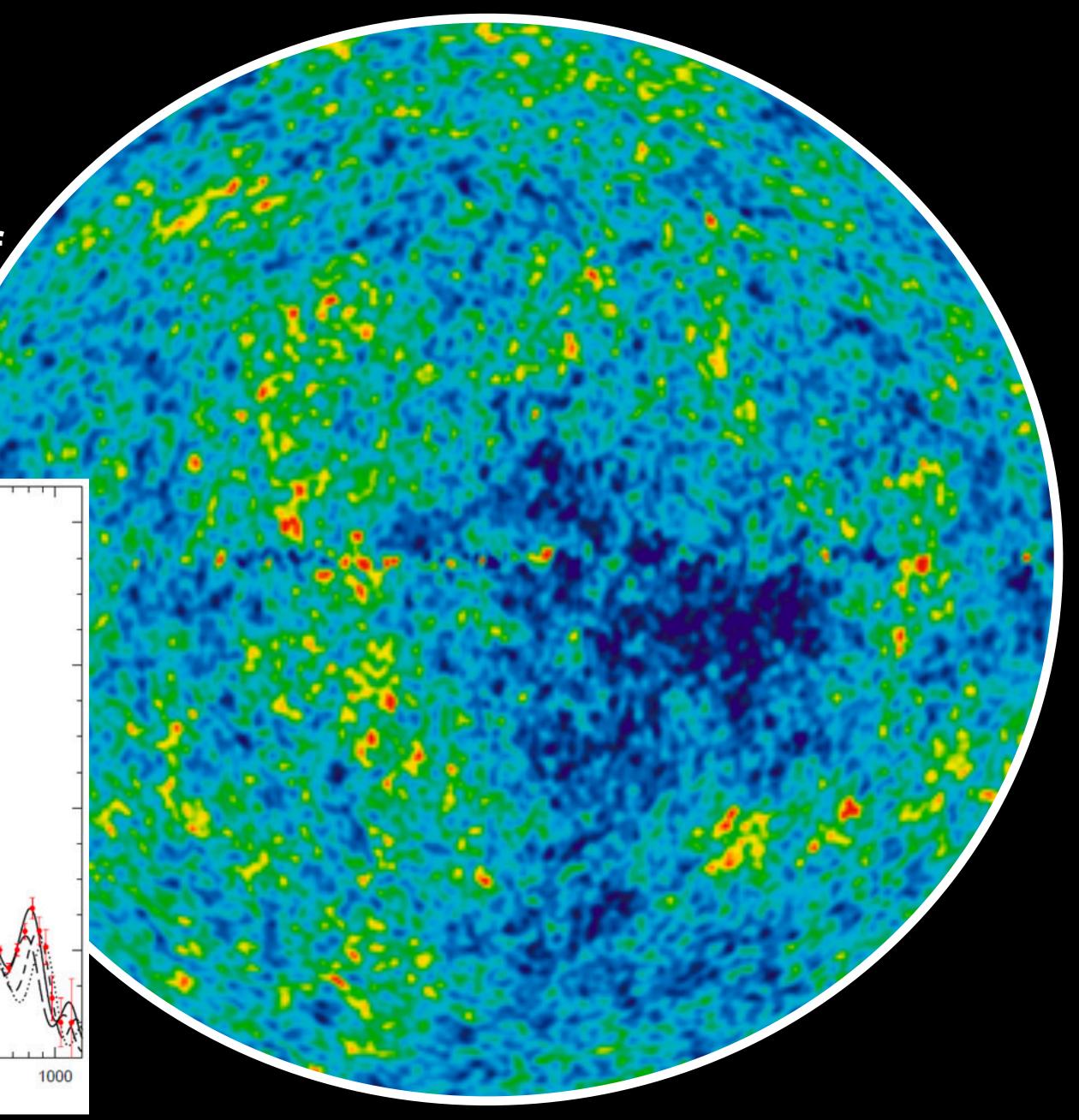
= 0.025 — — — -

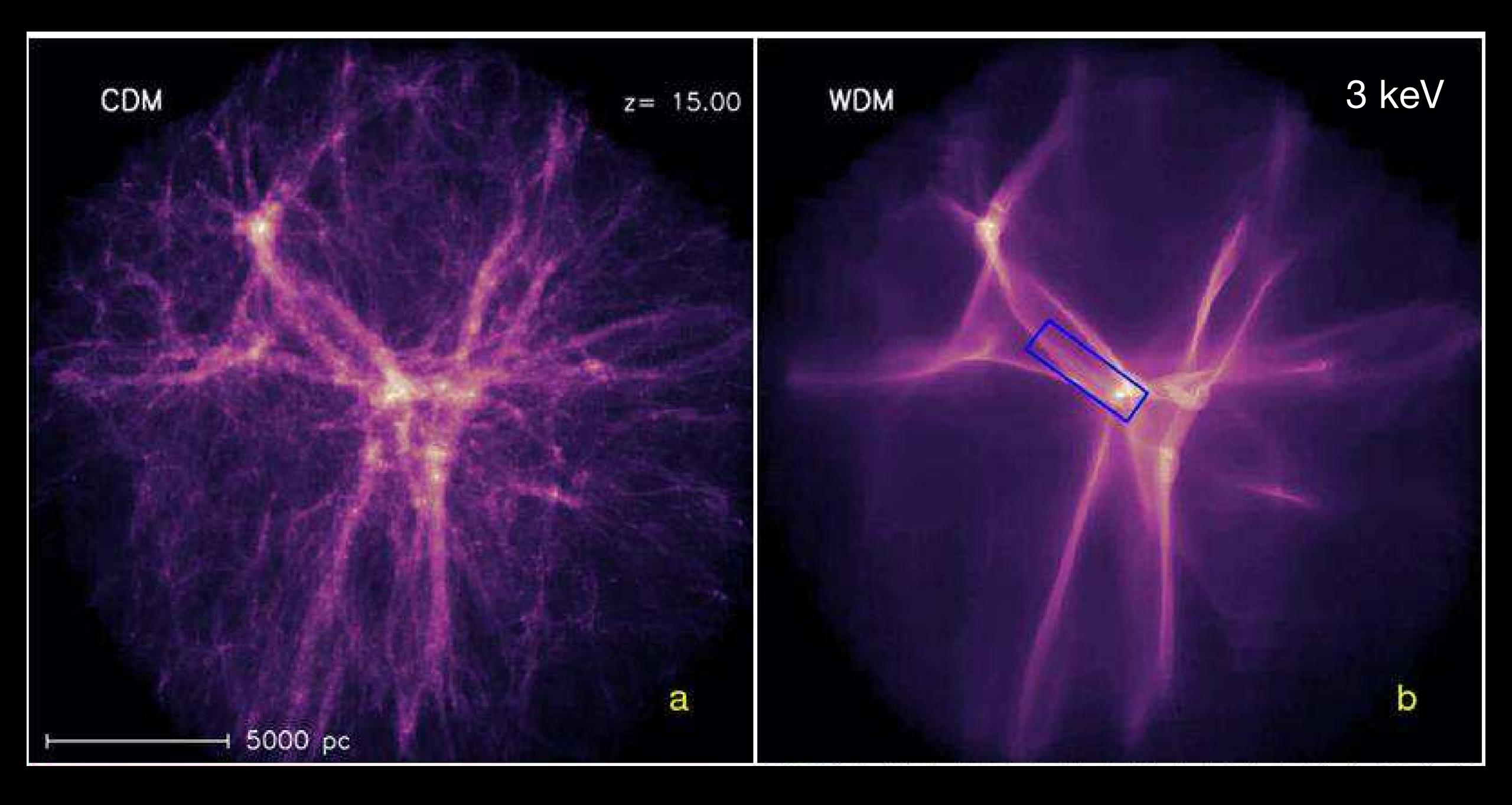
Multipole moment I

WMAP 7-year data .

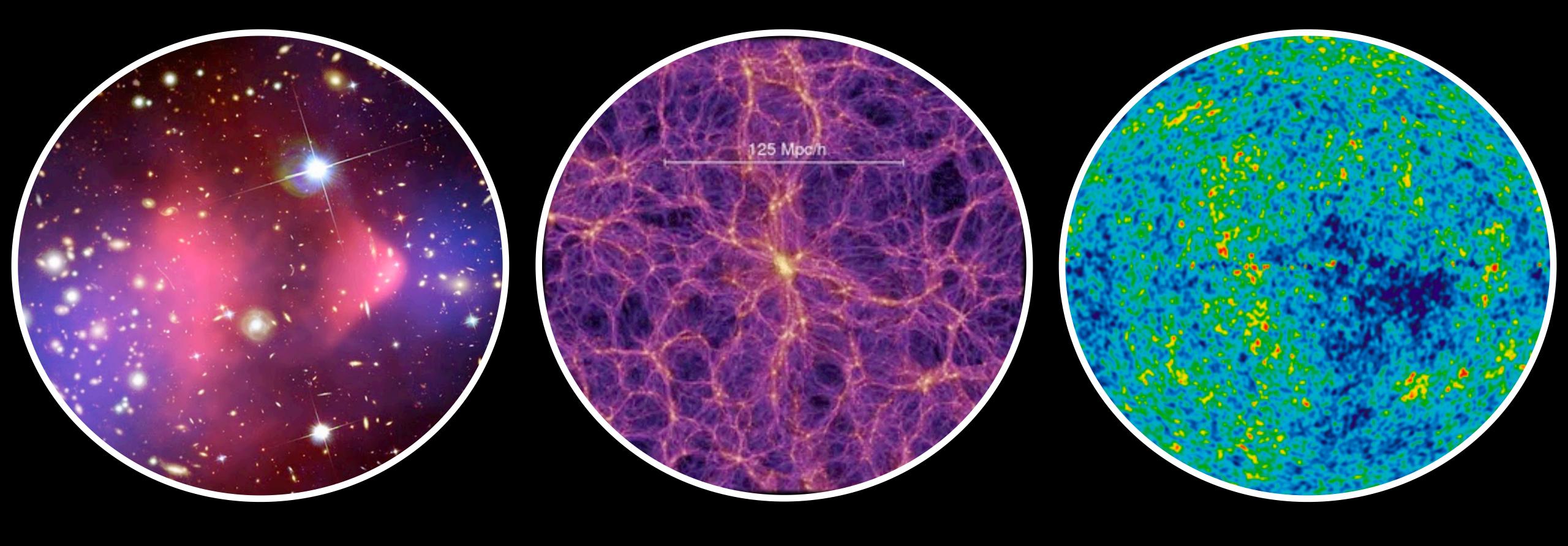
Large component that doesn't interact with

radiation is needed.



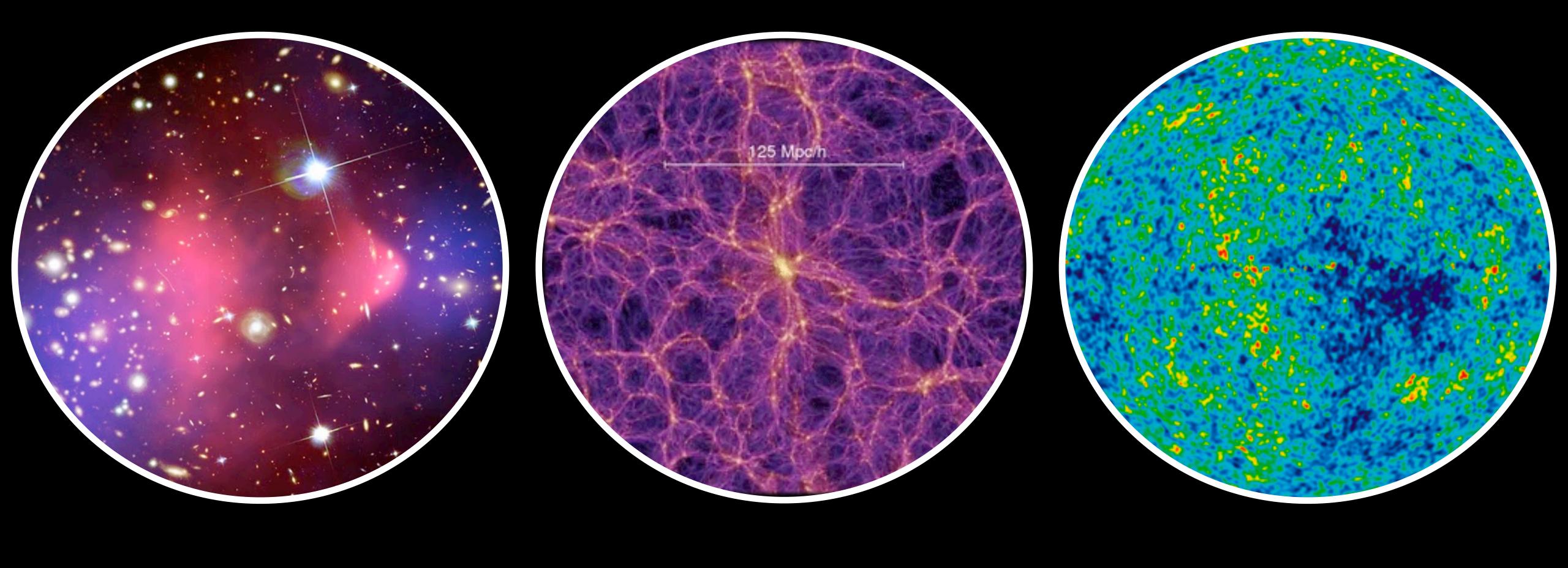


#### **The Present**



How a cosmologist views the evidence for dark matter.

#### **The Present**







#### **Maximum Mass of the Dark Matter Particle**



$$M_{\rm DM} \le M_{\rm dSph} = 2 \times 10^5 M_{\odot} = 10^{71} \text{ GeV}$$

#### Minimum Mass of the Dark Matter Particle

$$R_{\rm DM} \le R_{\rm dSph}$$
  $\lambda = \frac{h}{p}$   $R_{\rm DM} = \frac{1.2 \, {\rm eV} \, \mu {\rm m}}{p}$   $p = m_{\chi} (10 \, {\rm km \, s^{-1}})$   $m_{\rm min} \approx 3 \times 10^{-22} \, {\rm eV}$ 

 $30 \text{ pc} \approx 10^{26} \mu m$ 

#### Minimum Mass of the Dark Matter Particle

#### Pauli-Exclusionary Principle

$$f(\vec{x}, \vec{p}) = \mu$$

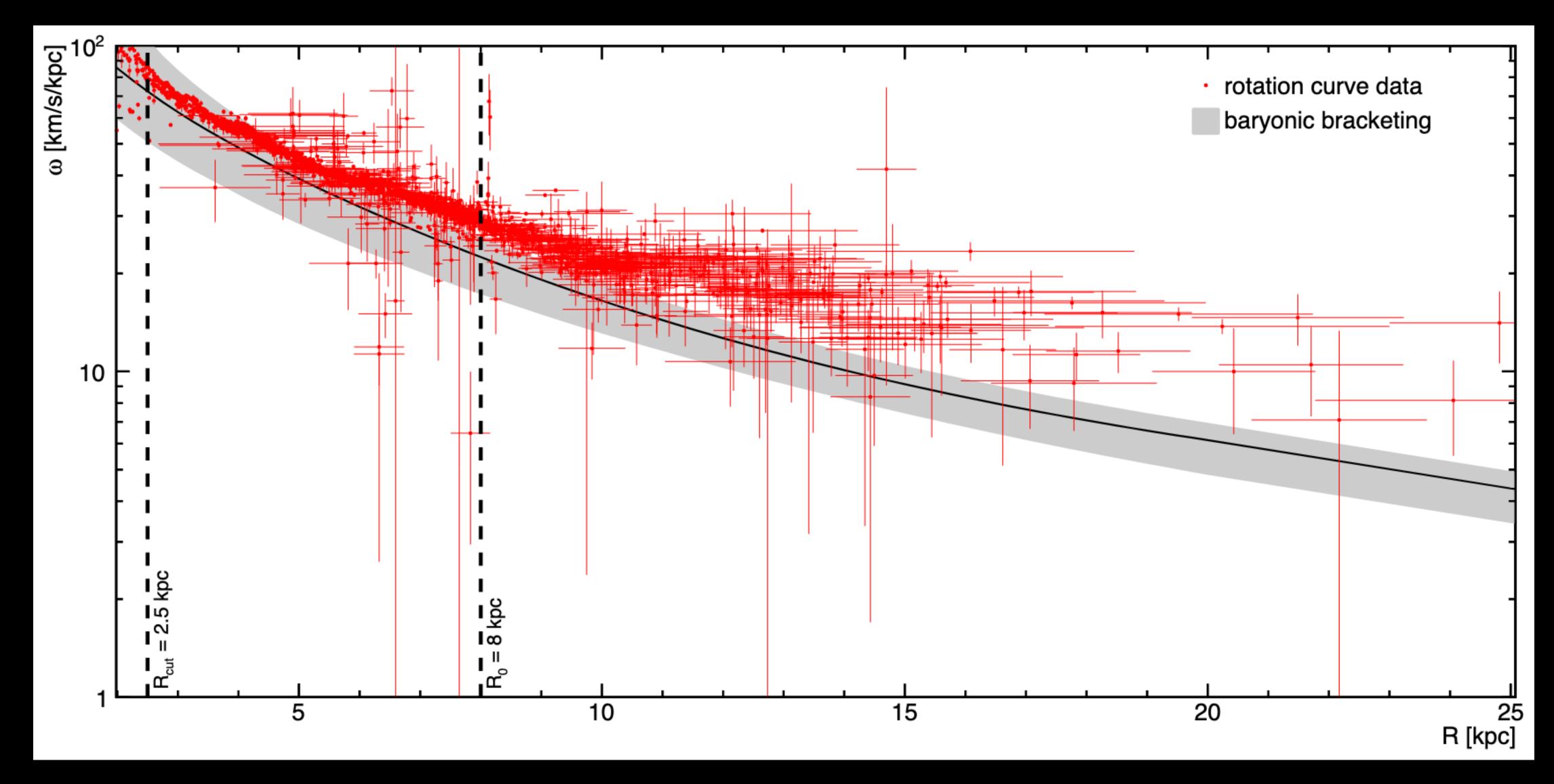
$$n_x \le \int_0^{p_{\text{max}}} \frac{d^3p}{(2\pi)^3} f(\vec{x}, \vec{p}) = \mu p_{\text{max}}^3$$

$$m_x^3 v_{\text{esc}}^3 \ge n_\chi = \frac{2 \times 10^{11} M_{\odot}}{4\pi (35 \text{kpc})^3 m_\chi}$$

If dark matter is a fermion

$$m_{\chi} > 100 \text{ eV}$$

## **Local Observations**

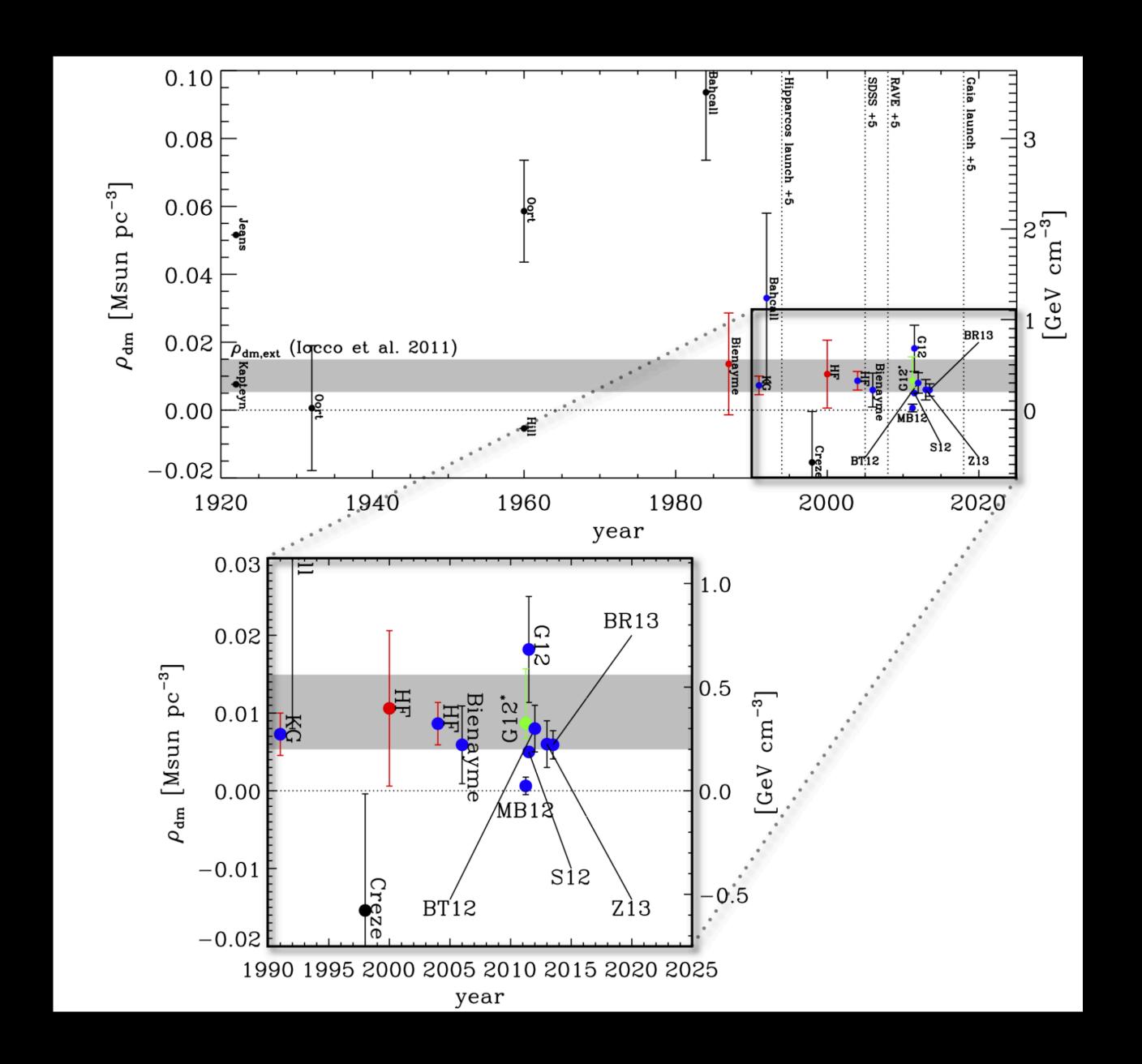


See:

Pato, locco (2015)

## **Local Observations**

- Many observations of local density
- Stellar rotation curves, height distributions, GAIA data
- o Still 50% uncertainties.



## **Local Observations**









## Local Observations

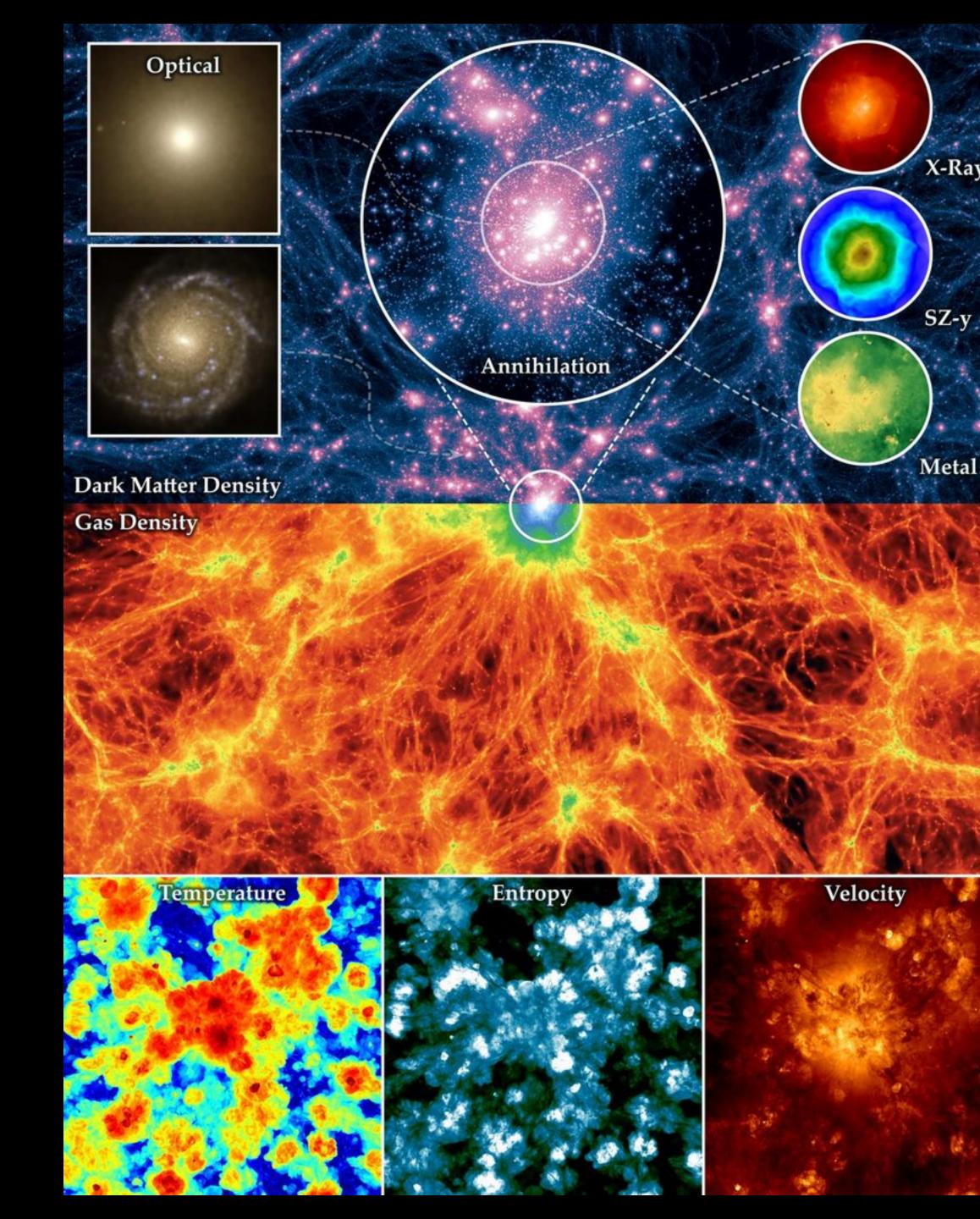
Navarro-Frenk White (NFW) Profile

$$\rho_{\chi} = \rho_0 \left(\frac{r}{r_s}\right)^{-1} \left(1 + \frac{r}{r_s}\right)^{-2}$$

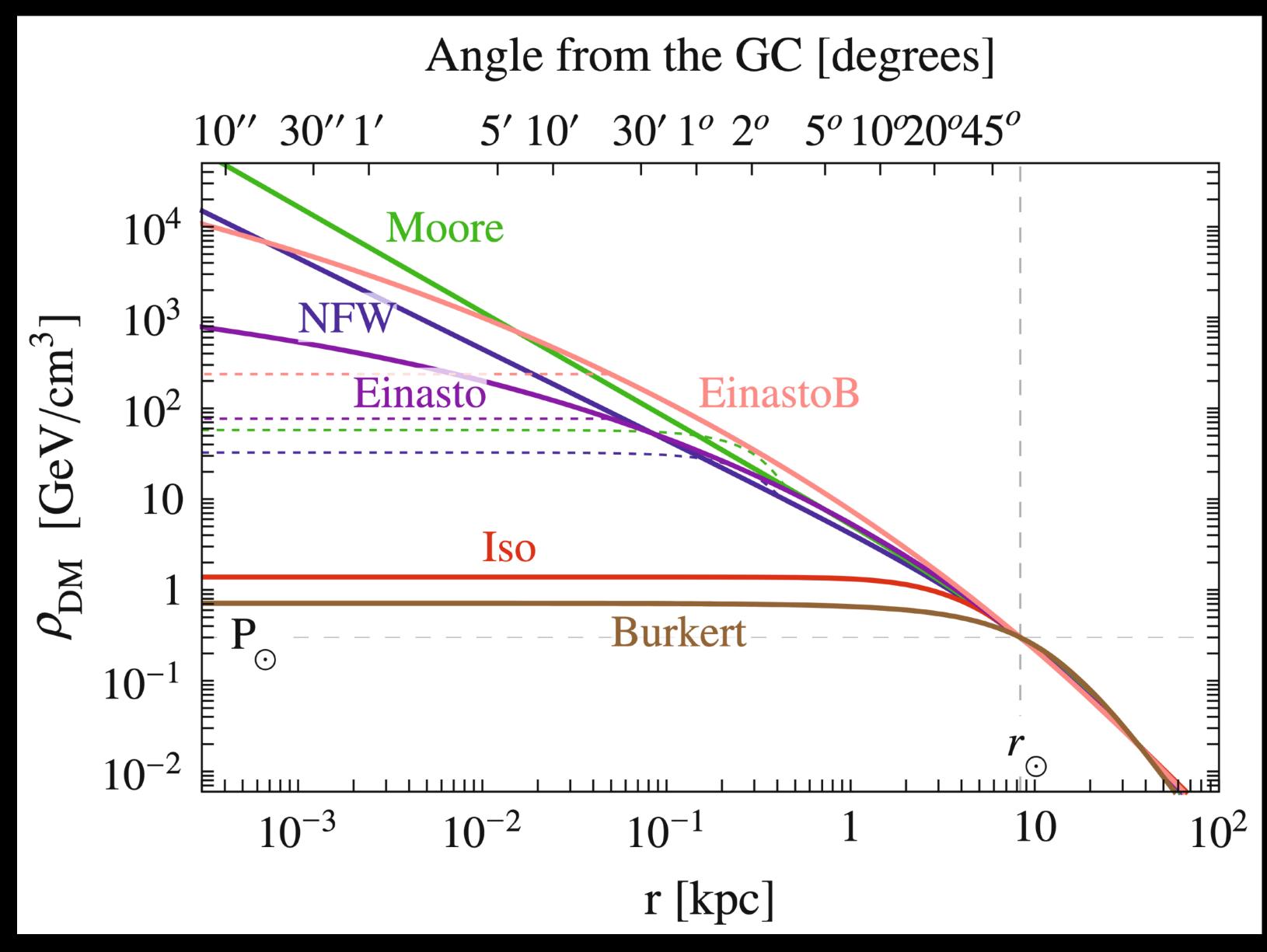
Generalized NFW Profile

$$\rho_{\chi} = \rho_0 \left(\frac{r}{r_s}\right)^{-\alpha} \left(1 + \frac{r}{r_s}\right)^{-3+\alpha}$$

• Einasto Profile  $\rho_{\chi} = \rho_0 \exp \left[ -\frac{2}{\alpha} \left( \frac{r^{\alpha}}{r_s} \right) - 1 \right]$ 

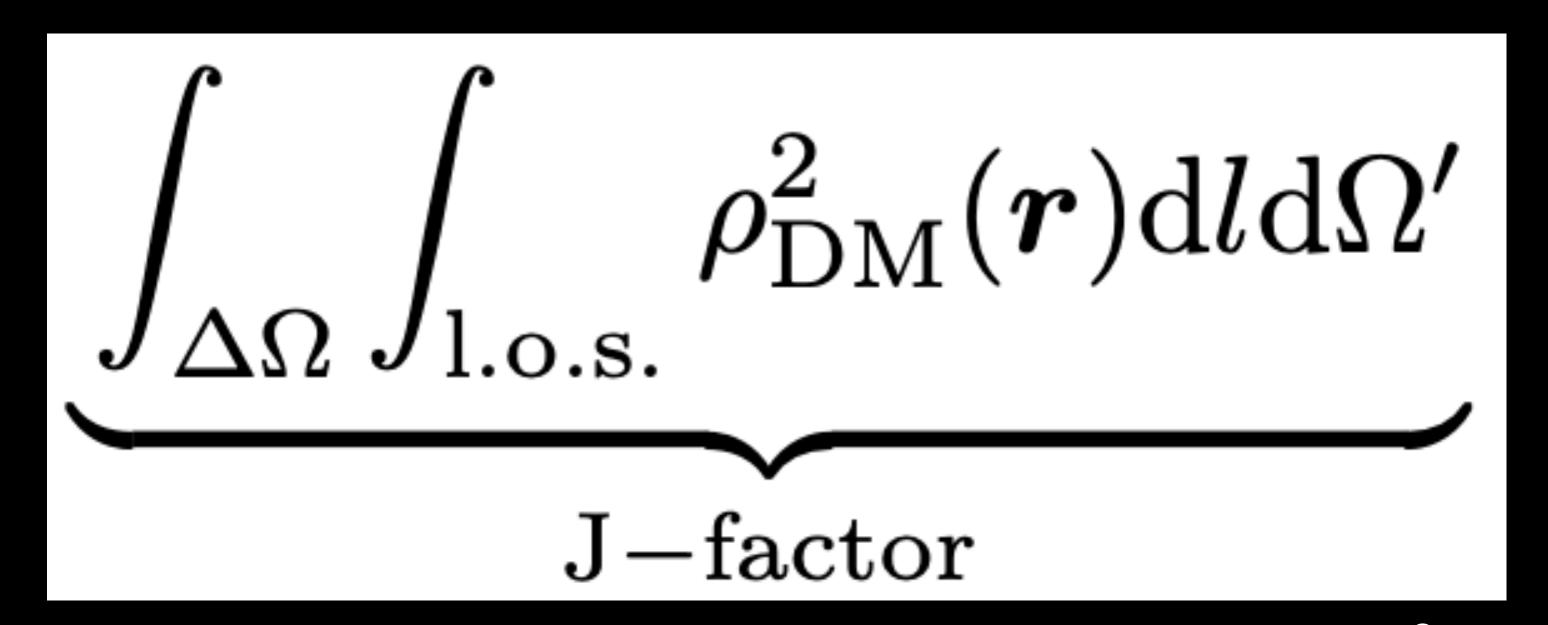


## **Local Observations**



## J-Factors

O Remember, we cannot constrain distances astrophysically.

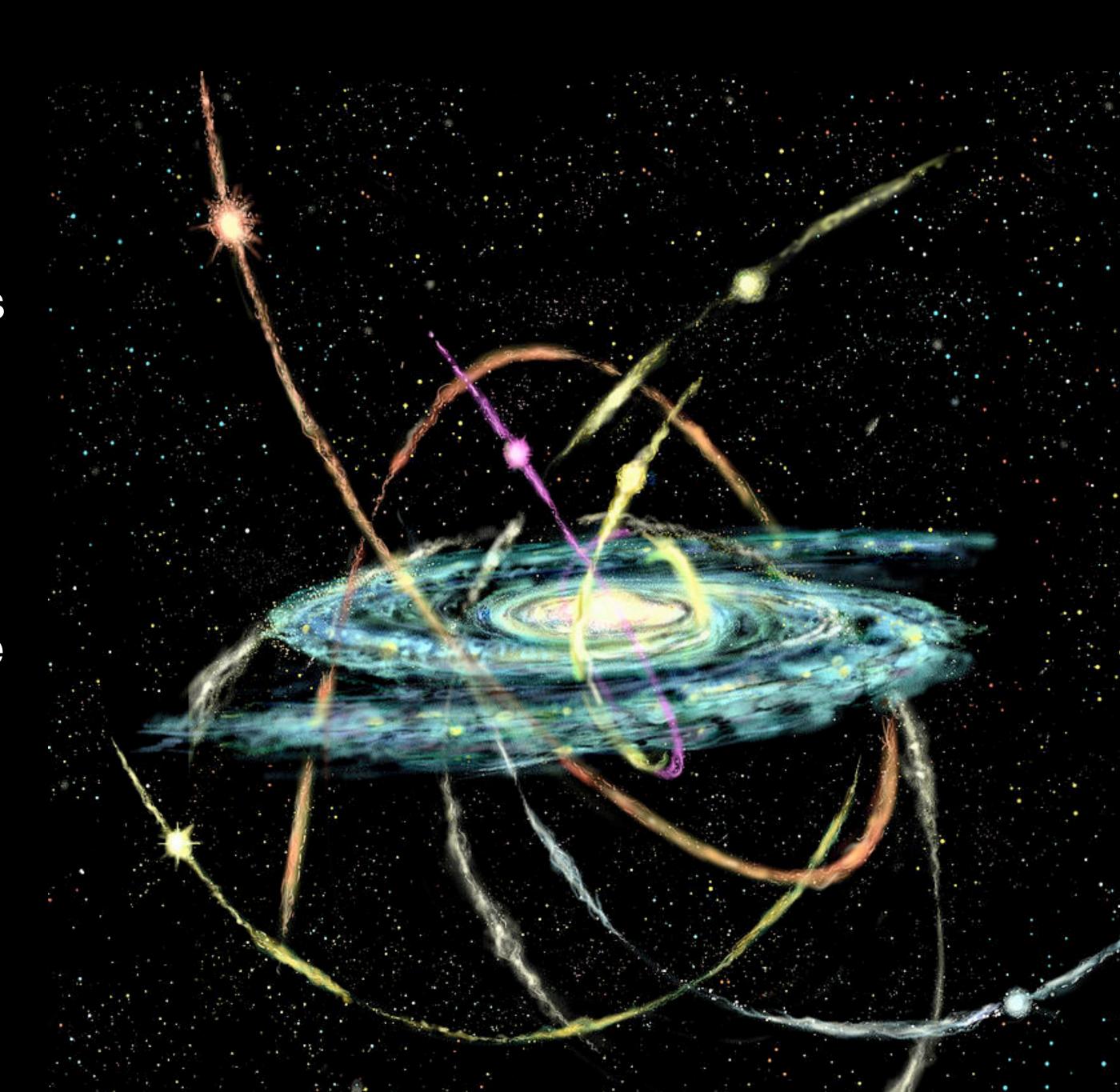


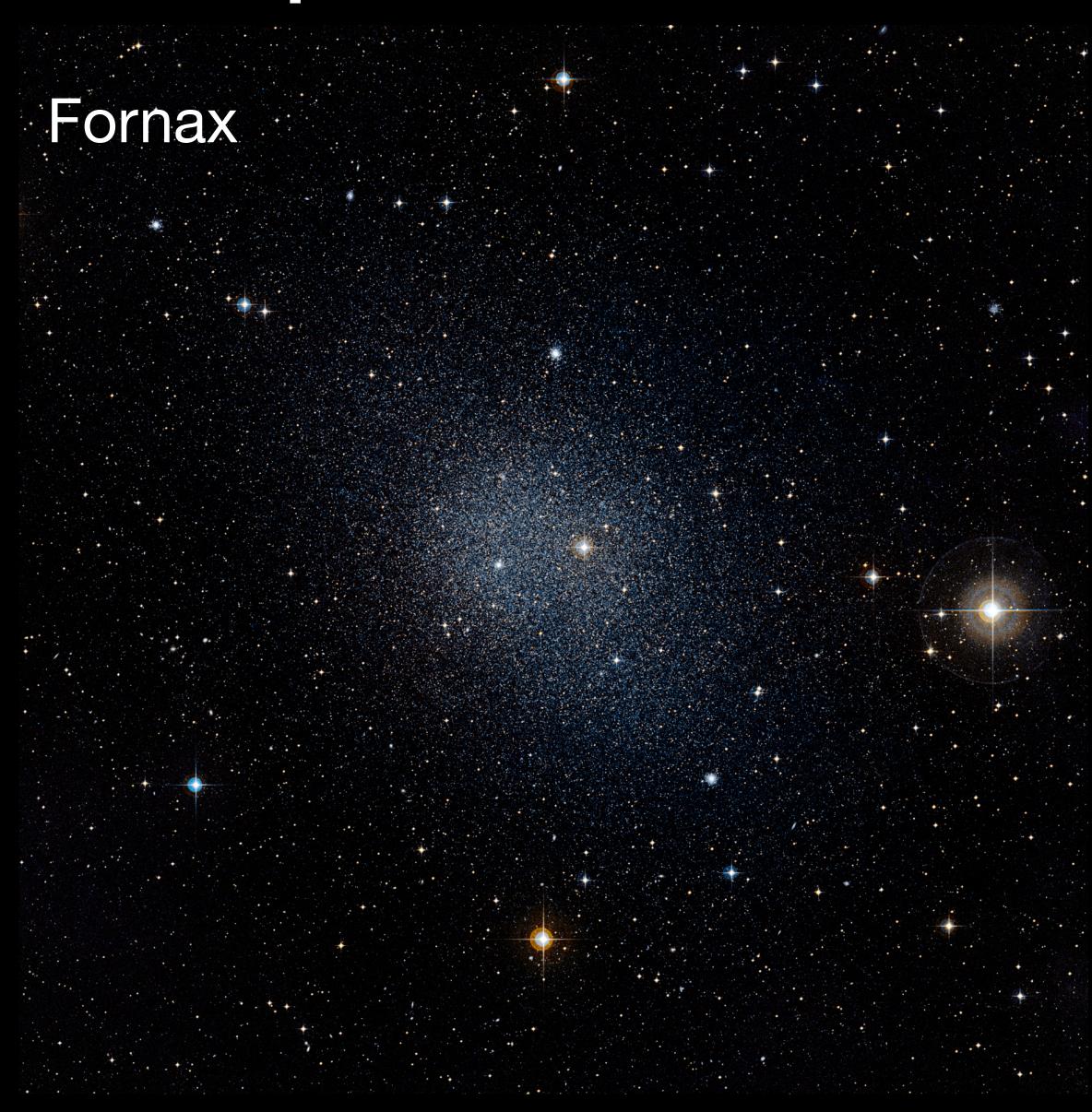
- O For decaying dark matter, this will depend on  $\rho$  and not  $\rho^2$  (and is normally called the D-factor for Decay).
- Large uncertainties: The majority of the signal depends on regions near the centers of structures, where the dark matter density is maximally uncertain.

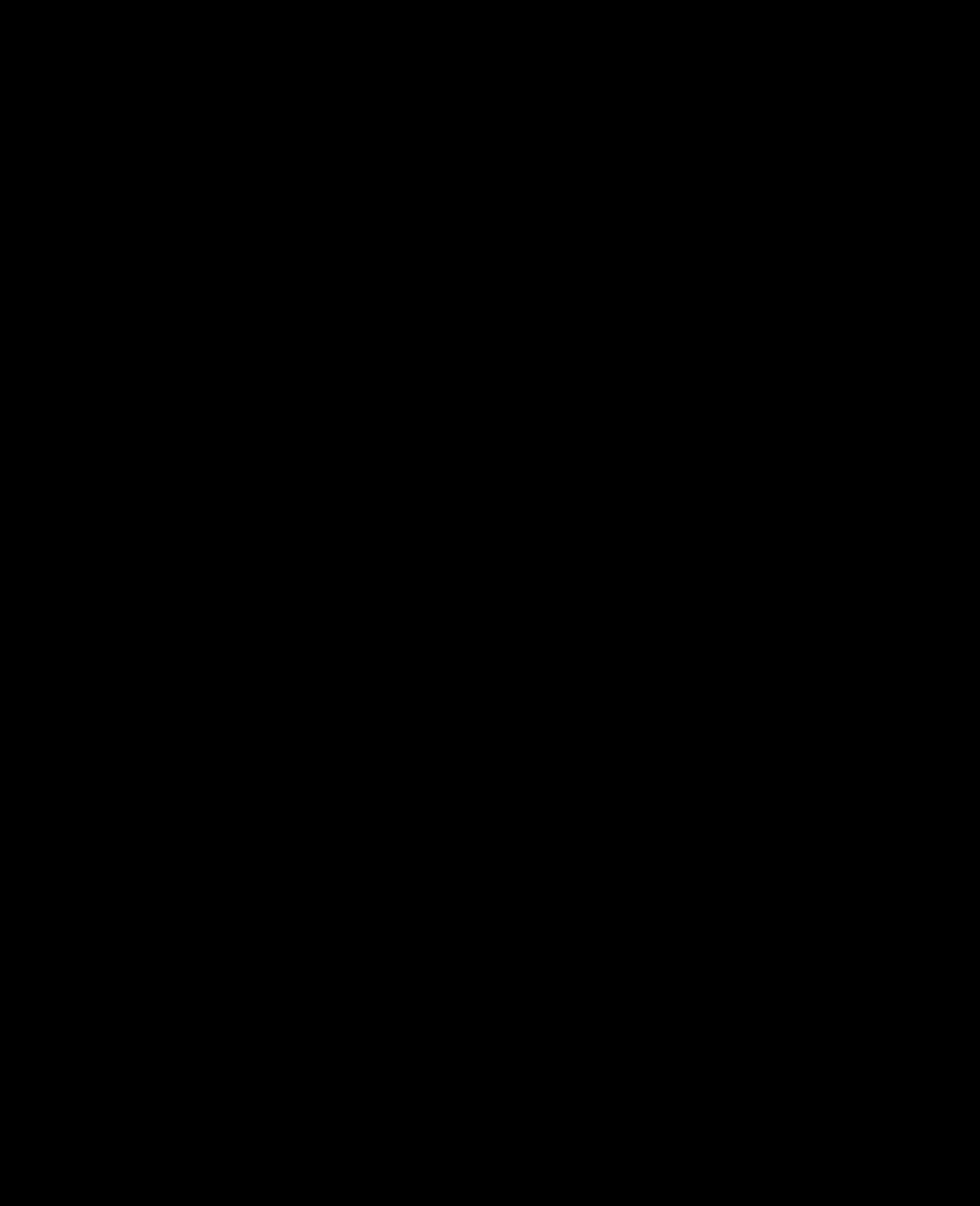
## **Dwarf Spheroidal Galaxies**

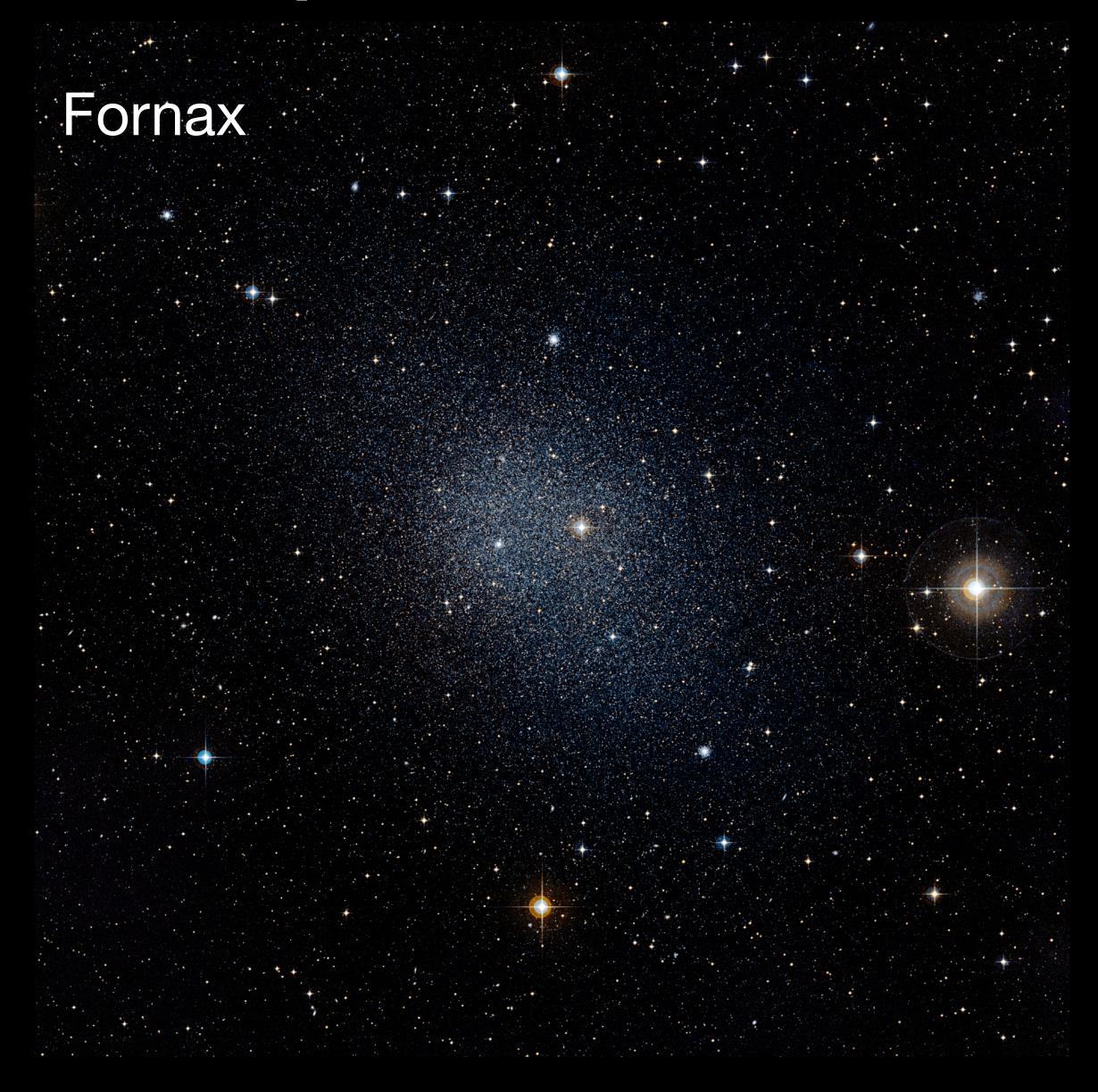
- Dark Matter dominated structures.
  - Stellar rotation curves almost always dark matter dominated.

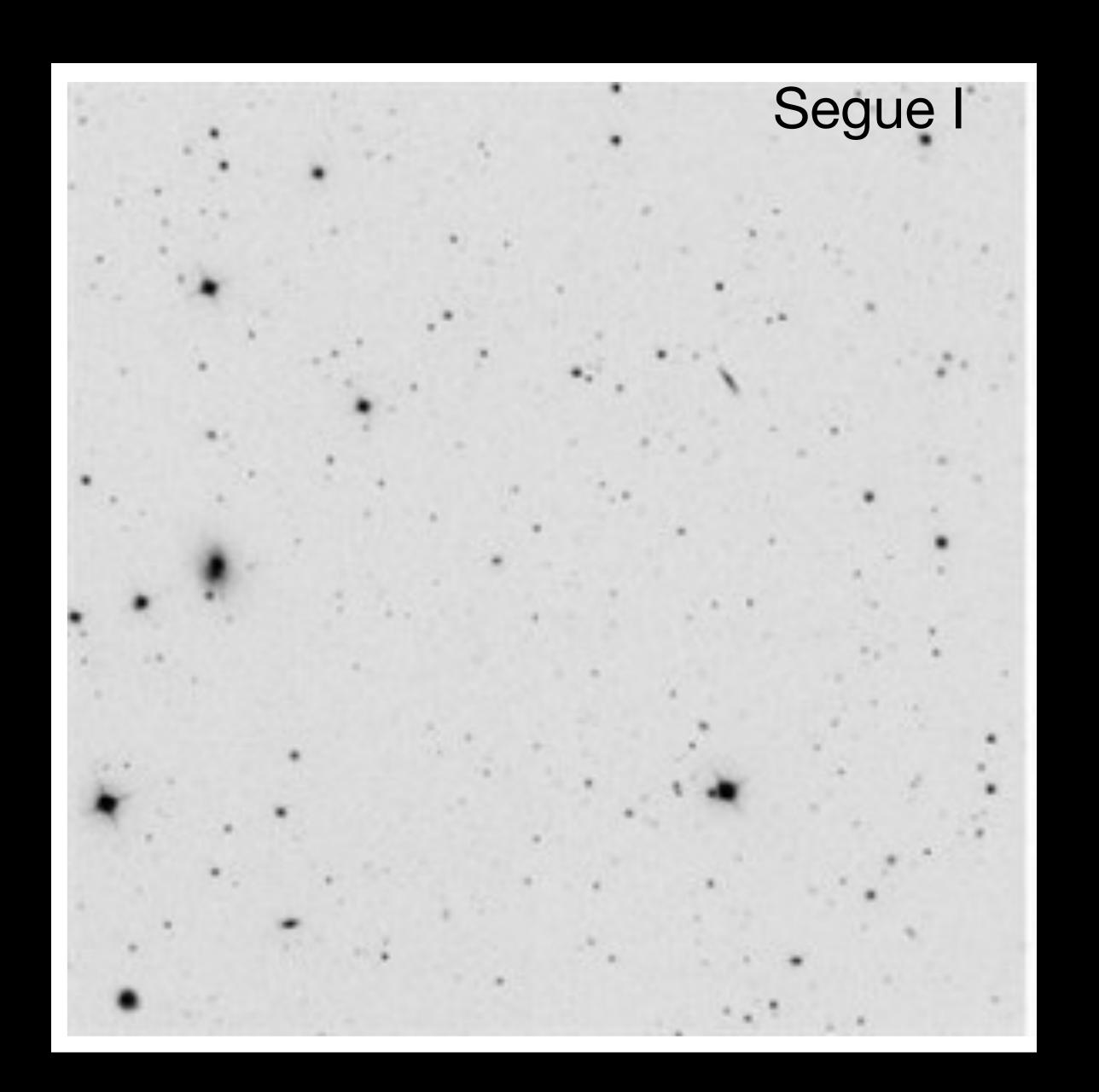
 Single stellar population with negligible baryonic feedback, SNe, etc.

