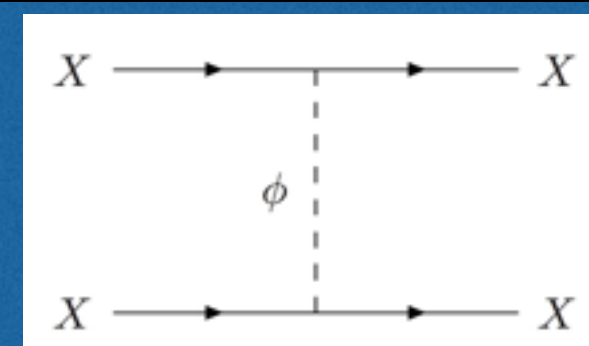
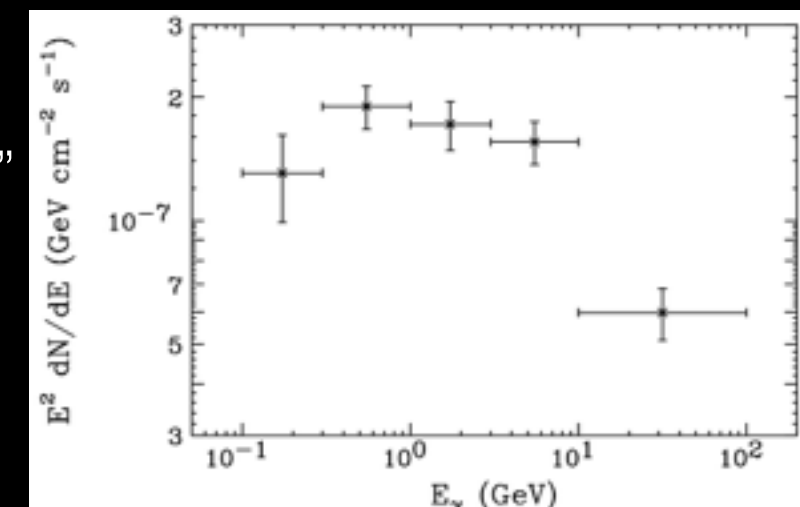
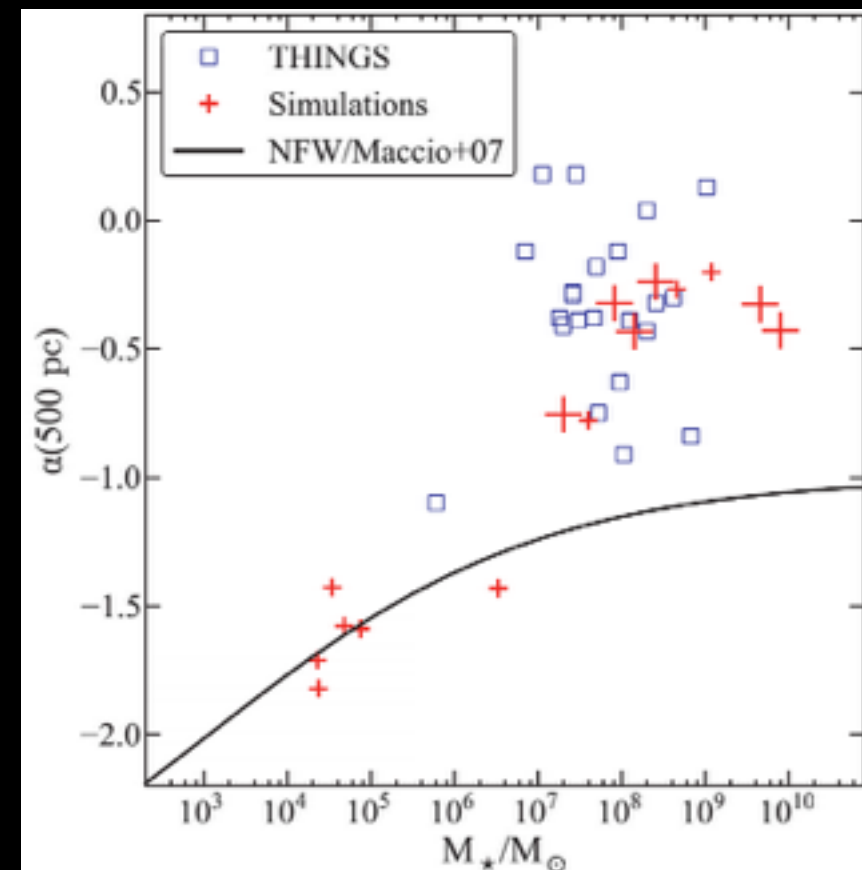