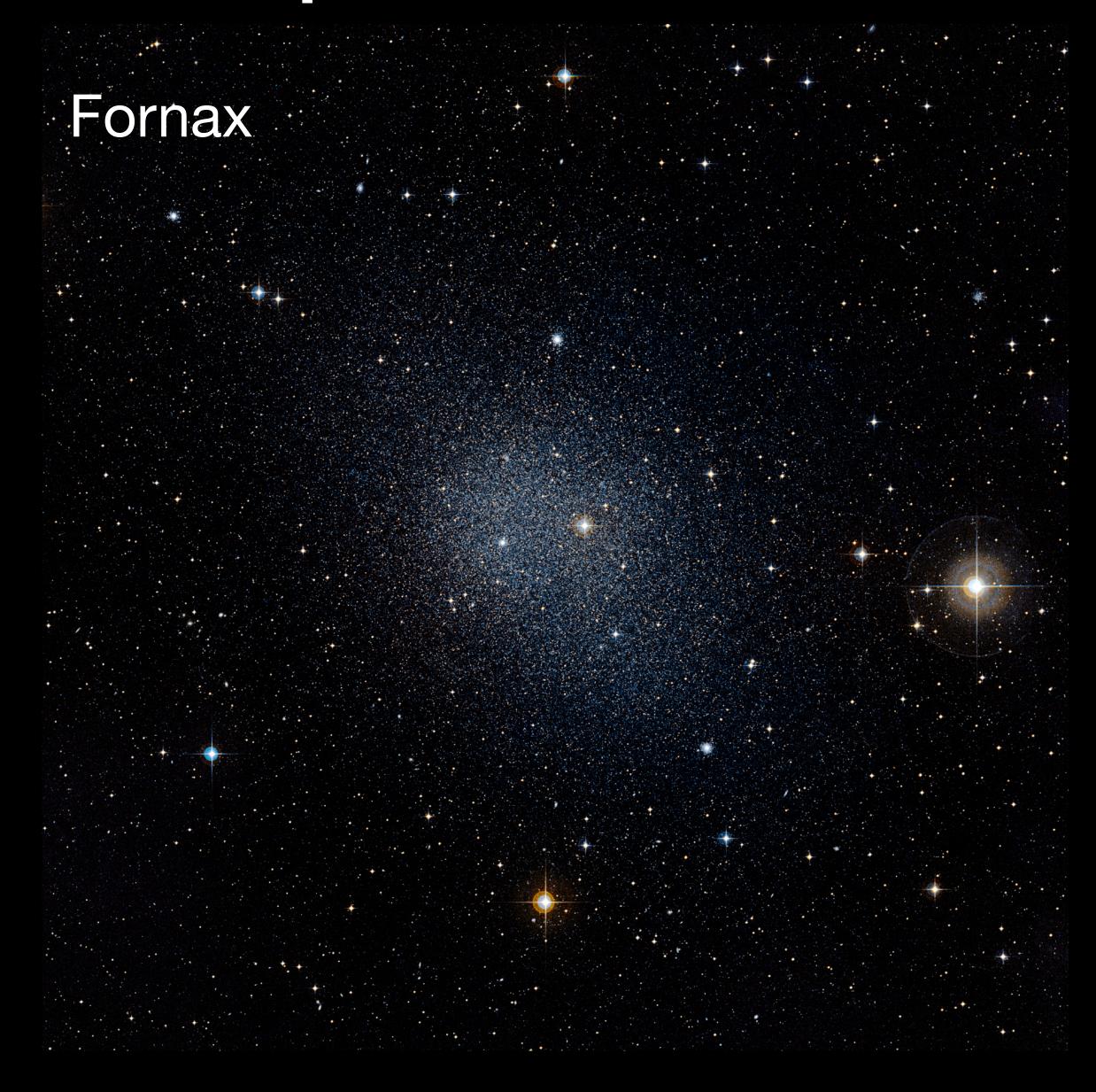


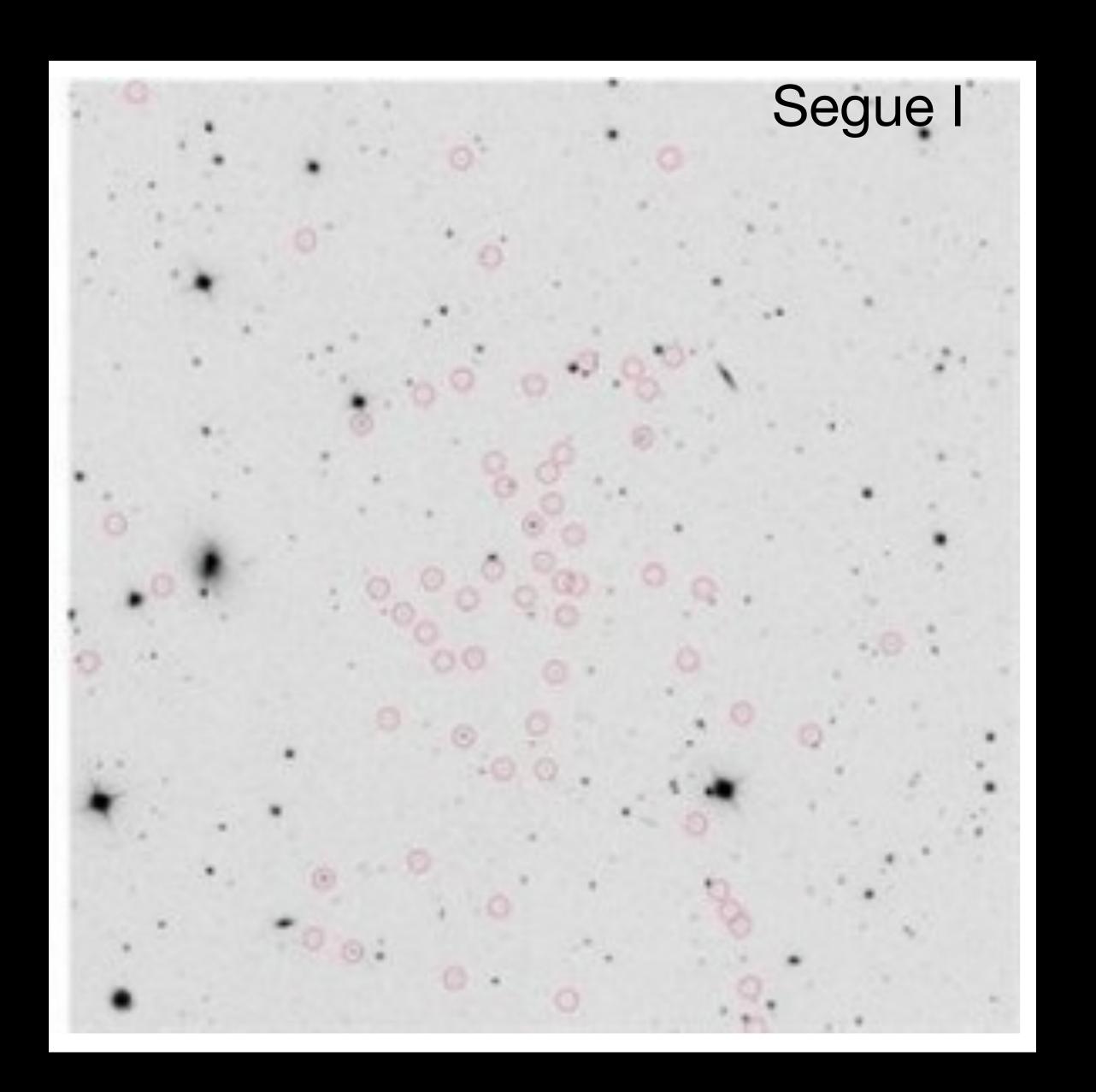


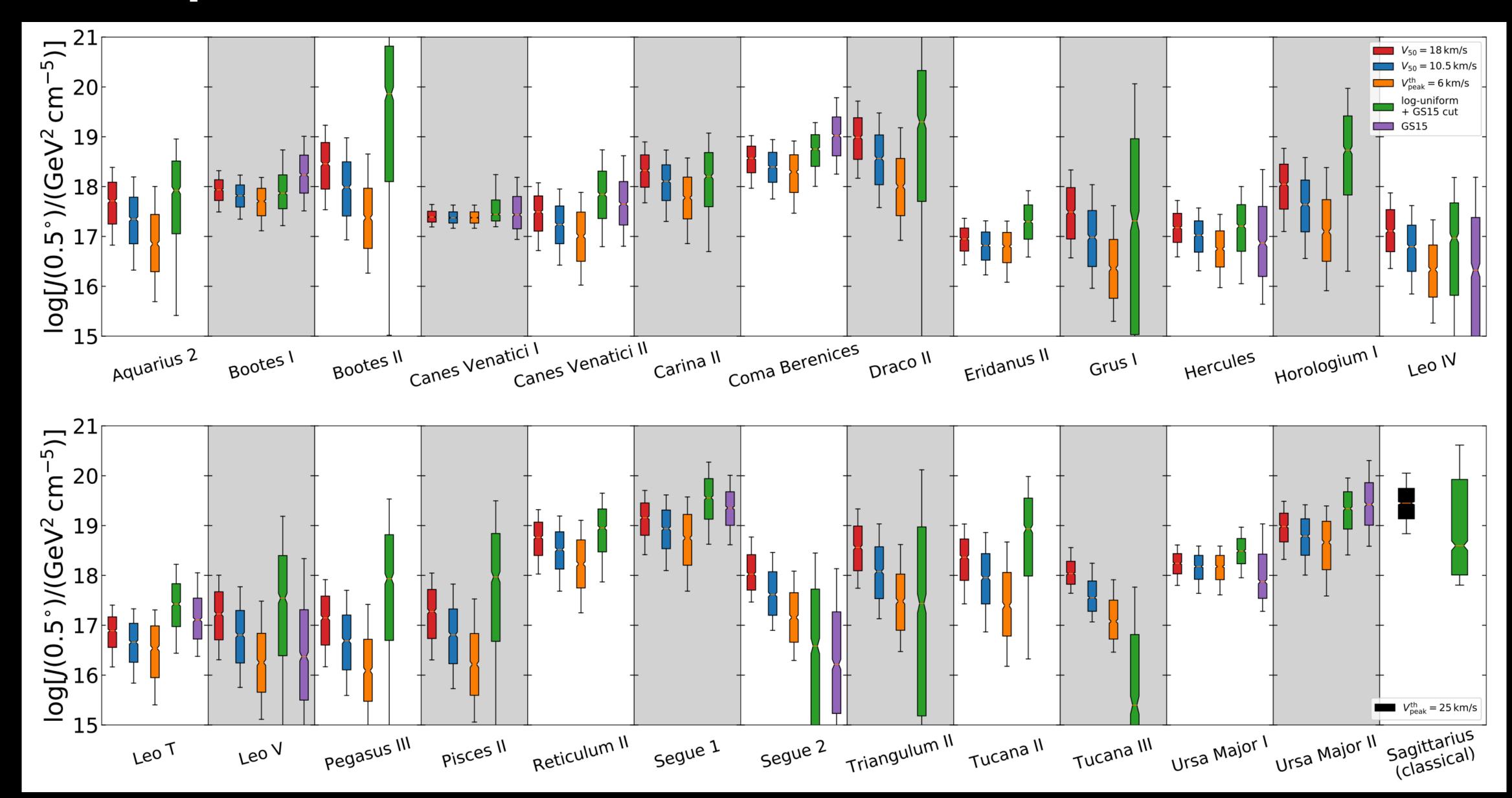






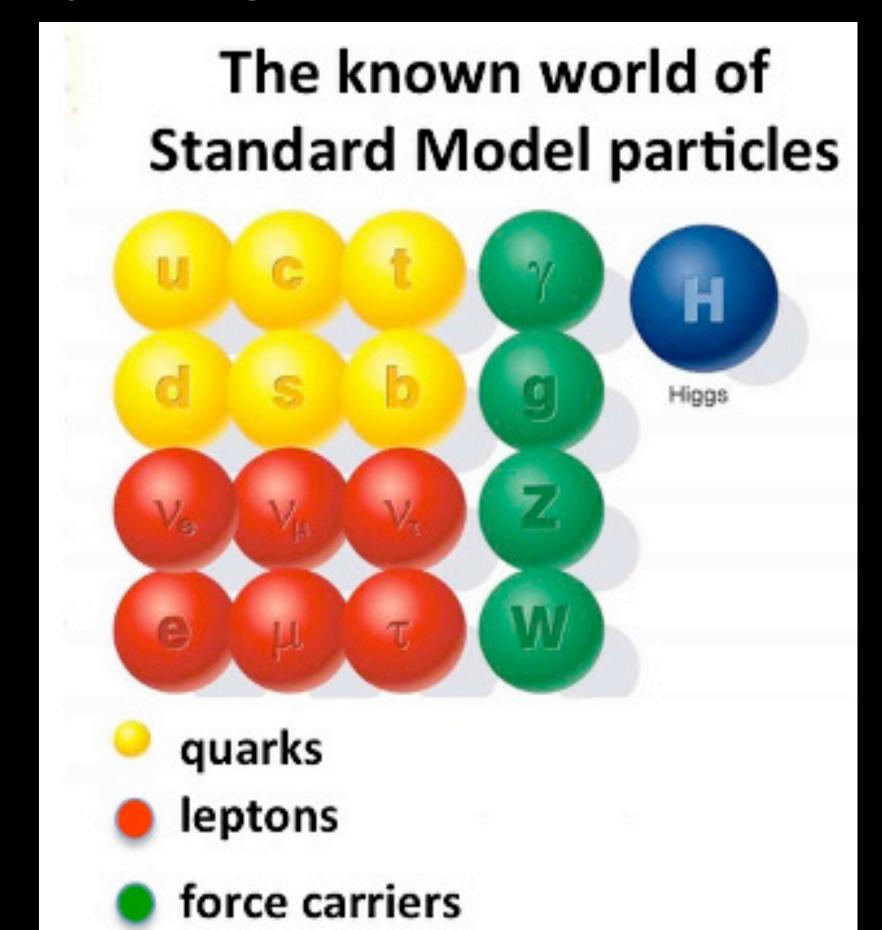


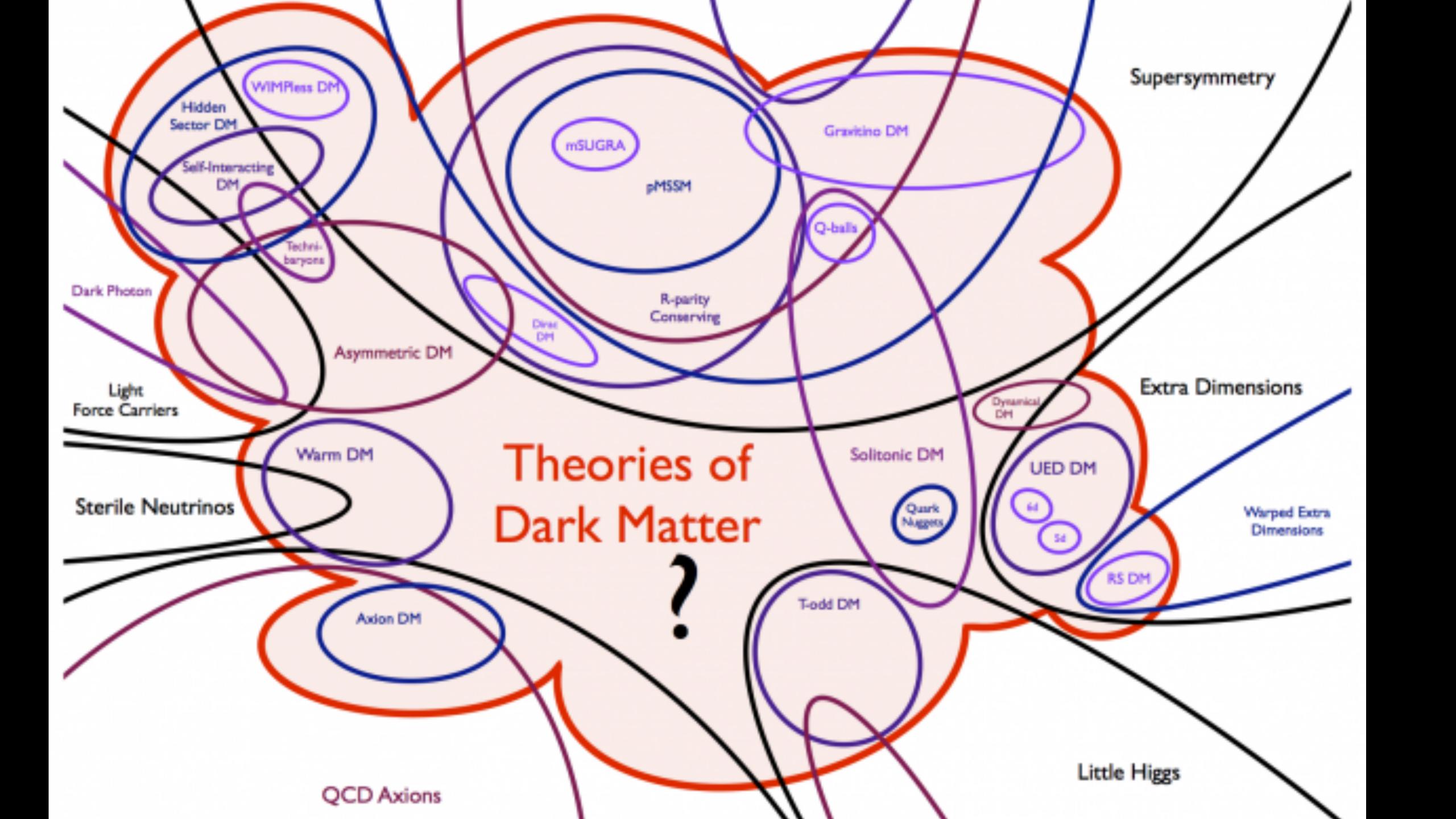




# Properties of the Dark Matter Particle

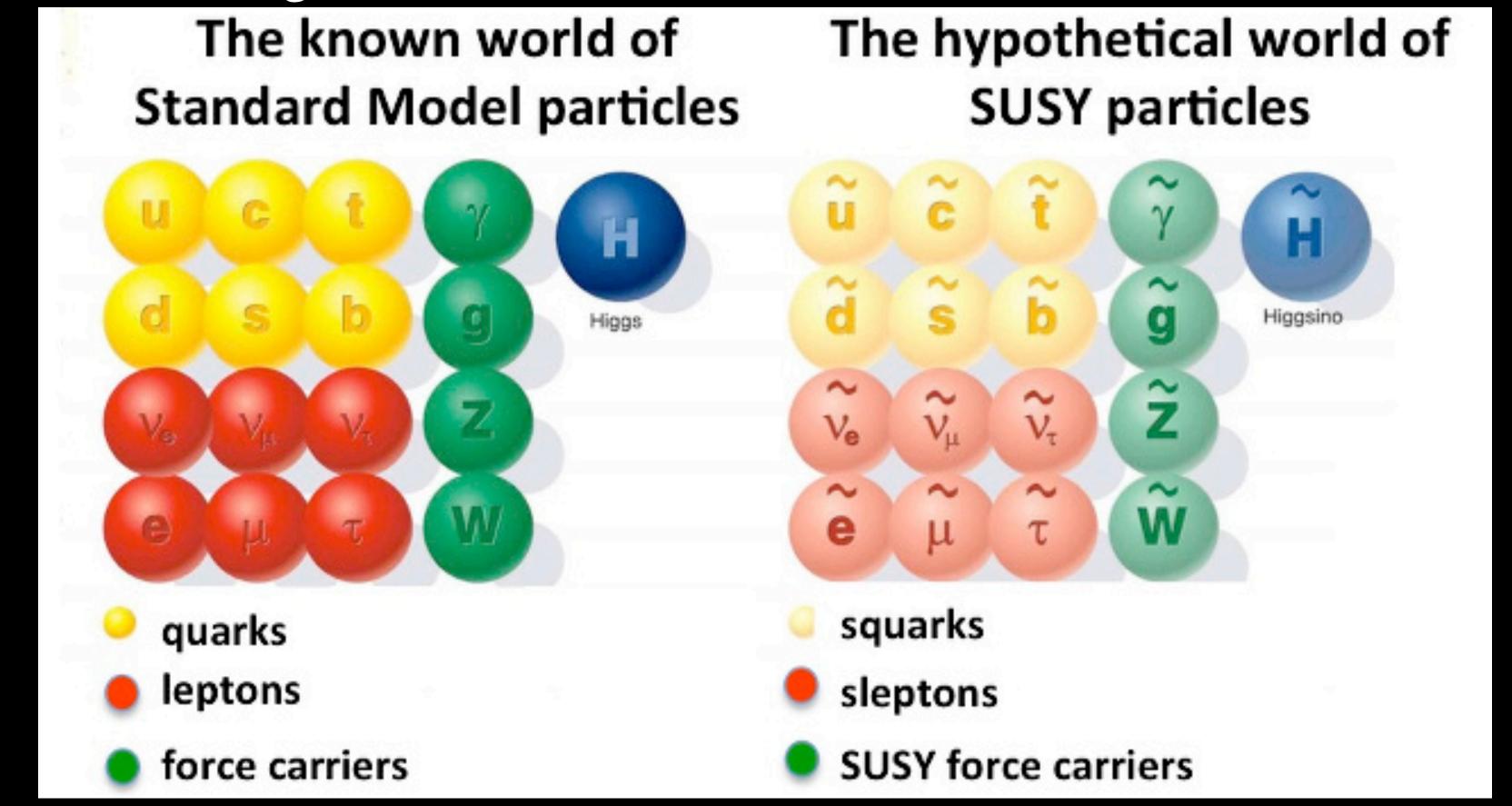
- o Three Key Particle Properties of Dark Matter
  - Stable
  - Does not Interact with Light (Dark)
  - Does not move relativistically throughout the Universe (Cold)

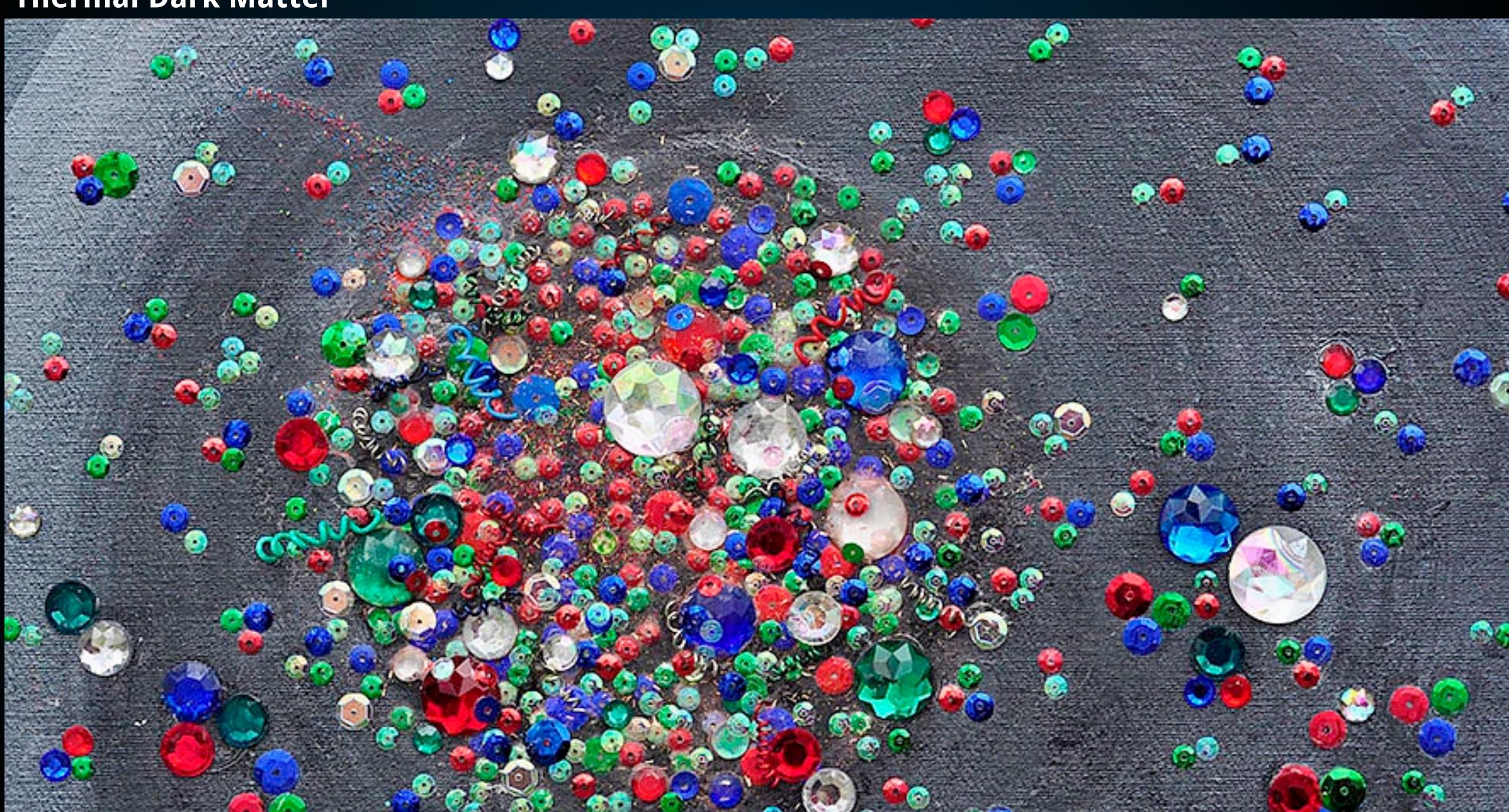




# WIMP (and WIMP-like) Dark Matter

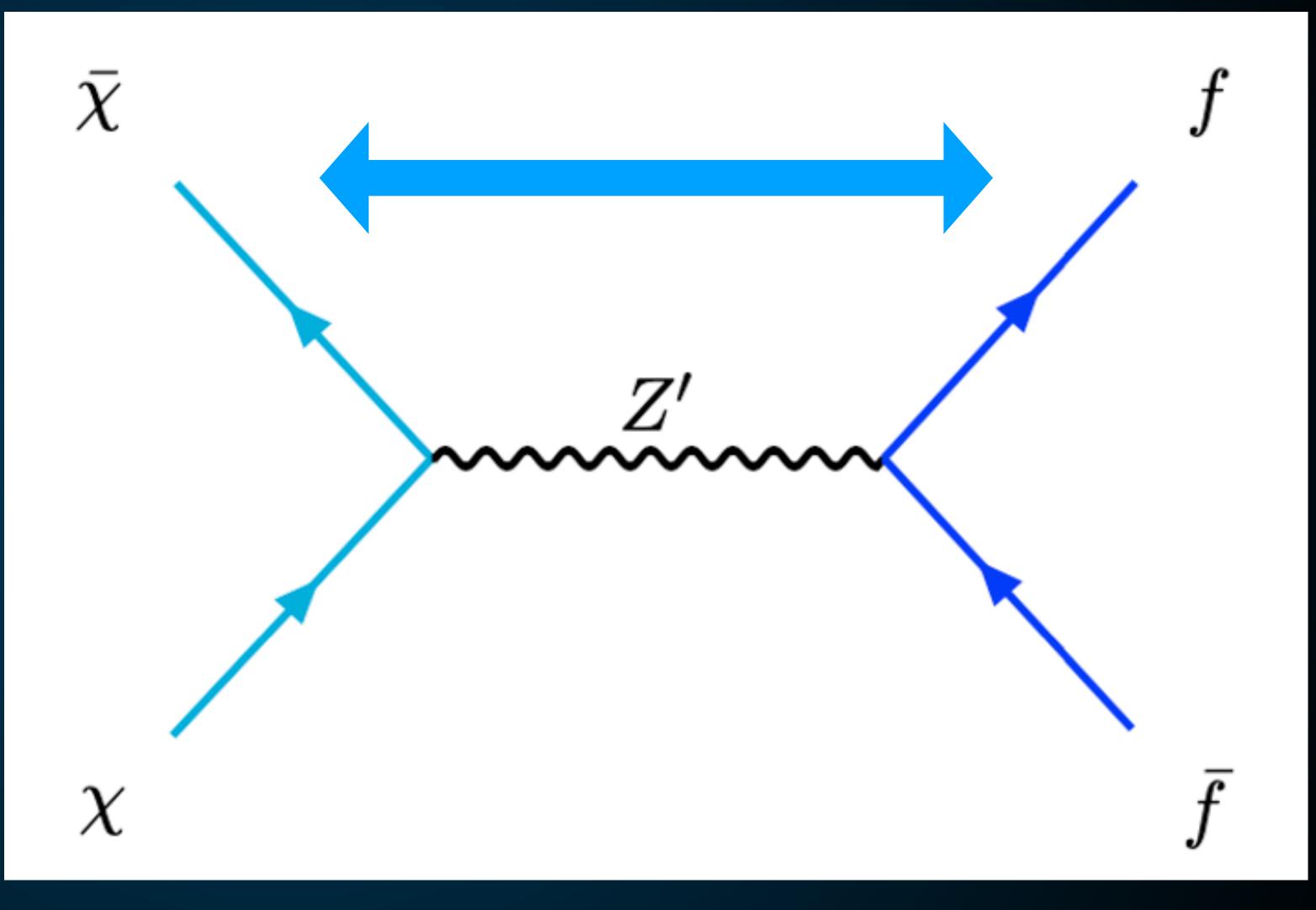
- o What about the lightest supersymmetric particle?
  - o R-parities can prevent decay into standard model particles (stable)
  - Likely Heavy (cold)
  - Would not interact with light if it is neutral (dark)





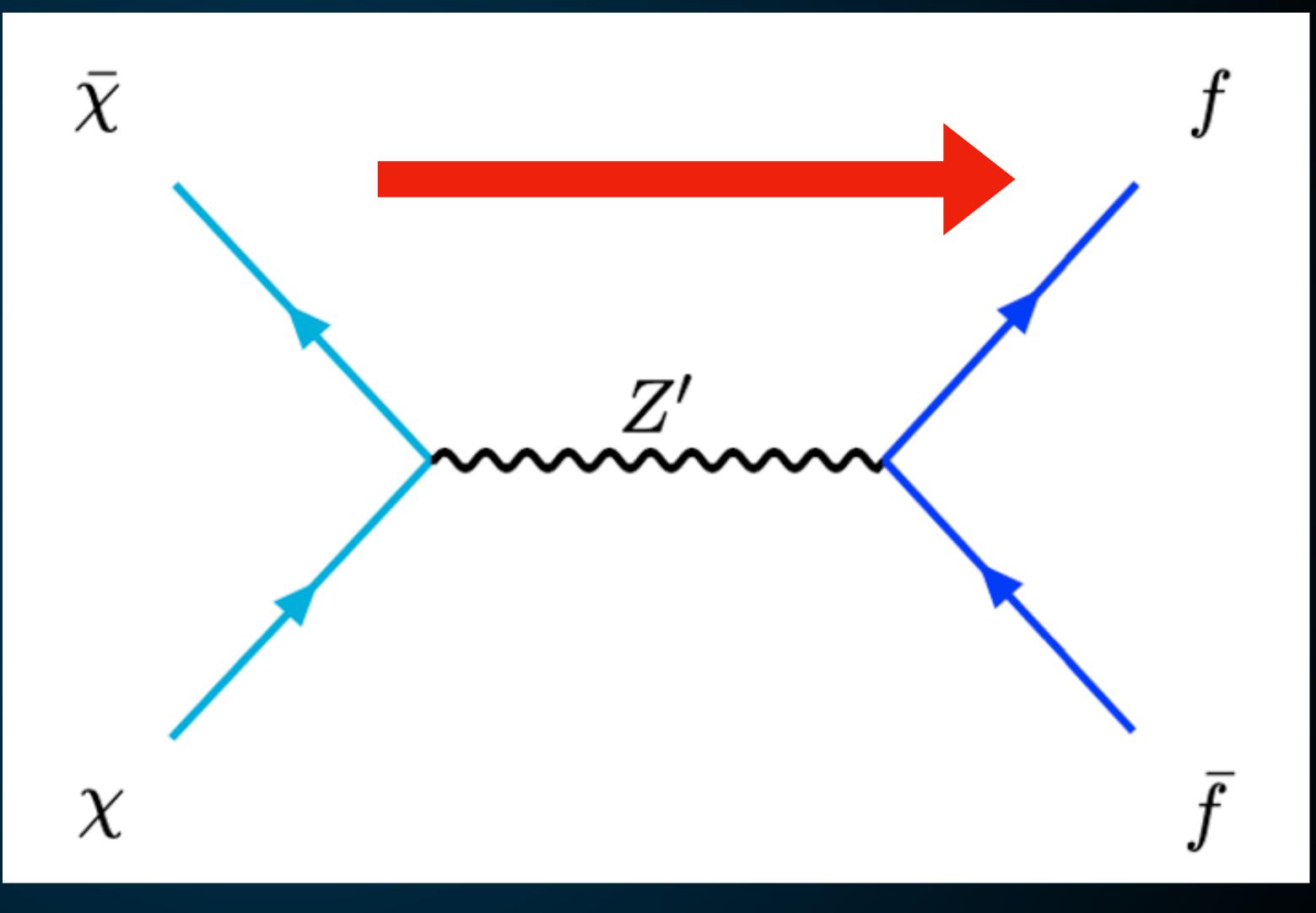
## **Thermal Dark Matter**





## **Thermal Dark Matter**





# Thermal Freeze-Out/WIMP Miracle

- o Three periods:
  - o Thermal equilibrium with baryons
  - Annihilating away
  - Hubble expansion wins

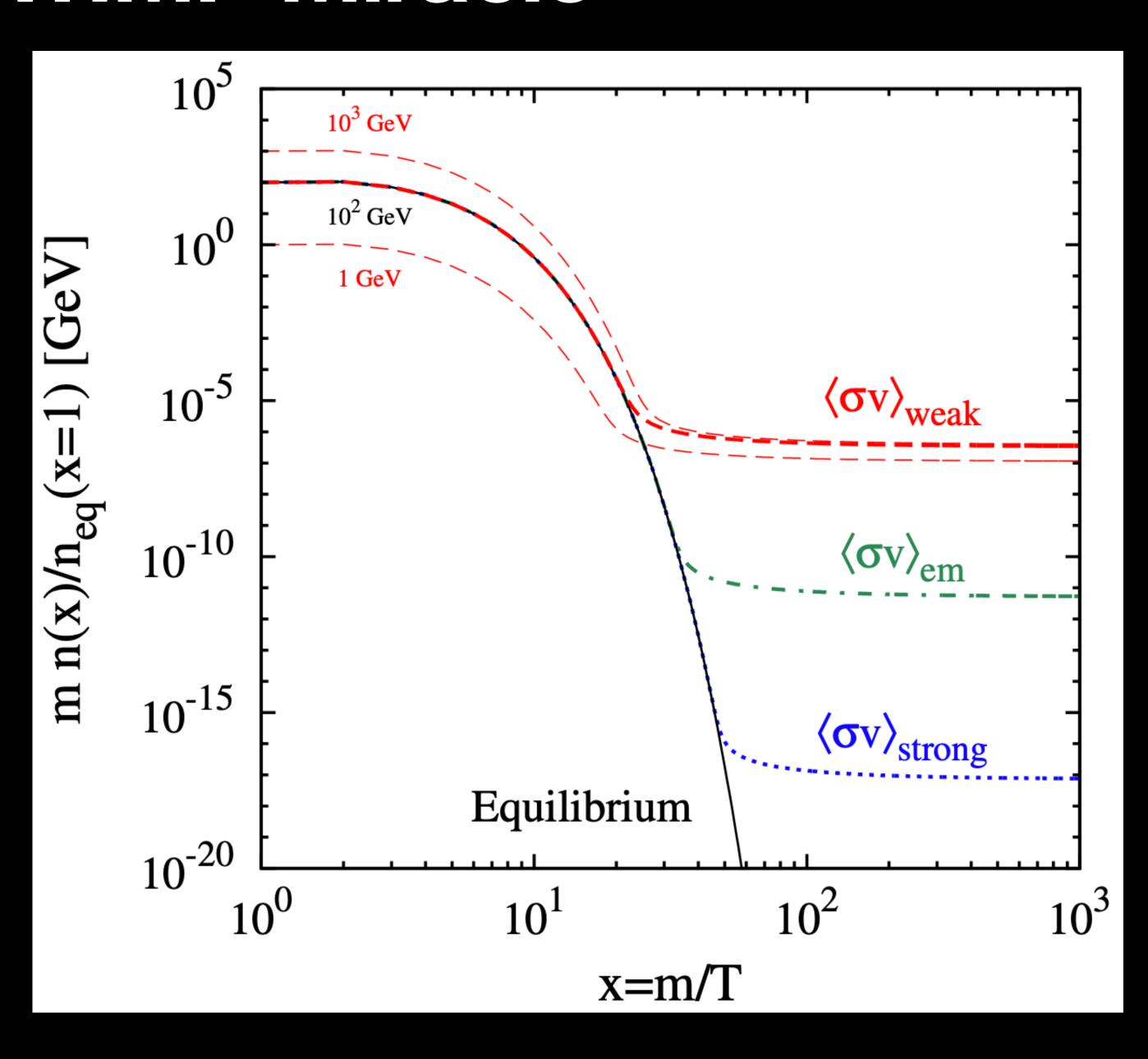
$$\frac{dn}{dt} + 3Hn = -\frac{n^2}{2} \langle \sigma v_{\rm rel} \rangle \times 2$$

$$\frac{dn}{dt} + 3Hn = (n_{\rm eq}^2 - n^2)\langle \sigma v_{\rm rel} \rangle$$

o For a cold particle during freeze out

$$n \sim (m_{\chi}T)^{3/2} \exp\left(-\frac{m_{\chi}}{T}\right)$$

If interaction has weak force strength current DM density = Present Density!



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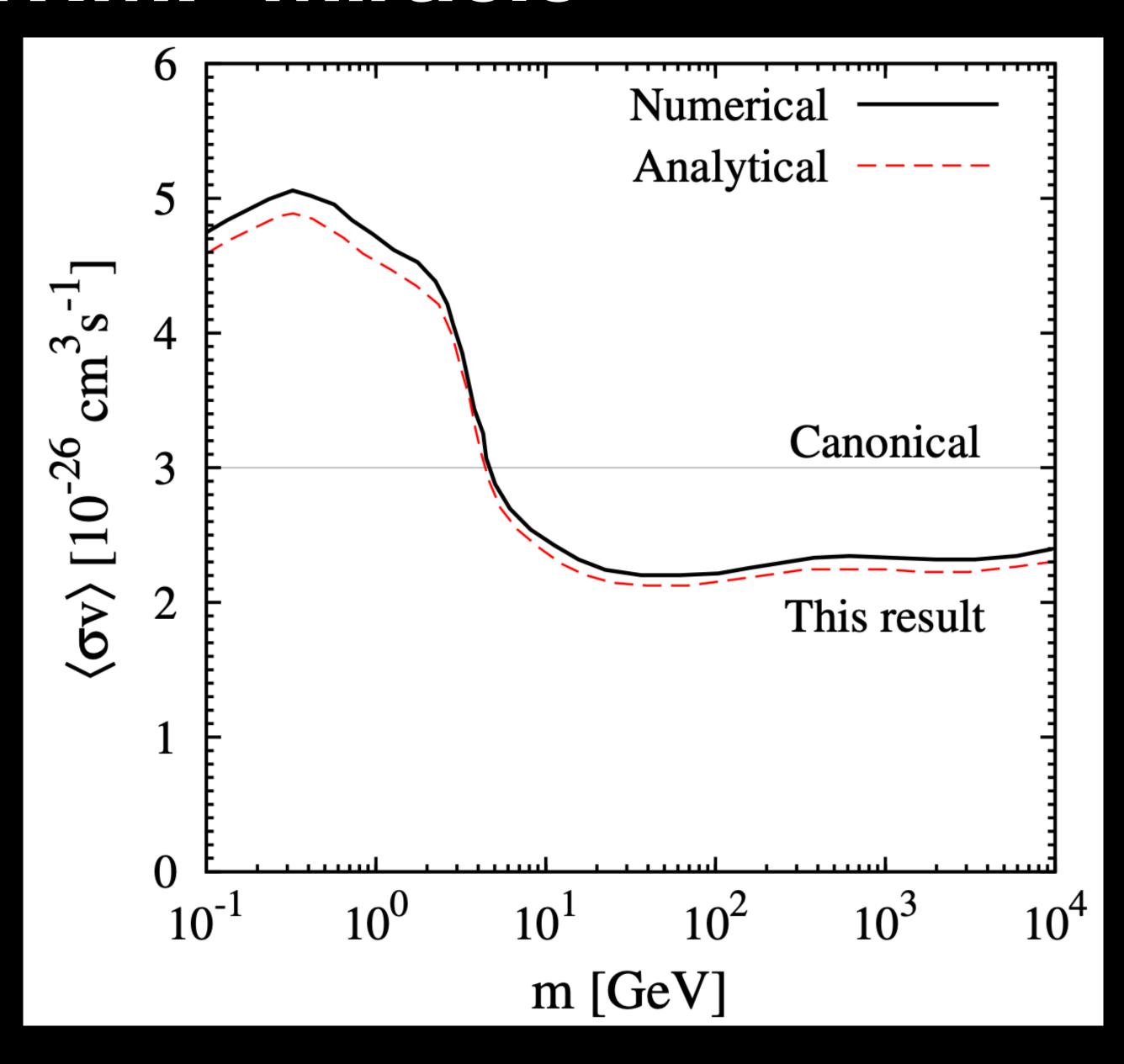
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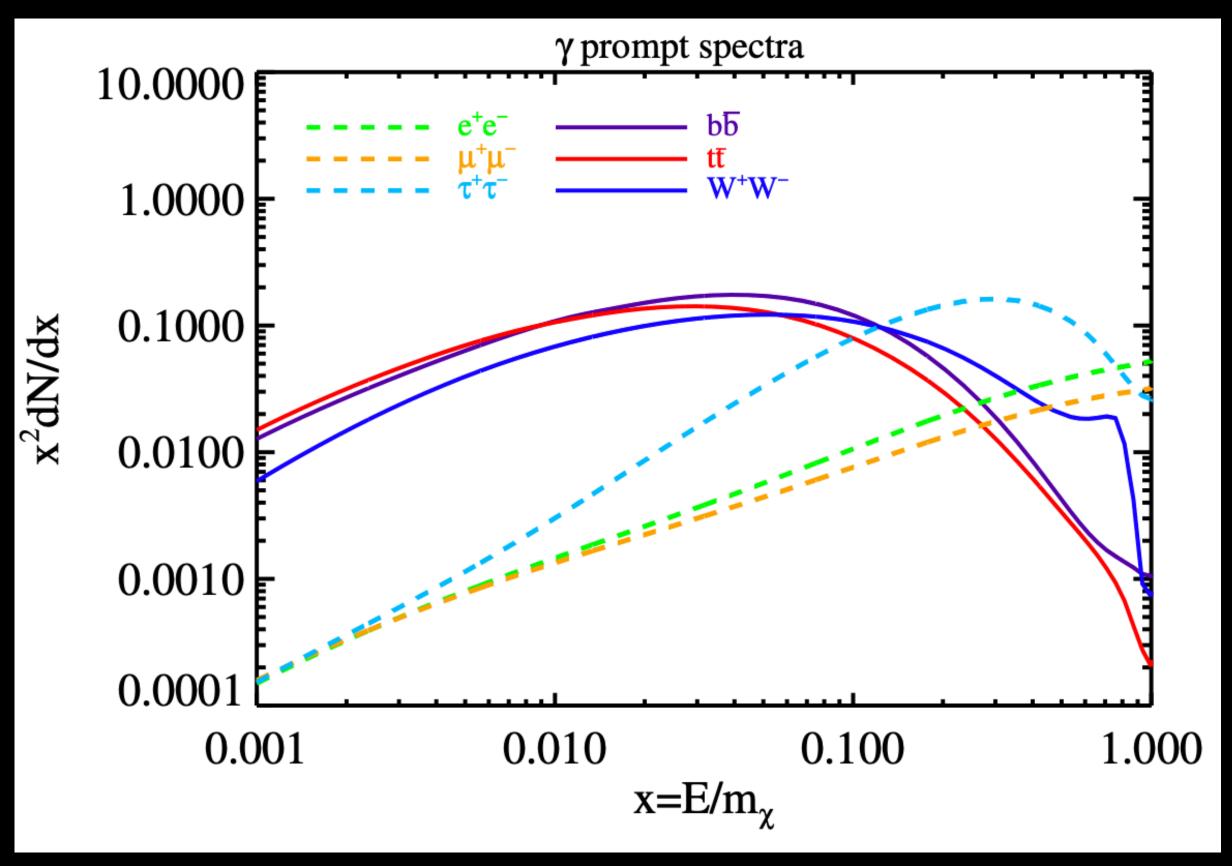
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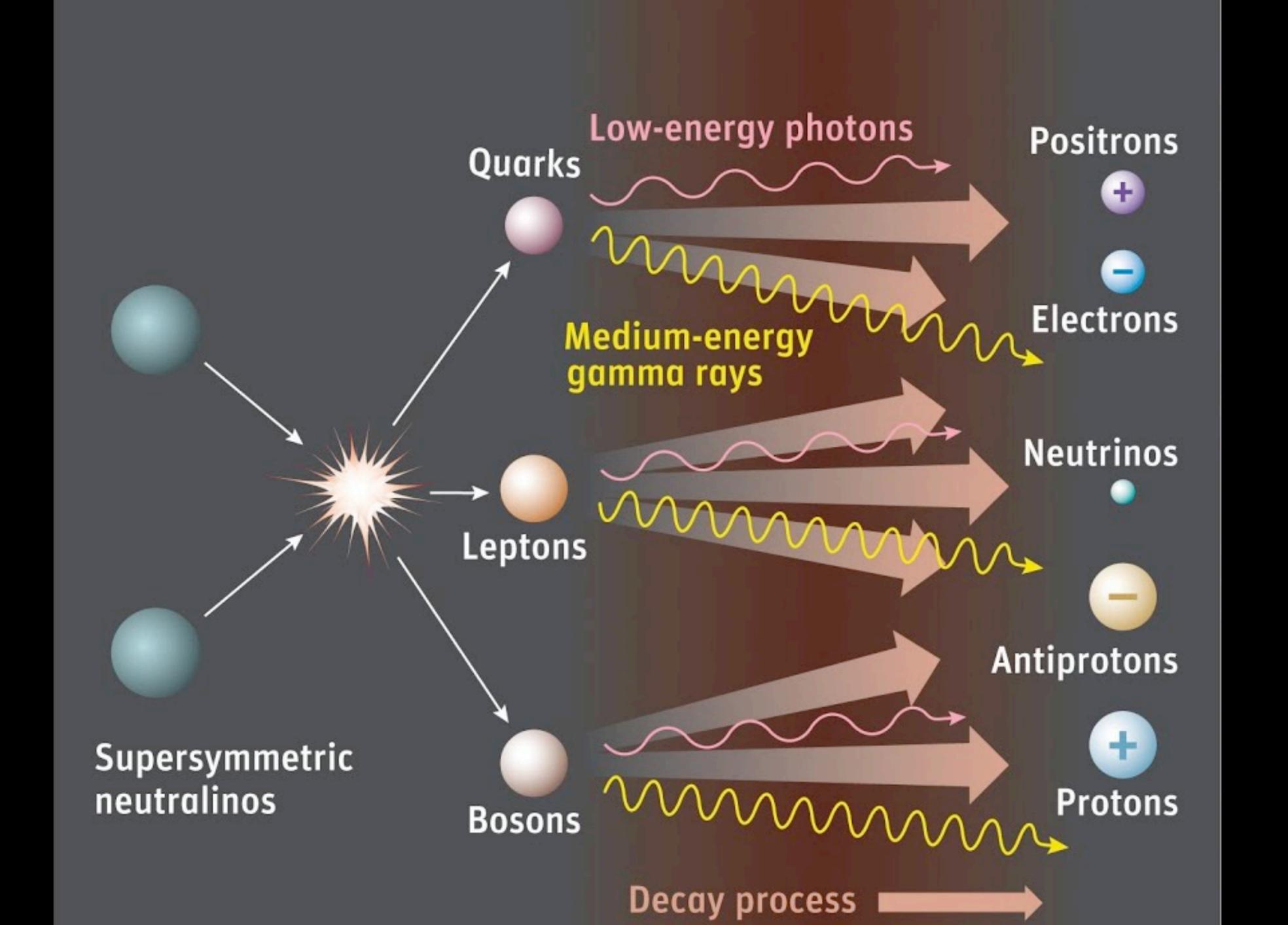


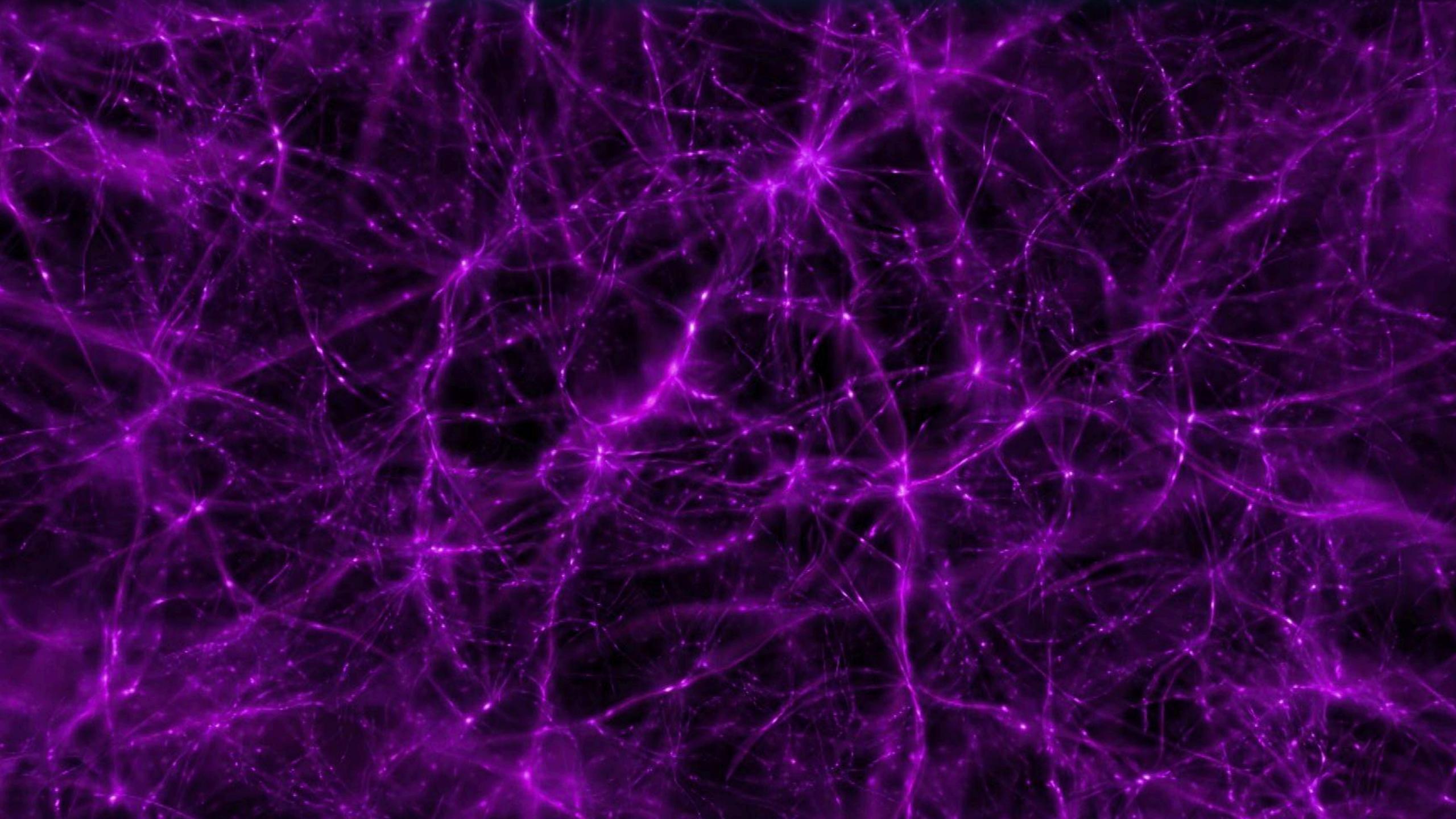
## Thermal Freeze-Out/WIMP Miracle

- Owhat does the dark matter annihilate into?
- o Any final states are possible:
  - Higgs-motivated
    - To heaviest quarks, bottom/top pairs
  - Leptophilic
    - Tau/muon/electron pairs
  - o Bosons
    - W/Z or combinations
- After fist baryonic particles are created, shower is formed following standard model physics to create gamma-rays, electrons/positrons/protons/neutrons