# The Indirect Detection of Self-Interacting Dark Matter



Tim Linden

KICP Fellow

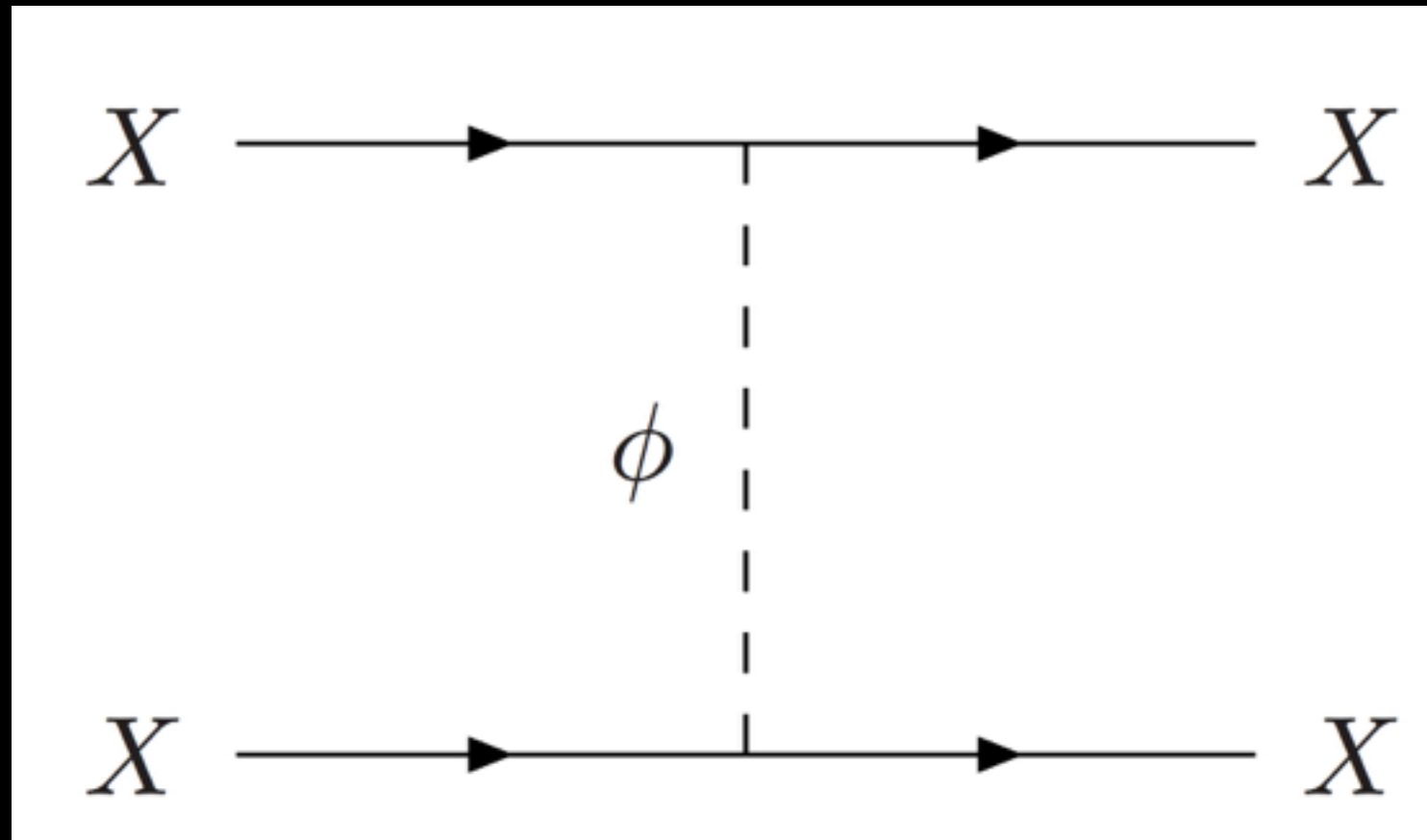


Based on Two Soon to be Submitted Papers:

“Density Profile of Self-Interacting Dark Matter in the Presence of Baryons”  
Kaplinghat, Linden, Keeley, Yu (2013)

“Indirect Searches for Self-Interacting Dark Matter”  