Gaskins (1604.00014)







## Calculation of Expected Flux

- Now we have all of the ingredients
  - o J-factor
  - o Annihilation Rate
  - Annihilation Final State/Spectrum

$$\phi_s(\Delta\Omega) = \underbrace{\frac{1}{4\pi} \frac{\langle \sigma v \rangle}{2m_{\rm DM}^2} \int_{E_{\rm min}}^{E_{\rm max}} \frac{{\rm d}N_{\gamma}}{{\rm d}E_{\gamma}} {\rm d}E_{\gamma}}_{\rm particle~physics} \times \underbrace{\int_{\Delta\Omega} \int_{\rm l.o.s.} \rho_{\rm DM}^2(\boldsymbol{r}) {\rm d}l {\rm d}\Omega'}_{\rm J-factor}$$

$$\langle \sigma v \rangle \approx 2 \times 10^{-26} \text{ cm}^3 \text{s}^{-1}$$

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$$\rho_0 = \frac{2 \times 10^{11} M_{\odot}}{\frac{4}{3} \pi \left(35 \text{ kpc}\right)^3} = \frac{2 \times 10^{68} \text{ GeV}}{3 \times 10^{69} \text{ cm}^3} \sim 0.1 \text{ GeV cm}^{-3}$$

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Pick:

$$m_{\chi} = 100 \text{ GeV}$$
  $\frac{E}{\text{ann}} = 200 \text{ GeV}$ 

$$<\sigma v>\approx 2\times 10^{-26}$$
 cm<sup>3</sup>s<sup>-1</sup>

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$$\Phi = \frac{E}{\text{ann }} \frac{\rho^2}{2m_{\chi}} < \sigma v > V$$

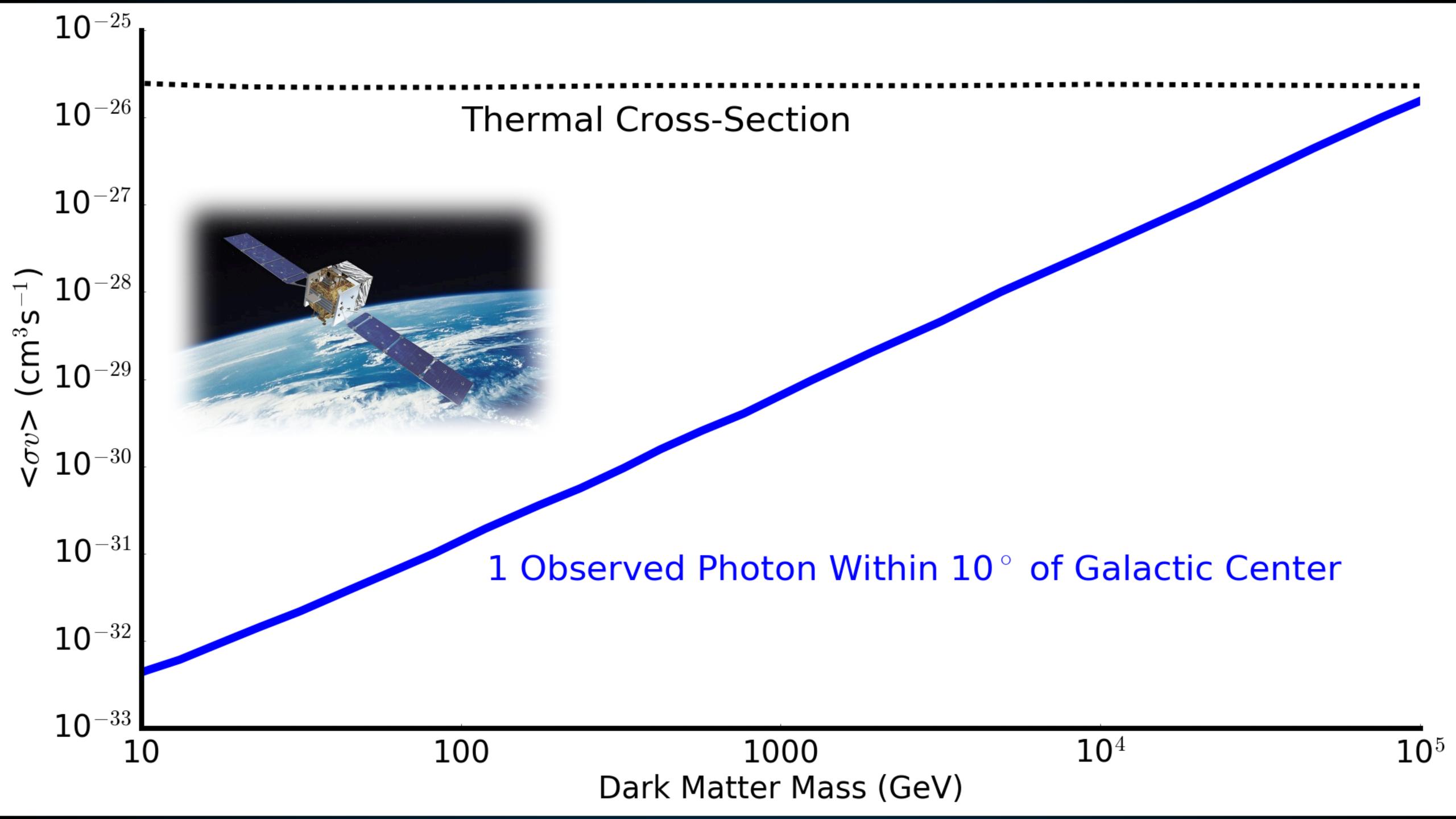
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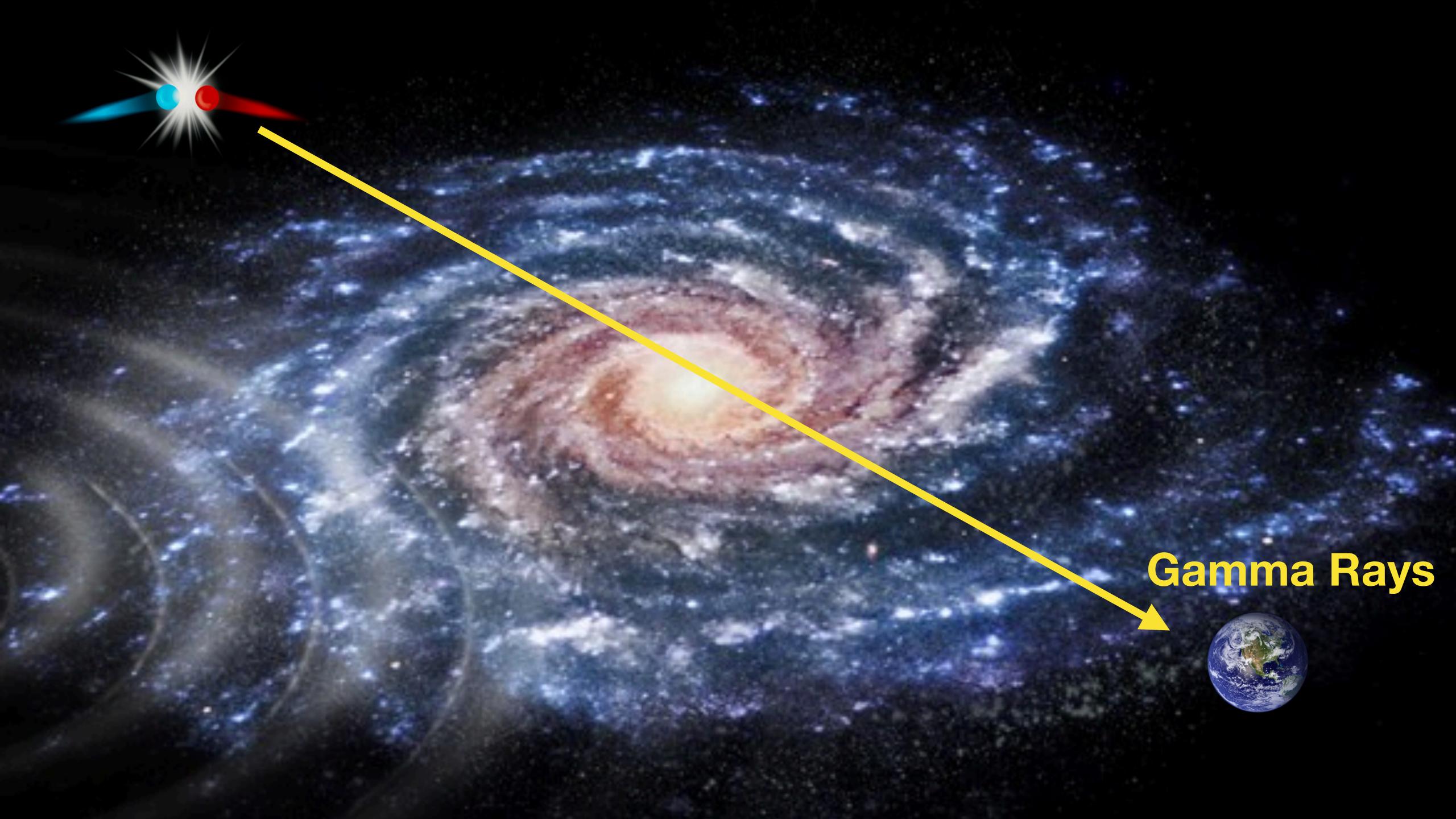
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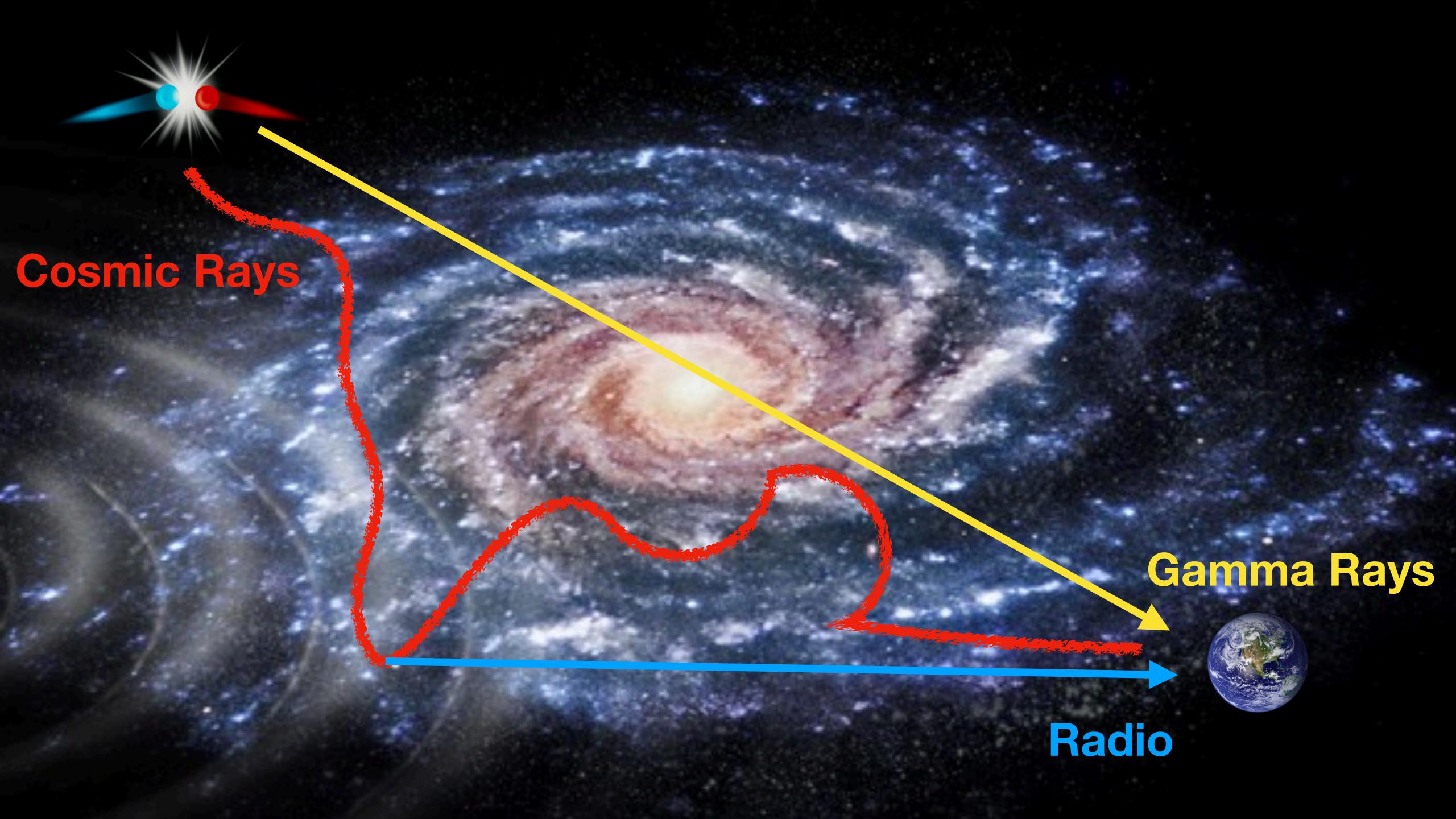
$$\Phi = \frac{E}{\text{ann }} \frac{\rho^2}{2m_{\chi}} < \sigma v > V = 200 \text{ GeV} \frac{0.1 \text{ GeV cm}^{-3}}{200 \text{ GeV}} \left(10^{-26} \text{ cm}^3 \text{s}^{-1}\right) 3 \times 10^{69} \text{ cm}^3 = 3 \times 10^{39} \text{ GeV s}^{-1}$$











### Gamma-Rays

Dwarf Spheroidal Galaxies
Galaxy Clusters
Milky Way Subhalos
Galactic Diffuse
Sun
Jupiter
Nearby Stars
Galactic Center Stars
Andromeda
Little Galaxies
Isotropic Gamma-Ray Background
Anisotropy Searches
Cusps

Morphology

## Cosmic-Rays

Positrons
Electron + Positron Spectrum
Antiprotons
Antineutrons
Antihelium
Cosmological Lithium Problem

#### **Antimatter**

### Neutrinos

The Sun
Direct Annihilation to Neutrinos
Hidden Sources
The Earth
Ultra-Heavy Dark Matter

Messengers

## Low-Energy

**Galactic Center Synchrotron Dwarf Galaxy Synchrotron Galaxy Cluster Synchrotron Diffuse Synchrotron** Sun **Jupiter Isotropic Background** X-ray background from Clusters **Anisotropy Searches Stellar Evolution Pulsar Evolution** Planetary Heating **Thermal Scattering Cosmic Microwave Background CMB** Absorption

**Targets** 

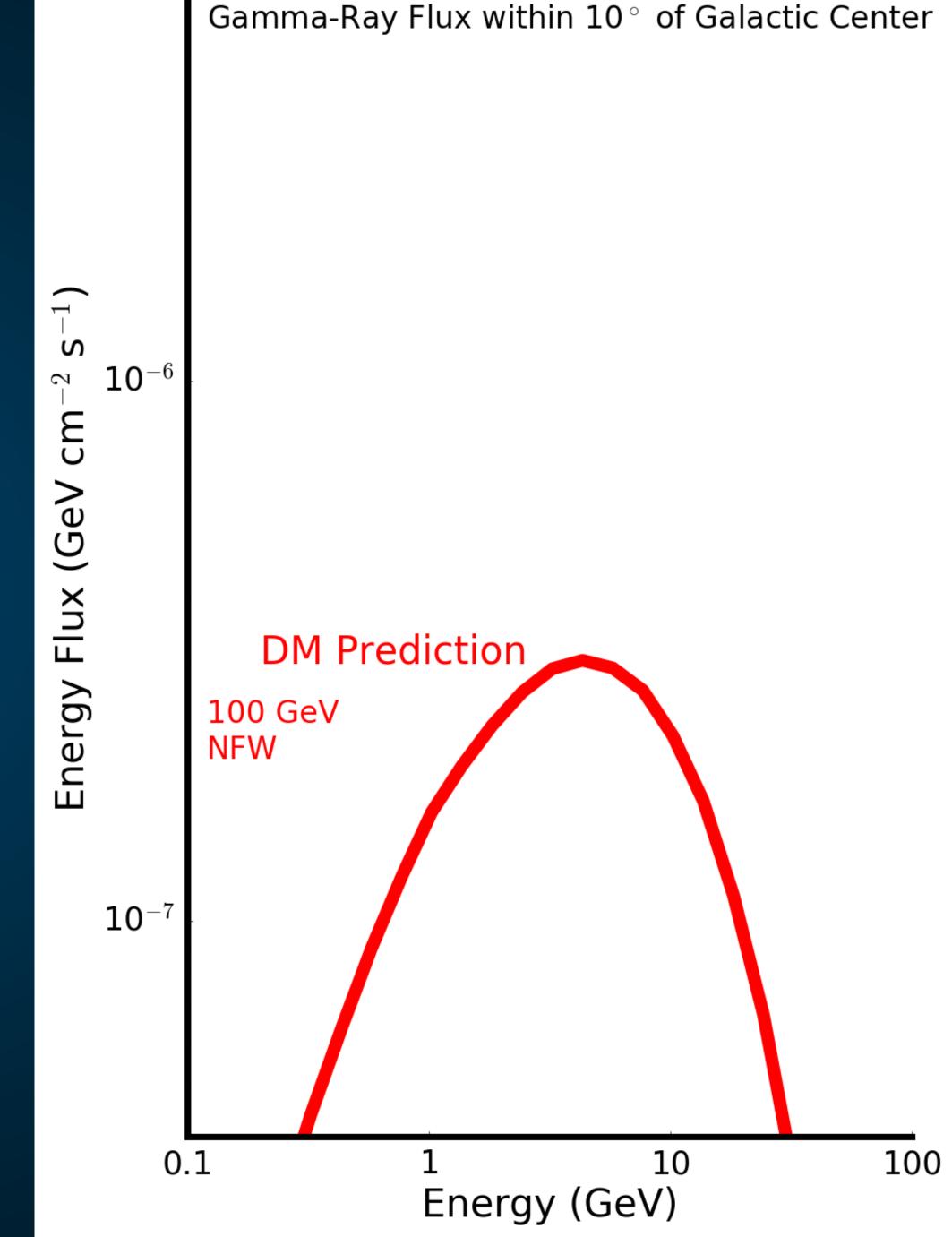
#### Thermal WIMPs and the Story of Tantalus

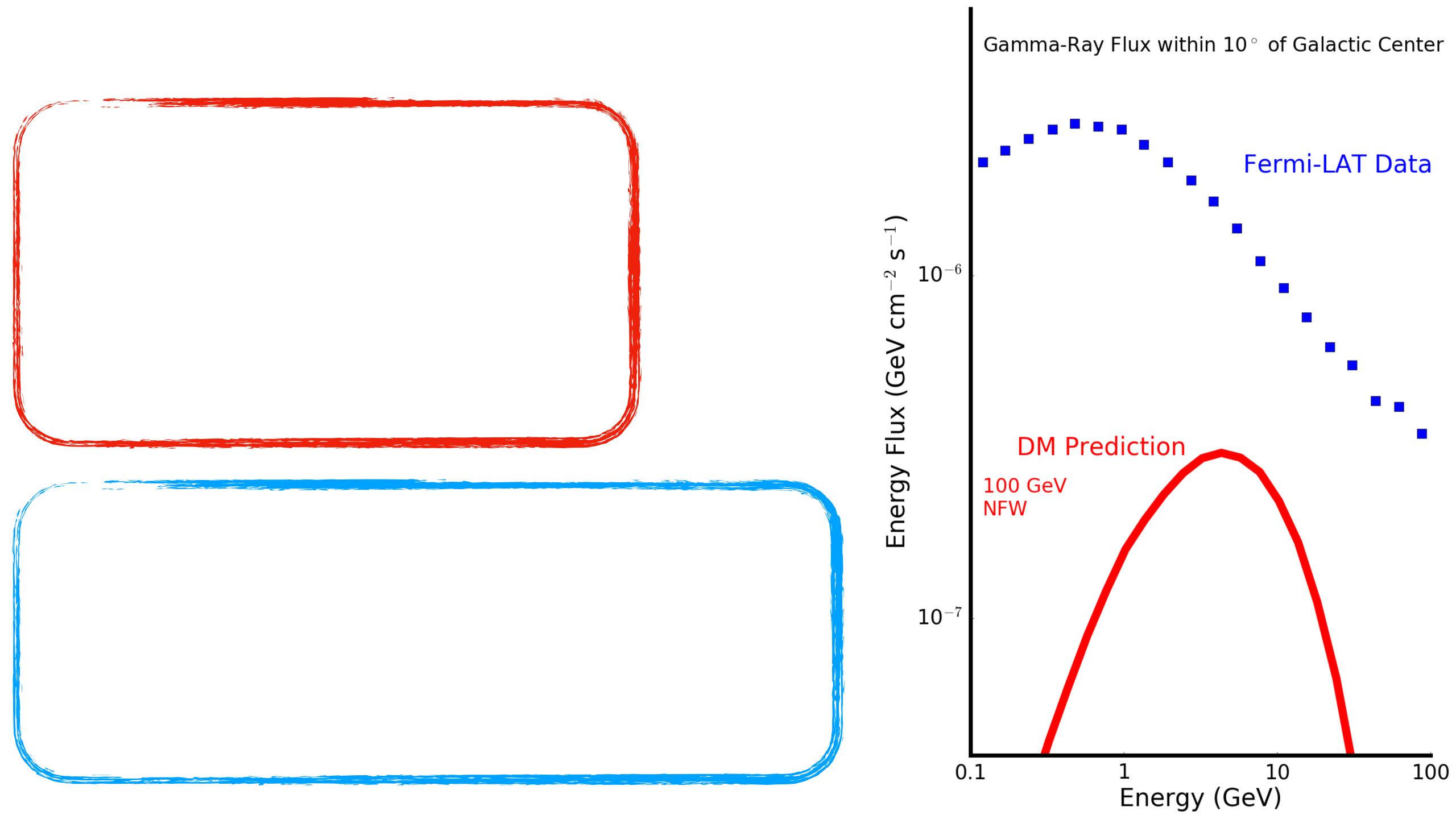
NFW Profile (Mass of Milky Way)

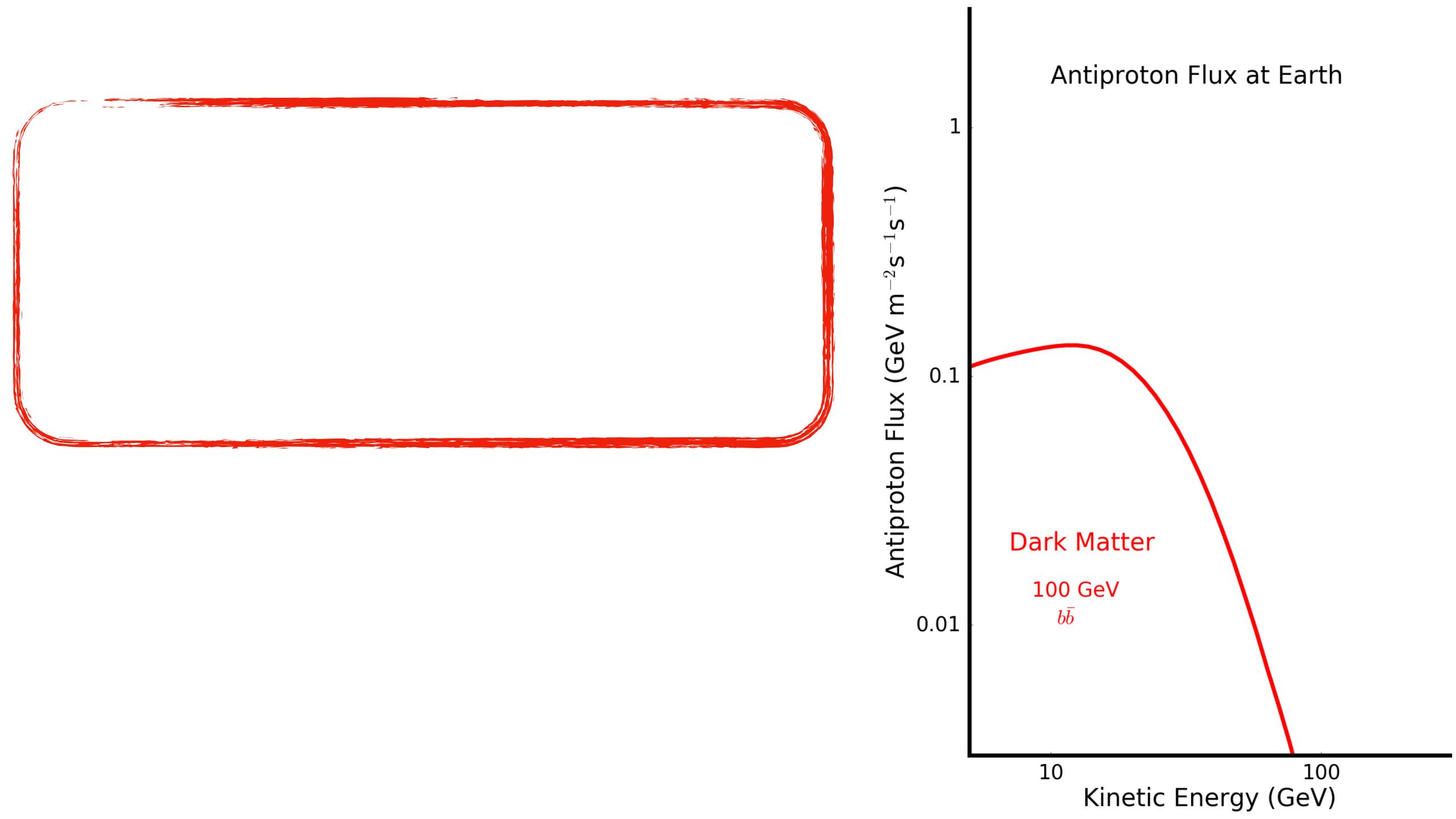
**Thermal Cross-Section (Early Universe)** 

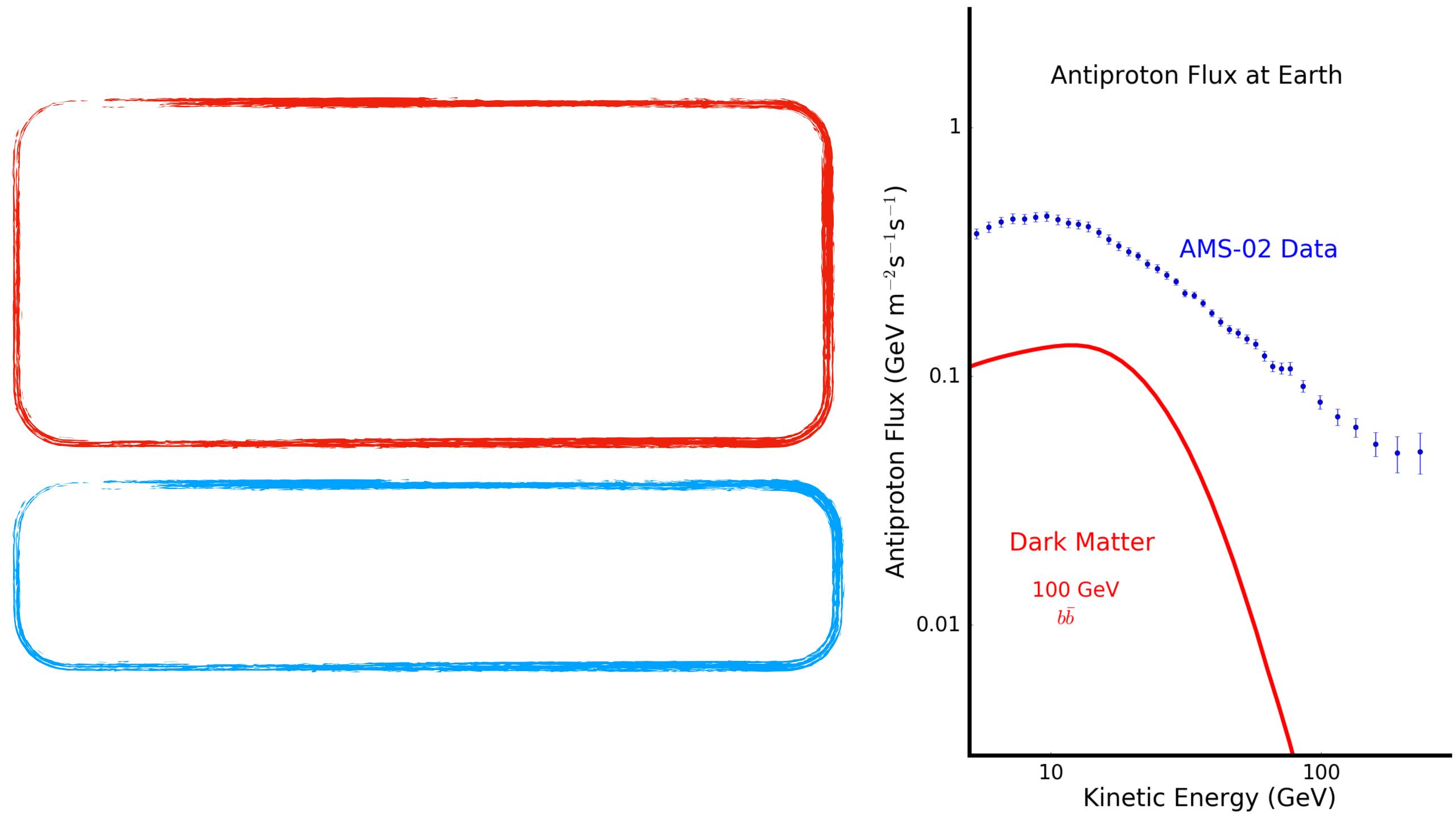
Dark Matter Mass (?)

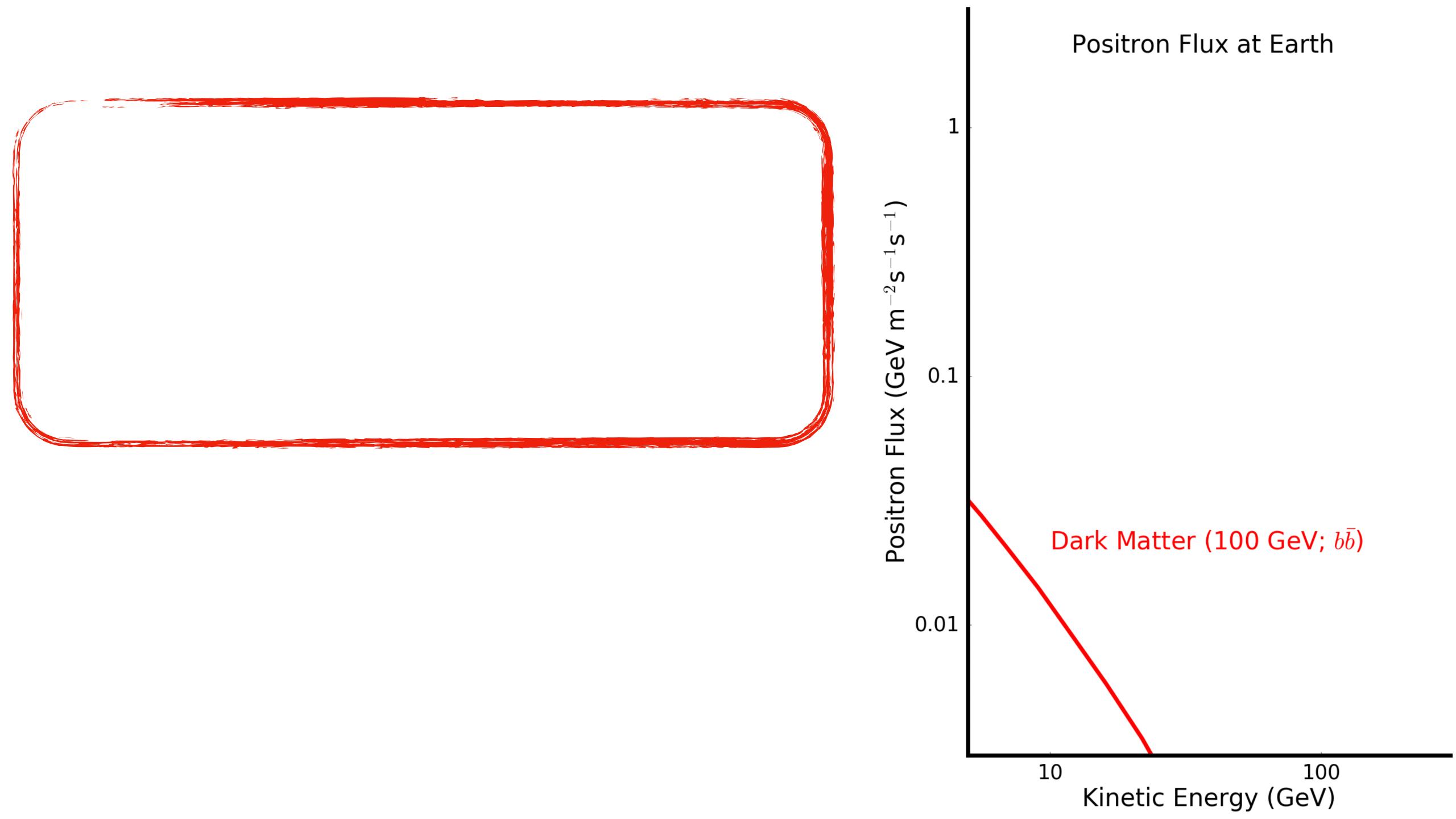
**Annihilation Final State (?)** 

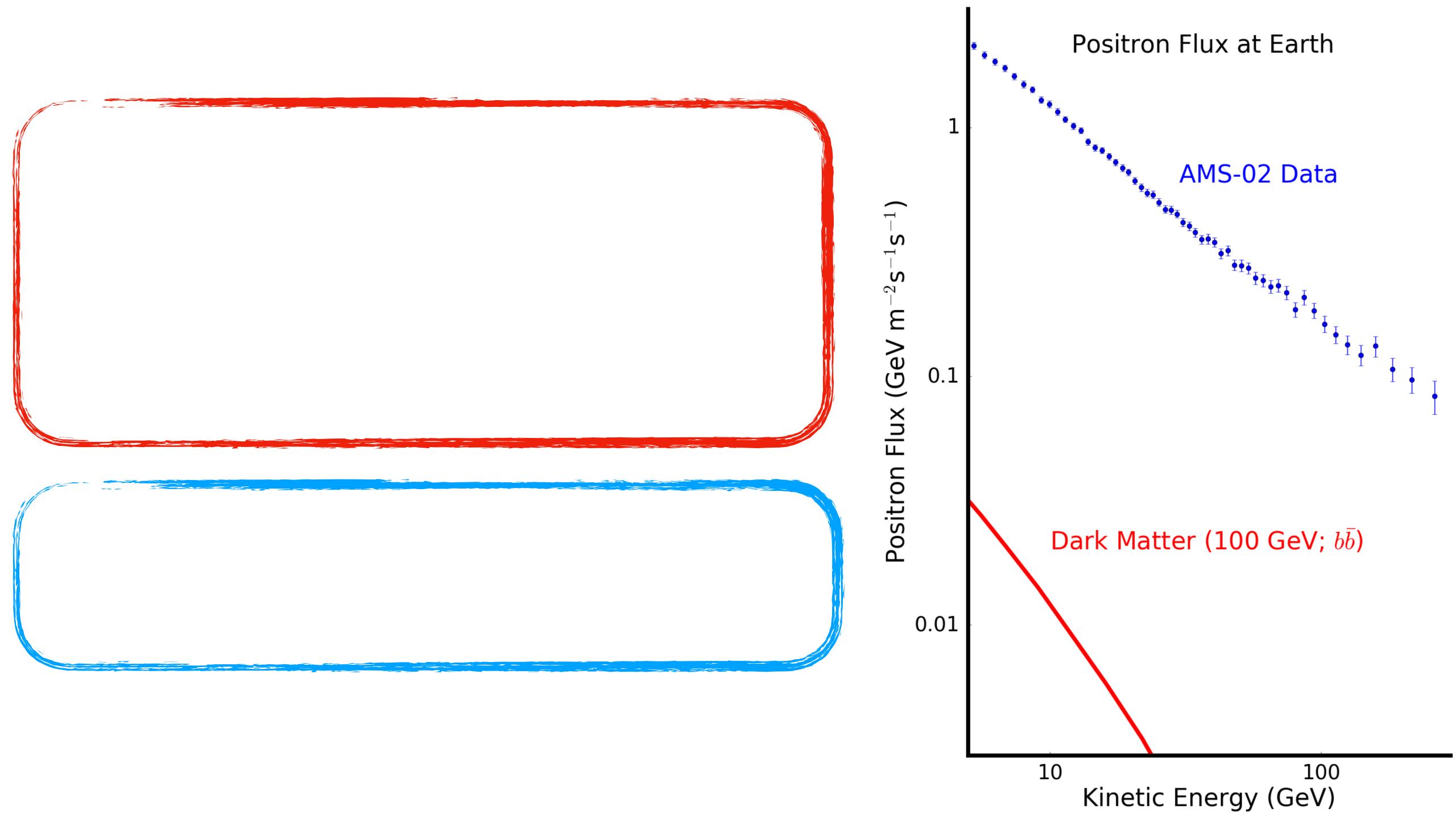


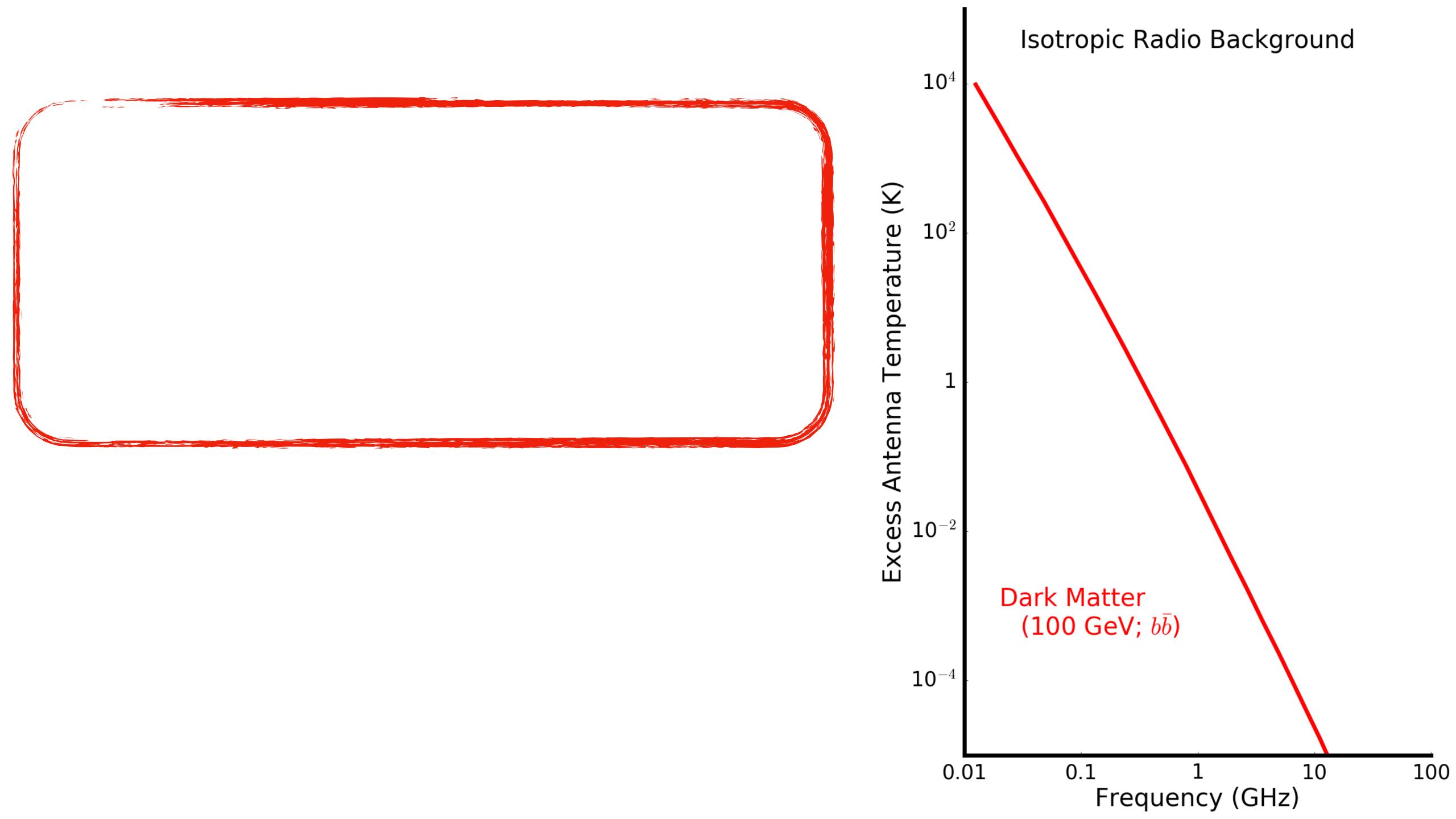


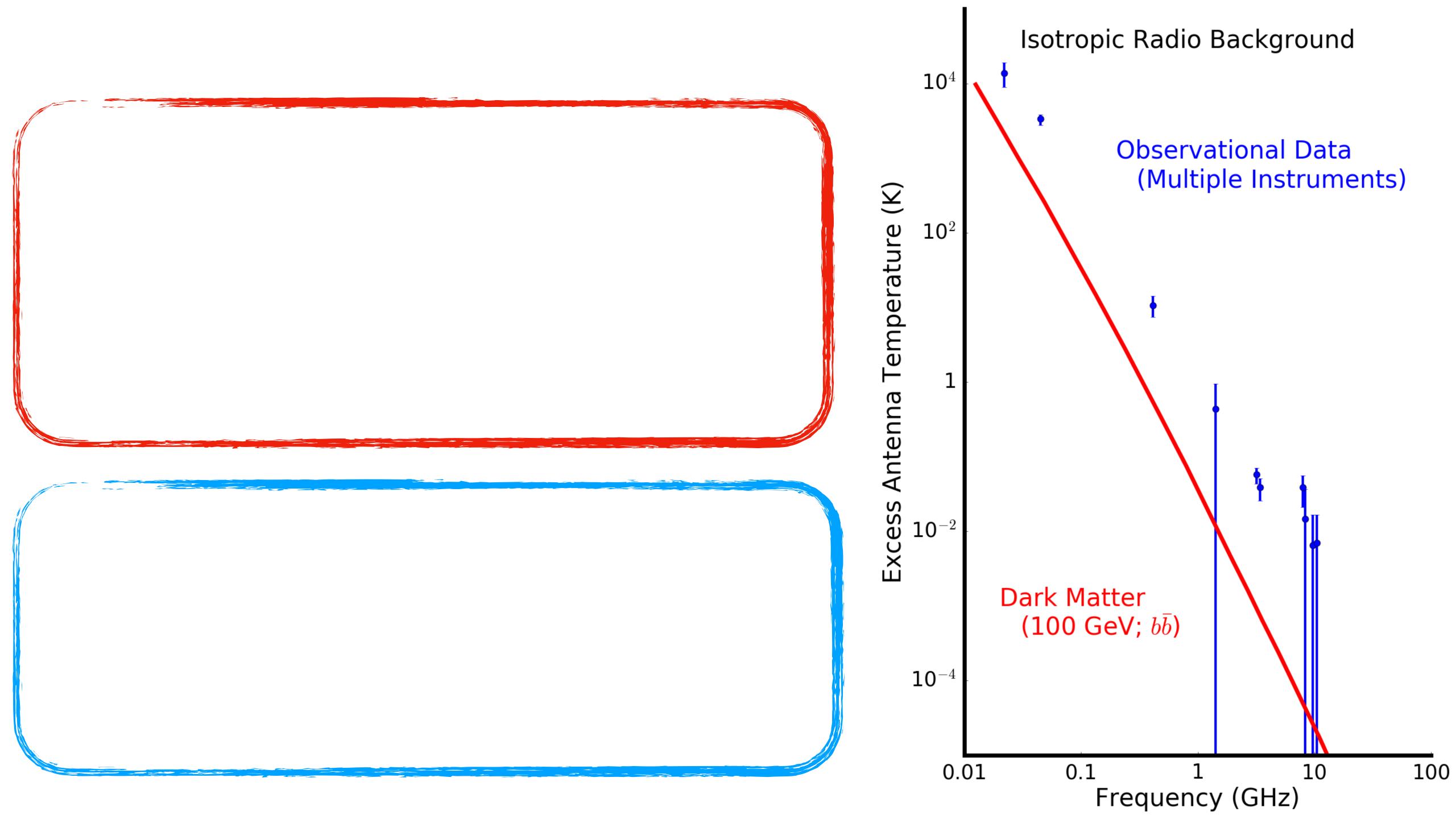




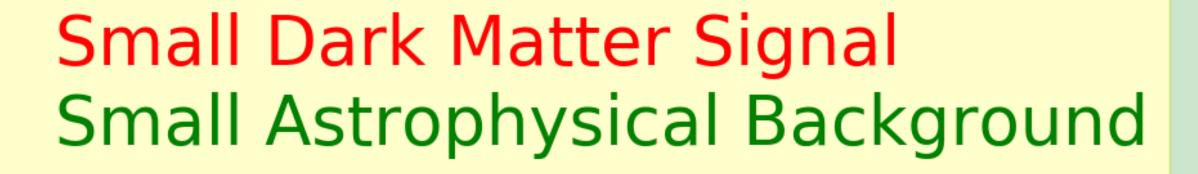












Small Dark Matter Signal Large Astrophysical Background

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Easy

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Fraction of Dark Matter Flux

Large Dark Matter Signal Small Astrophysical Background

Easy

# Easy

Small Dark Matter Signal Large Astrophysical Background

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Fraction of Dark Matter Flux

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Easy

# Easy

Small Dark Matter Signal Large Astrophysical Background

## Hard

Large Dark Matter Signal Large Astrophysical Background

Hard

Large Dark Matter Signal Small Astrophysical Background

Easy

Easy

Small Dark Matter Signal Large Astrophysical Background

Hard

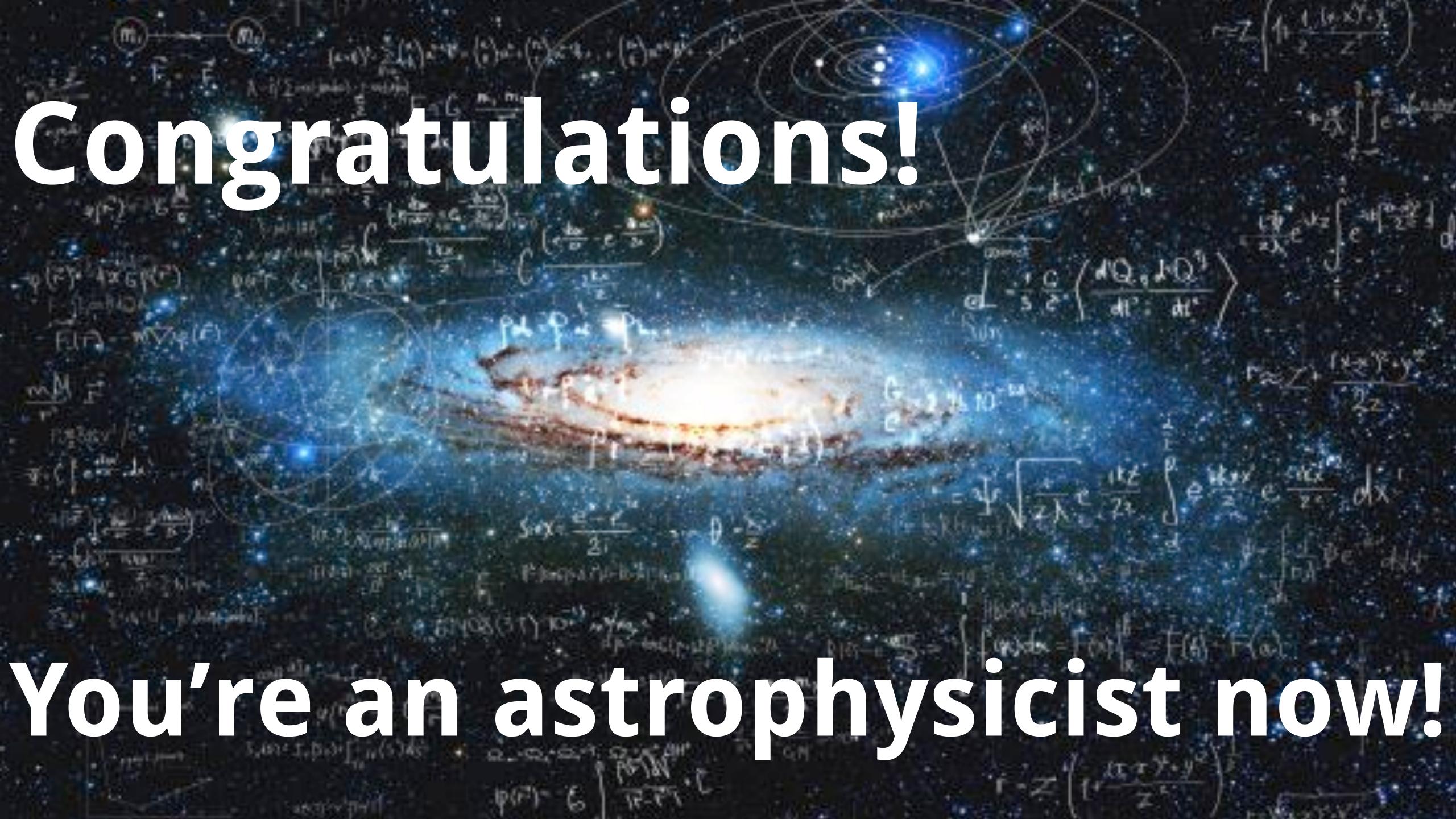
Large Dark Matter Signal Large Astrophysical Background

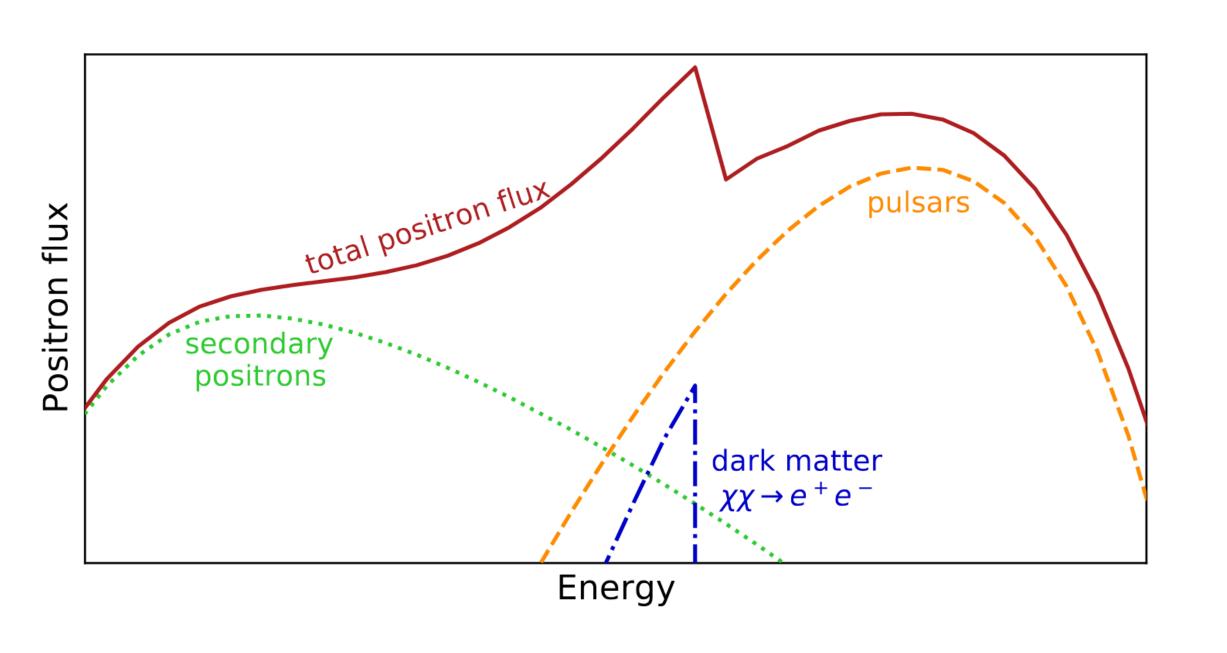
Fraction of Dark Matter Flux

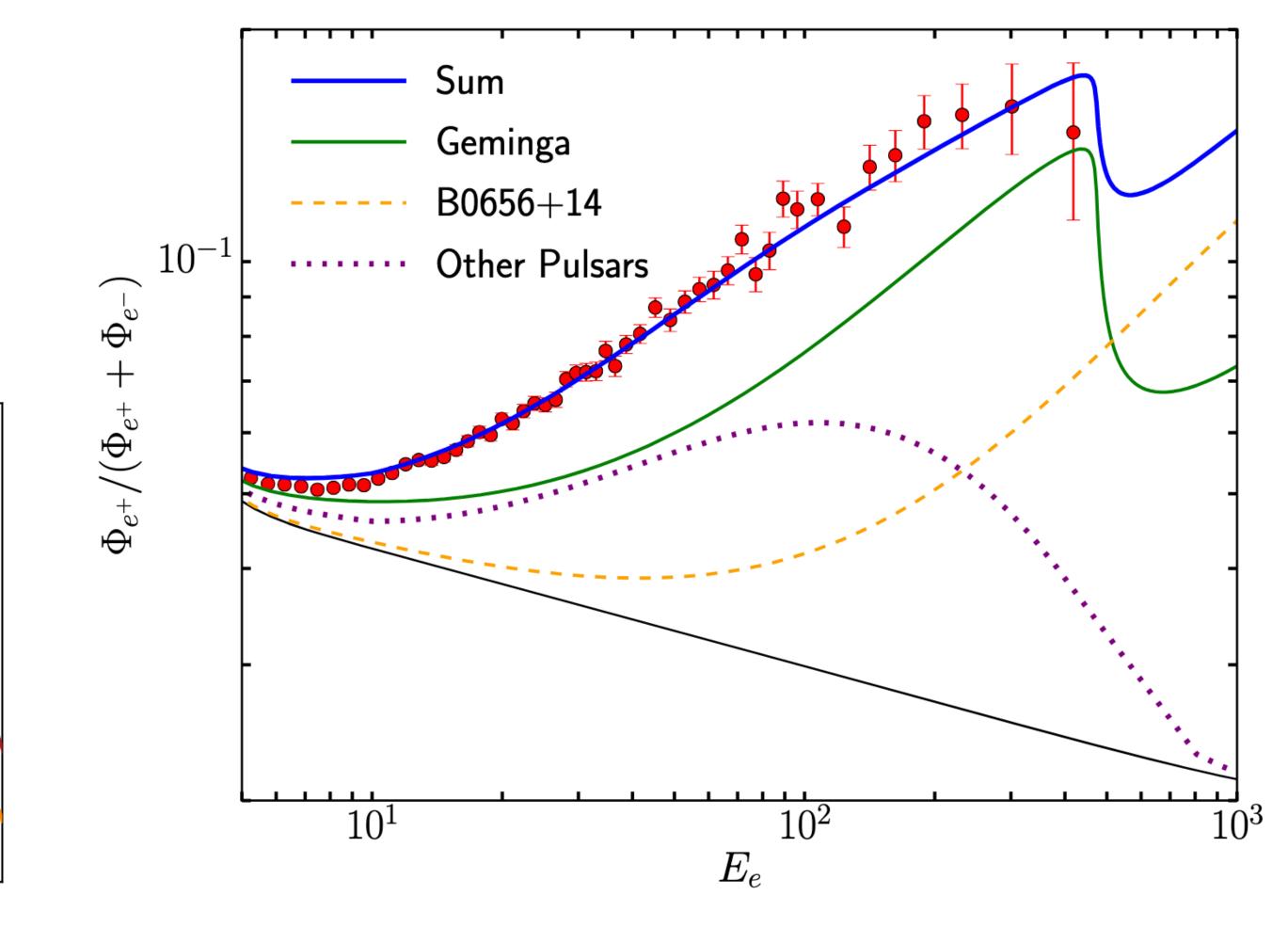
Fraction of Dark Matter Flux





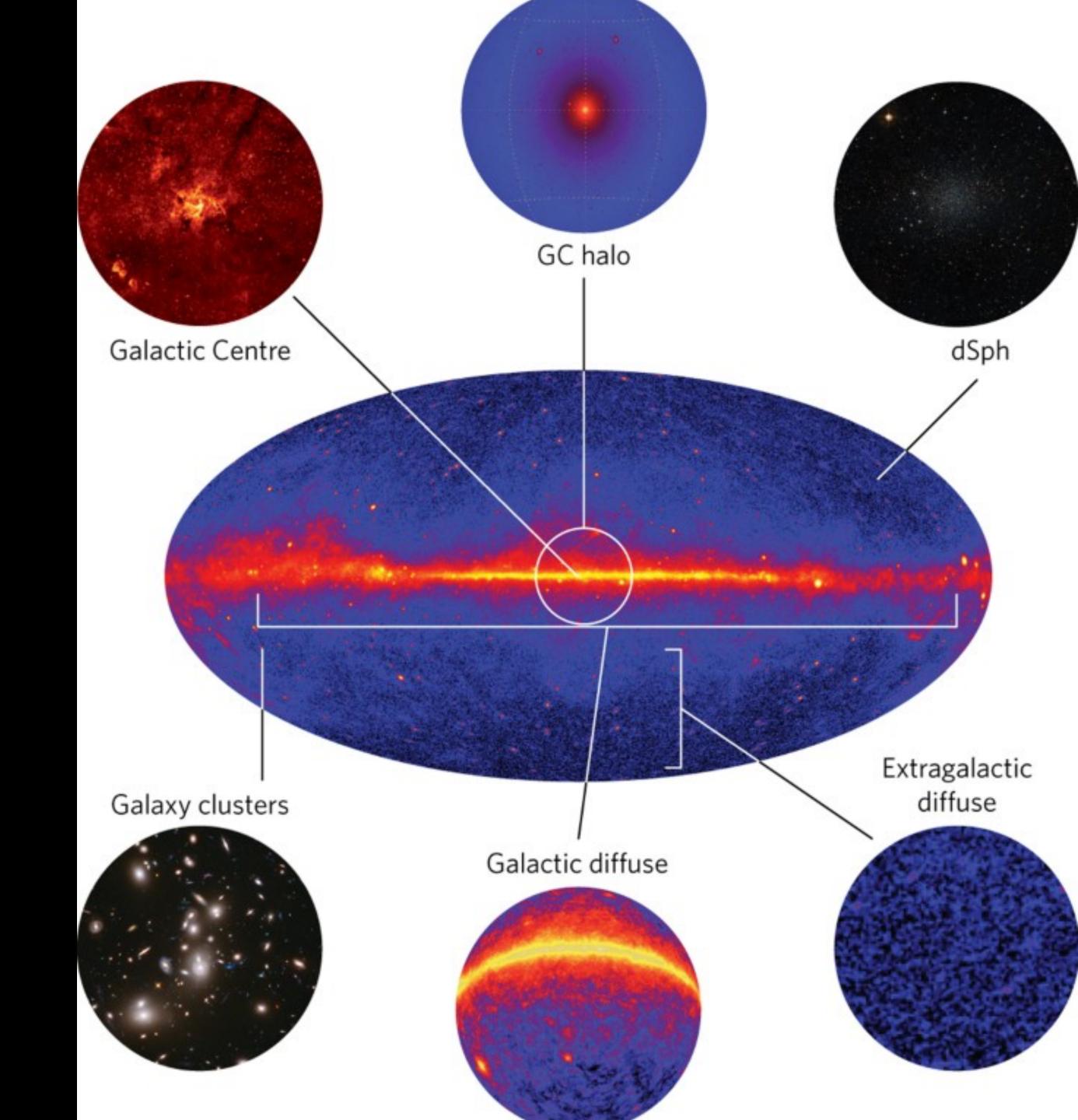






## Targets

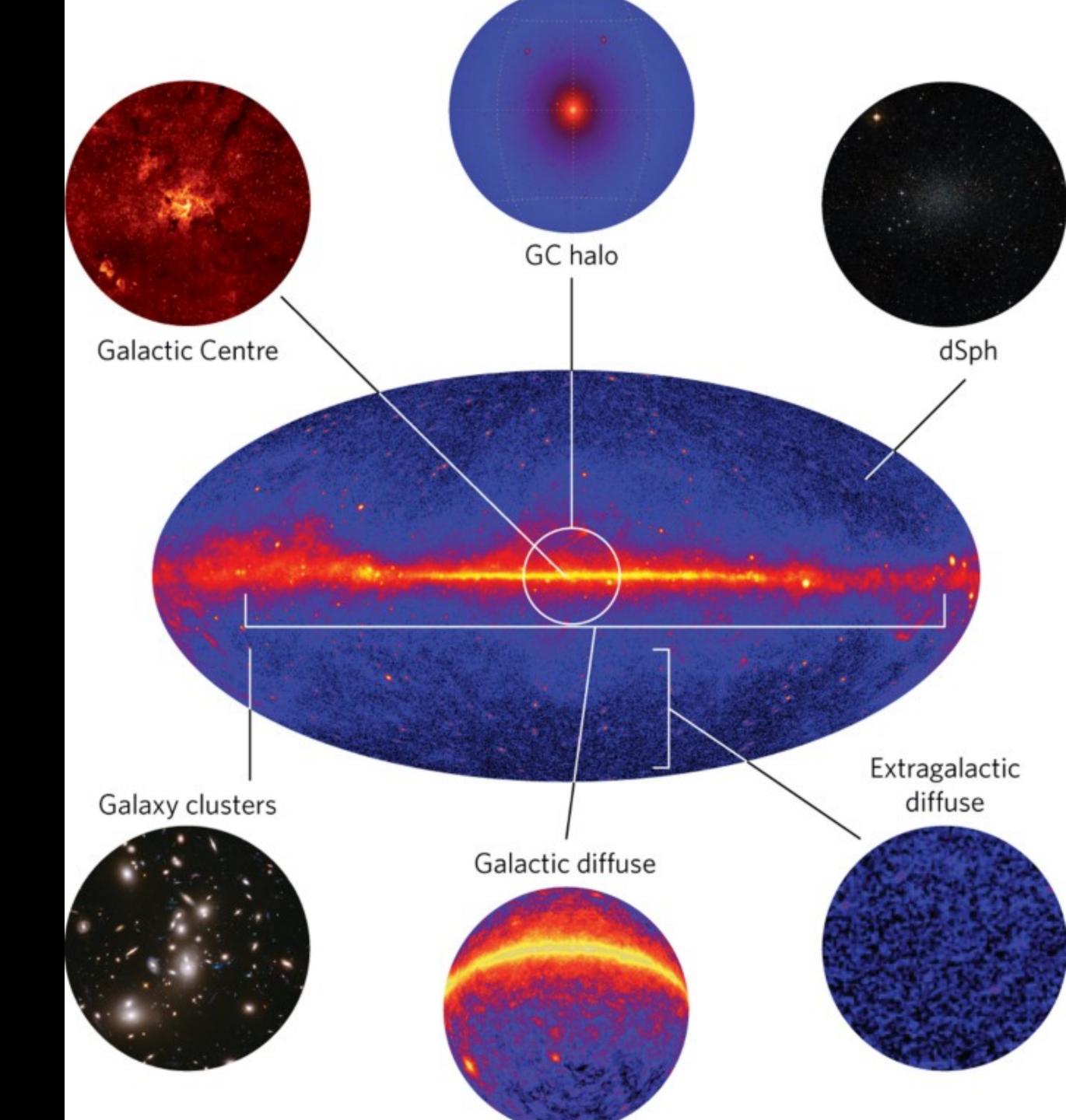
 2008: Very uncertain which dark matter targets would be most sensitive



## Targets

 2008: Very uncertain which dark matter targets would be most sensitive

 2023: Some consensus that most interesting targets are Galactic center and dSphs (though surprises always possible)



### Gamma-Rays

Galactic Center

Dwarf Spheroidal Galaxies

Galaxy Clusters

Milky Way Subhalos

**Galactic Diffuse** 

Sun

Jupiter

**Nearby Stars** 

Galactic Center Stars

Andromeda

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Isotropic Gamma-Ray Background

**Anisotropy Searches** 

Cusps

511 keV line

Morphology

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#### **Positrons**

**Electron + Positron Spectrum** 

**Antiprotons** 

Antineutrons

Antihelium

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**Galaxy Cluster Synchrotron** 

**Diffuse Synchrotron** 

Sun

**Jupiter** 

**Isotropic Background** 

X-ray background from Clusters

**Anisotropy Searches** 

**Stellar Evolution** 

**Pulsar Evolution** 

**Planetary Heating** 

**Thermal Scattering** 

Cosmic Microwave Background

**CMB** Absorption

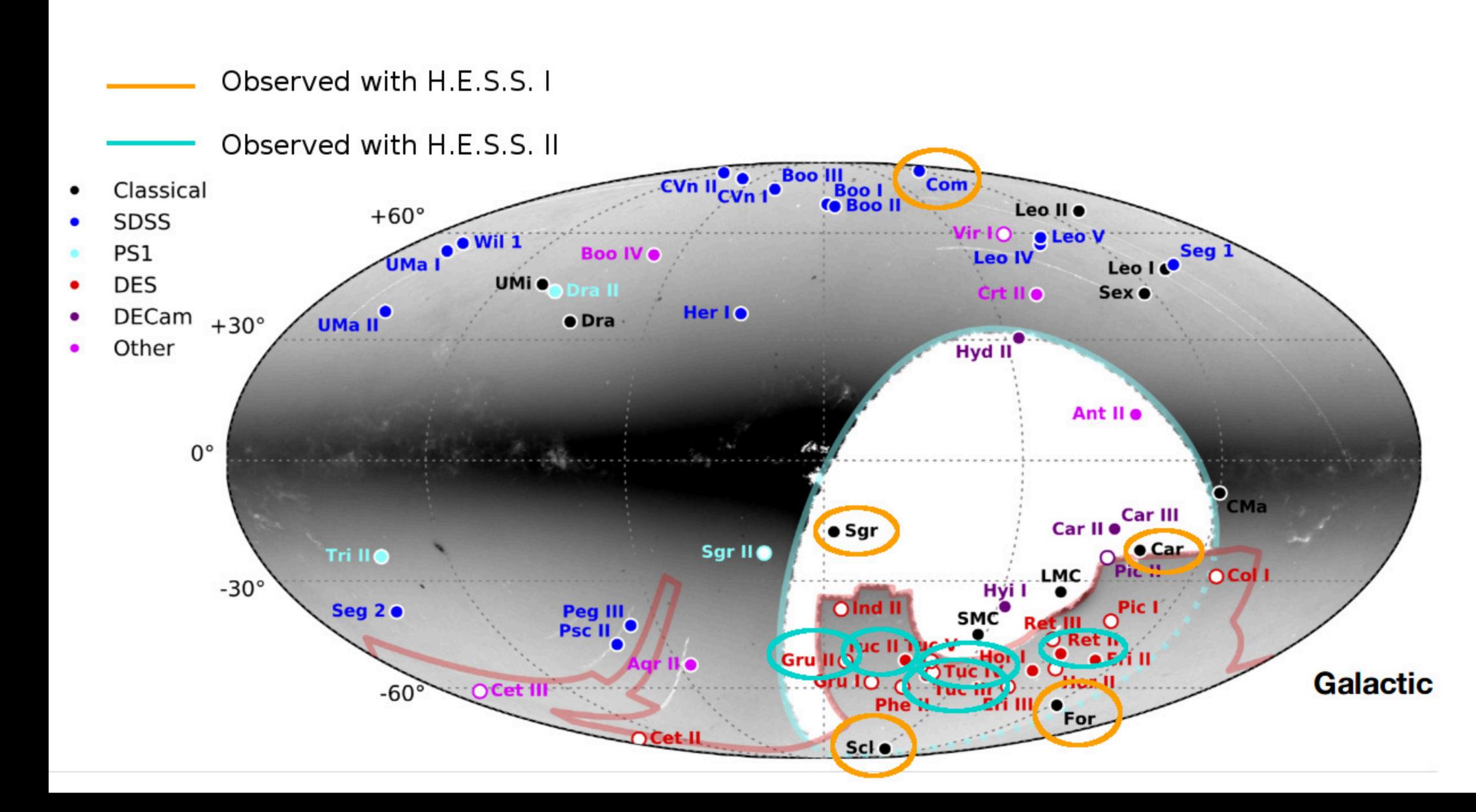
#### **Targets**

## Topics:

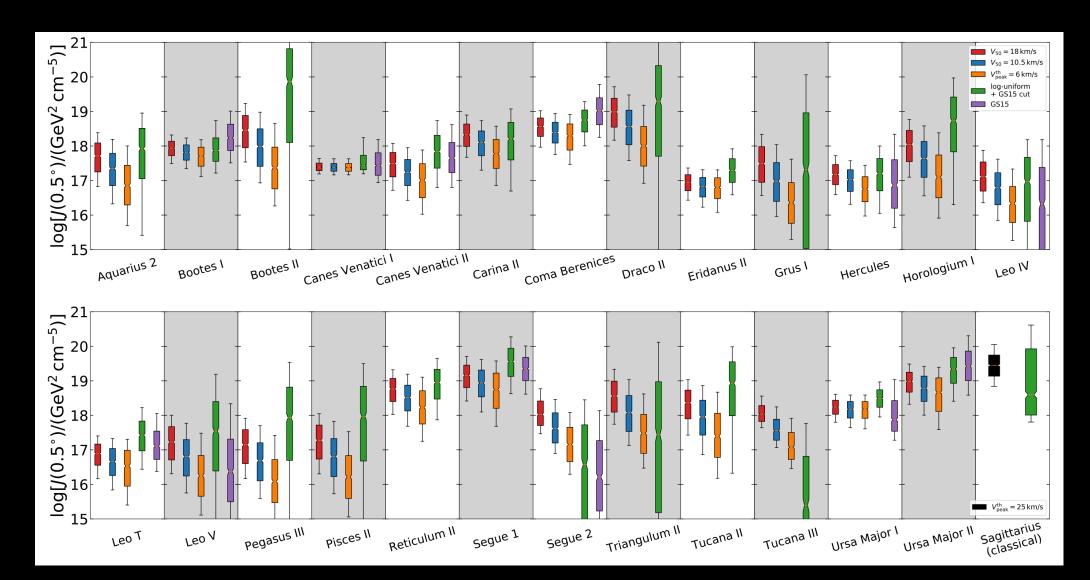
- o Thermal/WIMP-like:
  - o <u>dSphs</u>
  - o Galactic Center Excess
  - o Positron Excess
  - Antiproton Excess
  - Antinuclei
  - o Gamma-Ray Lines
  - Neutrinos

- Axions
- Sterile Neutrinos
- o Primordial Black Holes

## dSphs



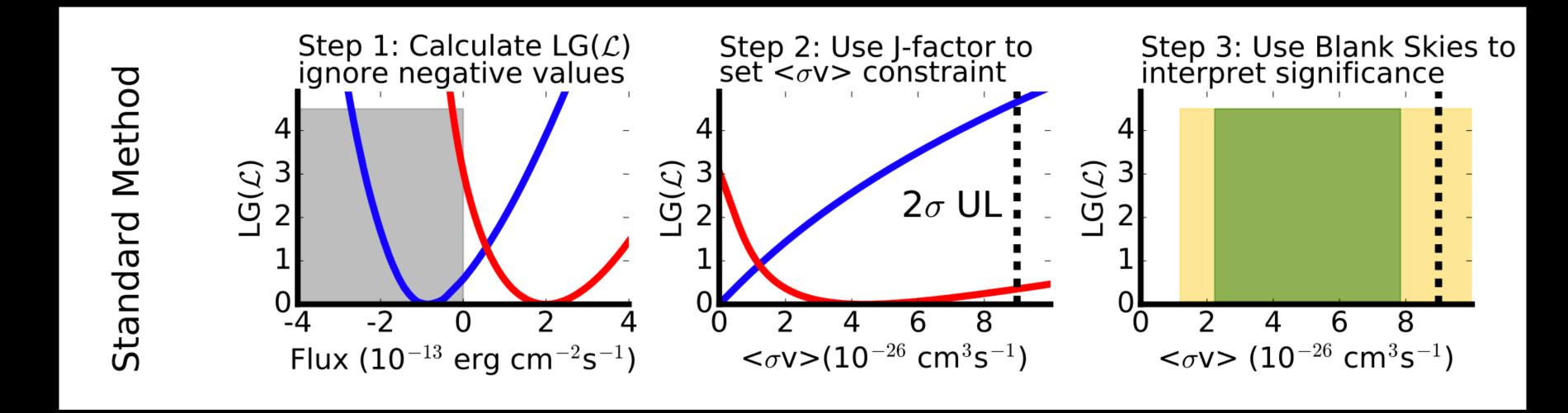
### dSphs



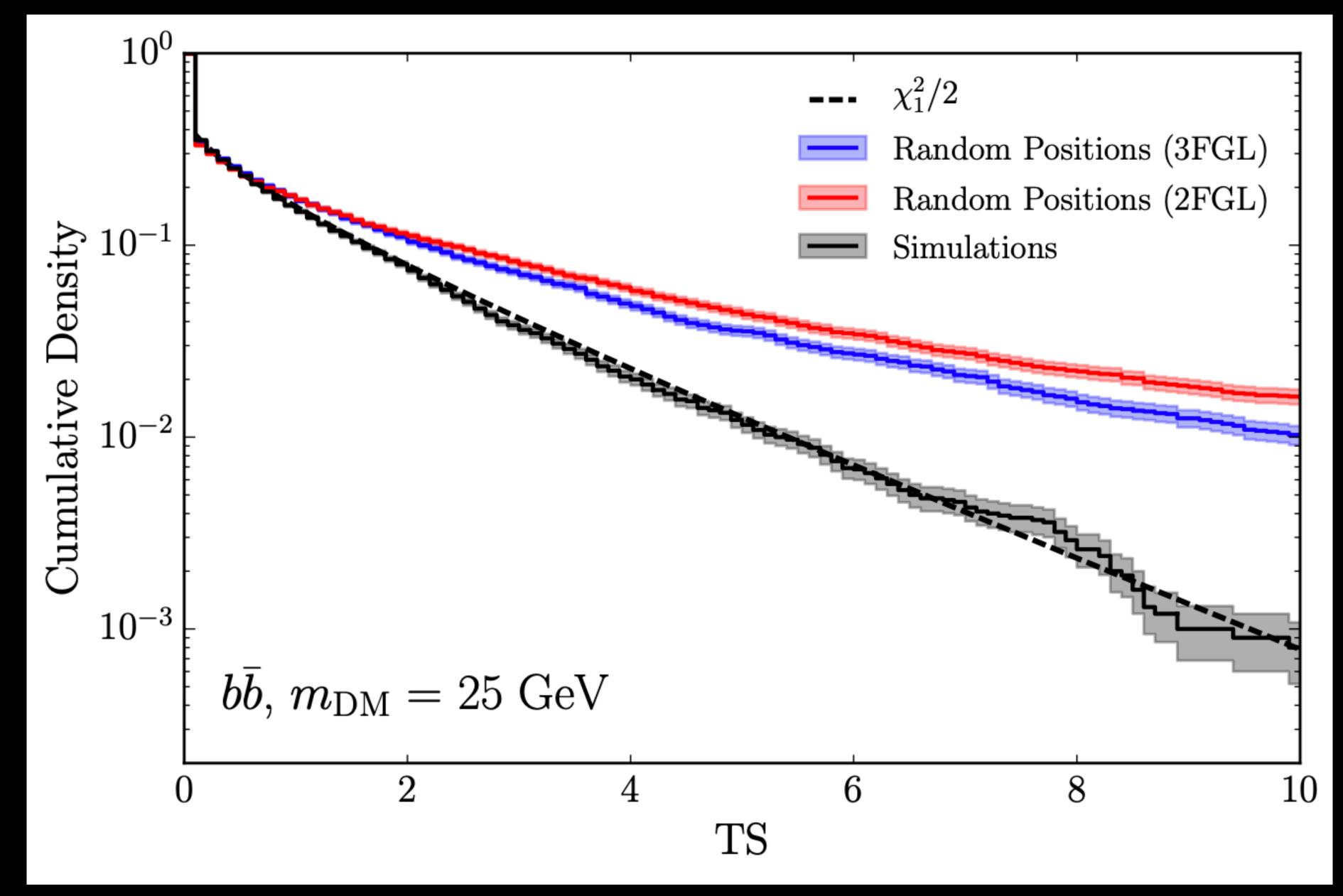
- -J-factors tell us the relative flux of each dSph
- -Can add the likelihood profiles of each source.
- -Use Blank Sky locations to test

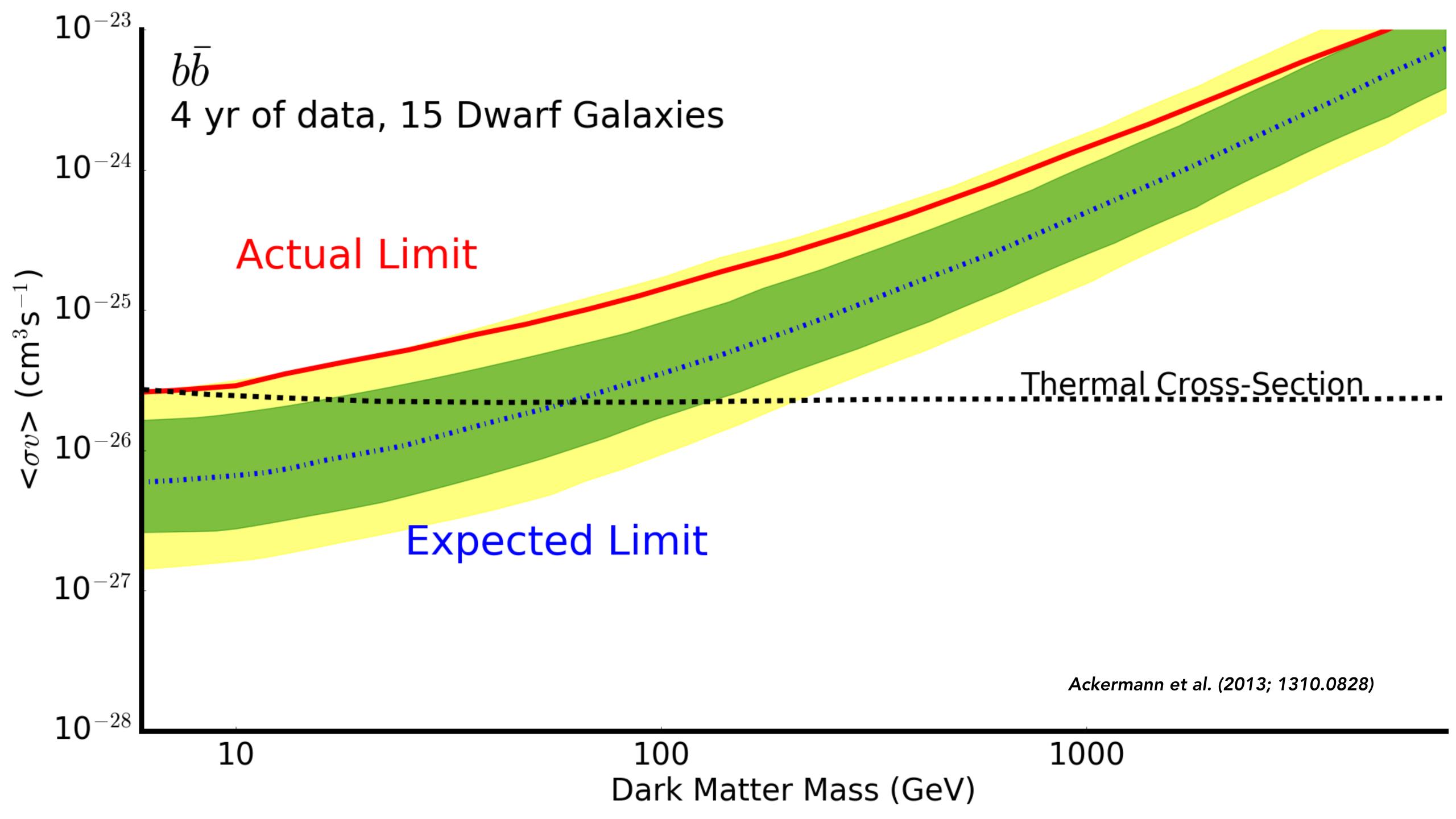
$$\mathcal{L}_{J}(J_{i} | J_{\text{obs,i}}, \sigma_{i}) = \frac{1}{\ln(10)J_{\text{obs,i}}\sqrt{2\pi}\sigma_{i}} \times e^{-(\log_{10}(J_{i}) - \log_{10}(J_{\text{obs,i}}))^{2}/2\sigma_{i}^{2}}$$

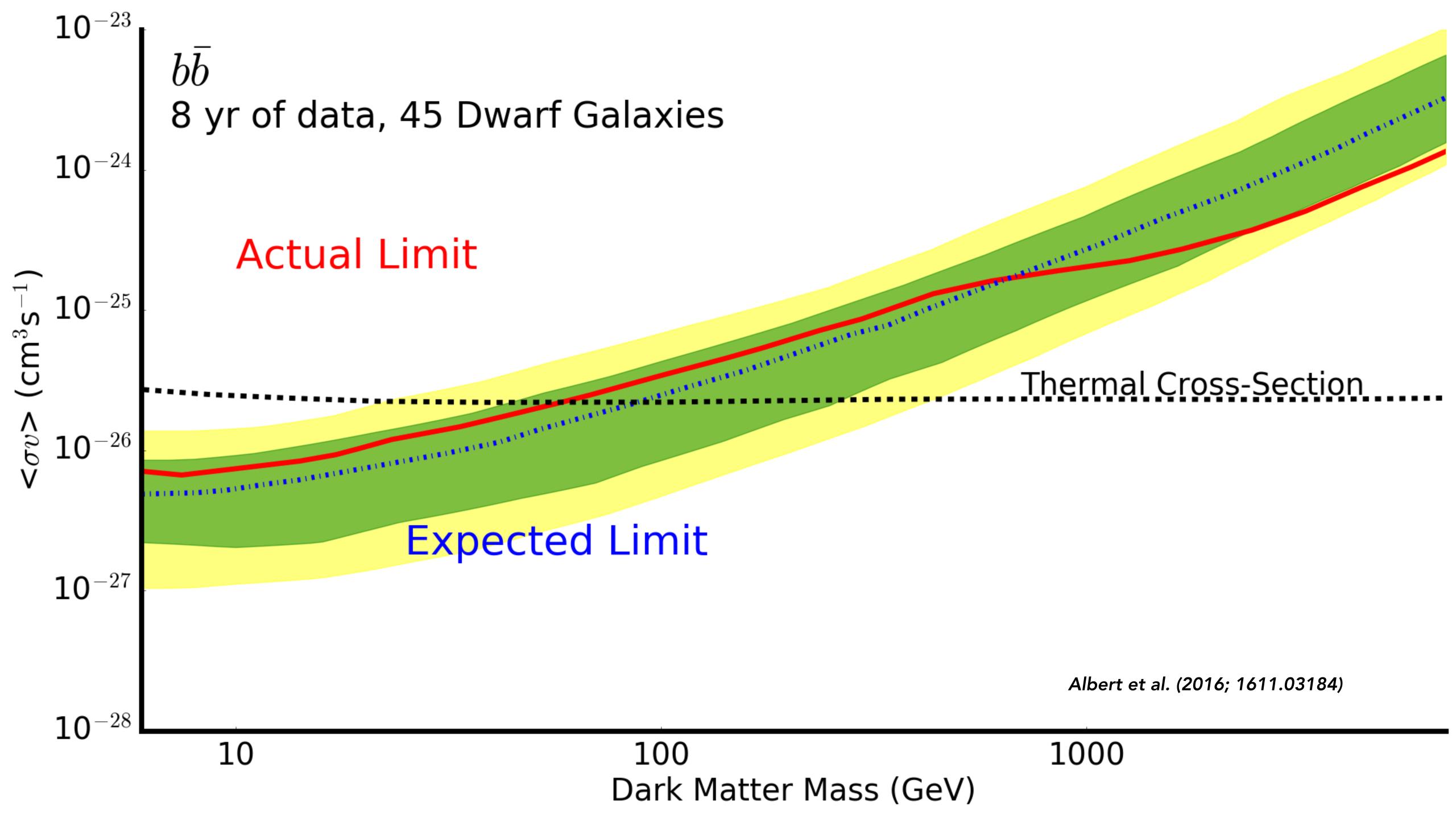
$$\mathcal{L}_i(oldsymbol{\mu},oldsymbol{ heta}_i\,|\,\mathcal{D}_i) = \prod_j \mathcal{L}_i(oldsymbol{\mu},oldsymbol{ heta}_i\,|\,\mathcal{D}_{i,j}).$$