Kaplinghat, Linden, Yu (2013)

# What are Self-Interactions?



Add a dark force carrier to the dark sector, which allows interactions between dark matter particles

In this case - a dark photon

# Cluster Constraints?

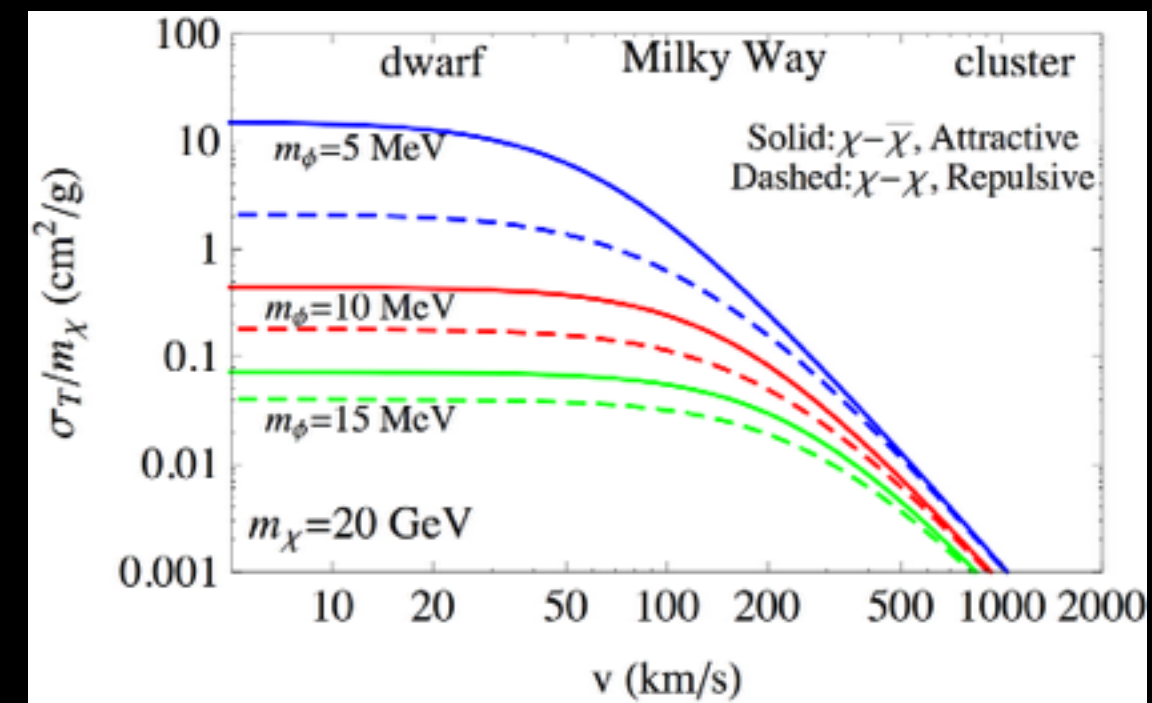
Bullet Cluster observations show that dark matter is collisionless — kind of?