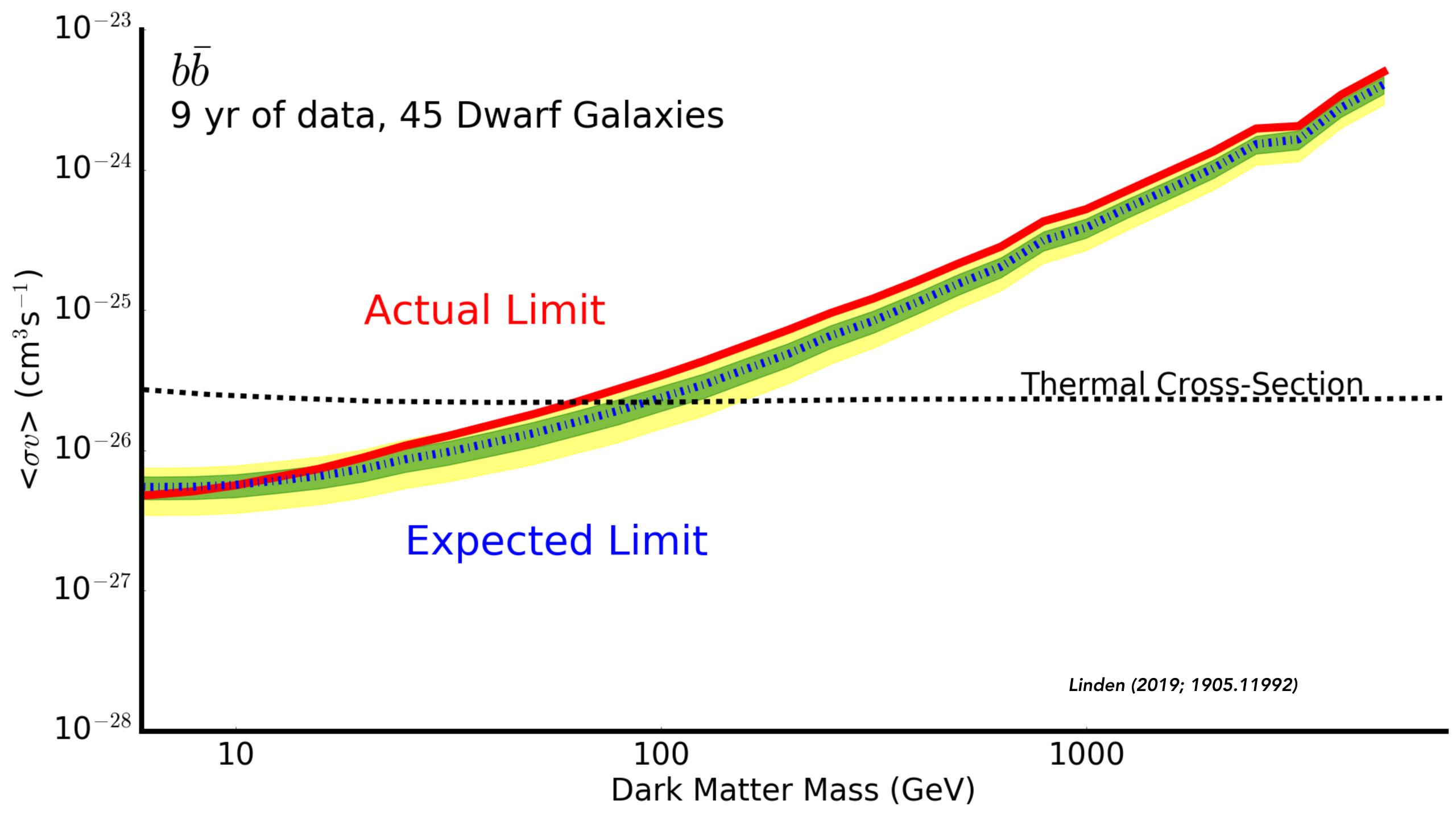


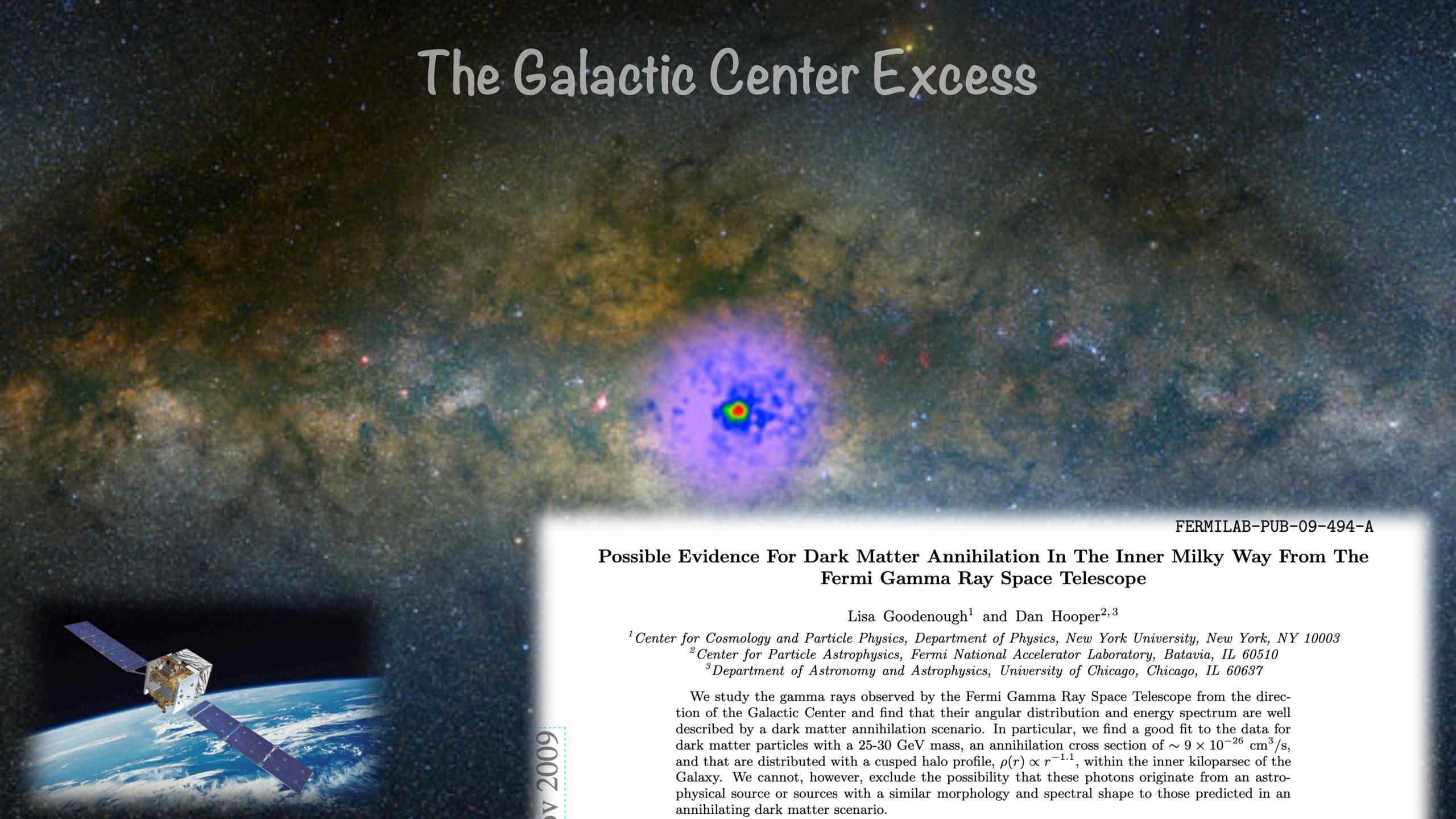
# dSphs

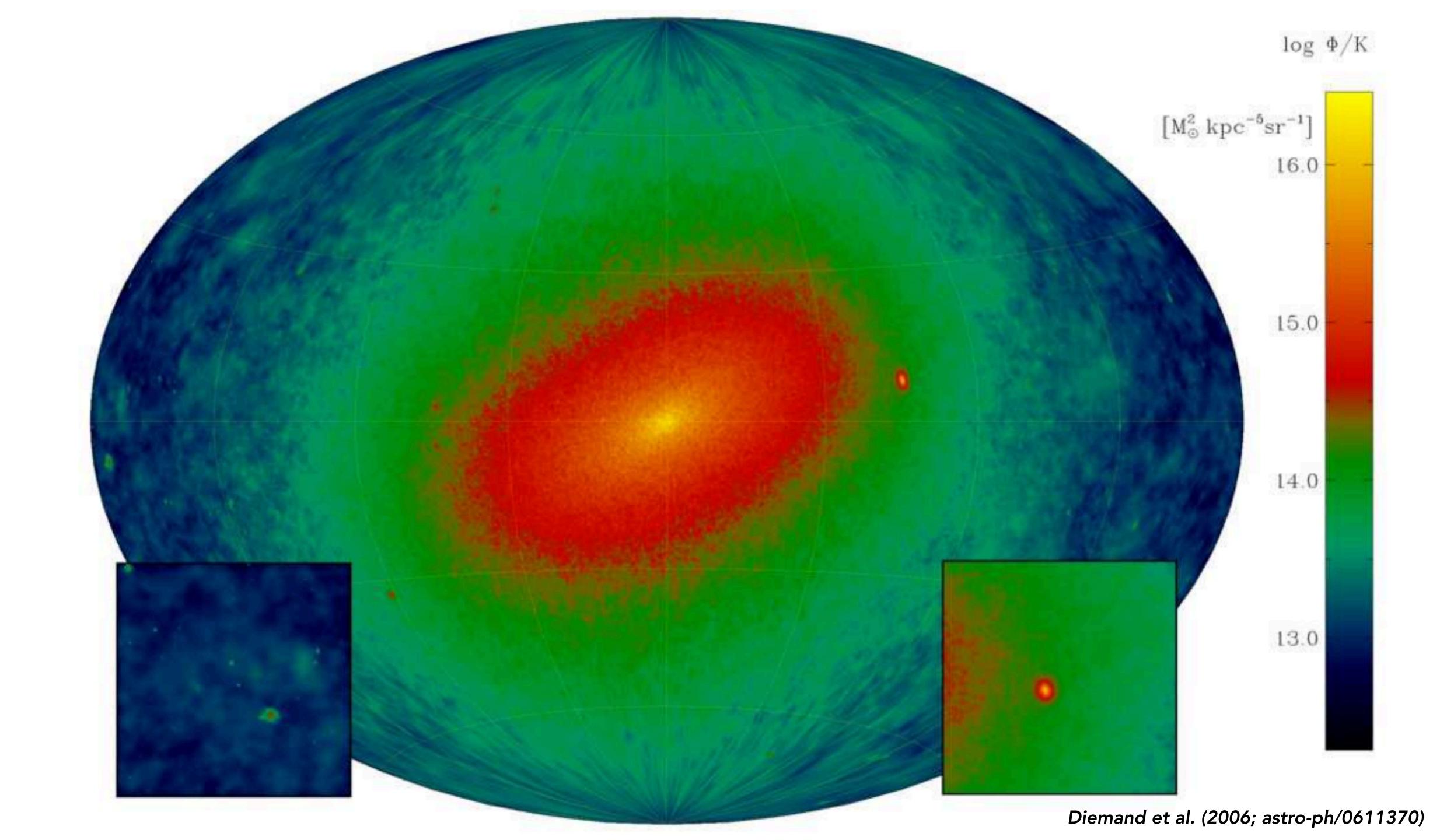


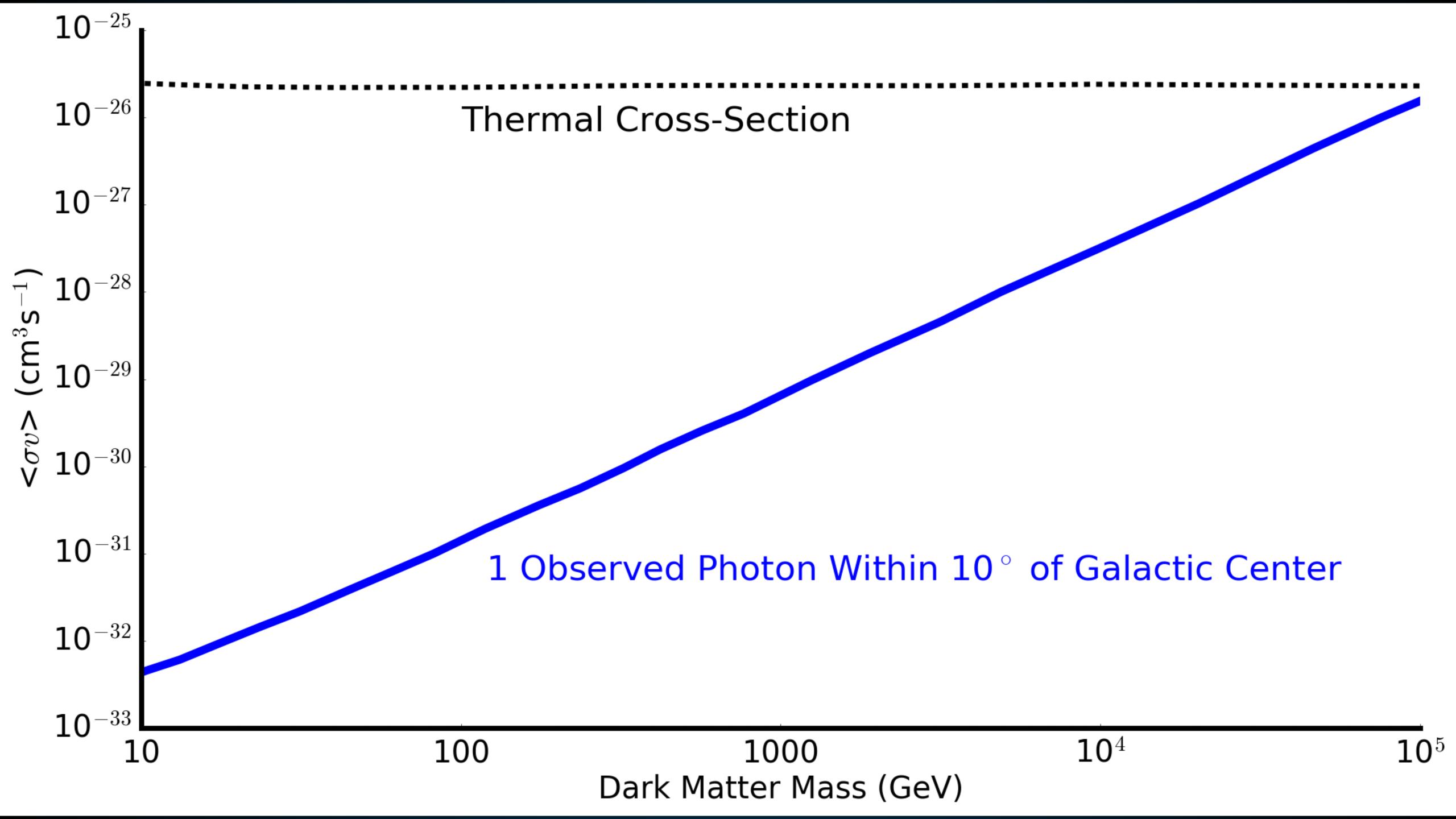


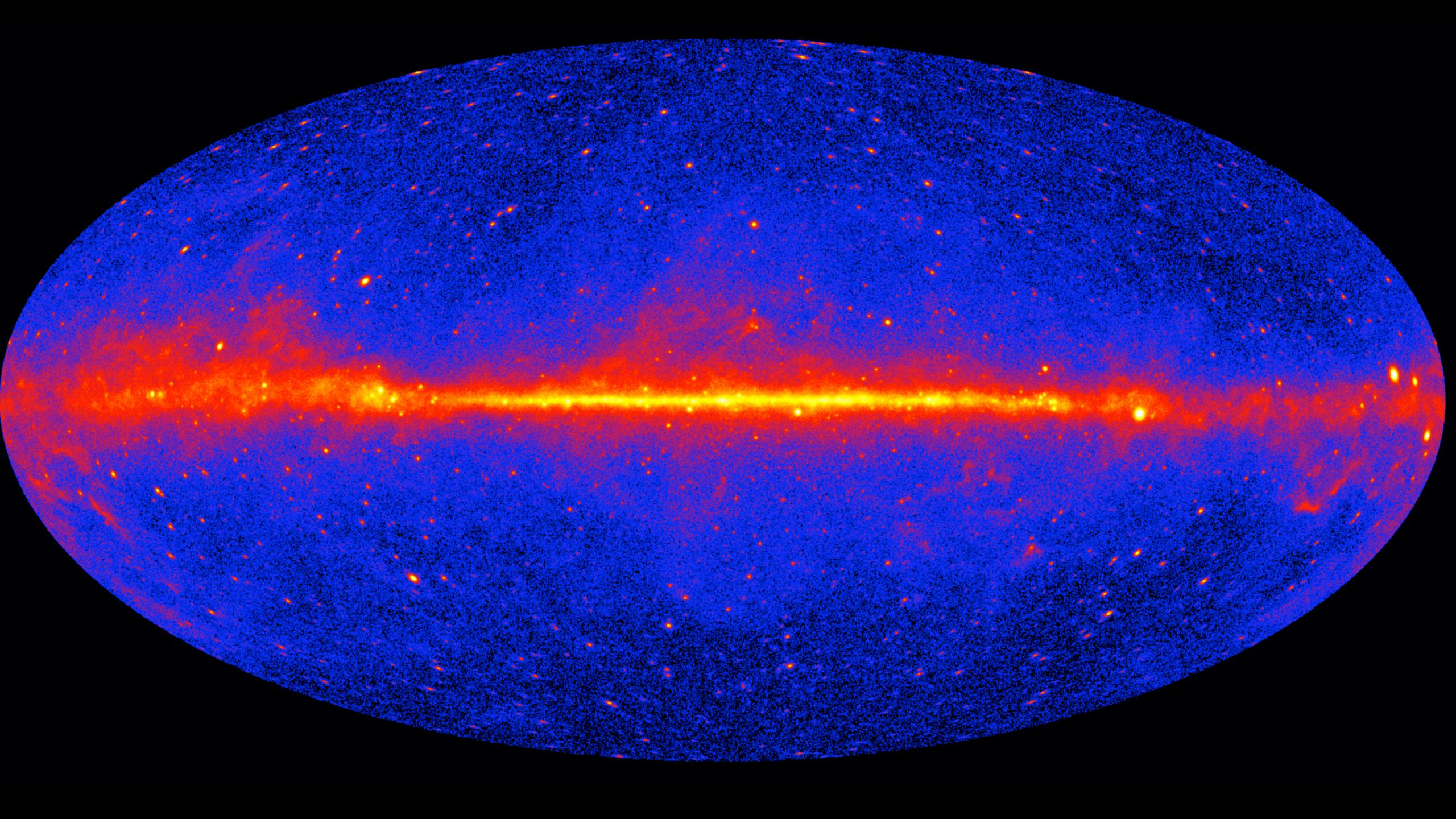


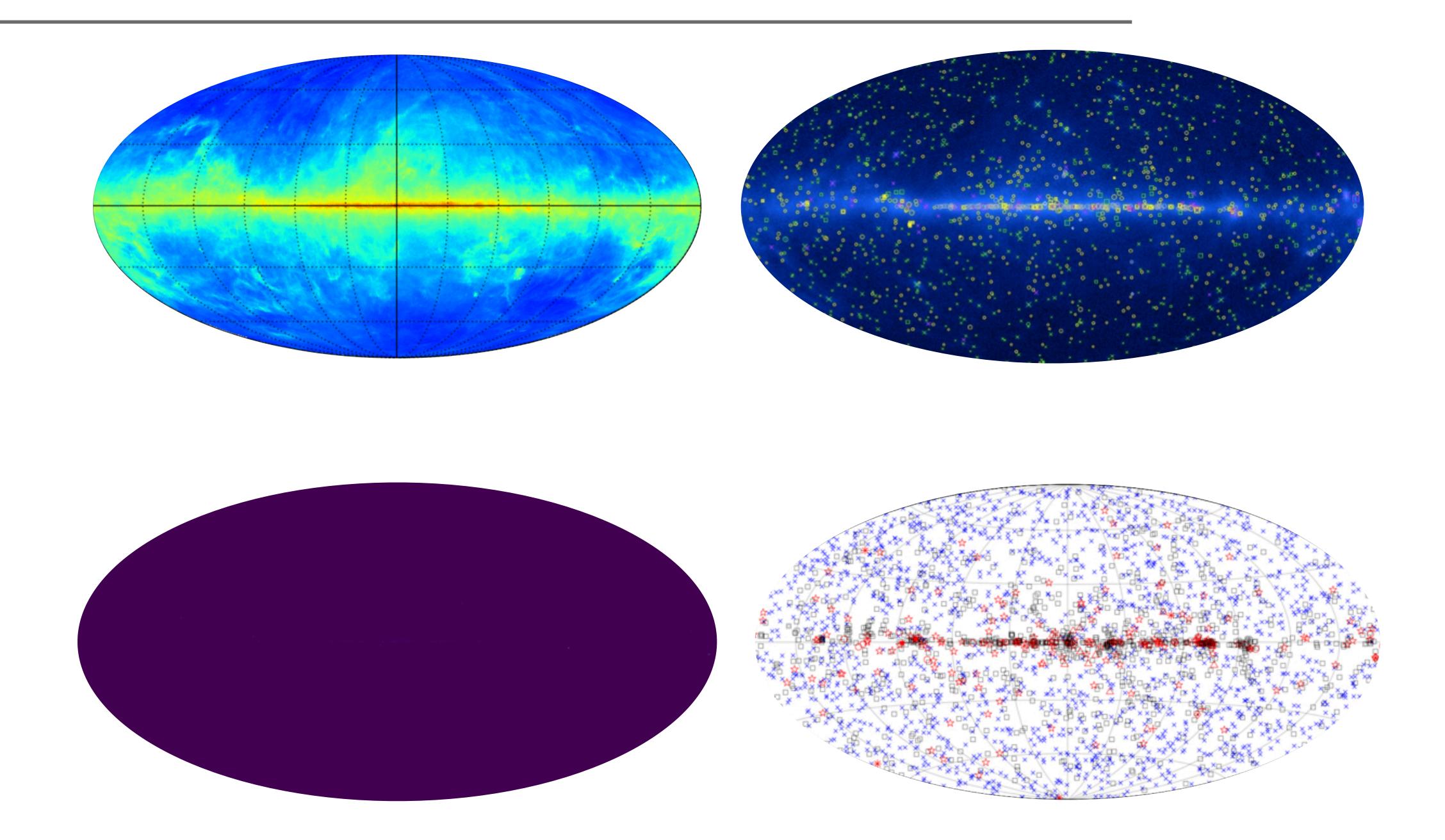


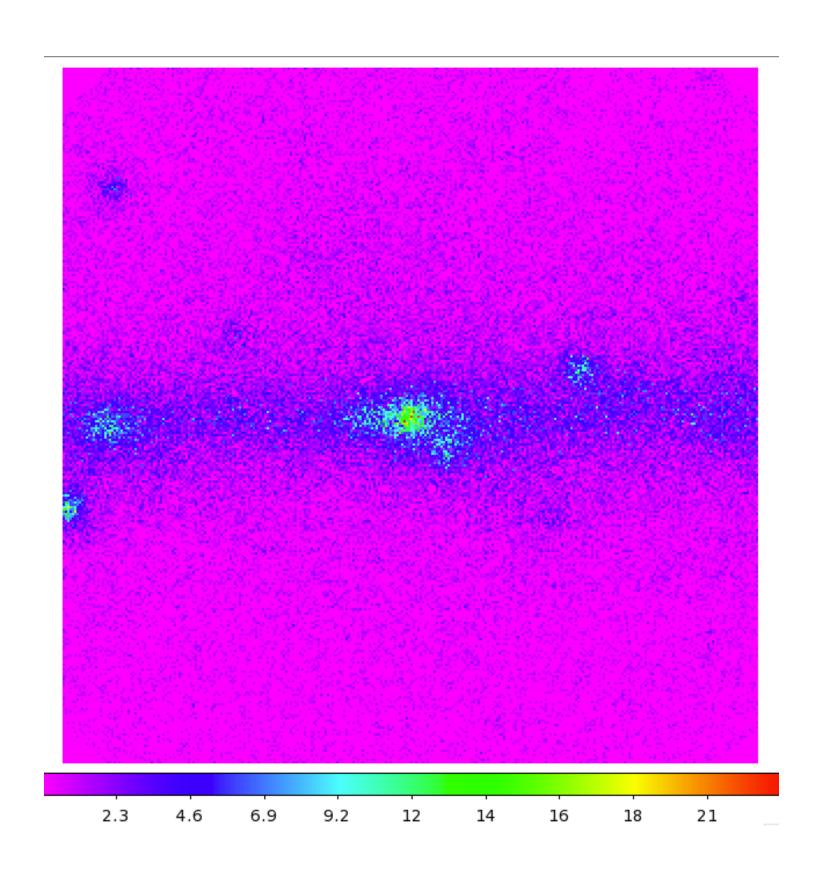


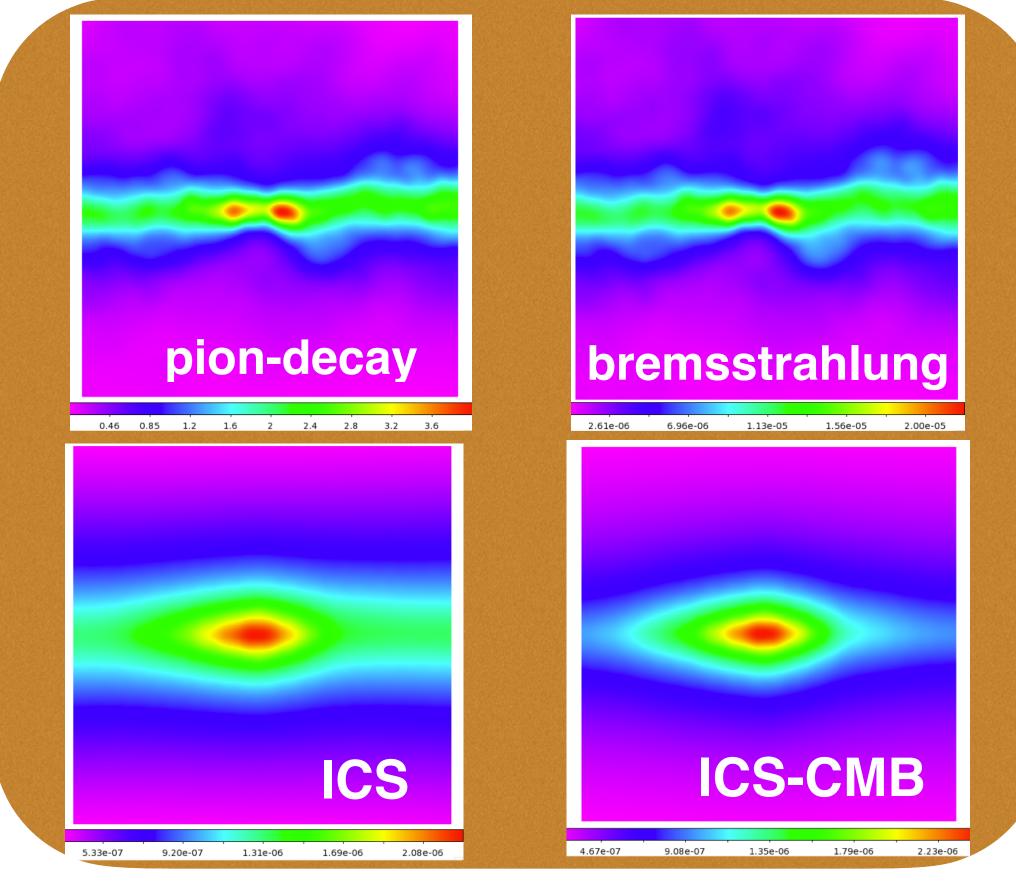


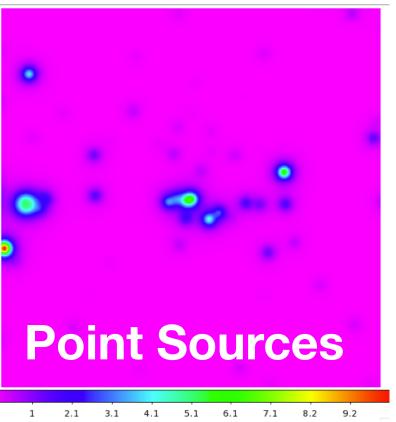


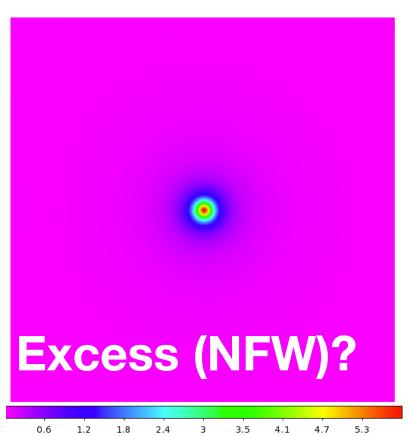








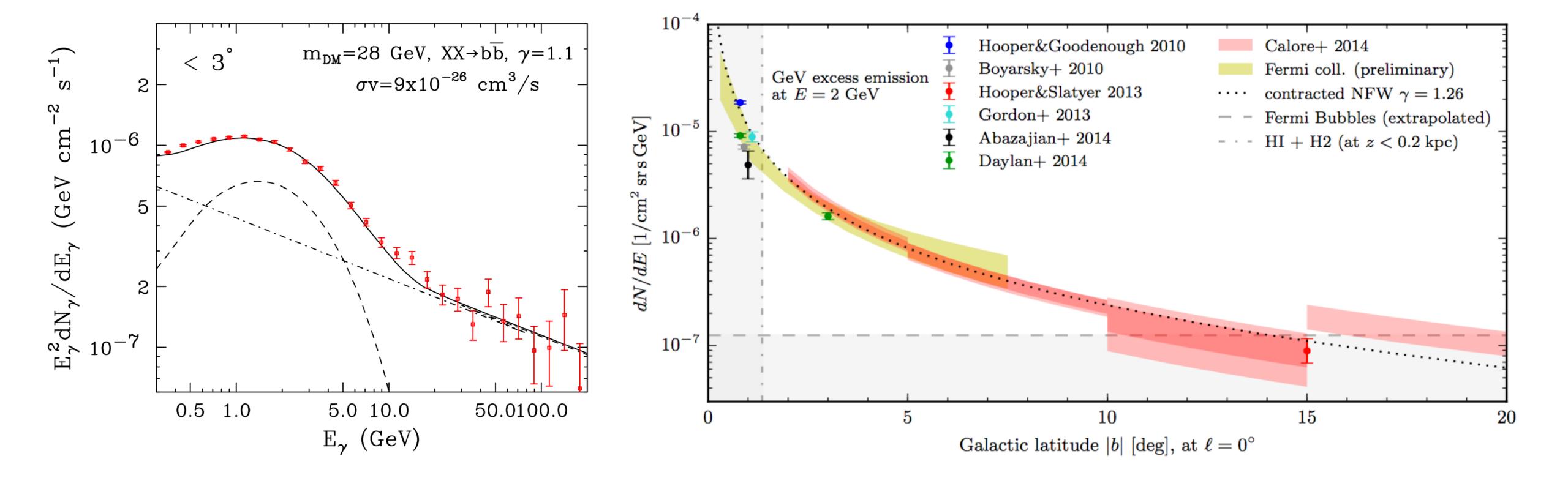


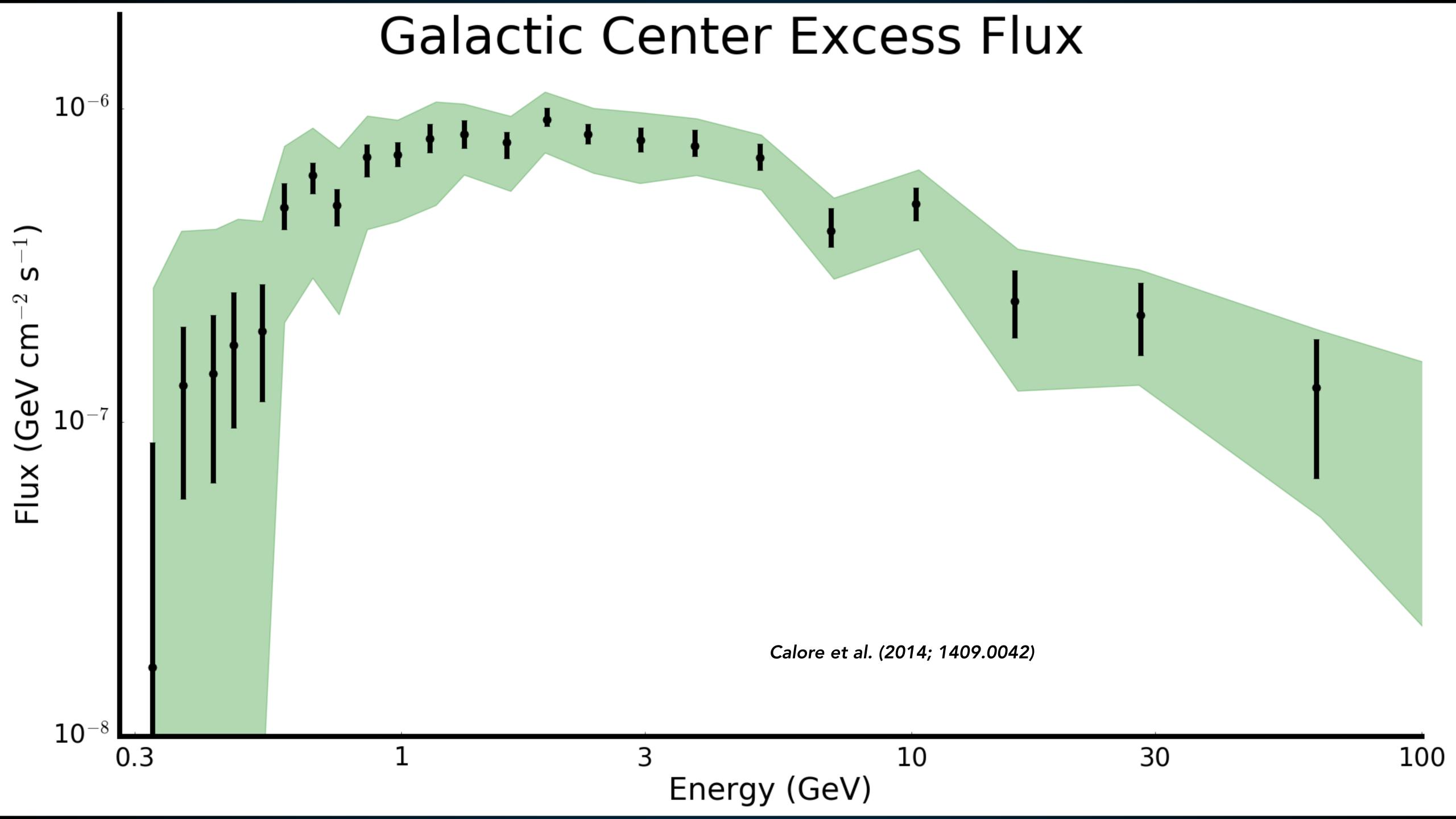


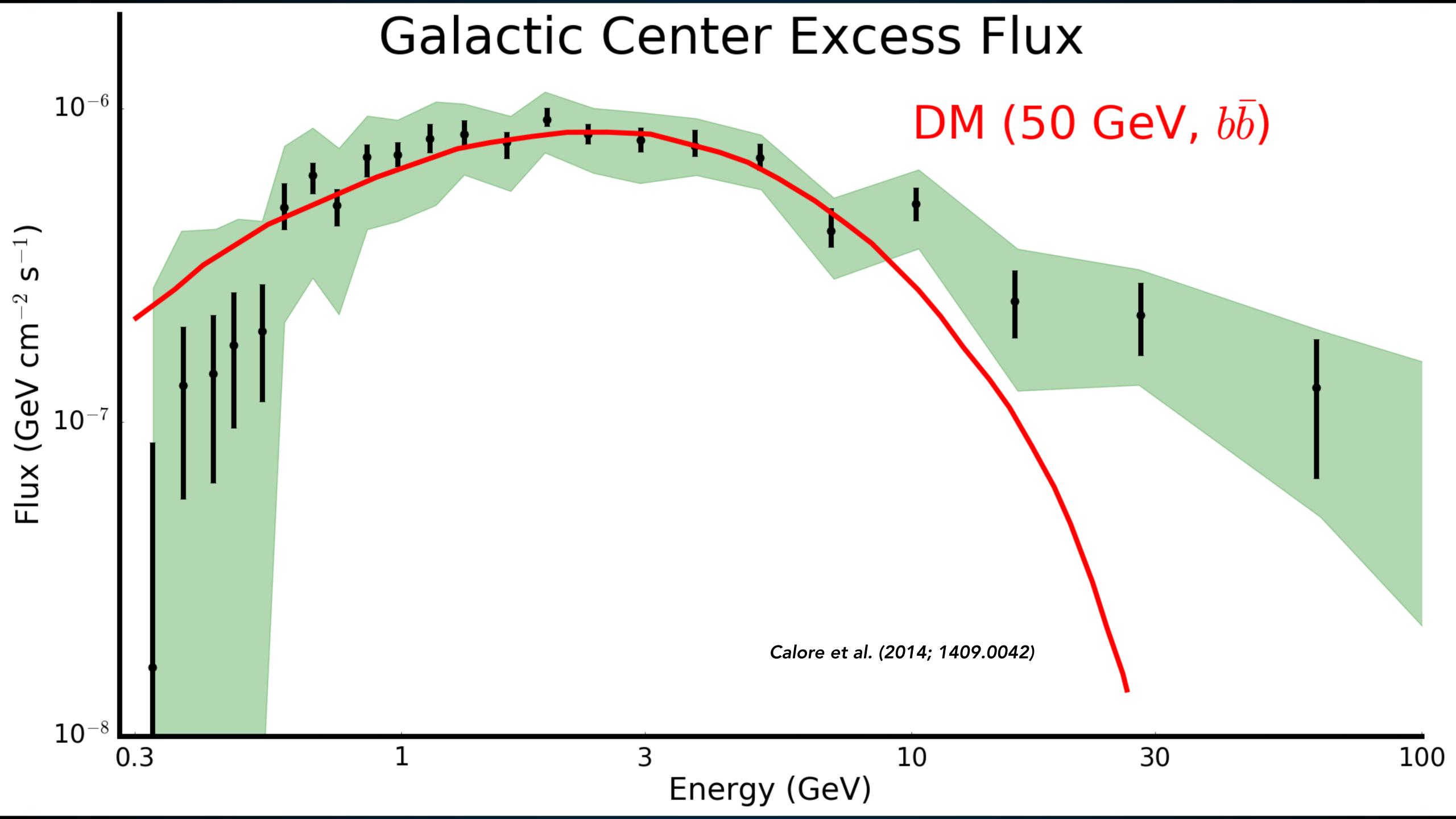
## Cosmic-Ray Transport

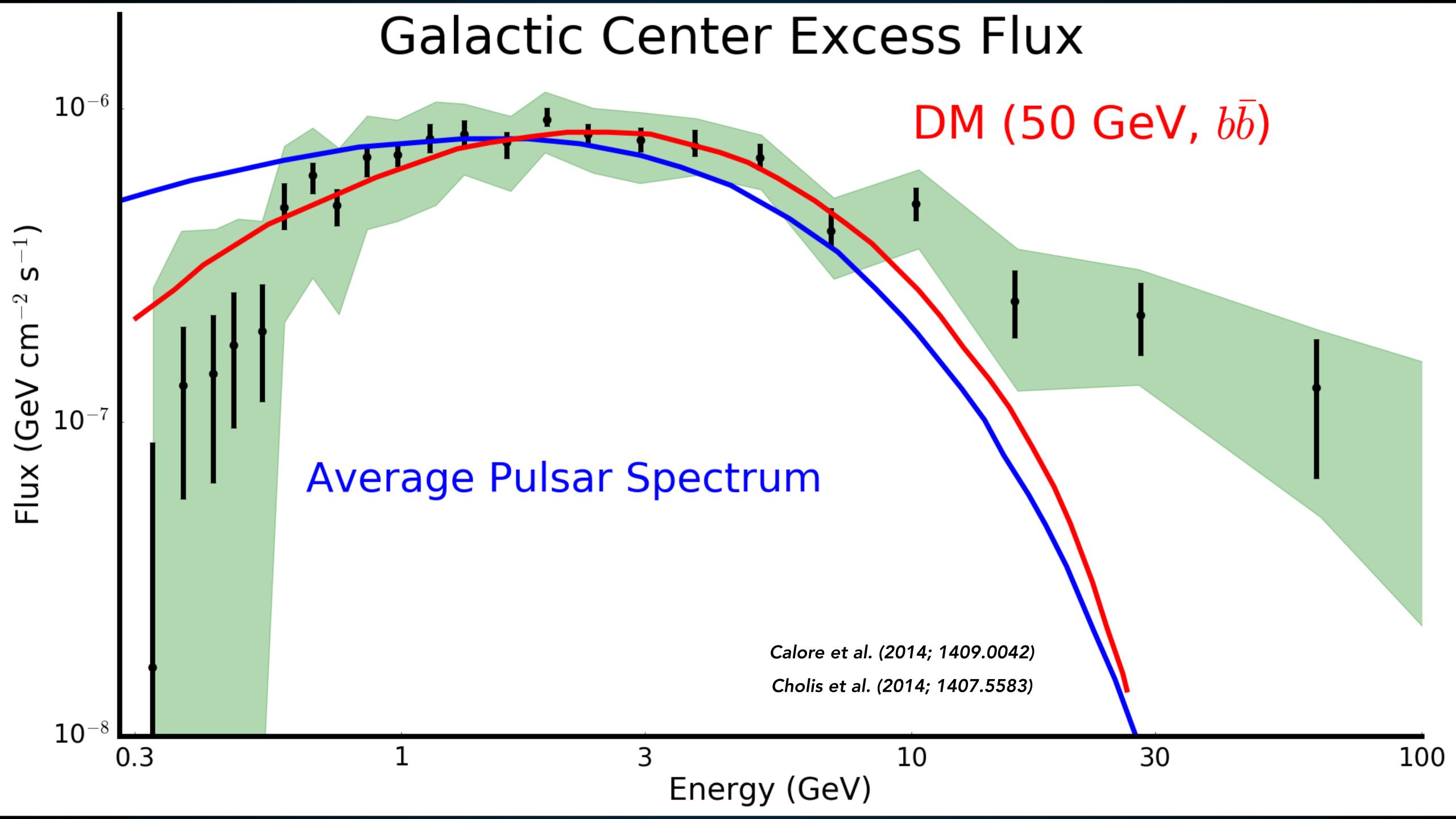
$$\frac{\overline{\partial \psi \left(\vec{r}, p, t\right)}}{\partial t} = \underbrace{\frac{\text{source}}{Q\left(\vec{r}, p, t\right)} + \overrightarrow{\nabla} \times (\overbrace{D_{xx}} \overrightarrow{\nabla} \psi)}_{\text{source}} - \underbrace{\frac{\text{convection}}{\overrightarrow{V} \psi}}_{\text{convection}} + \underbrace{\frac{\text{re-acceleration}}{\partial p} p^2 D_{pp} \frac{\partial}{\partial p^2} \psi}_{\text{convection}} - \underbrace{\frac{\partial}{\partial p} \left(\dot{p} \psi - \frac{p}{3} \left(\overrightarrow{\nabla} \times \overrightarrow{V}\right) \psi\right)}_{\text{energy loss}} - \underbrace{\frac{\psi}{\tau_f}}_{\text{ragmentation}} - \underbrace{\frac{\psi}{\tau_f}}_{\text{radioactive decay}}$$

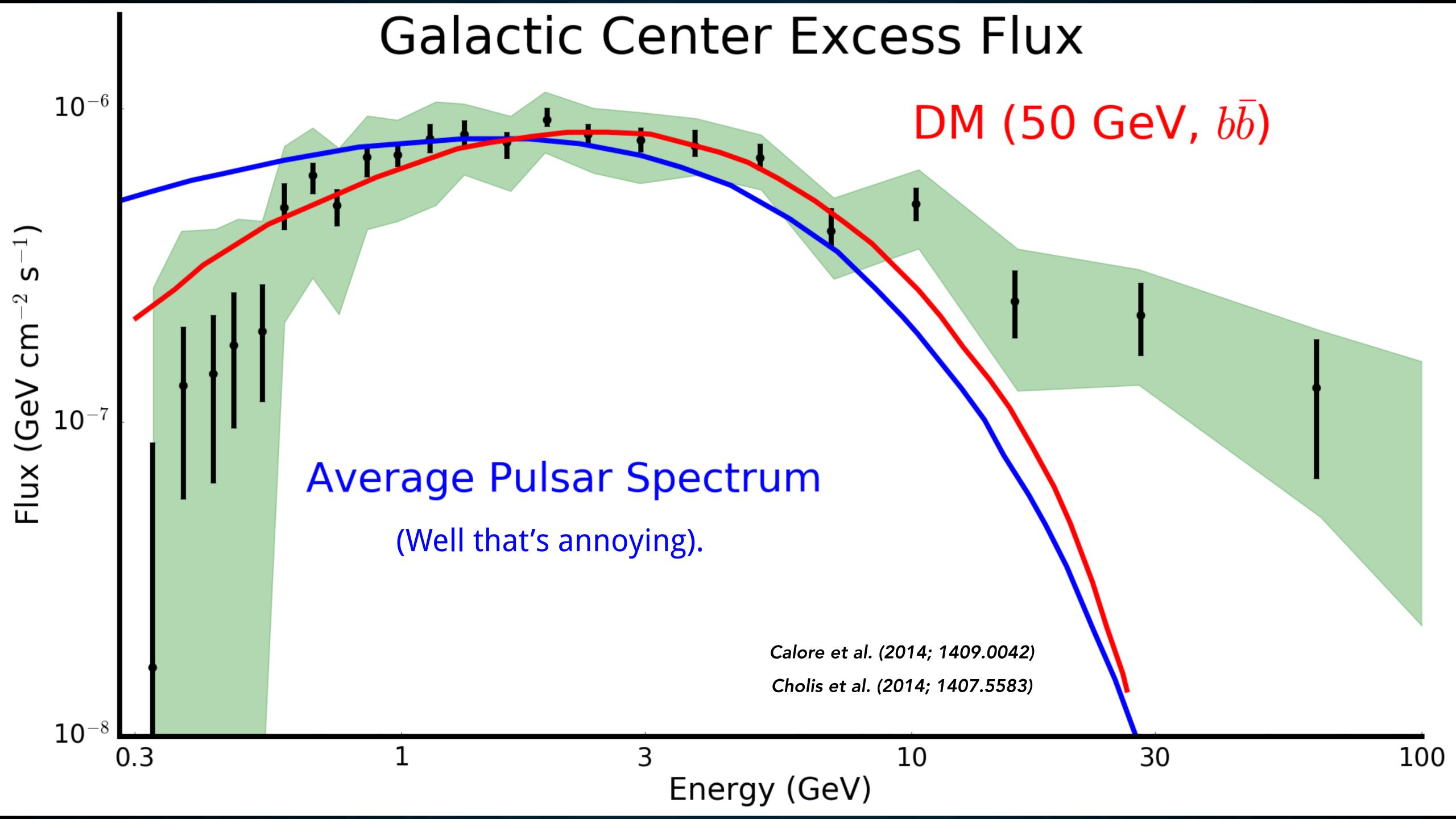
- Convection: Winds driven by injection of cosmic-rays and relativistic gas from the Milky Way
- Energy Losses: (Next Section)
- Fragmentation: Nuclei can be split by interactions
- Radioactive Decay: For radioactive nuclei