Actual constraint is about  $1 \text{ cm}^2/\text{g} \sim 2 \text{ barn}/\text{GeV}$

This is about 10 orders of magnitude above the expected WIMP annihilation cross-section  
(WIMP miracle implies  $\sigma \sim 10^{-10} \text{ barn}$ )

Moreover, the self-interaction cross-section may be velocity dependent

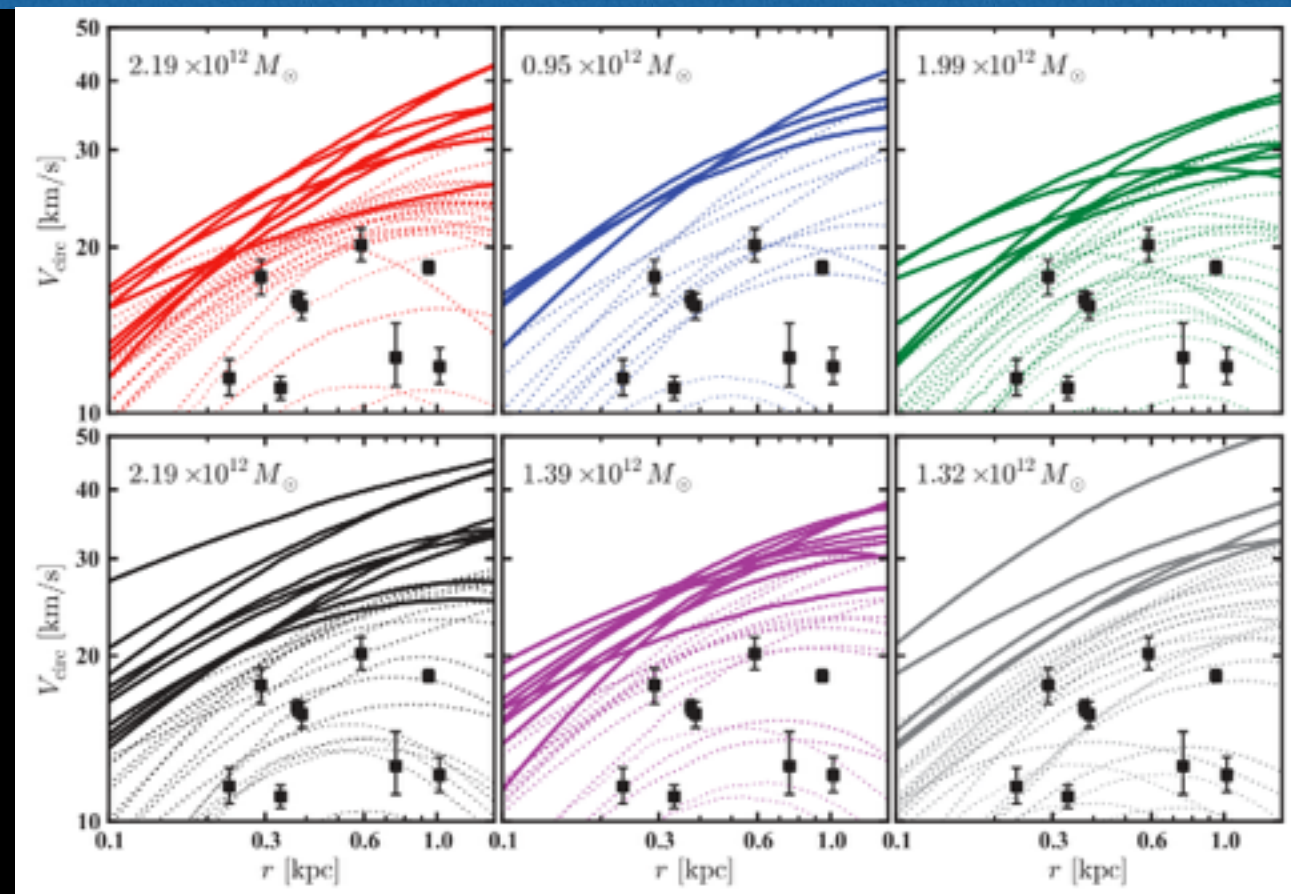




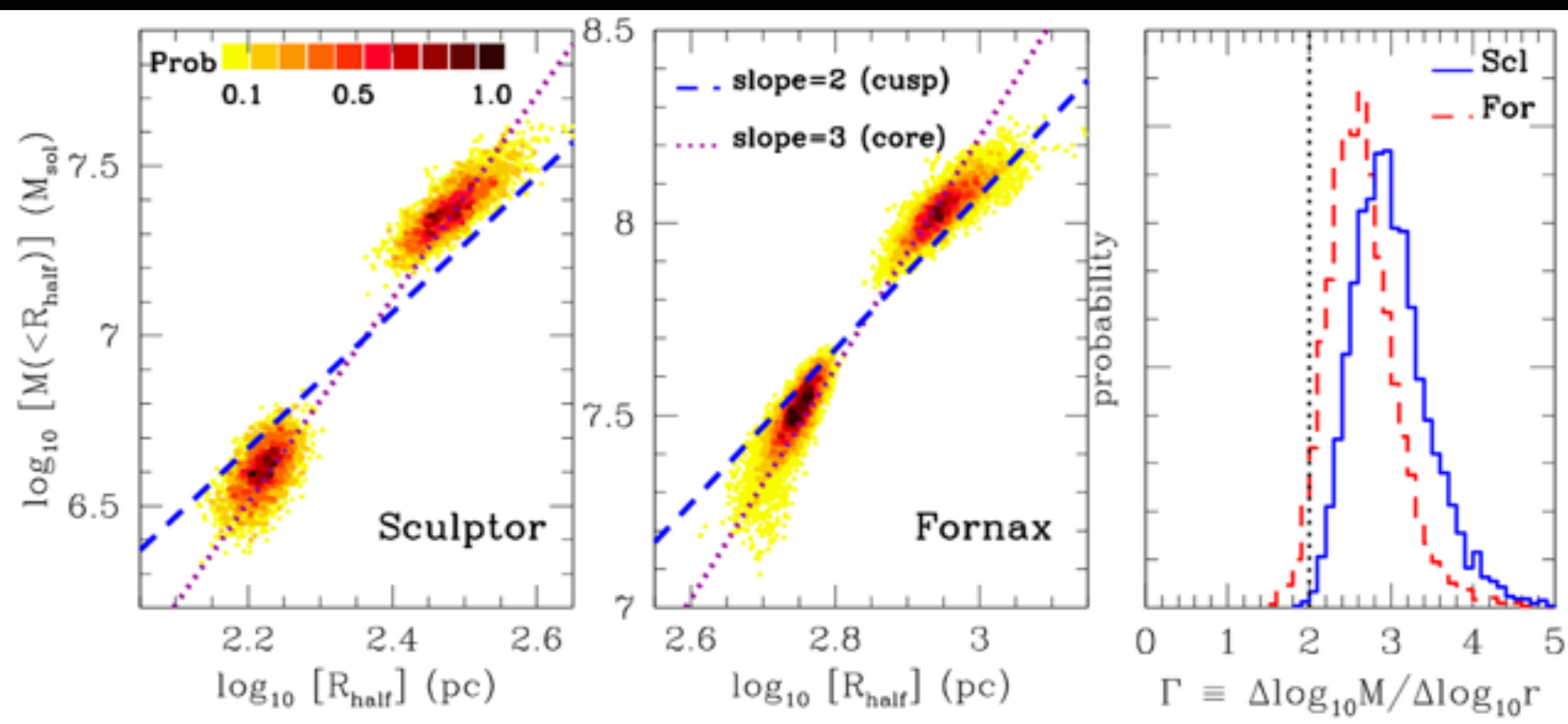
# Why Self-Interactions?

Cusped profiles (e.g. NFW) may not fit the observed density profiles of dwarfs

Stellar feedback is unlikely to solve the problem, as these are dark matter dominated systems

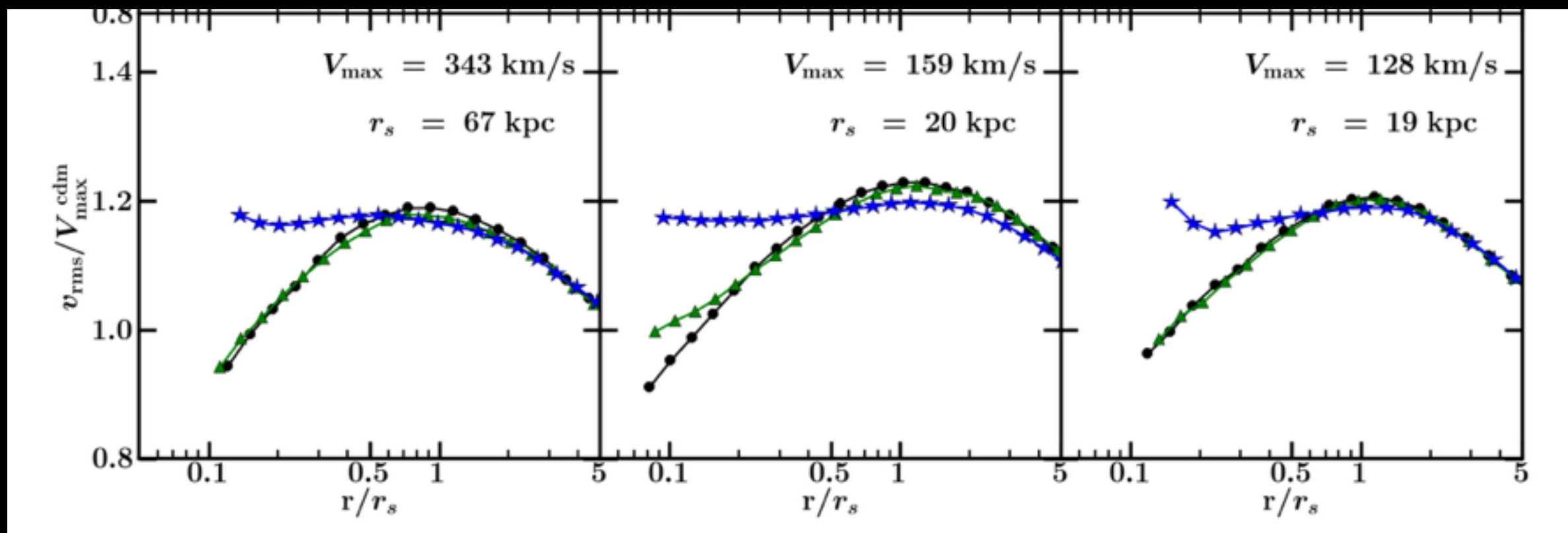


“Too Big to Fail” - Boylin-Kolchin et al. 2012



“Core-Cusp Problem”  
Walker & Penarrubia (2012)

# Why Self-Interactions?



Self-Interactions produce a DM core at approximately the point where one interaction is expected per Hubble time.