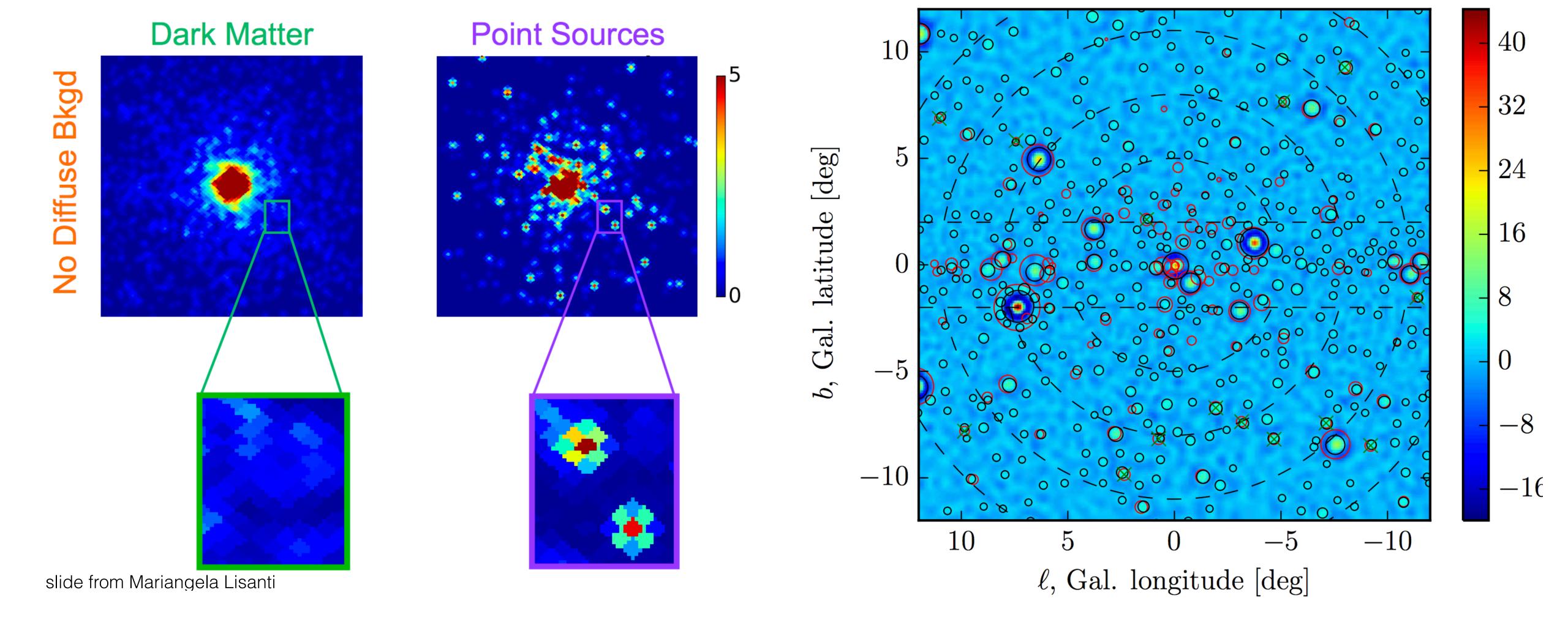


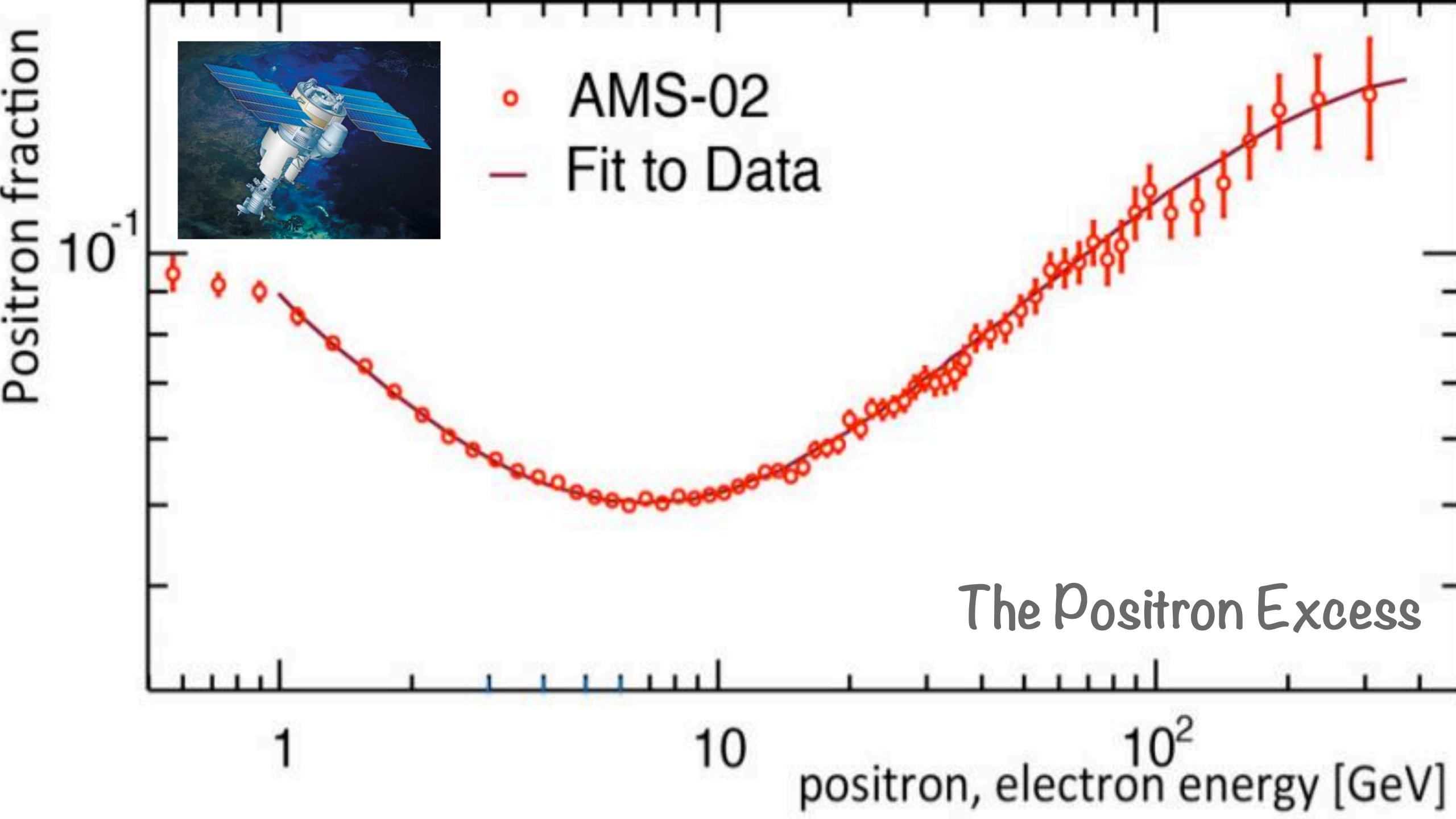


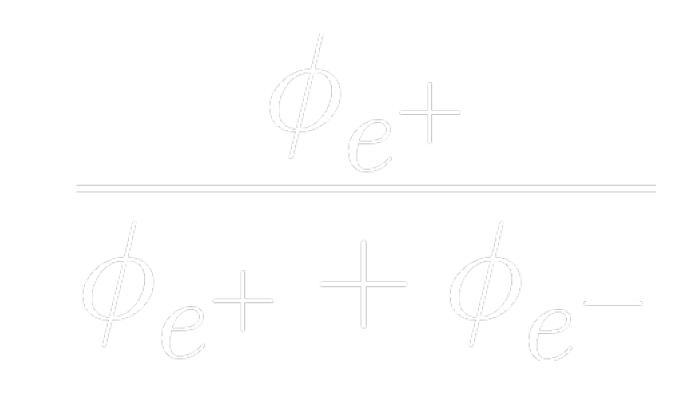


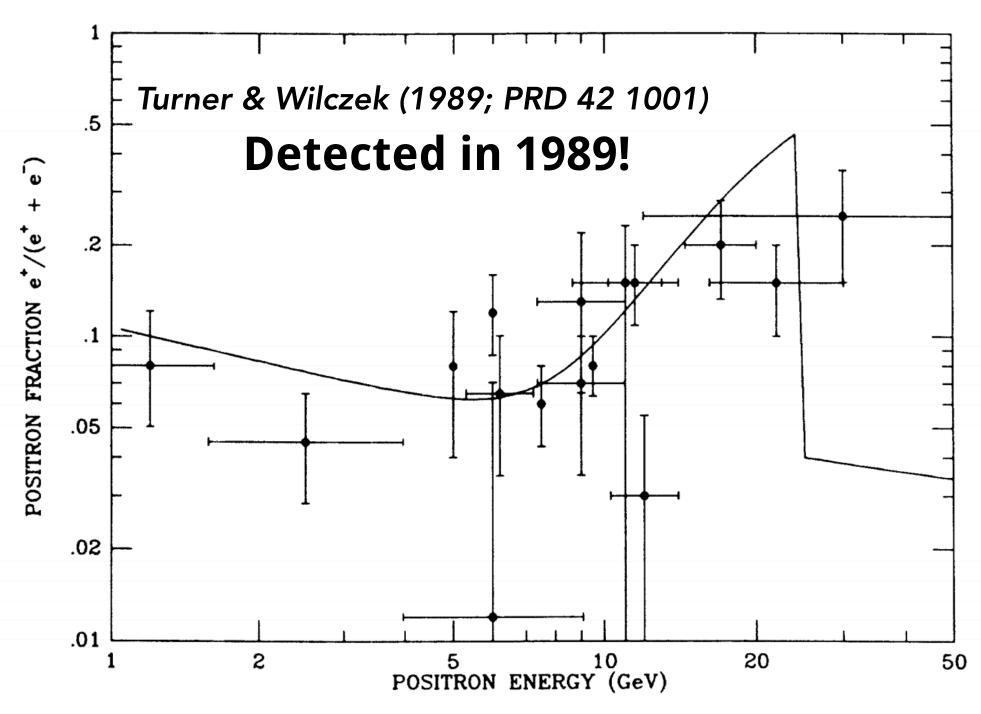


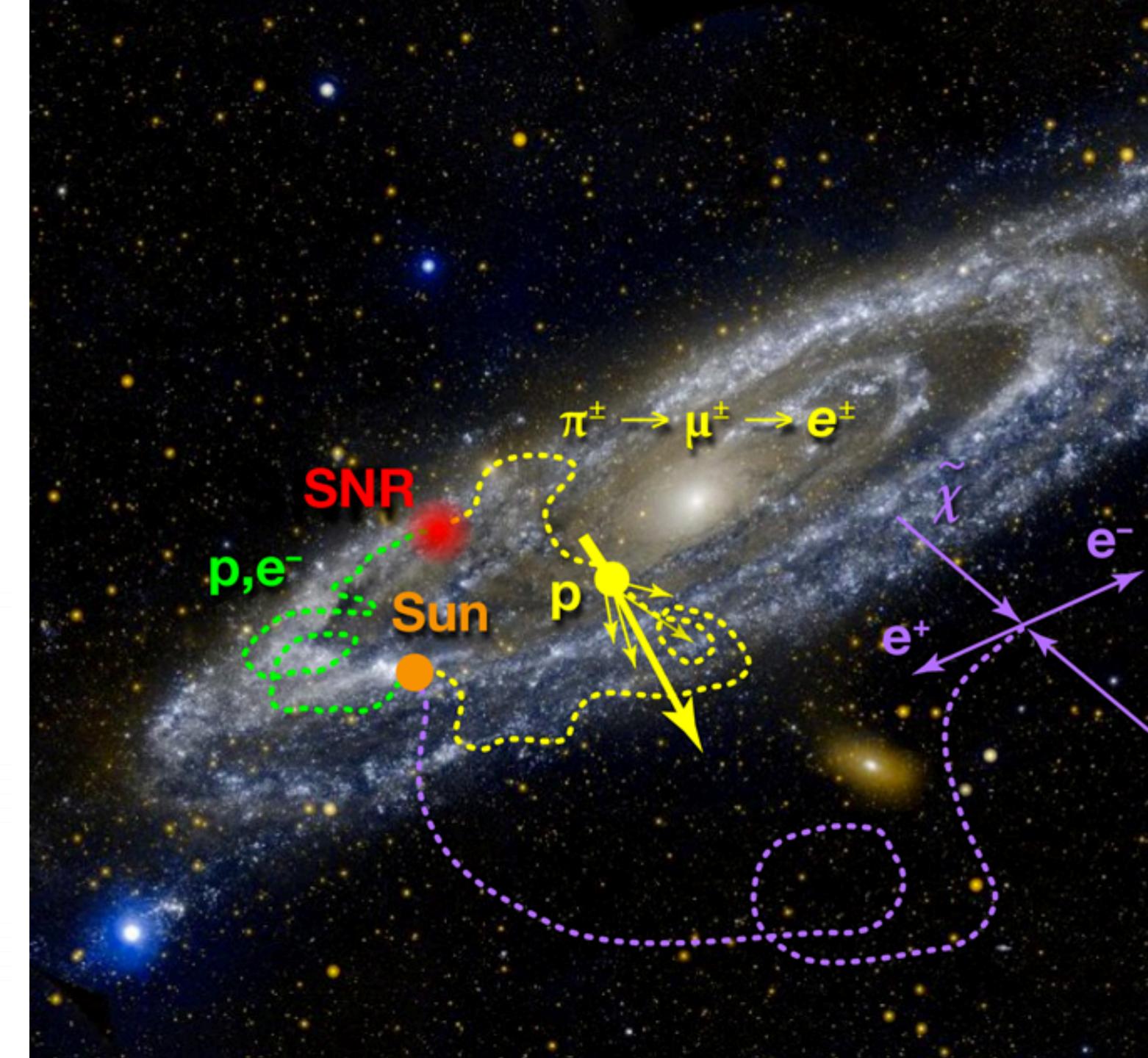


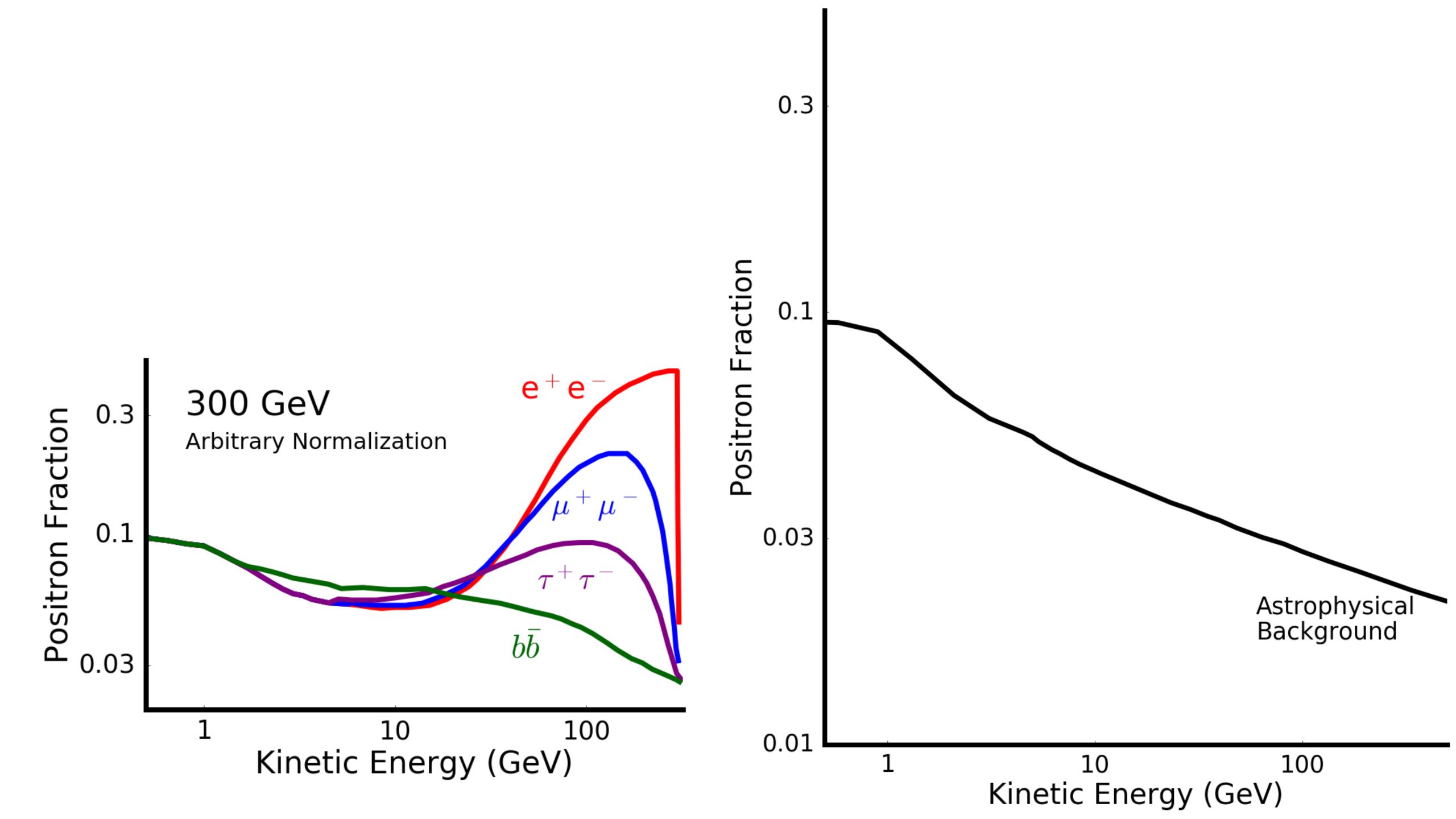


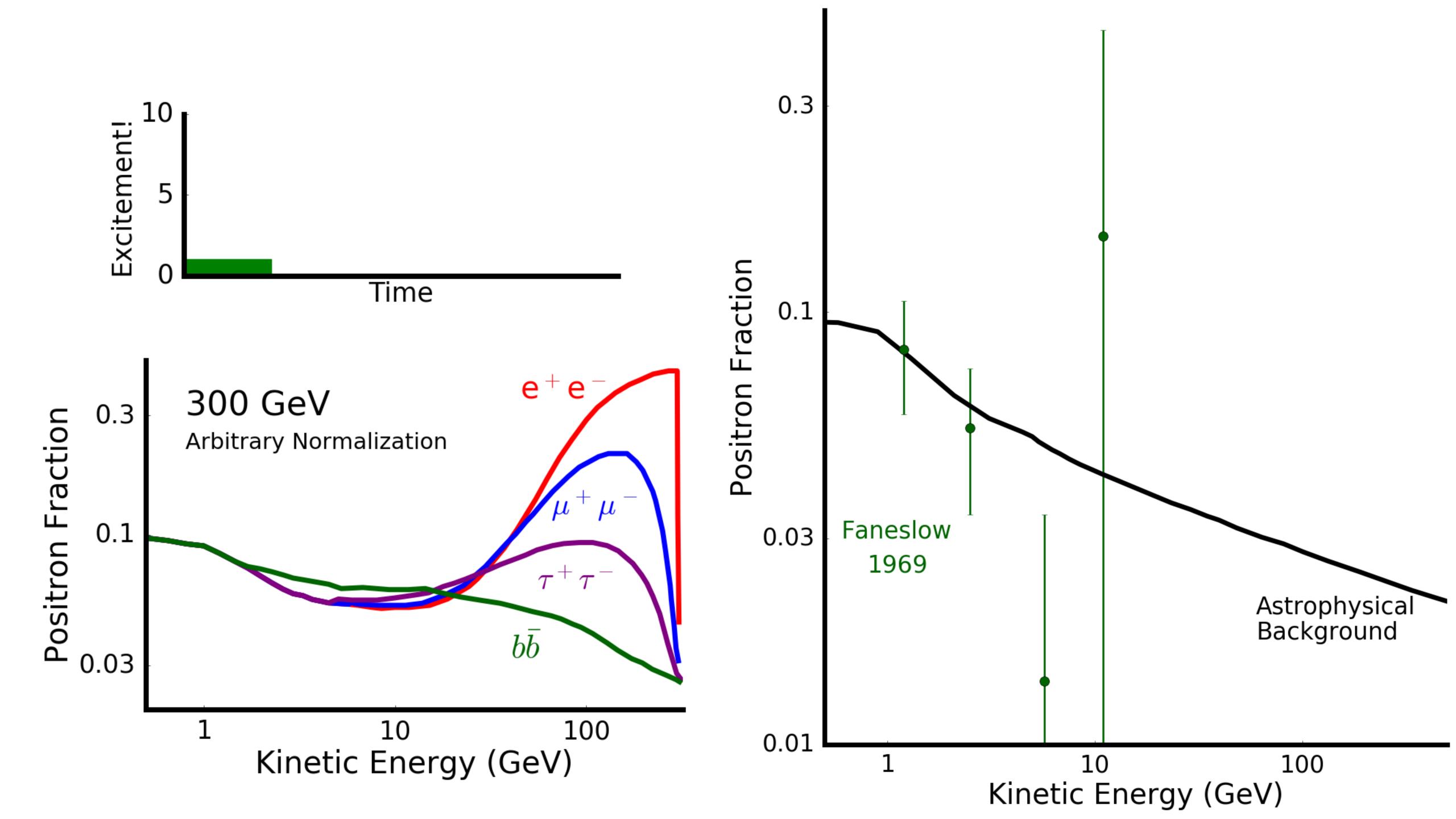


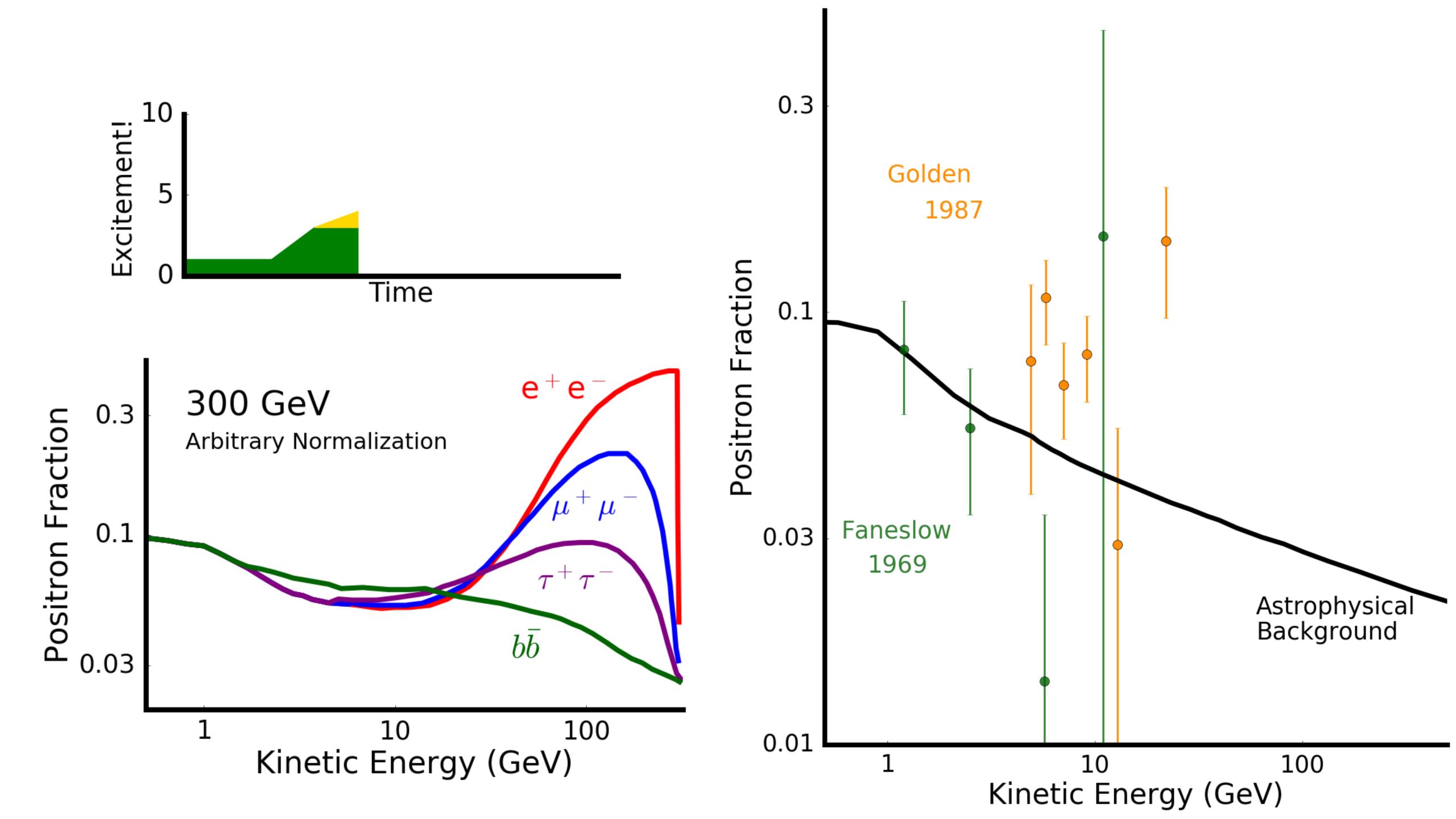


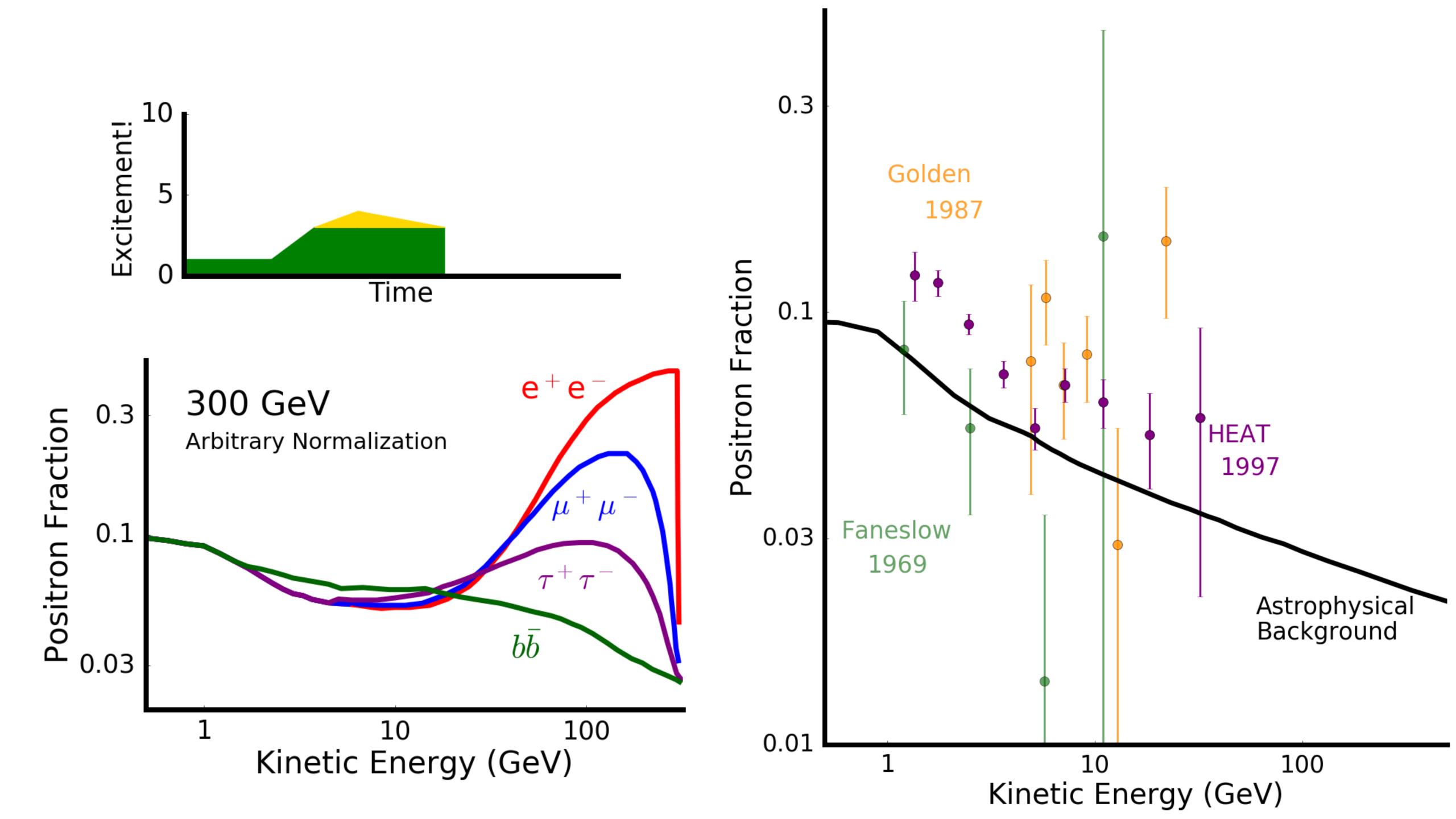


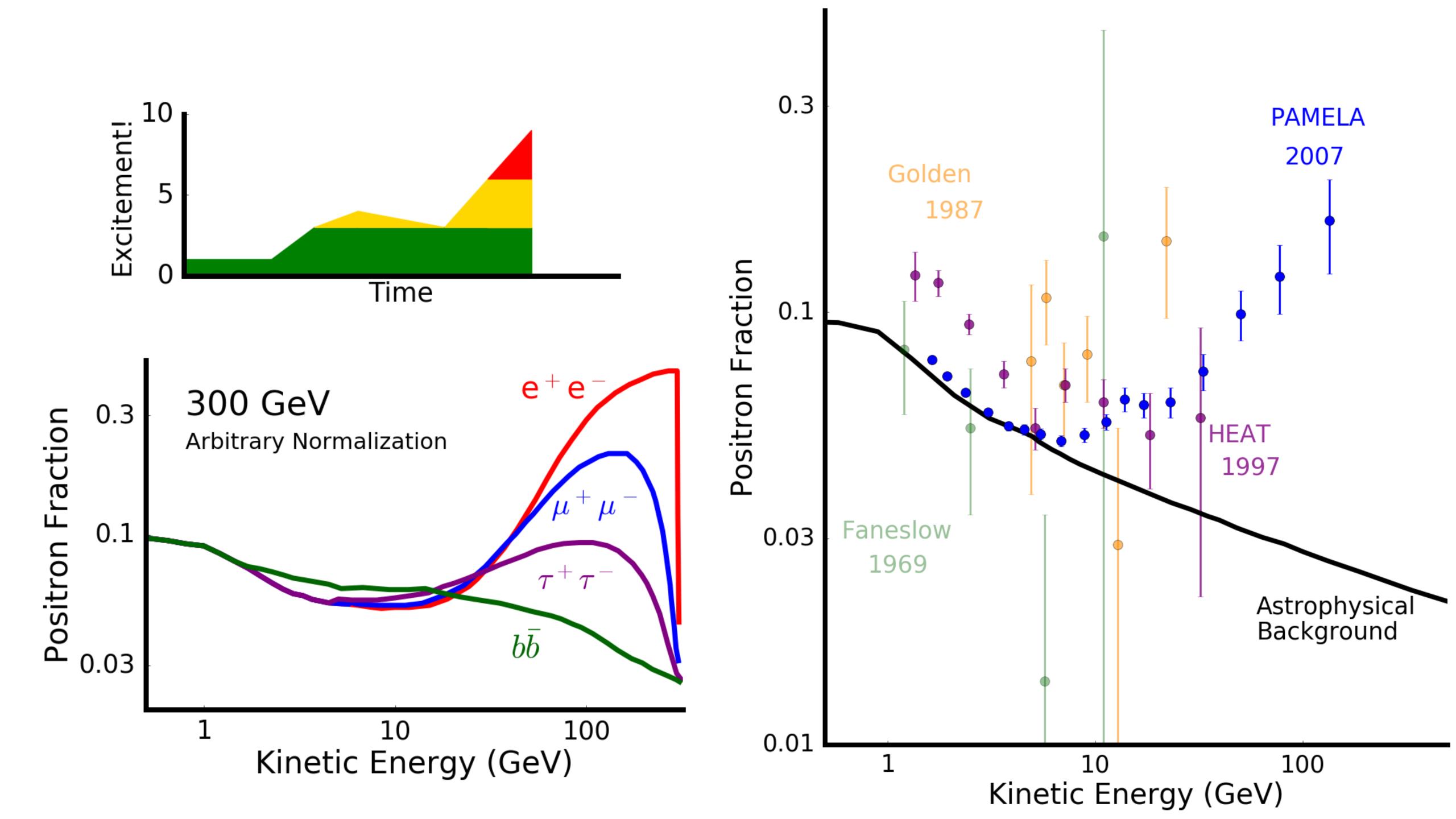


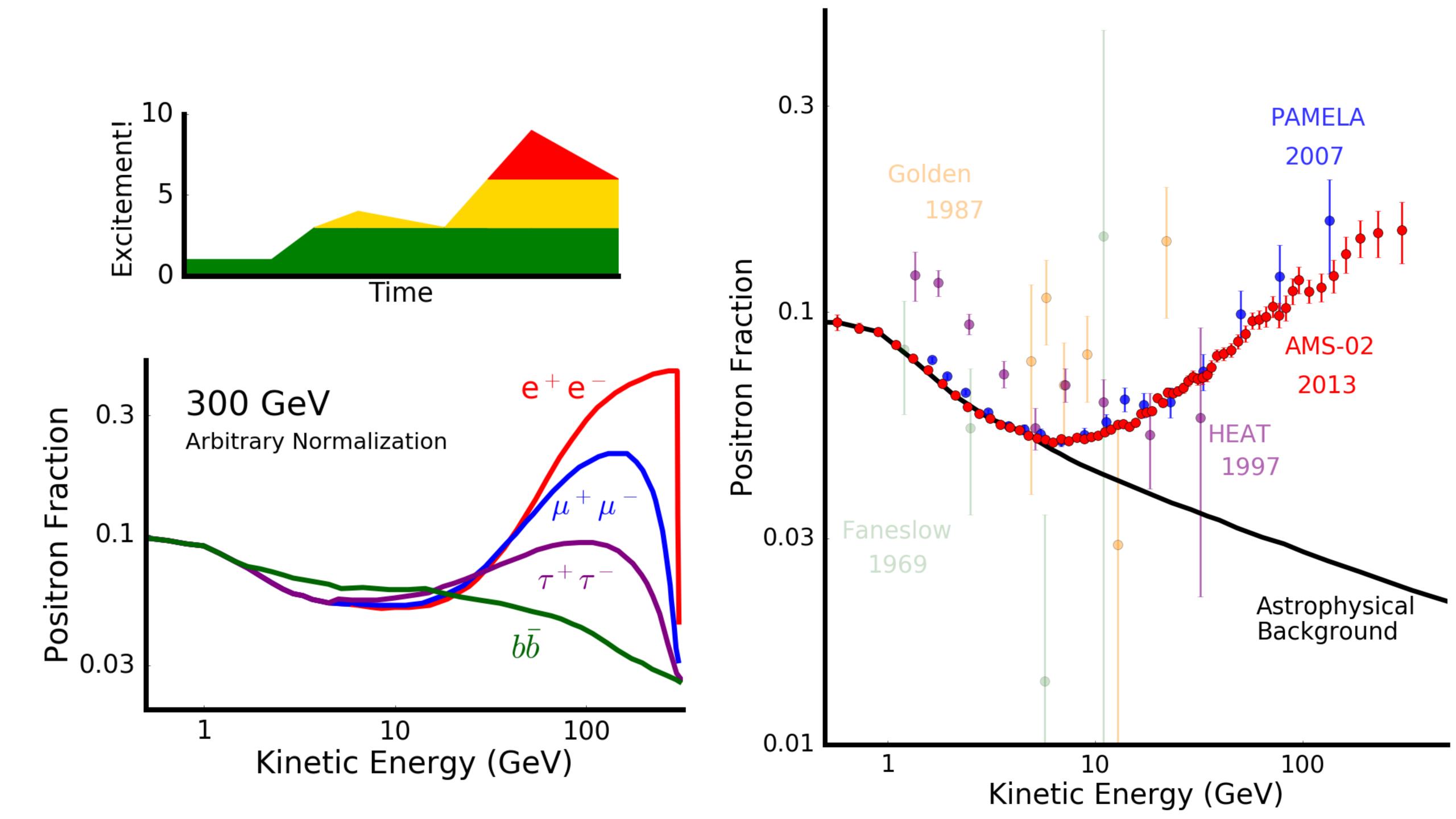


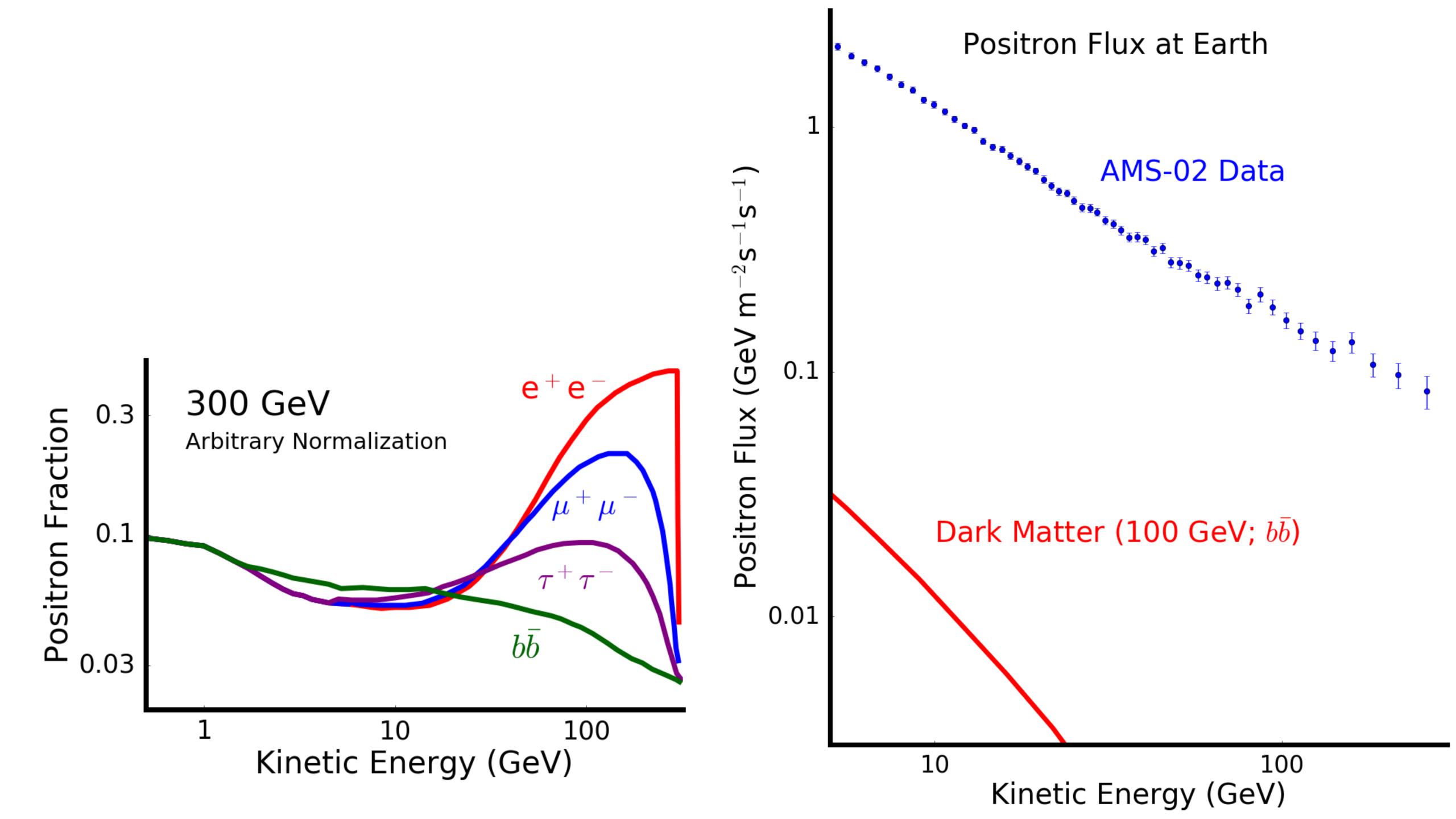


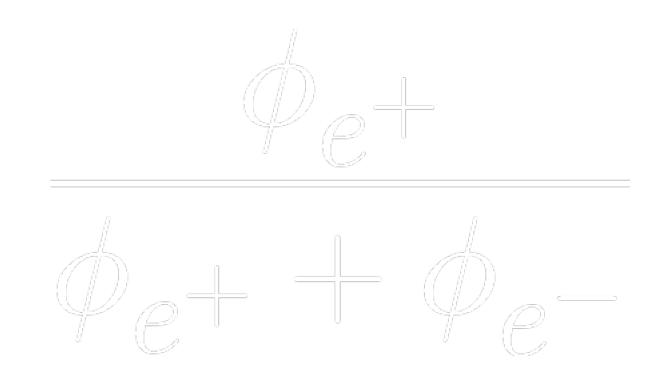


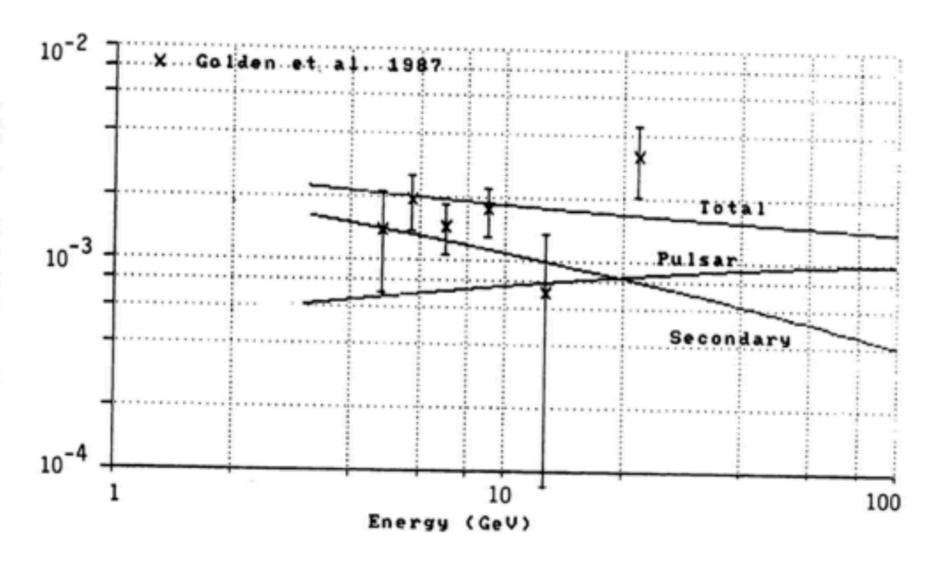


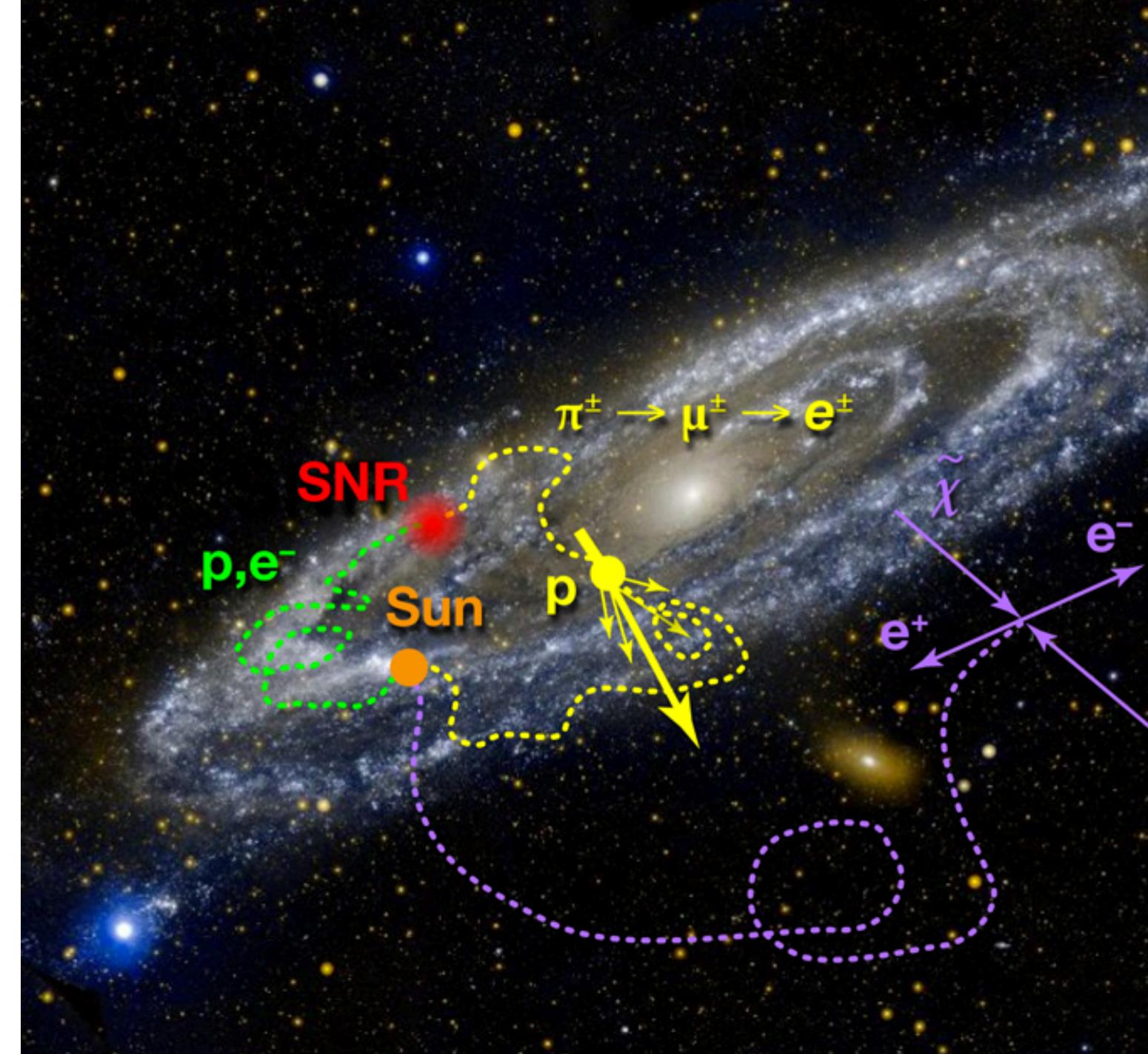


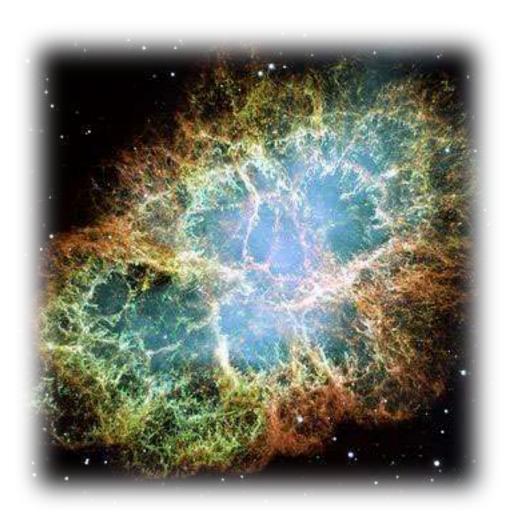


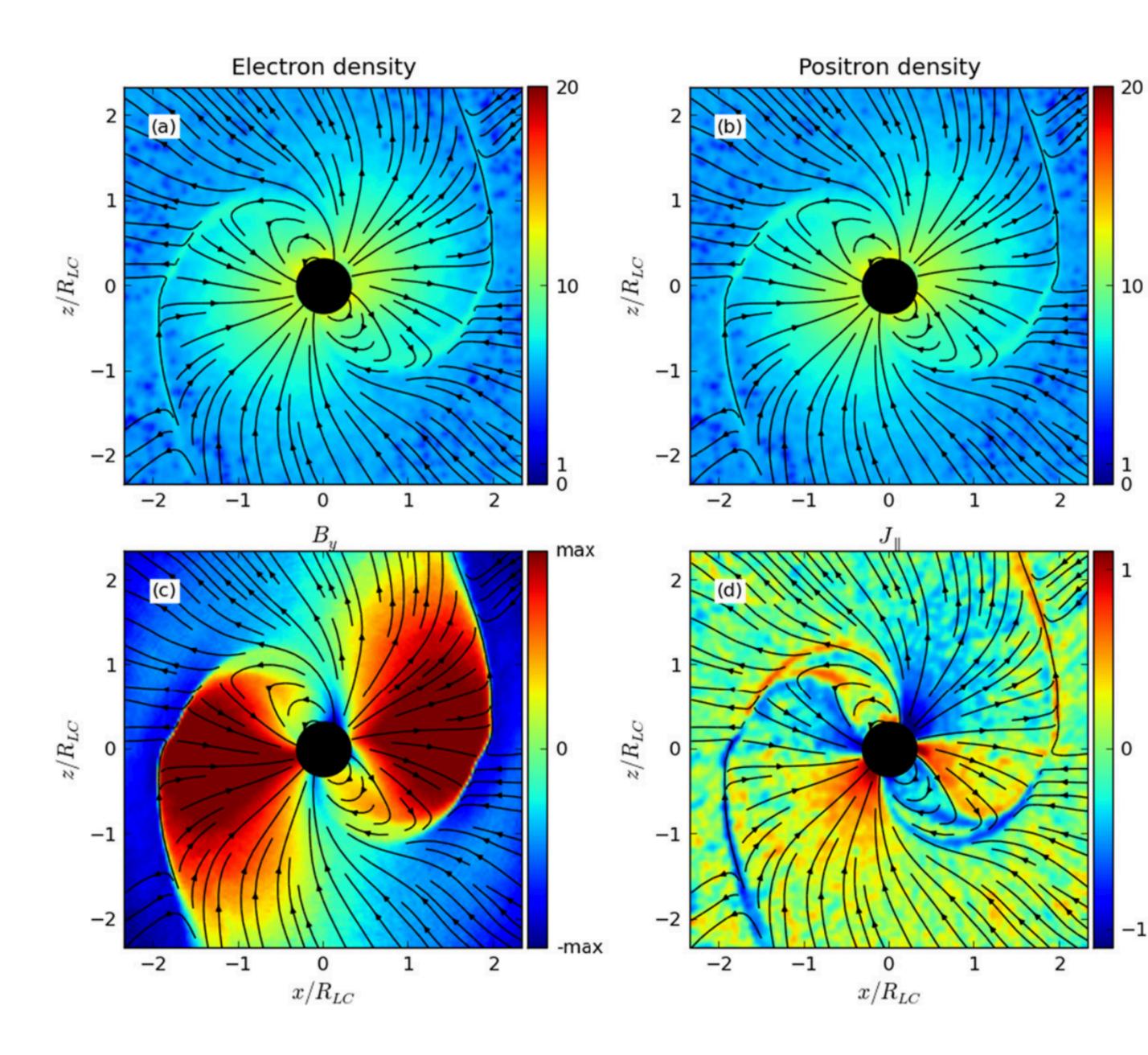


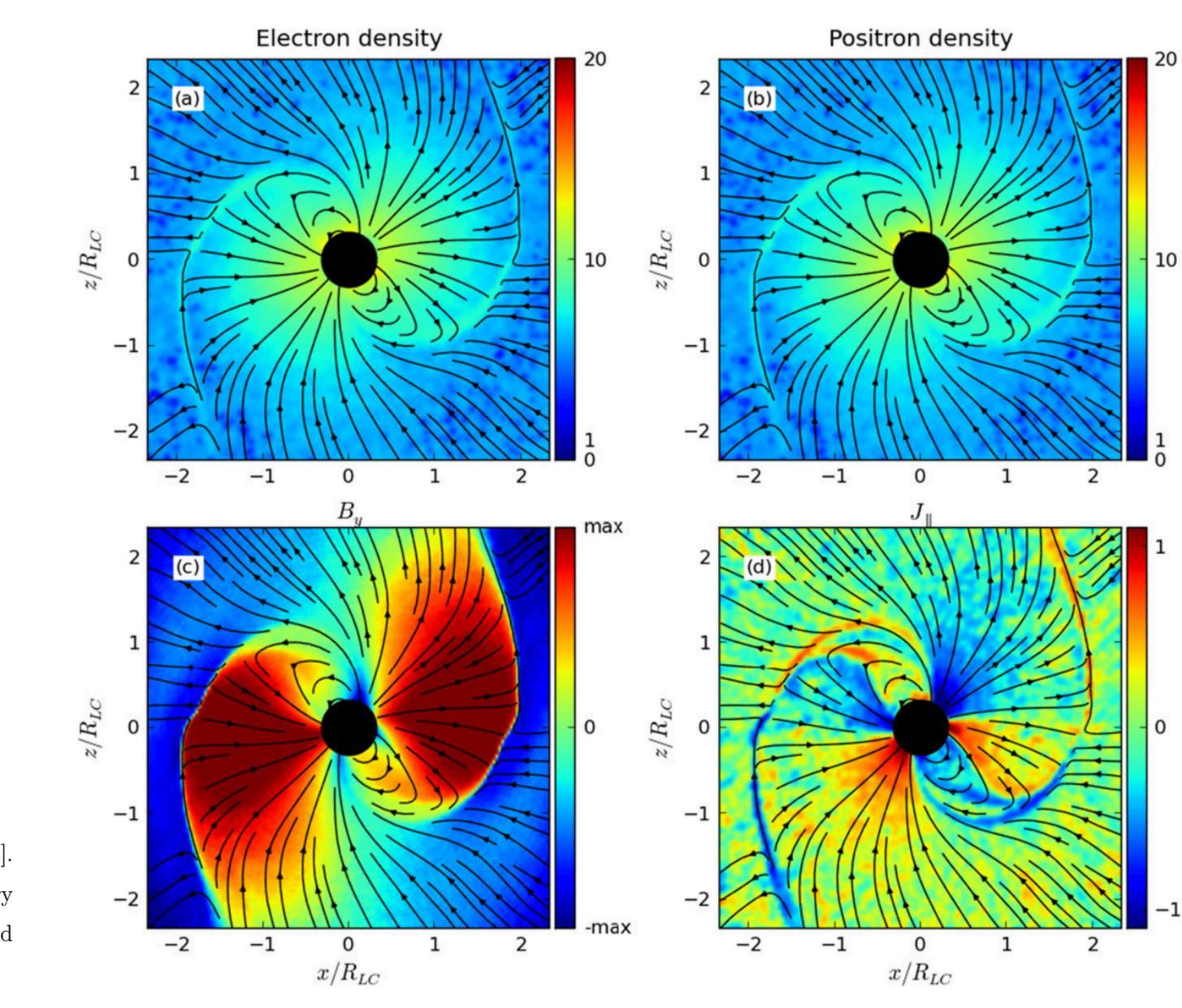












%. A quantitative discussion of plausible values for  $f_{e^{\pm}}$  was recently given in Ref. [38]. We shall not review their discussion here, but Ref. [38] argues (see in particular their very informative App. B and C) that in the context of a standard model for the pulsar wind nebulae, a reasonable range for  $f_{e^{\pm}}$  falls between 1% and 30%.



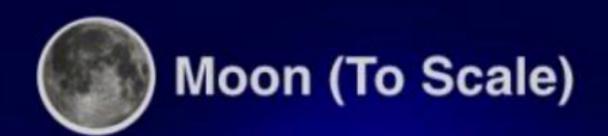
Geminga

SNR (hadronic/leptonic)

TeV Halo (escaped e<sup>+</sup>e<sup>-</sup>)

PWN (confined e<sup>+</sup>e<sup>-</sup>)



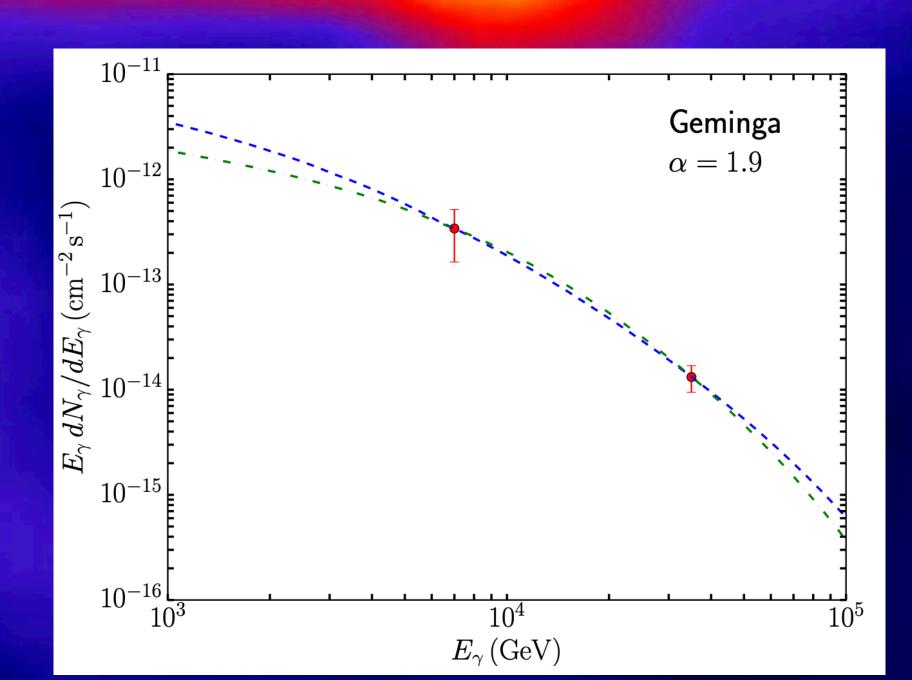


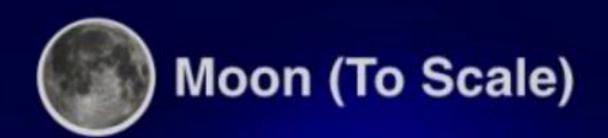
Geminga

SNR (hadronic/leptonic)

TeV Halo (escaped e<sup>+</sup>e<sup>-</sup>)

PWN (confined e<sup>+</sup>e<sup>-</sup>)



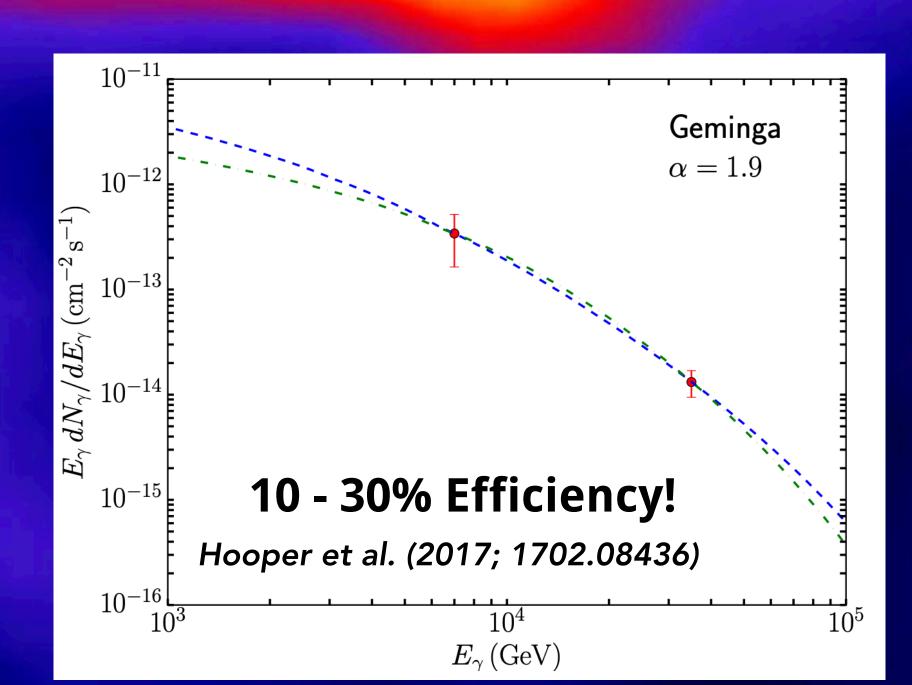


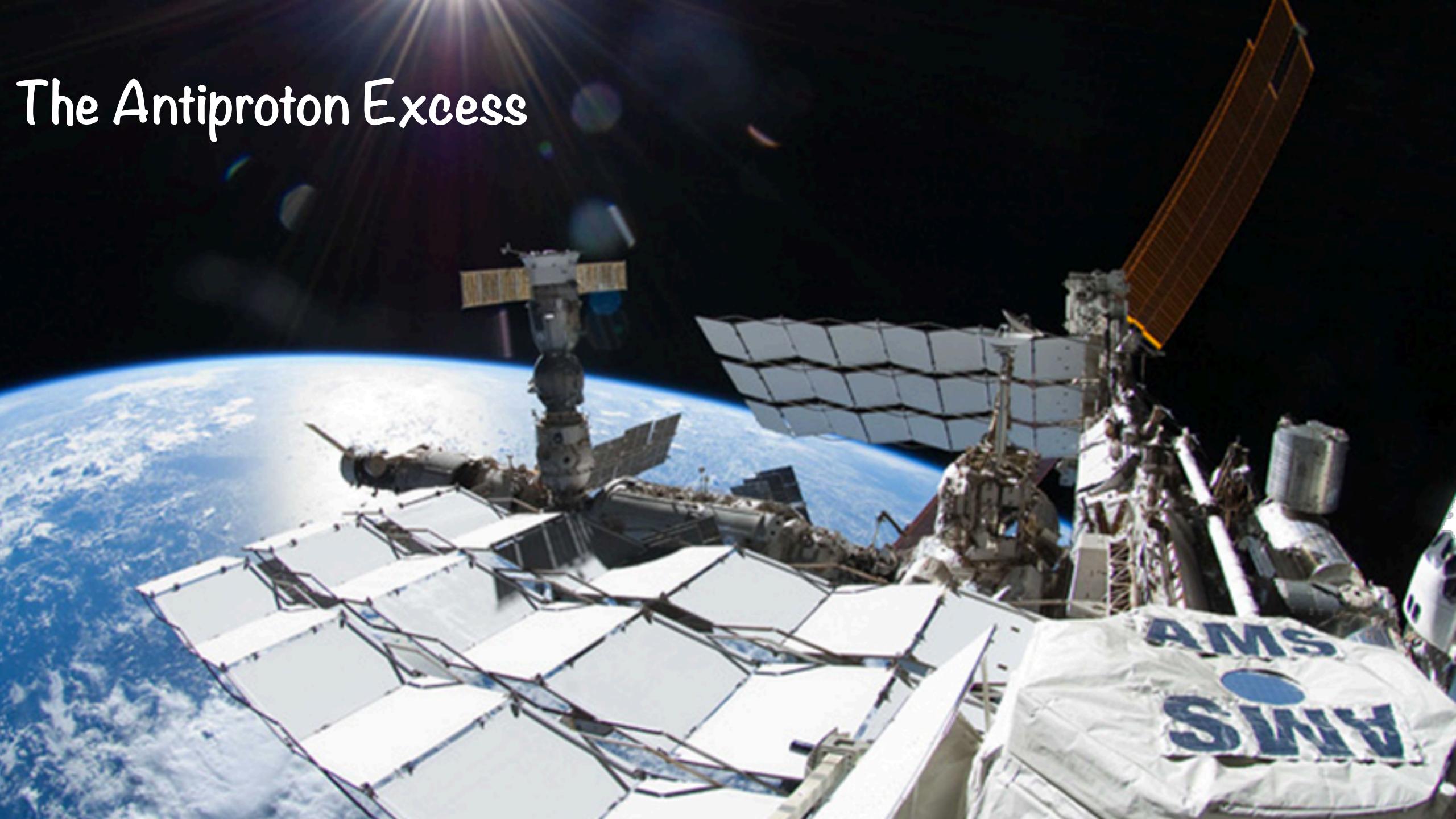
Geminga

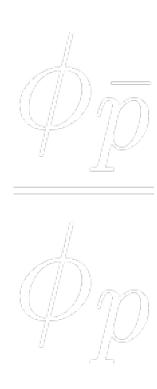
SNR (hadronic/leptonic)

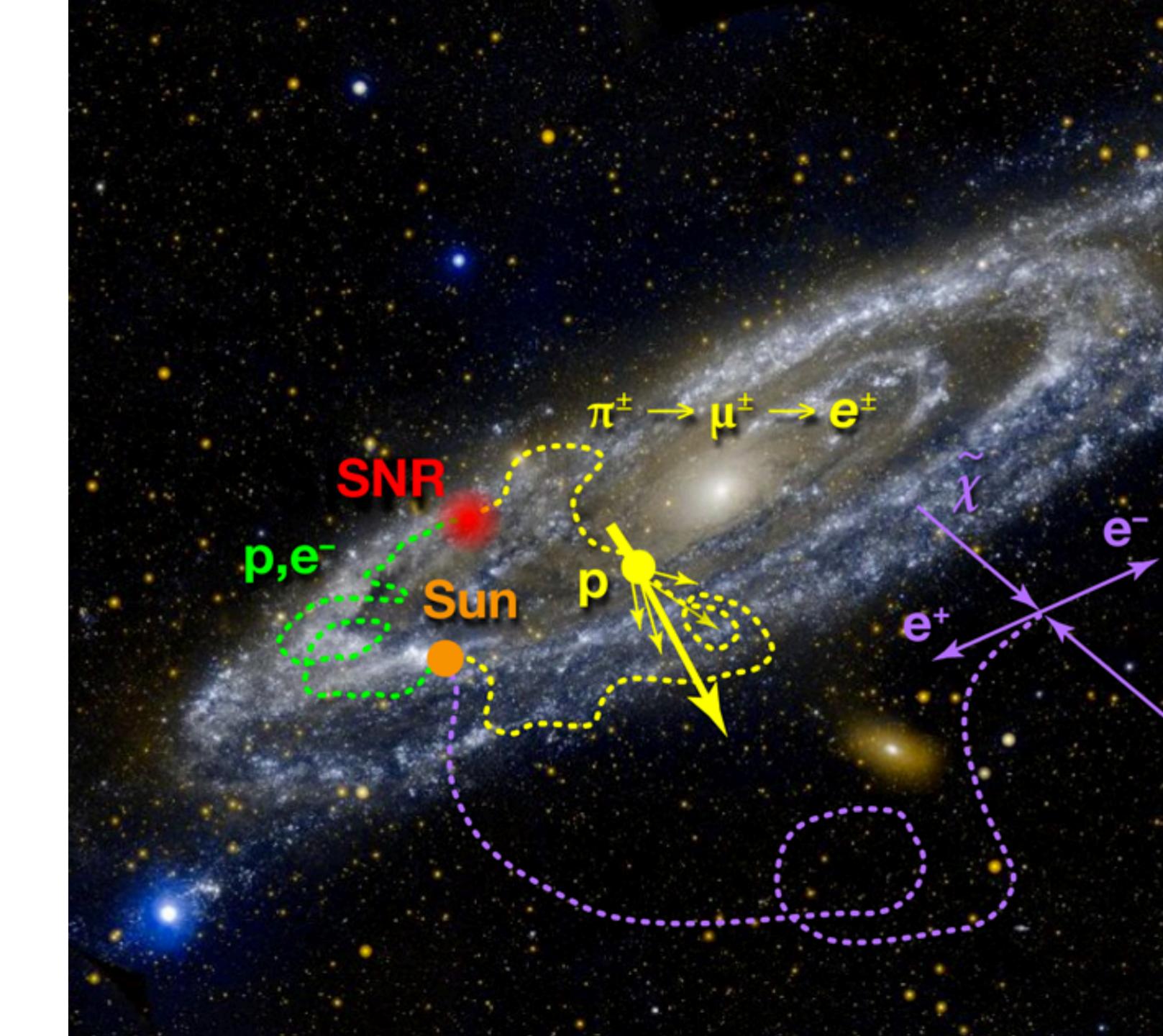
TeV Halo (escaped e<sup>+</sup>e<sup>-</sup>)

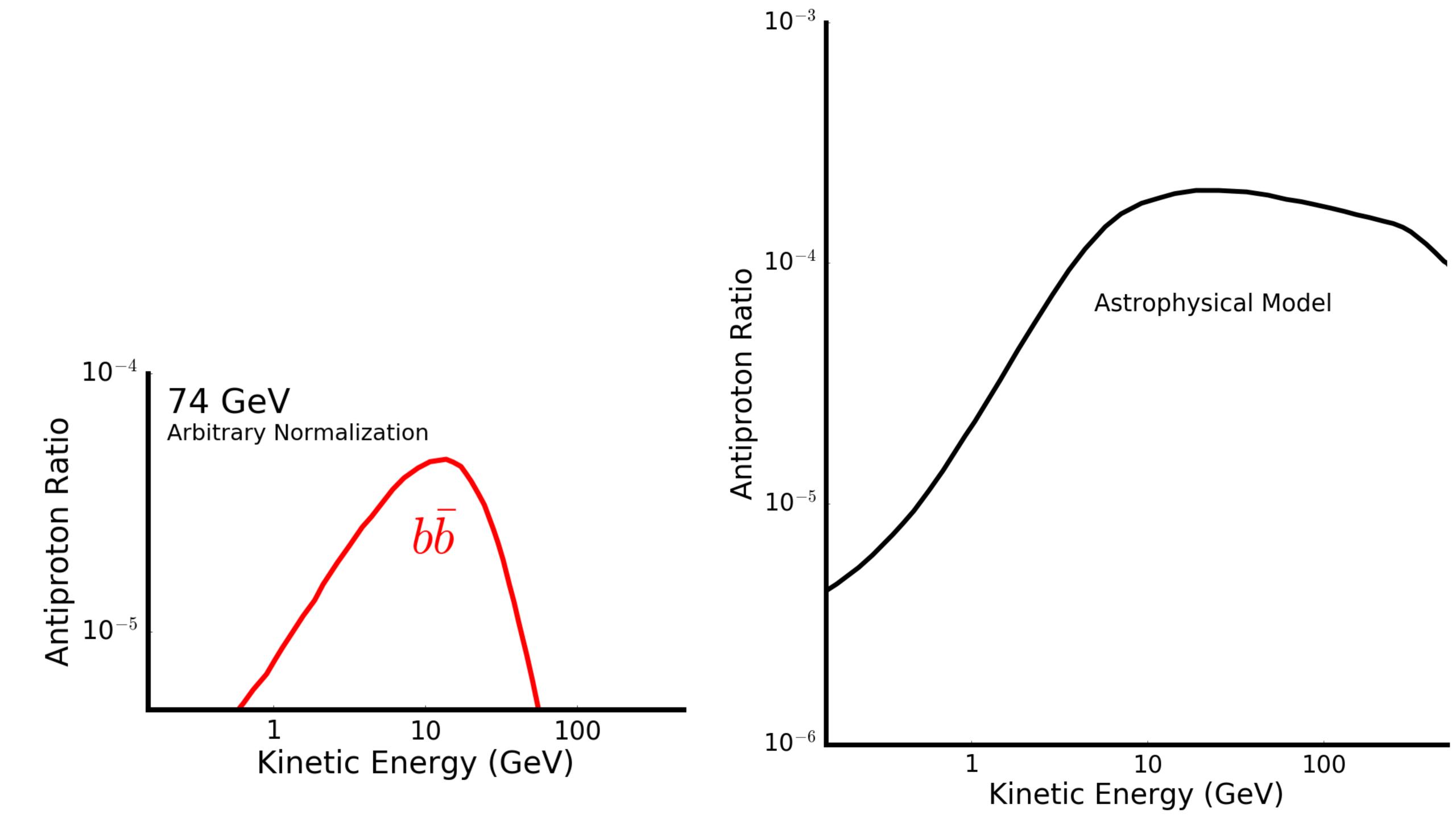
PWN (confined e<sup>+</sup>e<sup>-</sup>)

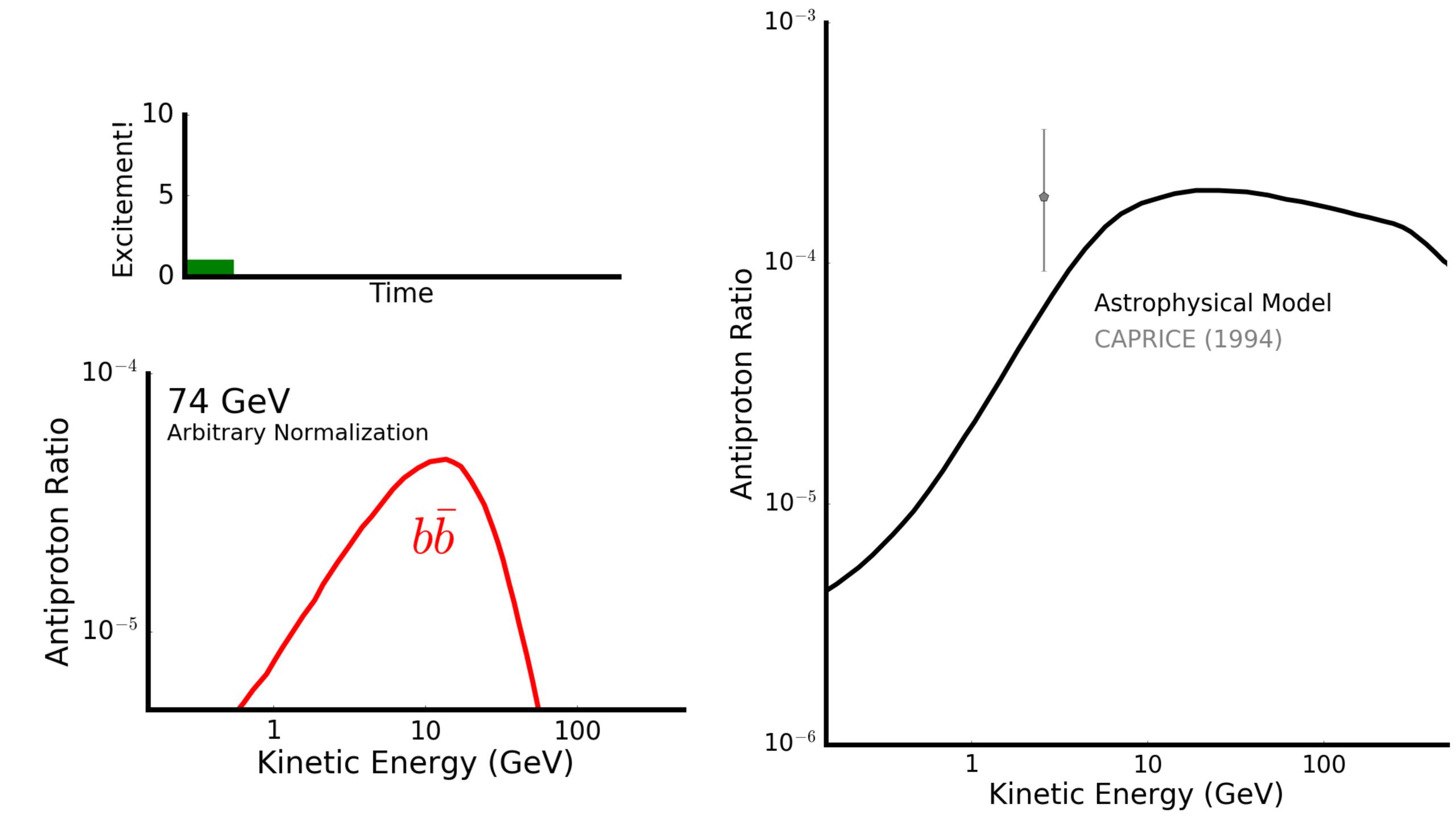


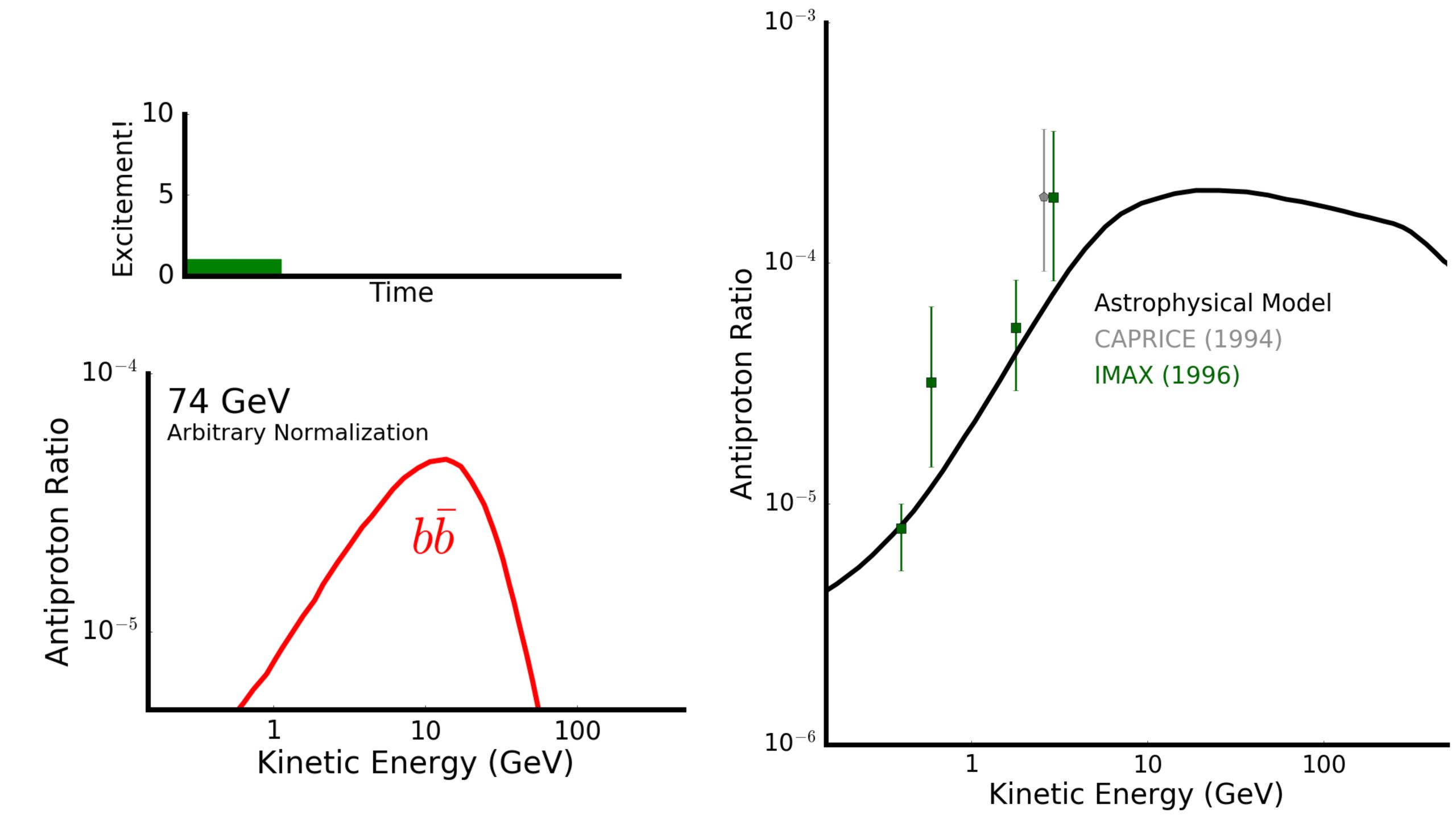


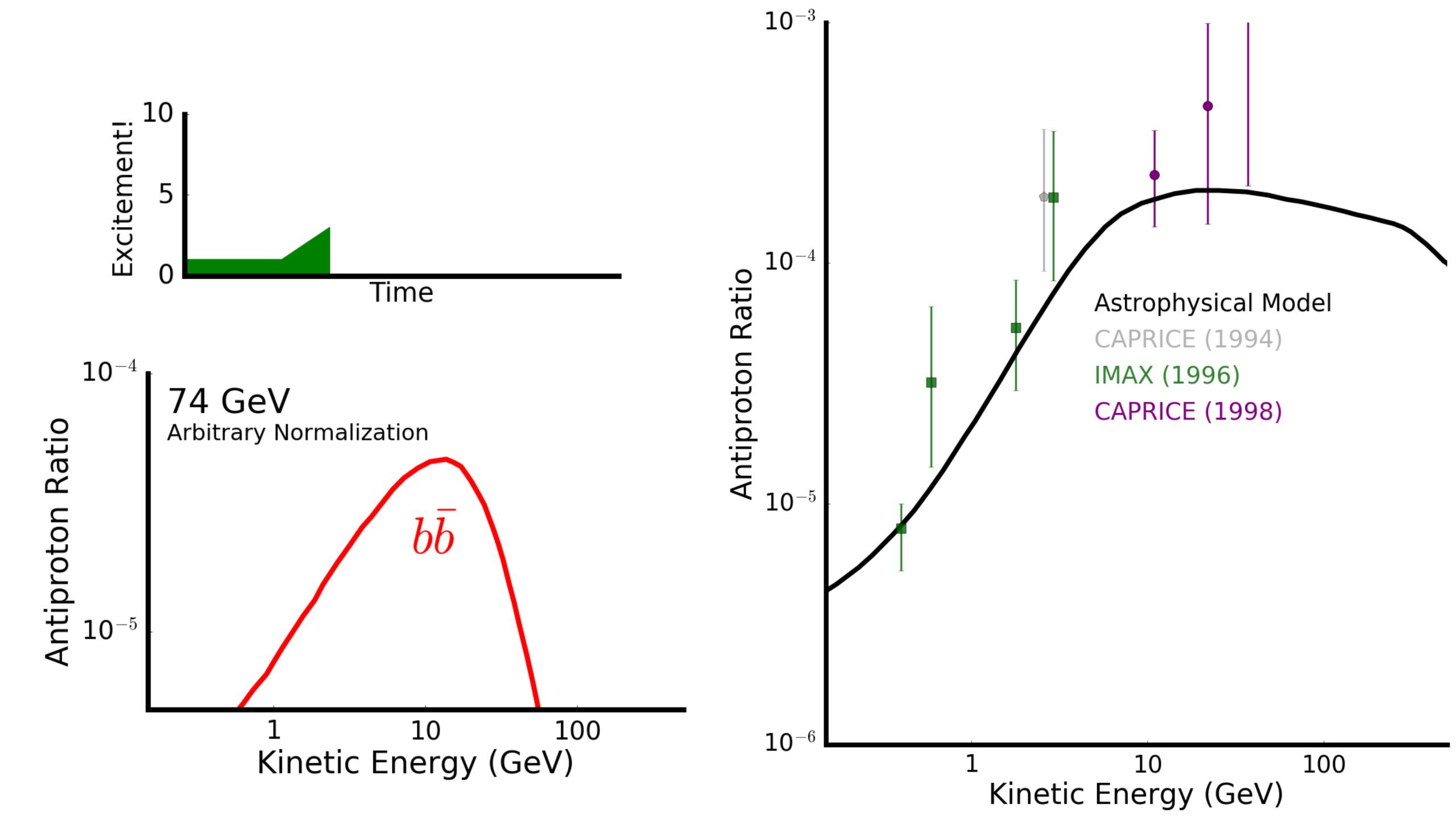


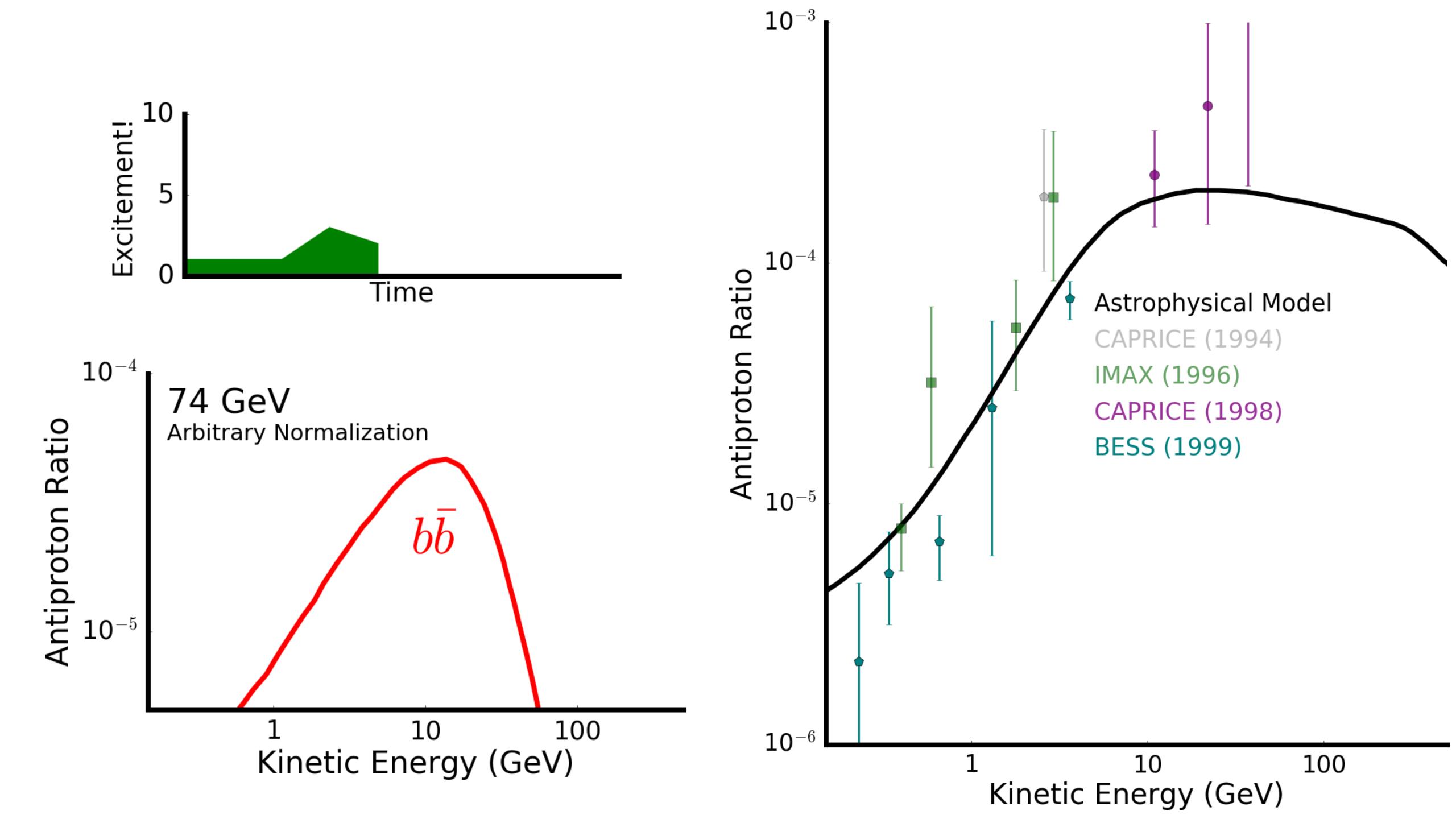


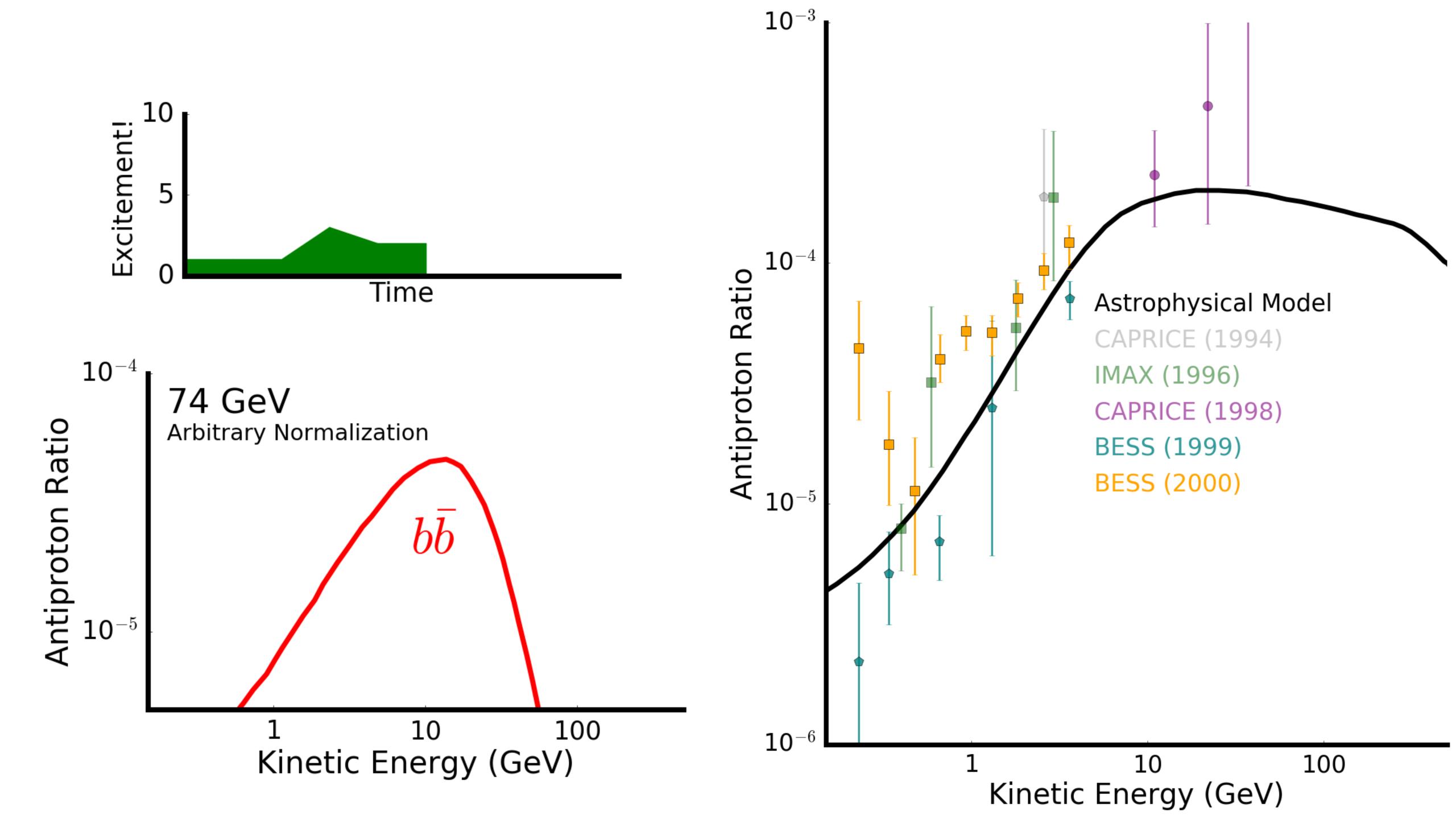


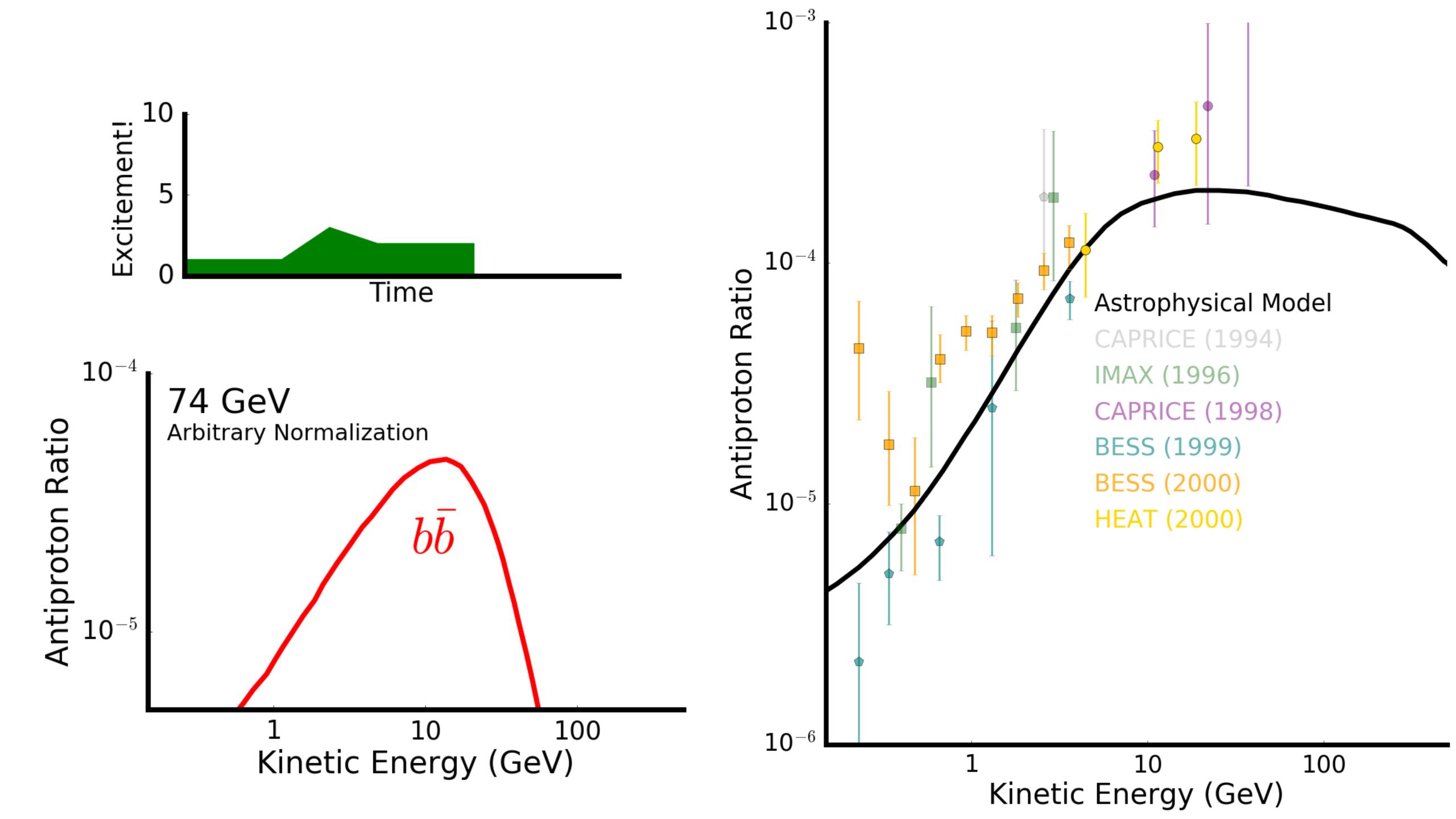


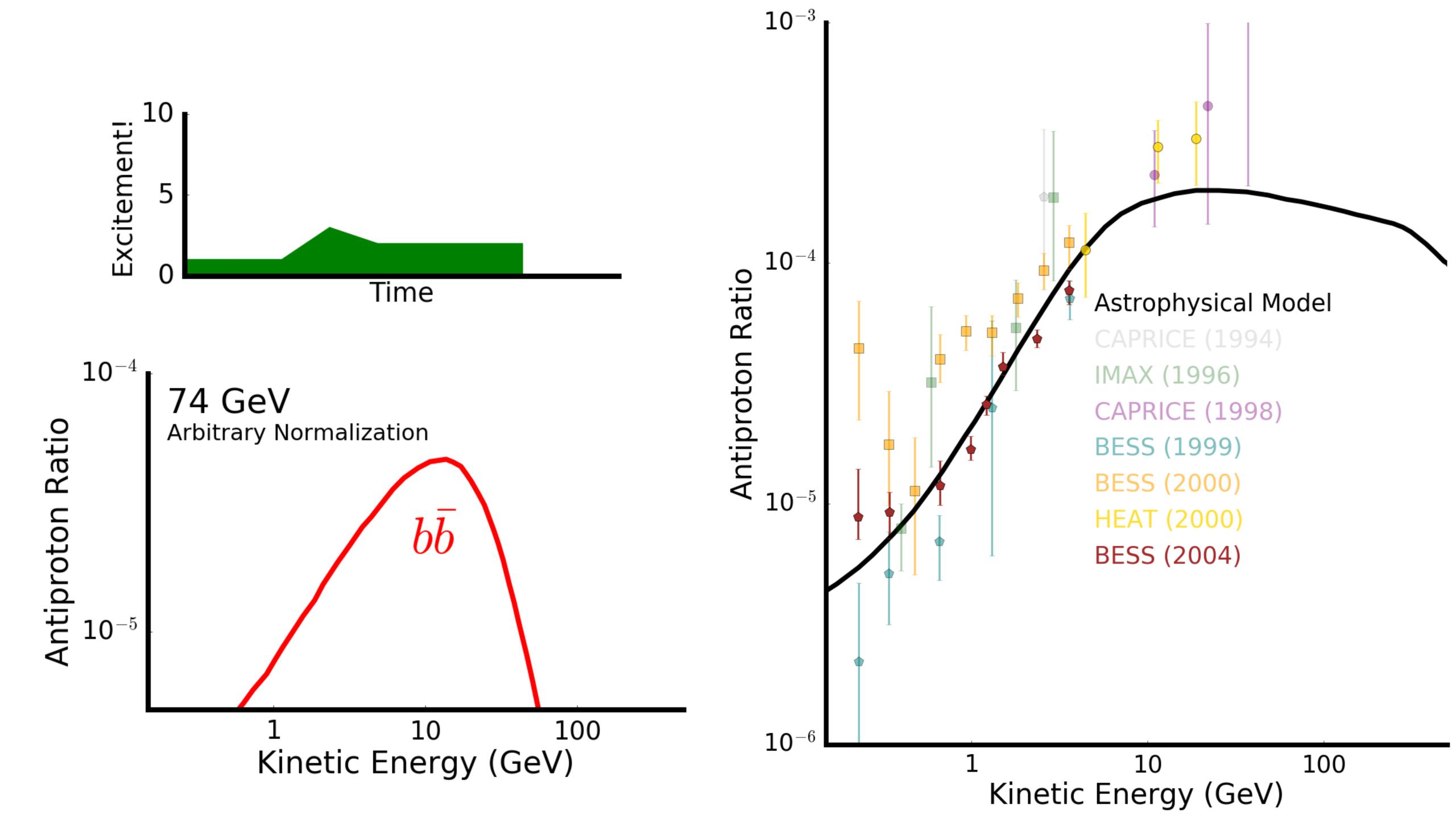


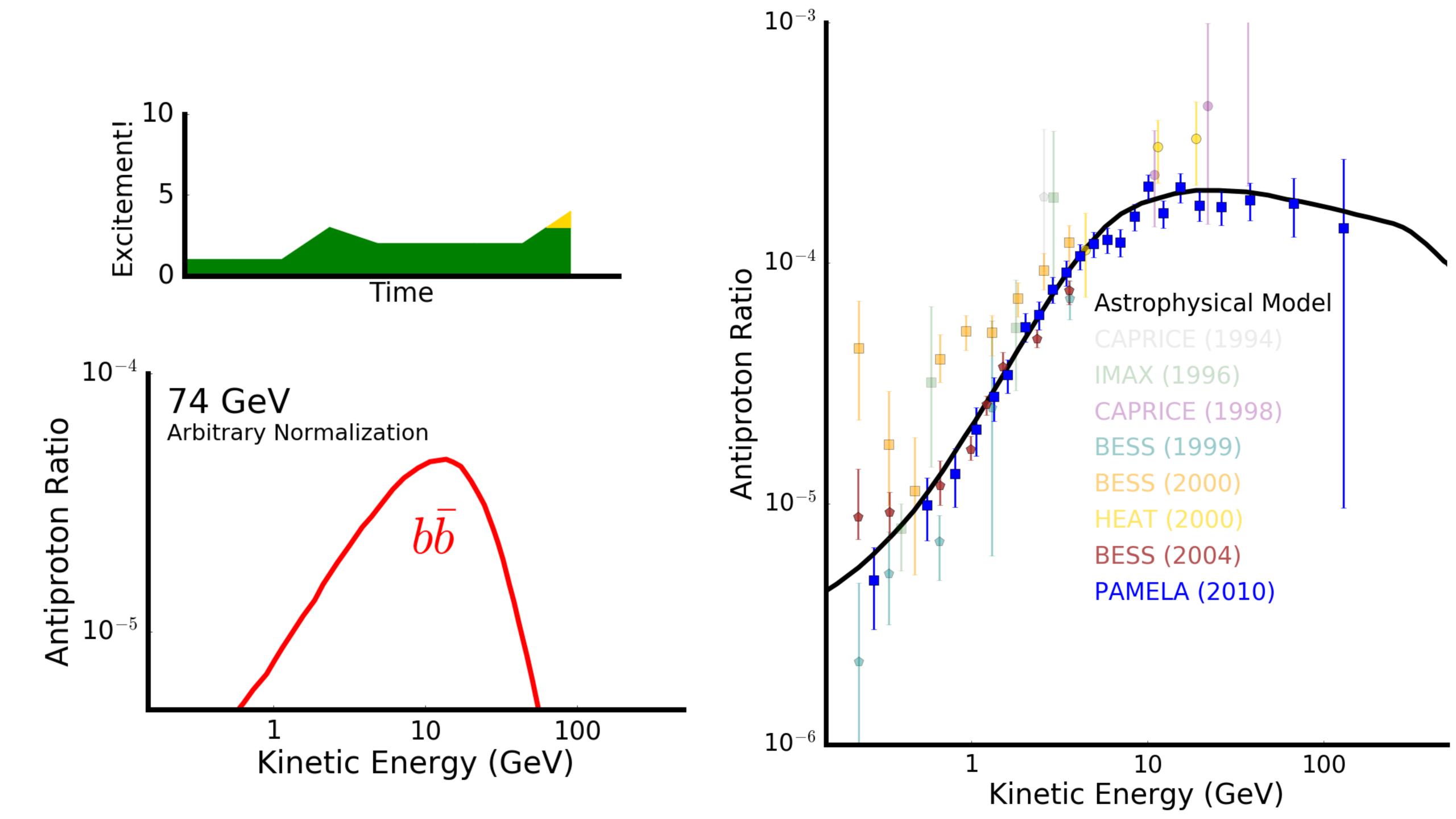


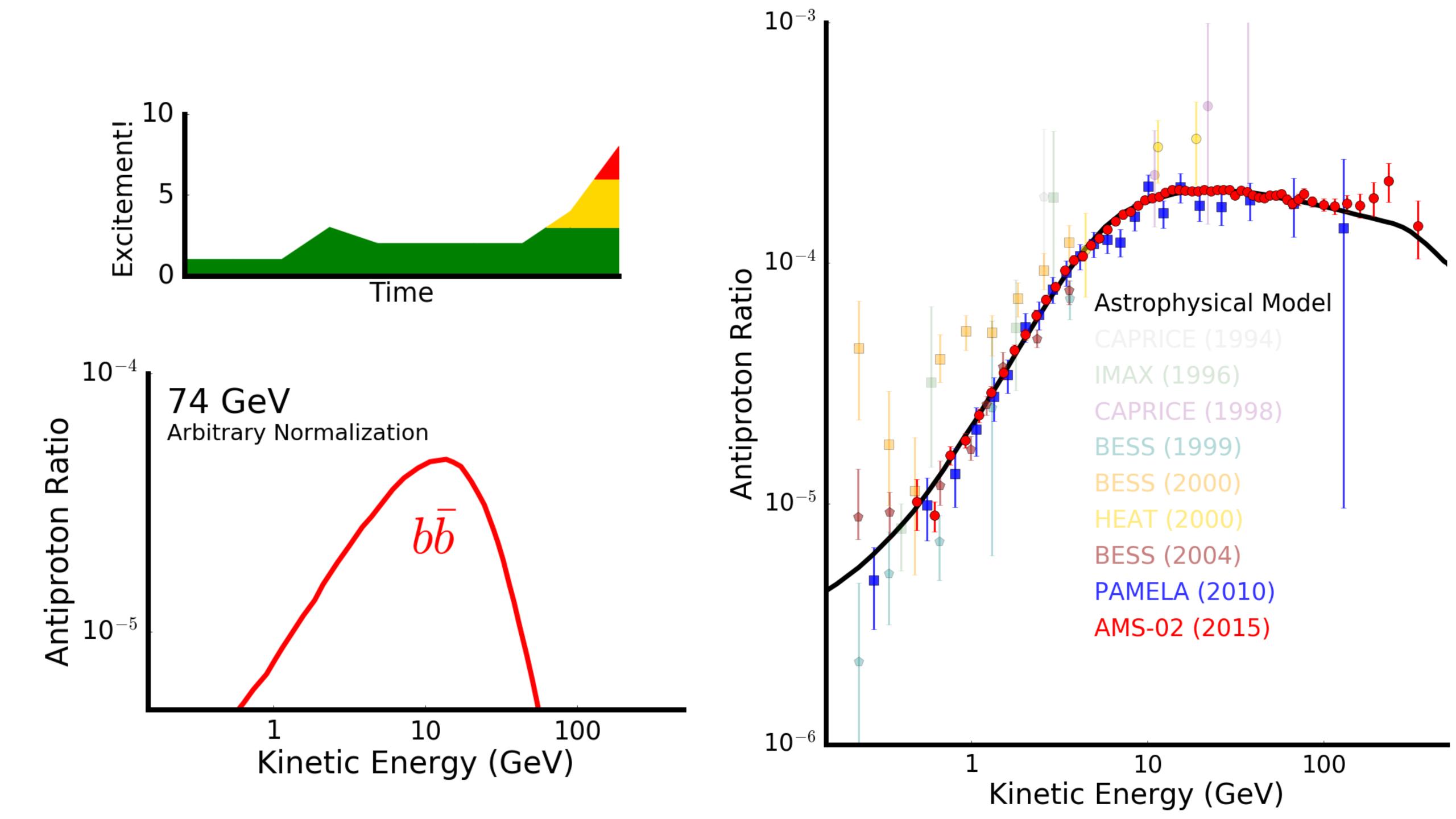


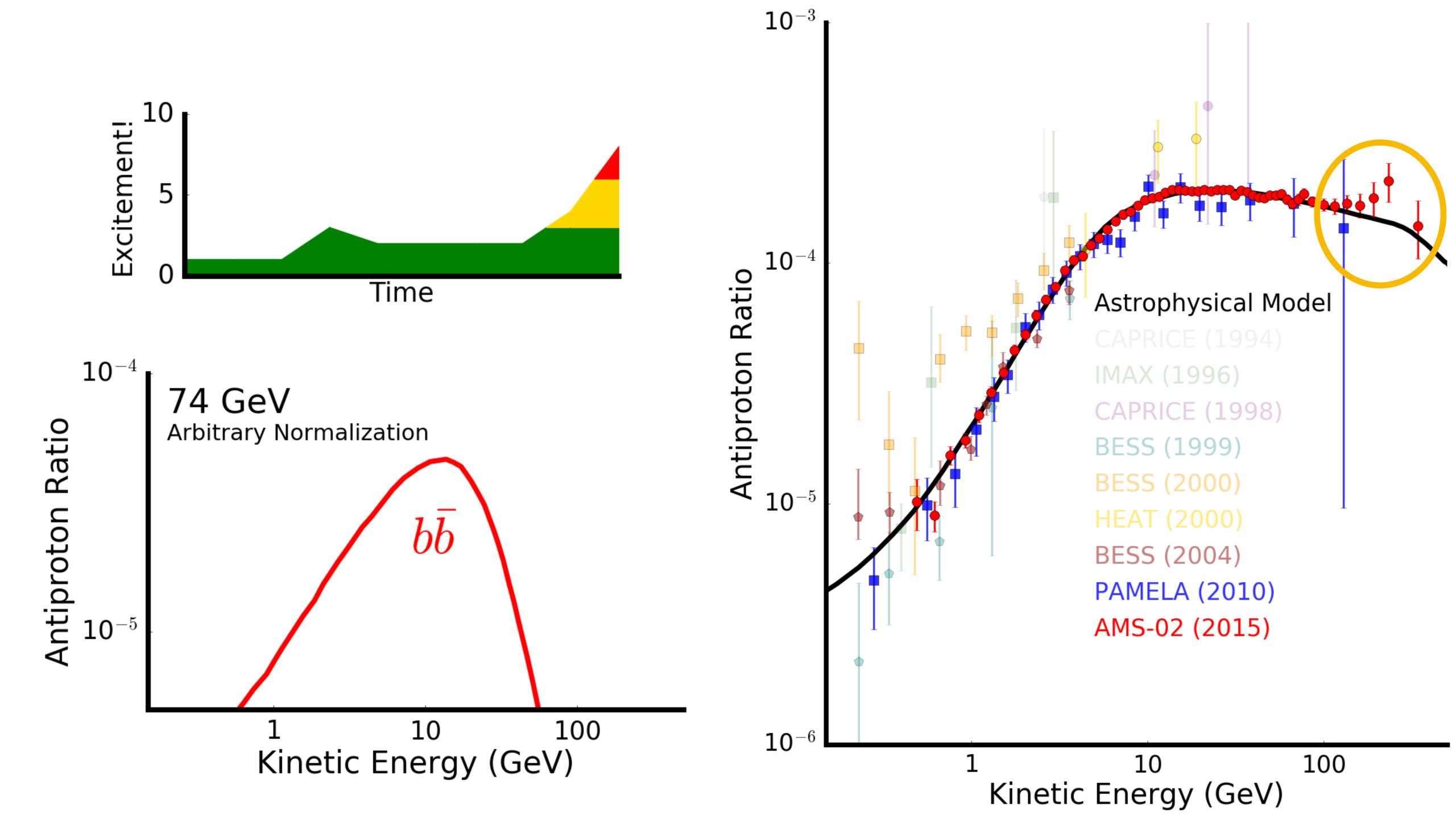


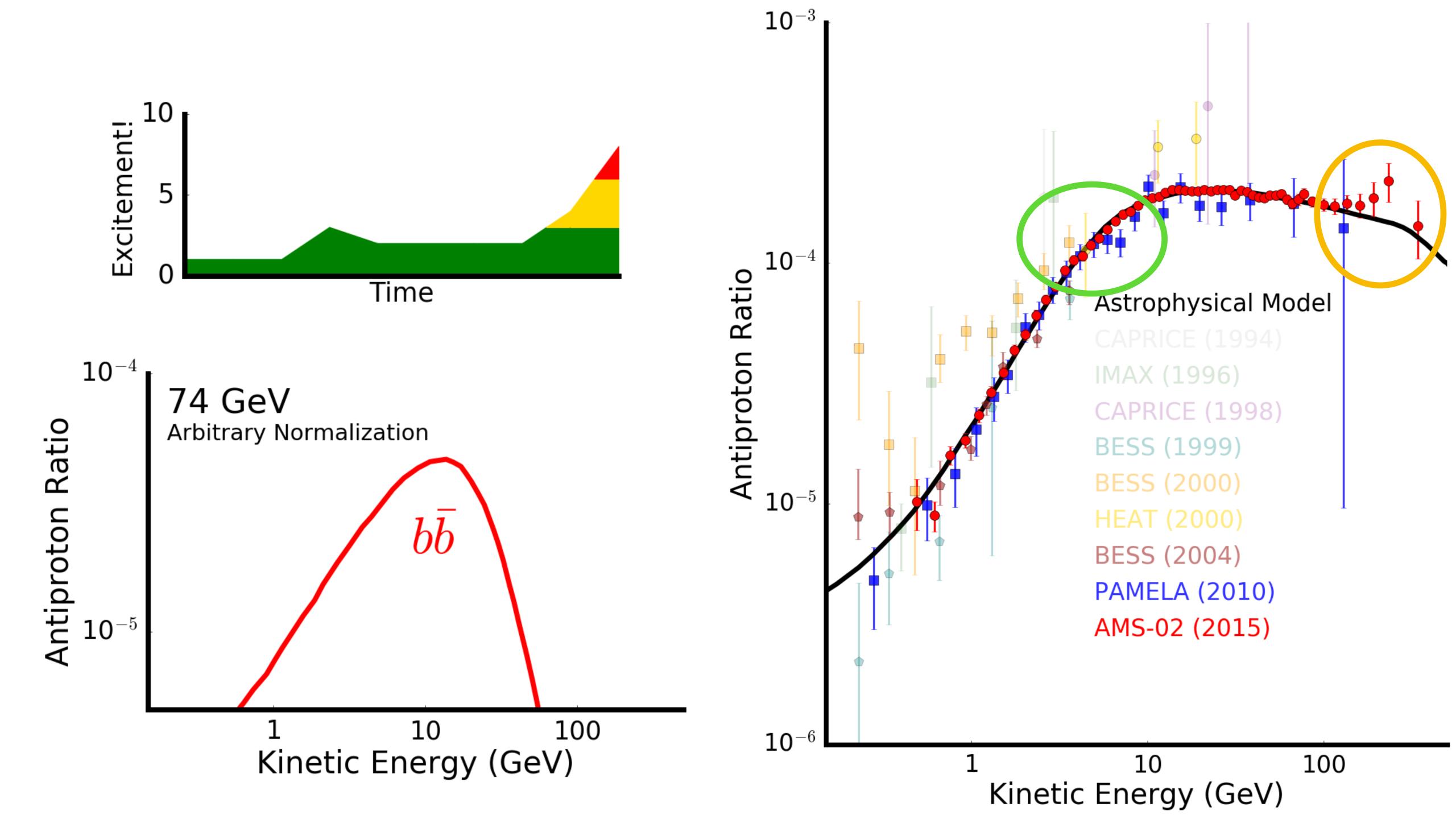


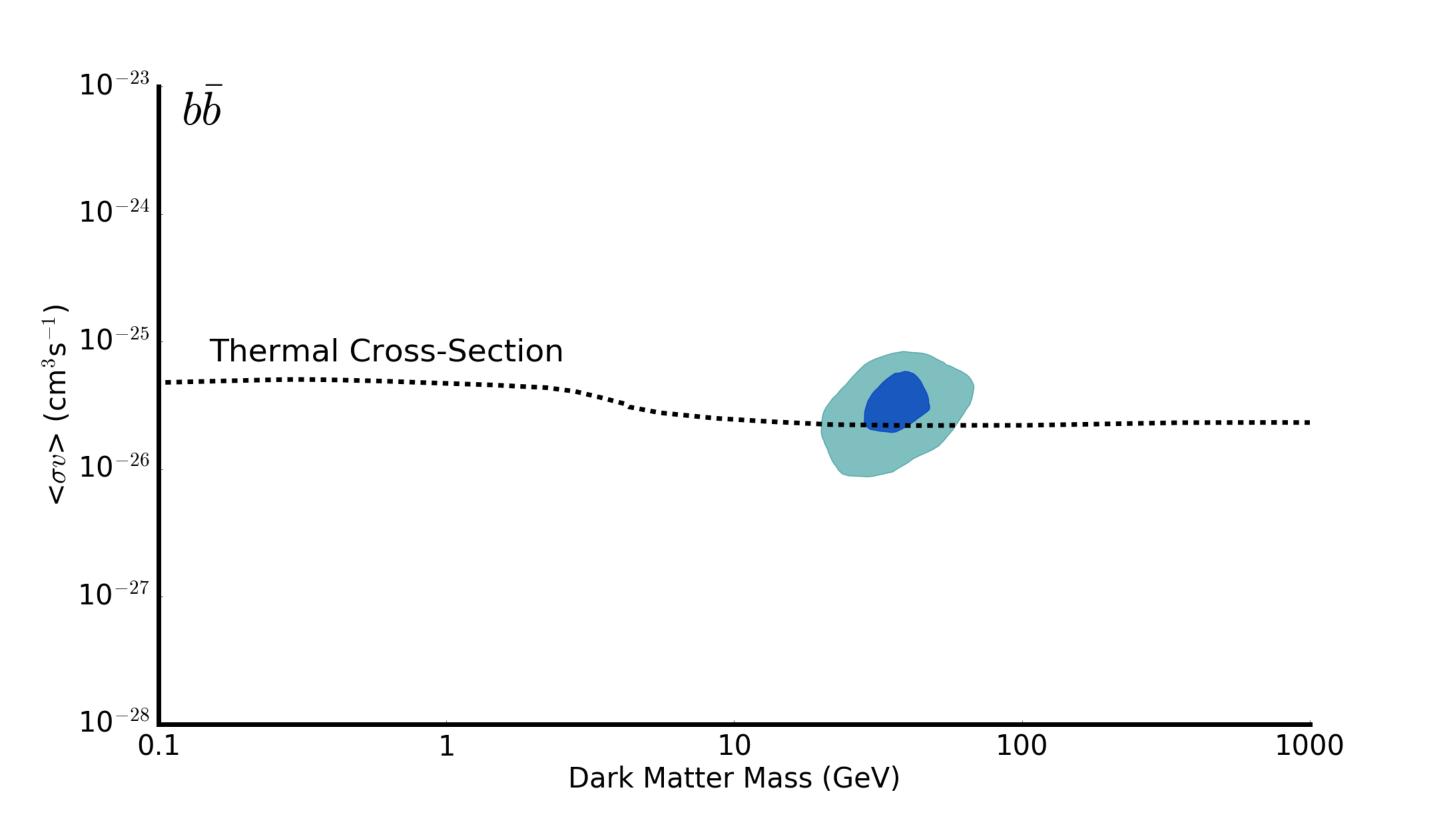


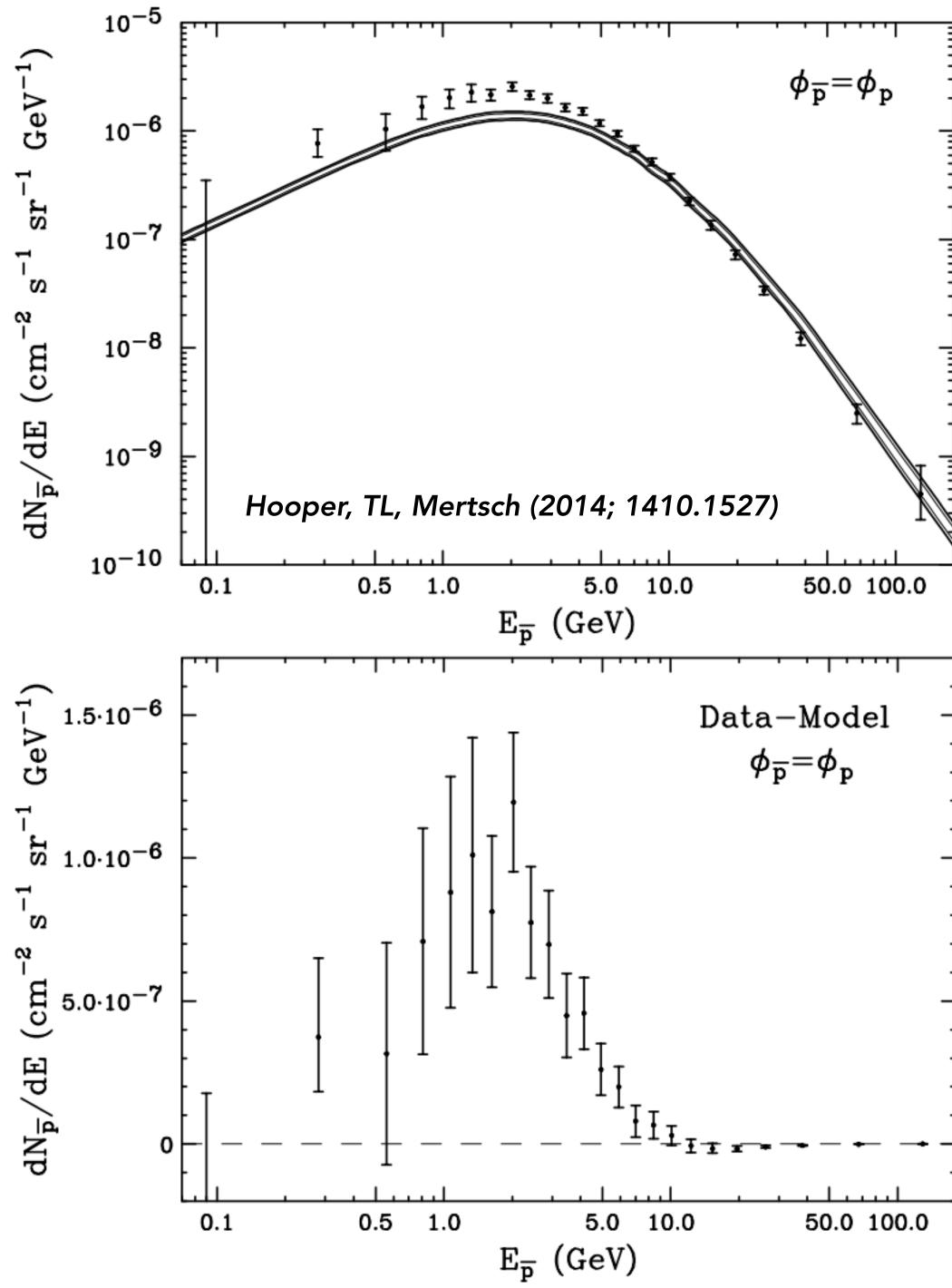


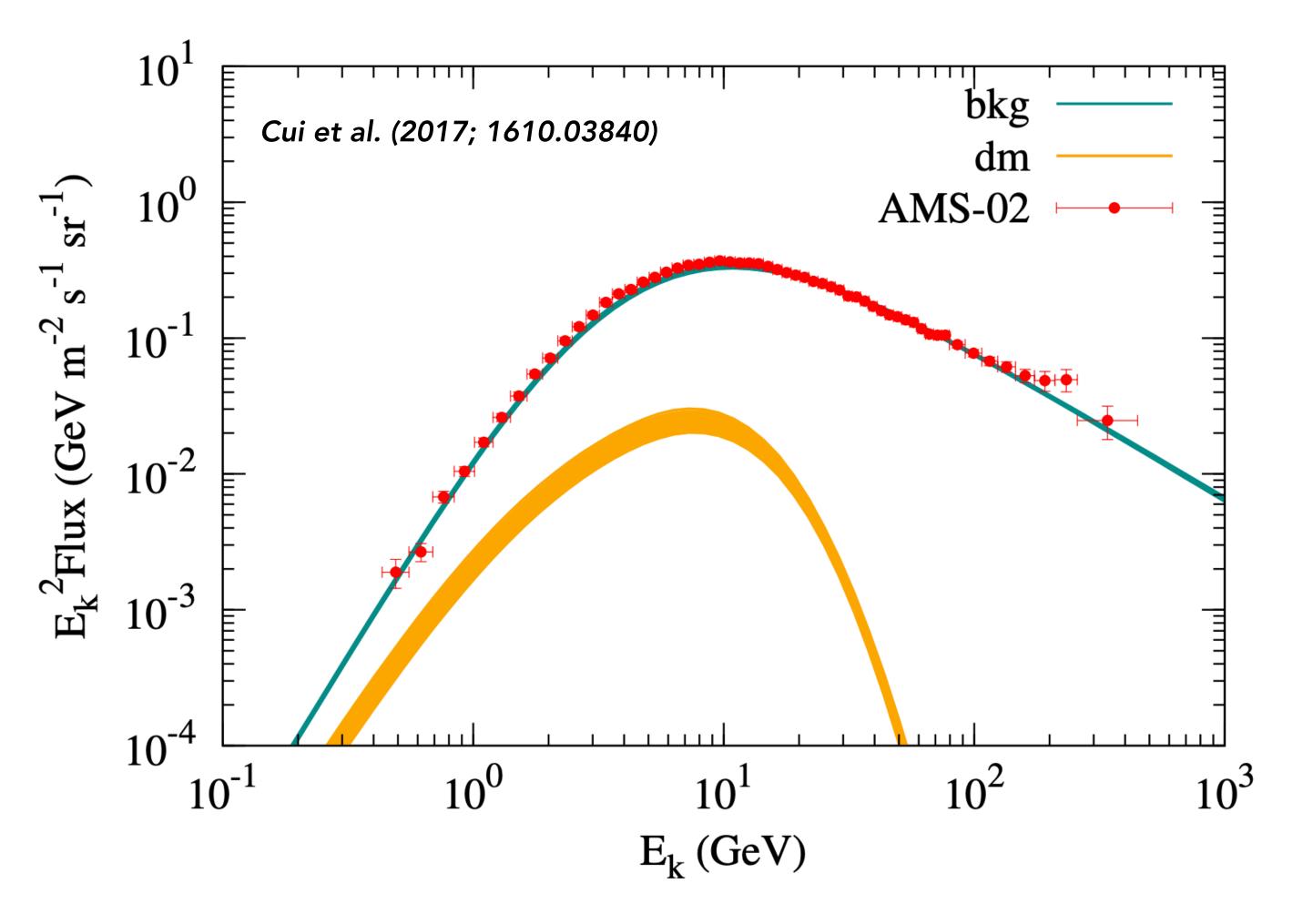


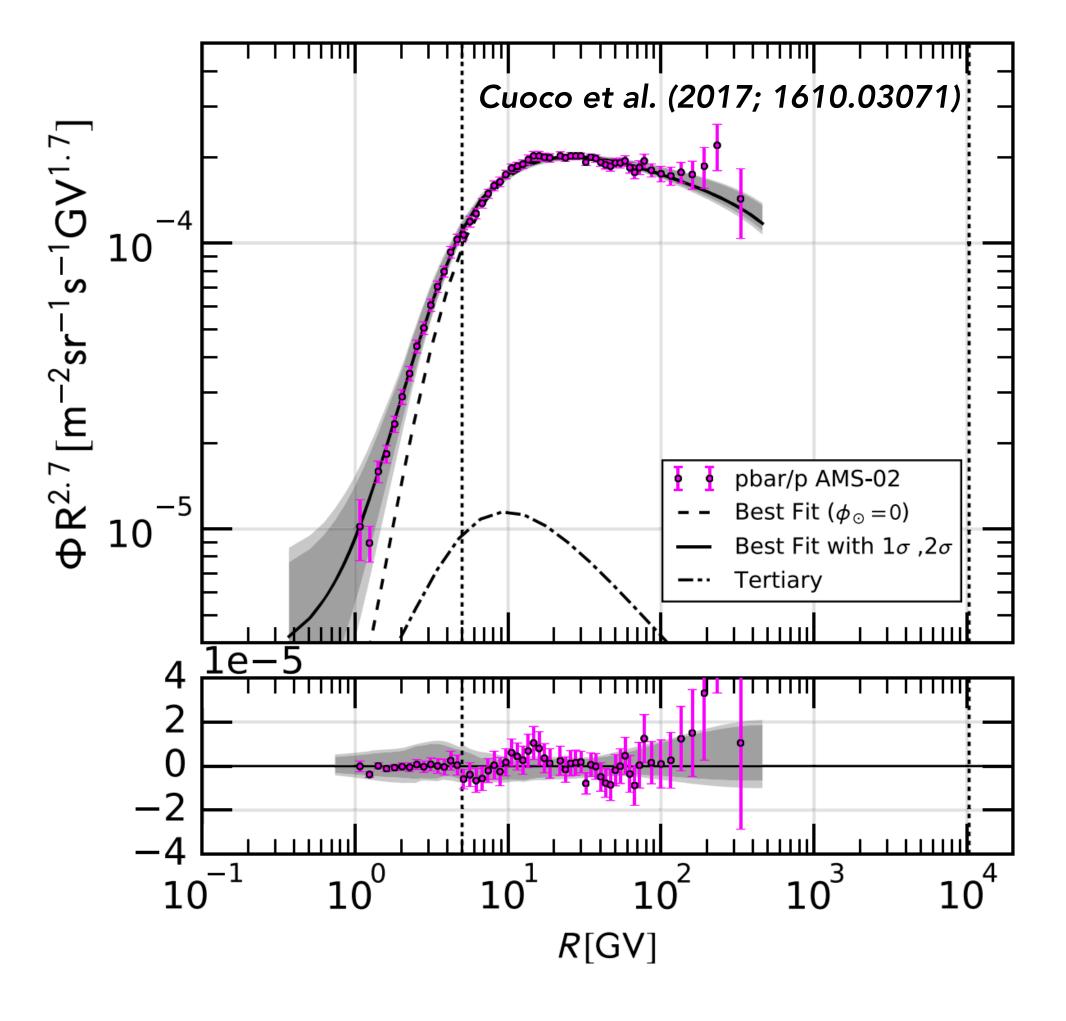




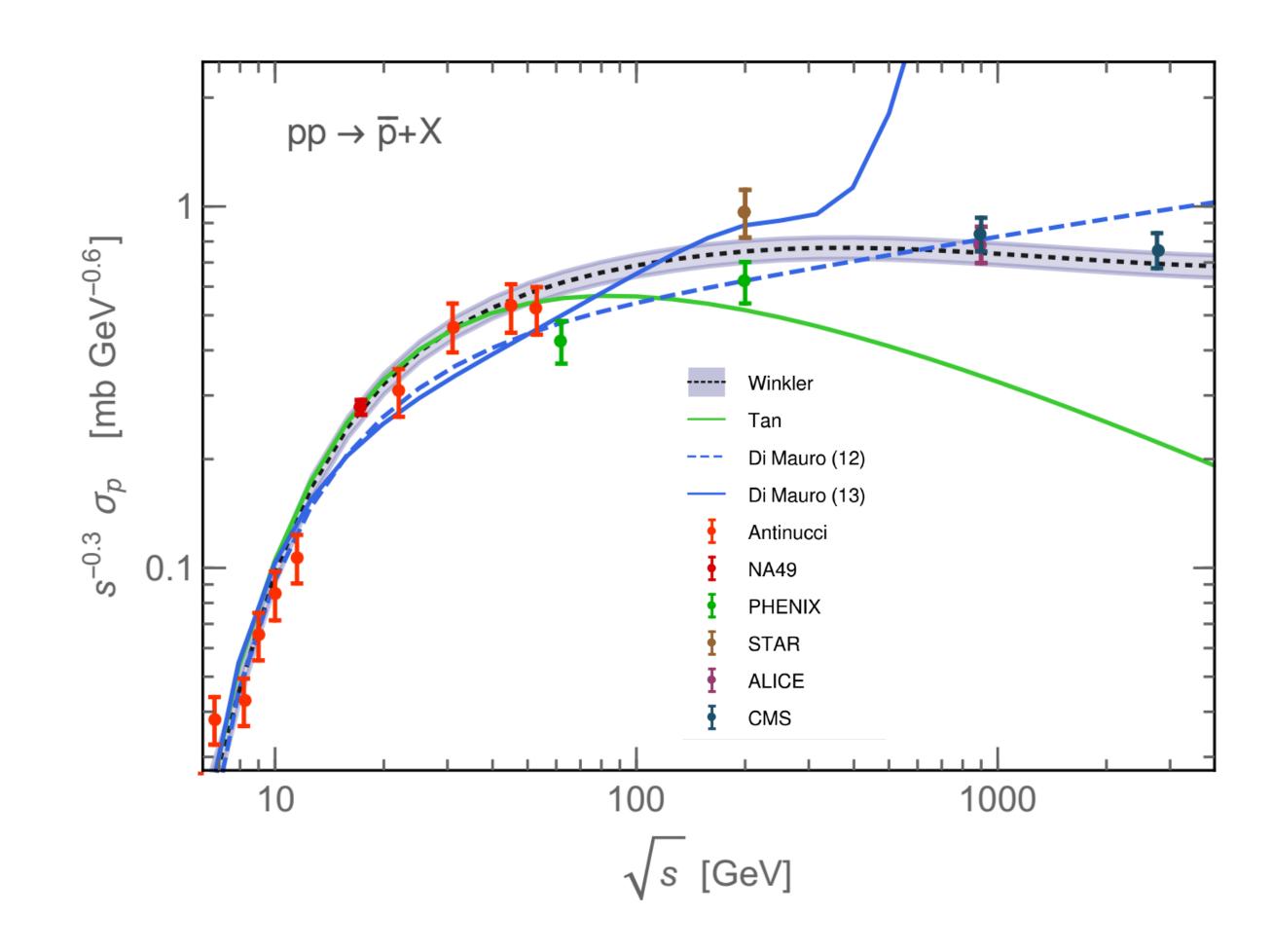








With great precision comes great responsibility:



**Galactic Primary to Secondary Ratios - Future AMS-02 Data! Inhomogeneous Diffusion - TeV Halos** 

**Solar Modulation - Voyager Data, Time-Dependent AMS-02 Data Antiproton Production Cross-Section - LHCb / Laboratory Experiments** 



#### Gamma-Rays

Galactic Center
Dwarf Spheroidal Galaxies
Galaxy Clusters
Milky Way Subhalos
Galactic Diffuse

Jupiter
Nearby Stars
Galactic Center Stars

Sun

Andromeda

**Little Galaxies** 

Isotropic Gamma-Ray Background Anisotropy Searches

Cusps 511 keV line

Morphology

#### Cosmic-Rays

**Positrons** 

**Electron + Positron Spectrum** 

**Antiprotons** 

**Antineutrons** 

**Antihelium** 

**Cosmological Lithium Problem** 

#### **Antimatter**

#### Neutrinos

The Sun

Direct Annihilation to Neutrinos
Hidden Sources
The Earth
Ultra-Heavy Dark Matter

Messengers

#### Low-Energy

Galactic Center Synchrotron

**Dwarf Galaxy Synchrotron** 

**Galaxy Cluster Synchrotron** 

**Diffuse Synchrotron** 

Sun

**Jupiter** 

**Isotropic Background** 

X-ray background from Clusters

**Anisotropy Searches** 

**Stellar Evolution** 

**Pulsar Evolution** 

**Planetary Heating** 

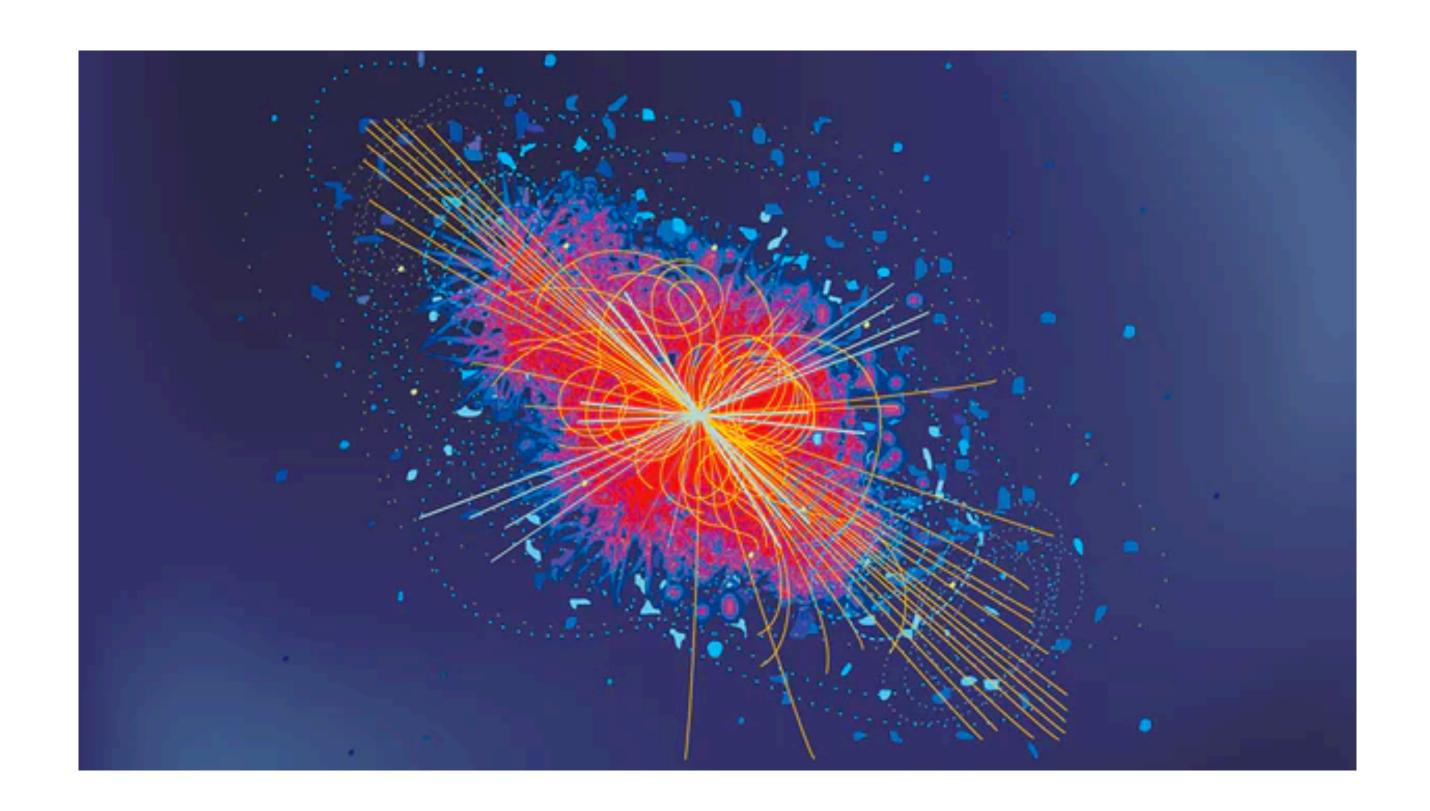
**Thermal Scattering** 

**Cosmic Microwave Background** 

**CMB** Absorption

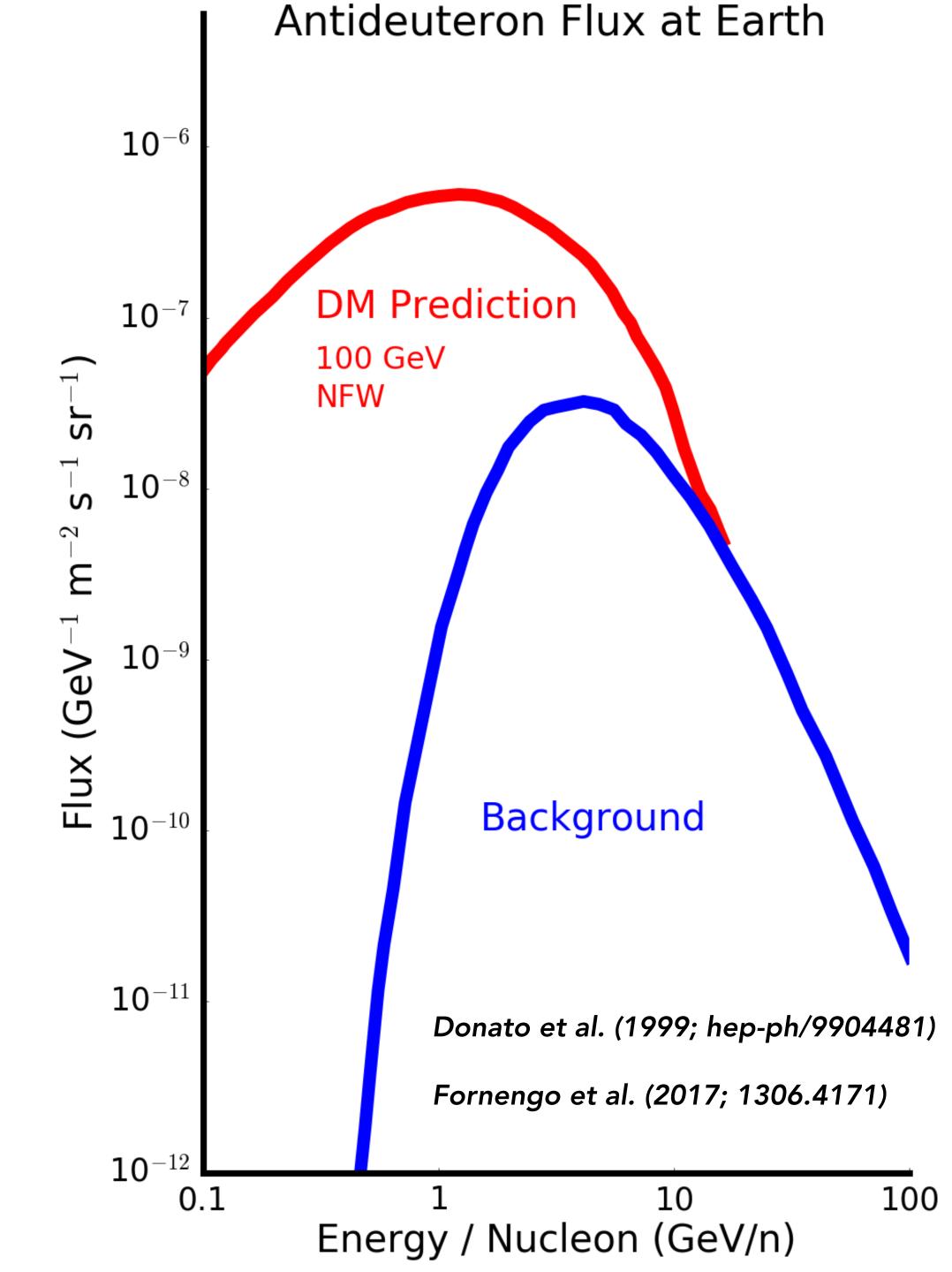
**Targets** 

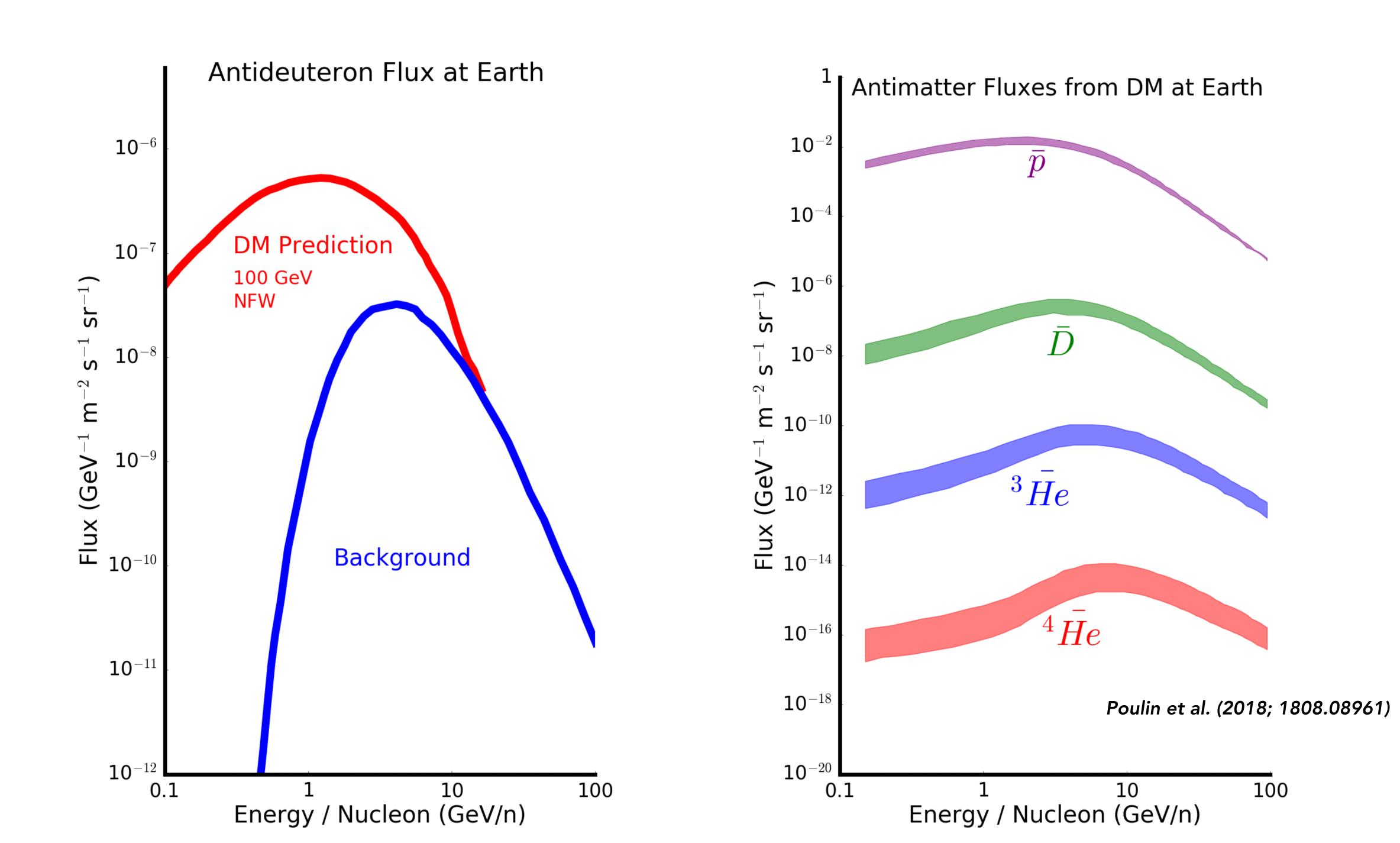




Astrophysical Antinuclei - Most be moving relativistically!

Dark Matter Antinuclei - Can be slow!





Small Dark Matter Signal Small Astrophysical Background

Hard

Large Dark Matter Signal Small Astrophysical Background

Easy

Easy

Small Dark Matter Signal Large Astrophysical Background

Hard

Large Dark Matter Signal Large Astrophysical Background

Small Dark Matter Signal Small Astrophysical Background

Acceptable?

Large Dark Matter Signal Small Astrophysical Background

Easy

Easy

Small Dark Matter Signal Large Astrophysical Background

Hard

Large Dark Matter Signal Large Astrophysical Background



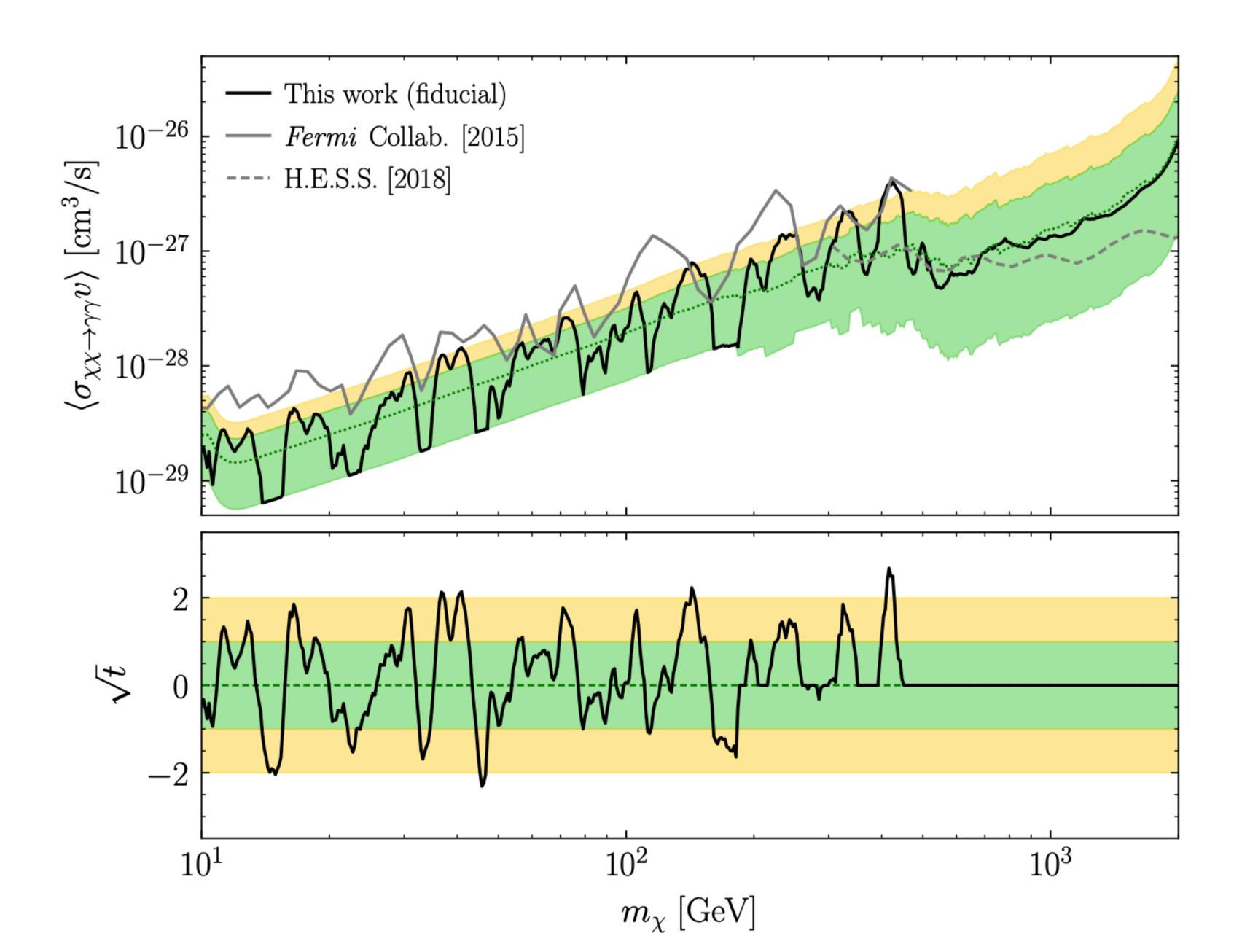
To date, we have observed eight events in the mass region from 0 to 10 GeV with Z=-2. All eight events are in the helium mass region.

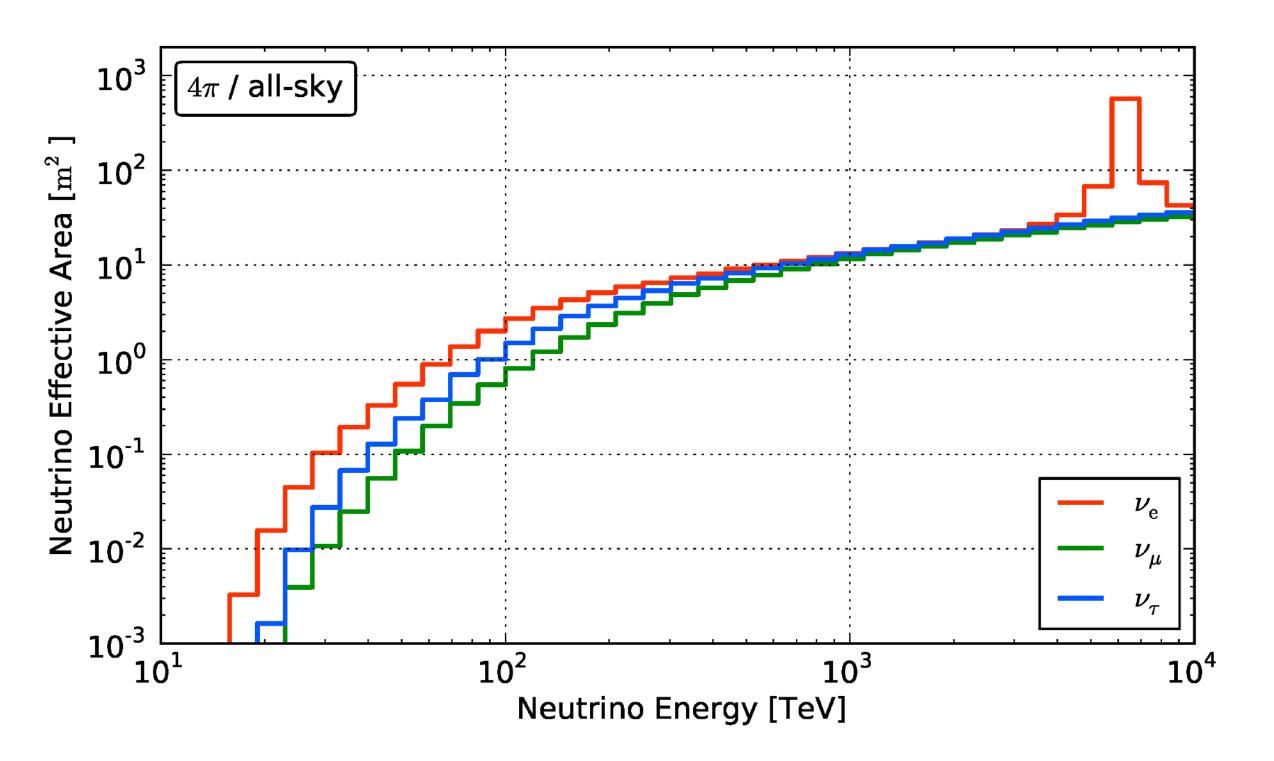
Currently (having used 50 million core hours to generate 7 times more simulated events than measured events and having found no background events from the simulation), our best evaluation of the probability of the background origin for the eight He events is less than  $3\times10^{-8}$ . For the two  ${}^4\text{He}$  events our best evaluation of the probability (upon completion of the current 100 million core hours of simulation) will be less than  $3\times10^{-3}$ .

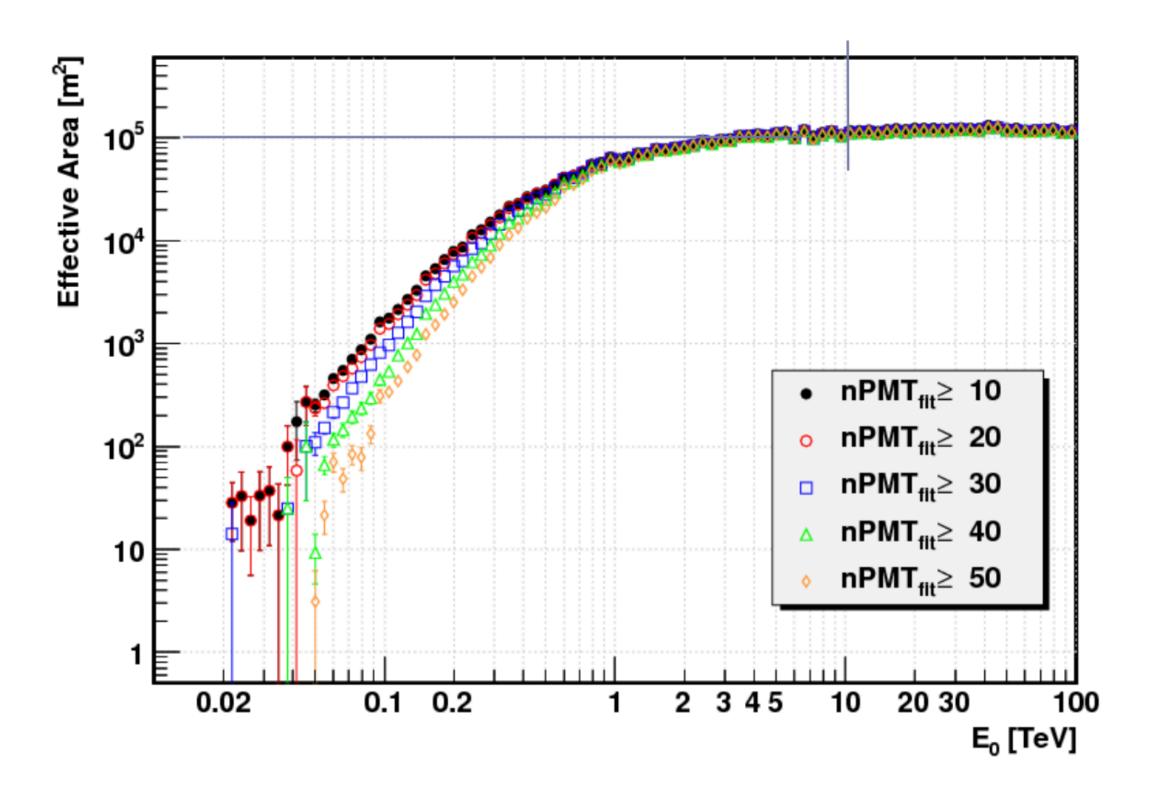
Note that for  ${}^4H\bar{}$ e, projecting based on the statistics we have today, by using an additional 400 million core hours for simulation the background probability would be  $10^{-4}$ . Simultaneously, continuing to run until 2023, which doubles the data sample, the background probability for  ${}^4H\bar{}$ e would be  $2\times10^{-7}$ , i.e., greater than 5-sigma significance.

X $X$ $X$			$\frac{\gamma}{\gamma}$
Signa 25 25 25 20 15 10 130 130 20 5	al significance (Solution 140) $E_{\gamma}$ [GeV]	OURCE)	Reg1 Reg2 - Reg3 Reg4 - Reg5

X $X$	\<	$\frac{1}{2}$
(L-S <sup>27</sup> 10 10 10 10 10 10 10 10 10 10 10 10 10	NFWc(γ=1.3) R3  Obs Limit P8 (5.8 yr stat+syst) Expected Limit Expected 68% Containment Obs Limit P7REP (3.7 yr, stat-only)  10 <sup>3</sup> 10 <sup>4</sup>	10 <sup>5</sup> m <sub>χ</sub> (MeV)



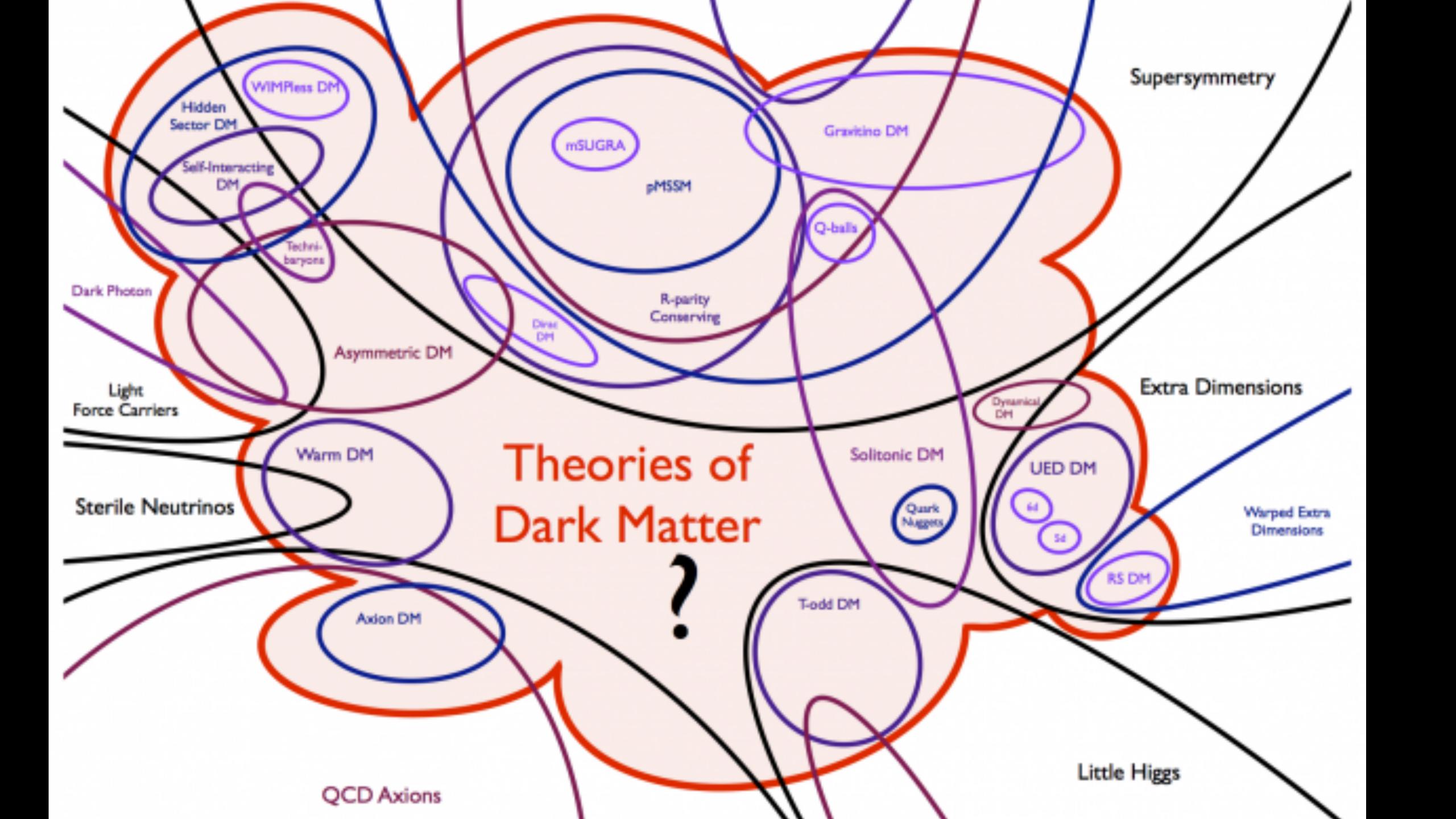




# Neutrino Searches

- o Celestial Body Focused Annihilation (e.g., Sun, GC Stars)
- O Direct Annihilation to Neutrinos (mass hierarchy??)
  - O Even then, cannot discount electromagnetic loops
- Exploiting gamma-ray attenuation.
- o Extremely heavy dark matter models.





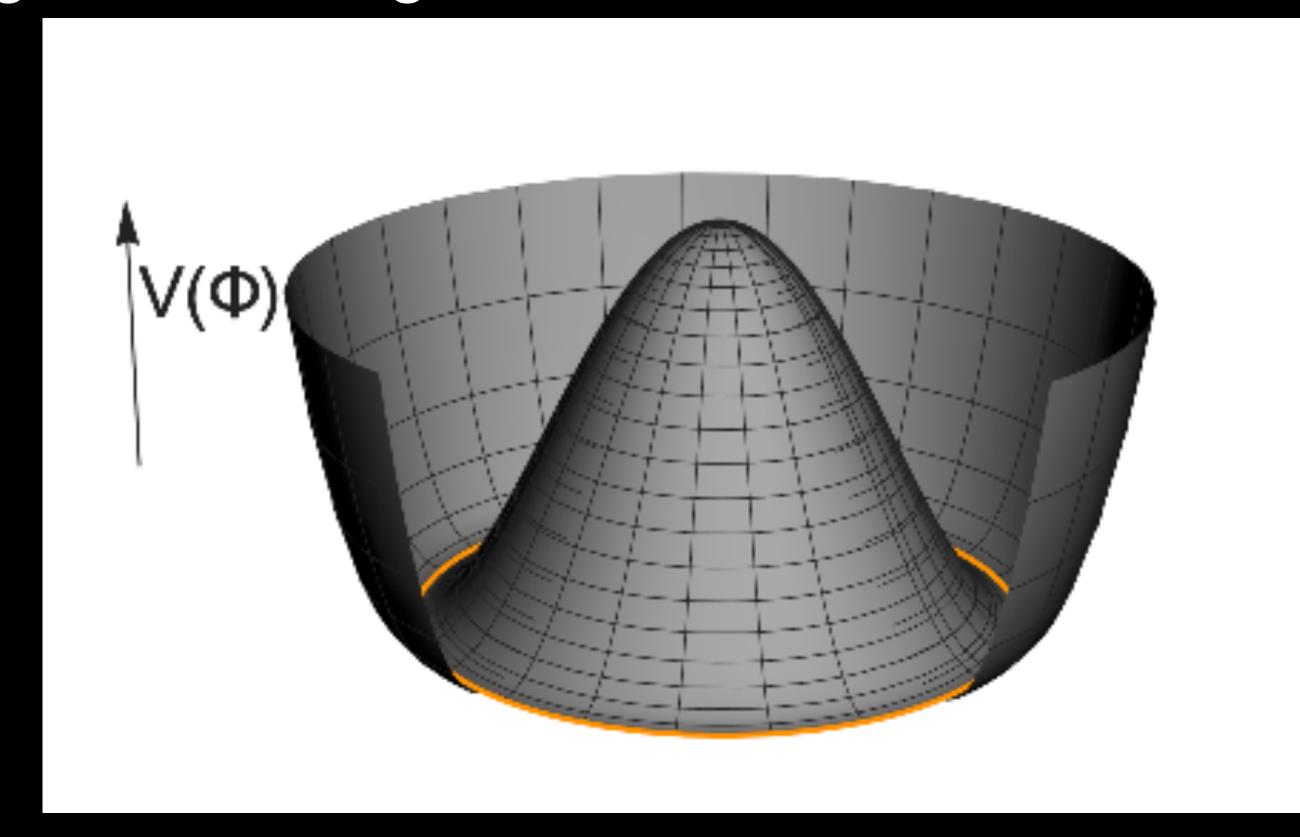
#### The Pecci-Quinn Mechanism

QCD Naturally Includes the following CP-violating term:

$$\mathcal{L} heta=ar{ heta}\,rac{g_s^2}{32\pi^2}\,G^a\mu
u\widetilde{G}^{a\,\mu
u}$$

o Turn theta into a dynamic variable:

$$\mathcal{L} \supset \left(rac{a(x)}{f_a} + ar{ heta}
ight) rac{g_s^2}{32\pi^2} \ G_{\mu
u}^a \widetilde{G}^{a\,\mu
u}$$



This naturally introduces a vacuum potential, which wants to fall to 0.

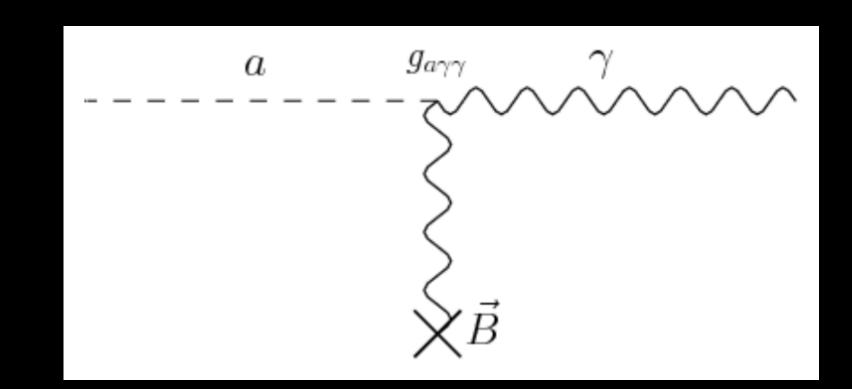
$$V\!(a) \sim \chi \, \left( 1 - \cos\!\left( rac{a}{f_a} + ar{ heta} 
ight) 
ight)$$

#### The Pecci-Quinn Mechanism

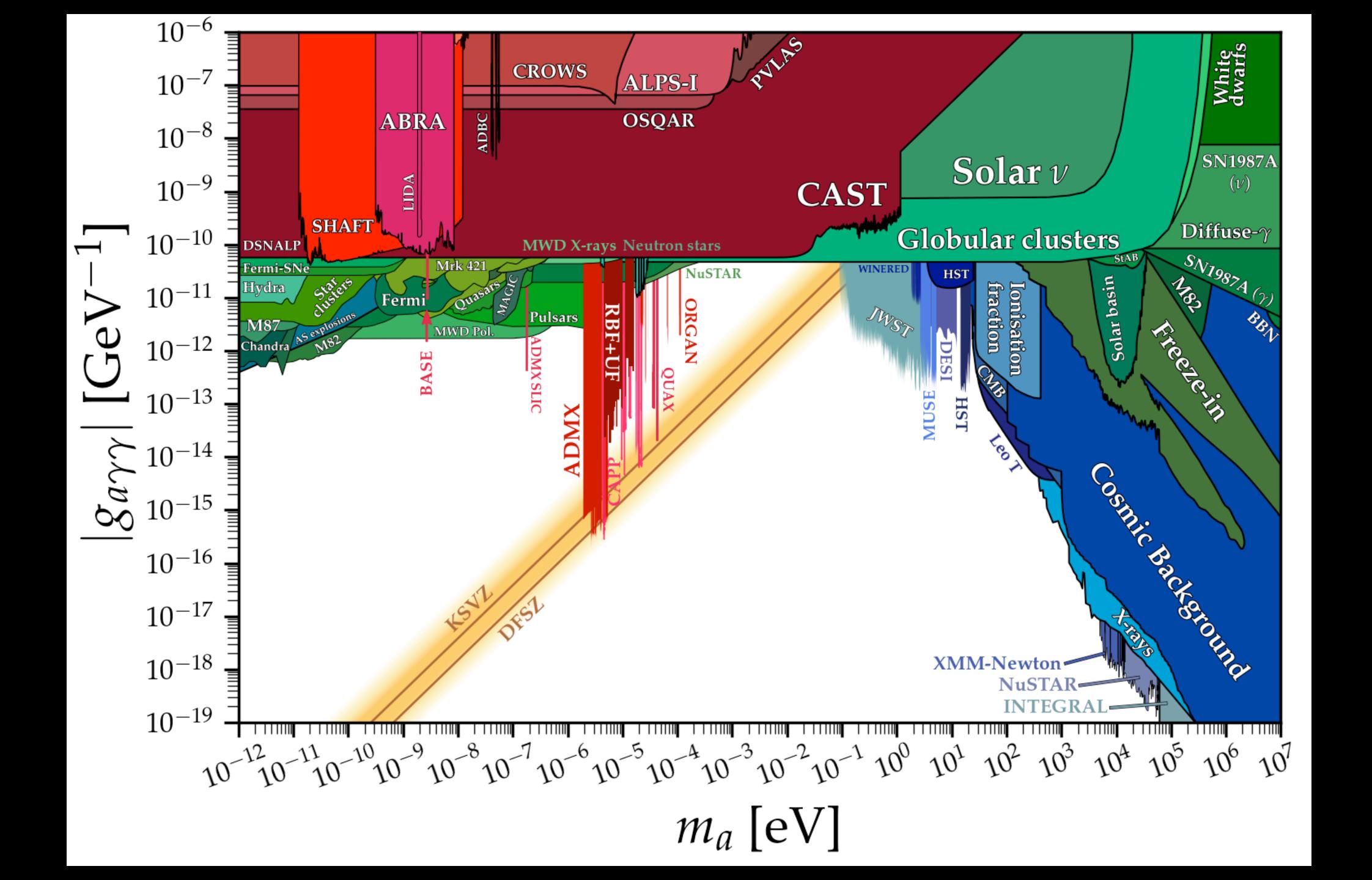
 The drawback (benefit?) is that this axion term must also couple to electromagnetic fields (photons).

$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma} a \vec{E} \cdot \vec{B},$$

$$g_{a\gamma} \simeq \frac{\alpha}{2\pi f_a} \sim 10^{-12} \text{ GeV}^{-1} \left(\frac{10^9 \text{ GeV}}{f_a}\right)$$



 This makes it possible to search for axion to photon conversion in the universe today.





electron neutrino



muon neutrino



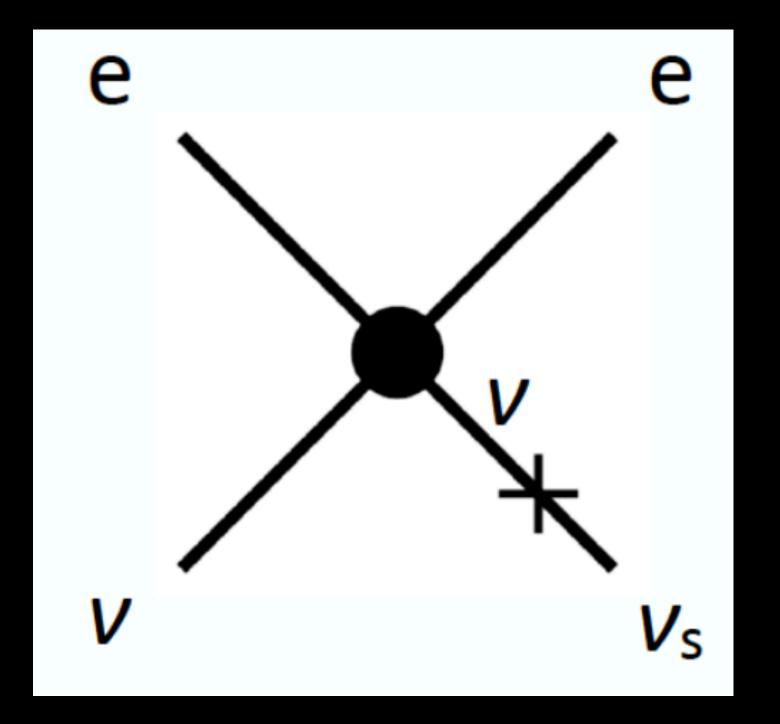
tau neutrino

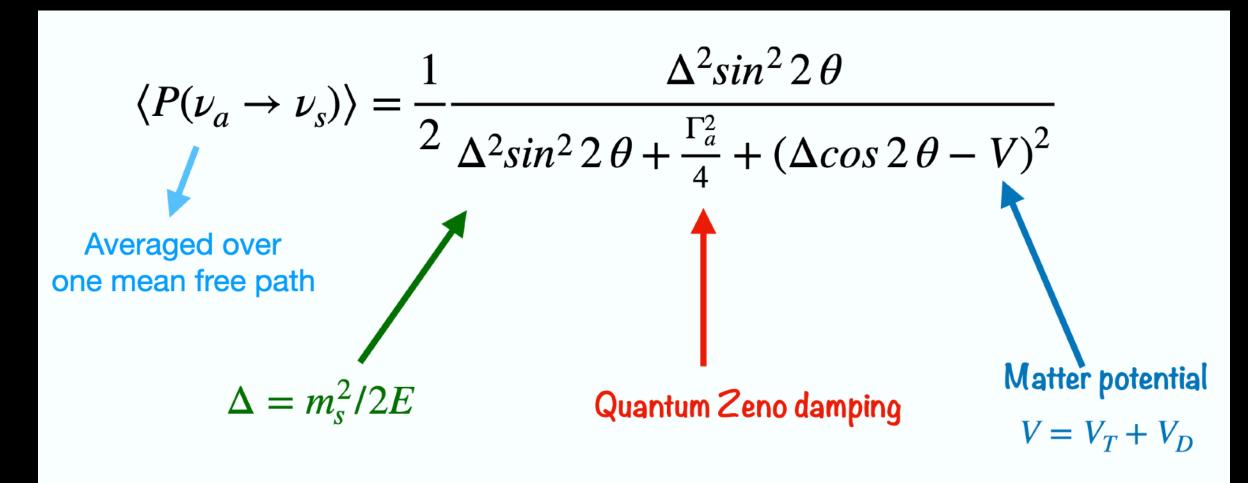


sterile neutrino

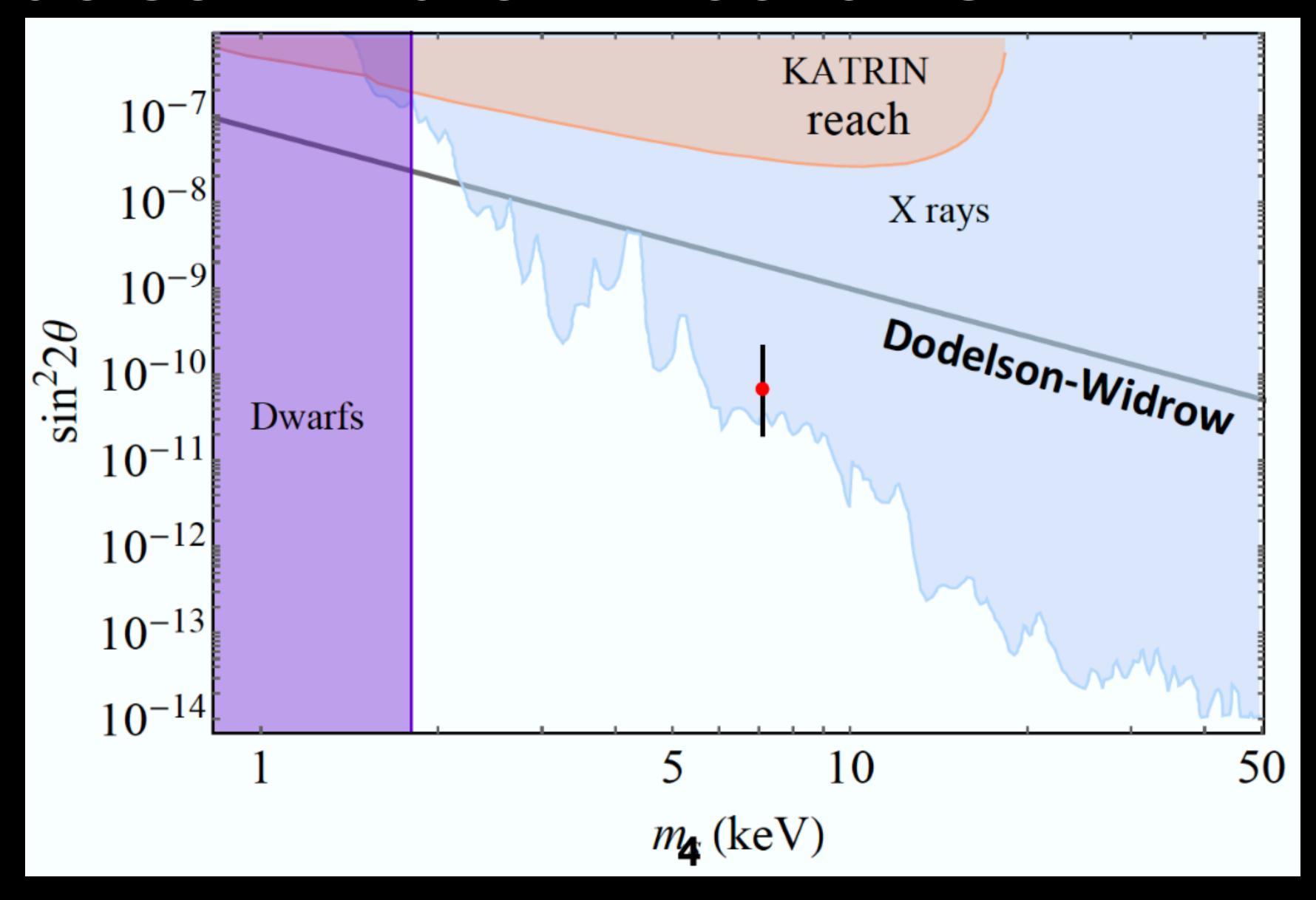
## The Dodelson-Widrow Mechanism

- Add a 4th generation of neutrinos (right-handed, solving the right/left handed neutrino problem).
- o Can also be correlated to leptogenesis.
- O Sterile neutrino cannot be in thermal equilibrium, but active neutrinos can be and can oscillate to sterile neutrinos with amplitude  $\sin^2(\theta_4)$

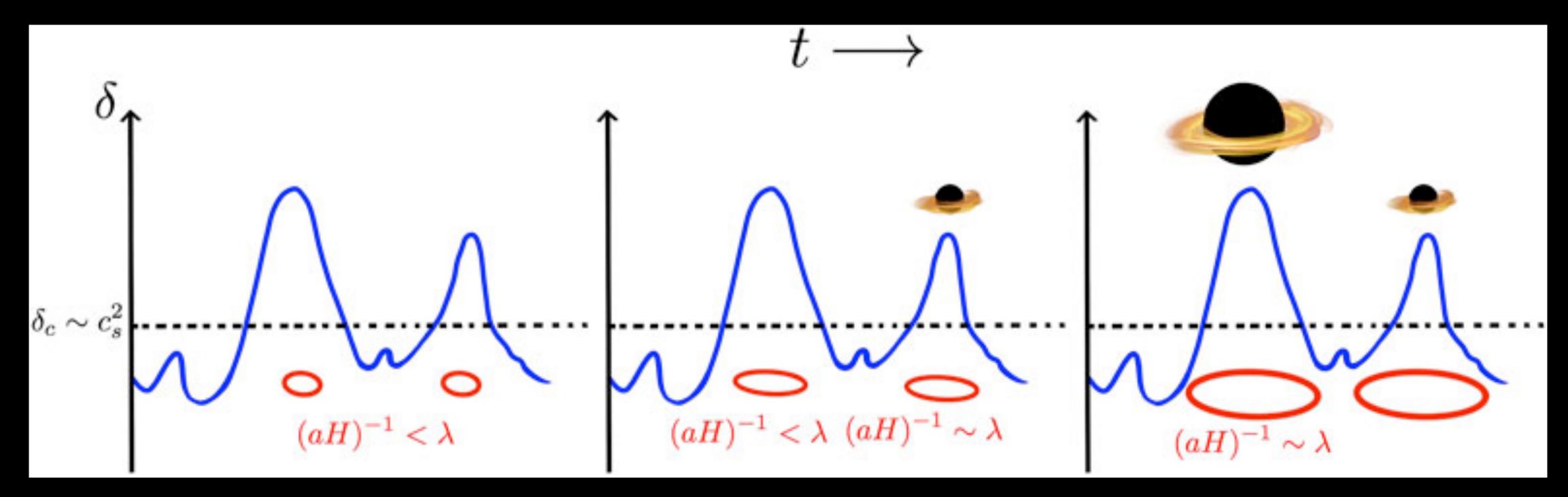




## The Dodelson-Widrow Mechanism



#### Primordial Black Holes



Density Perturbations in the Early Universe Could Collapse into BHs with masses that are related to the mass inside a Hubble radius at the time of their formation:

$$M_{
m H} = rac{4}{3} \, ar{
ho} \Big(rac{c}{H}\Big)^3 = rac{c^3}{2\,GH} \sim 10^{15} \, {
m g} \, \left(rac{t}{10^{-23} {
m s}}
ight)$$

o Formation details dependent on radiation/matter domination during that time.