$$\sigma_{SIDM} \left( \frac{\rho(r)}{M_{DM}} \right) v_{rms} = \Gamma(r) > 1$$

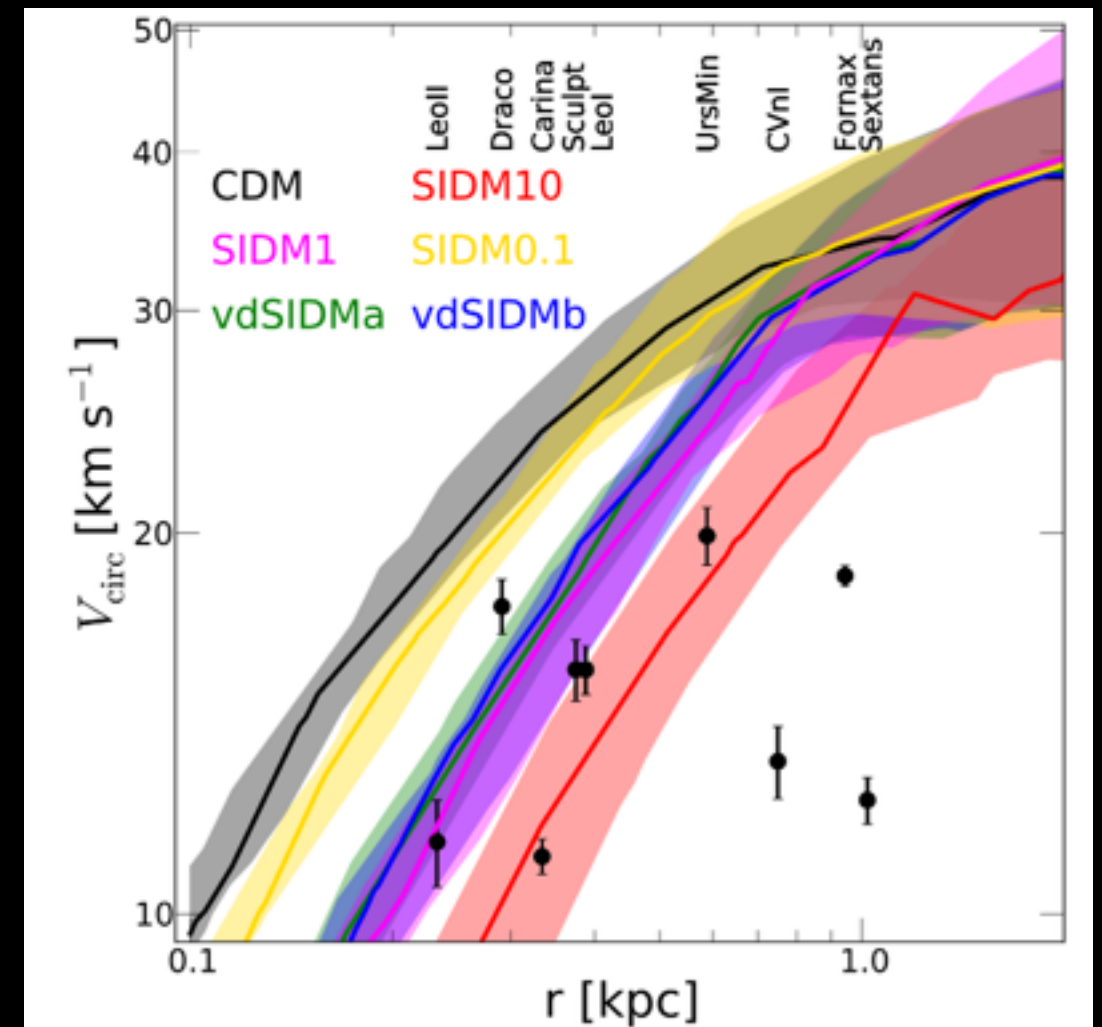
Inside this region, interactions move high temperature (high  $v_{rms}$ ) DM into the center and move low temperature gas outside, creating a constant density, constant temperature core

# Why Self-Interactions?

DM less peaked at GC  $\rightarrow$  higher mass clusters have smaller circular velocities inside the core

This solves too-big-to fail and the core-cusp problem

Requires a self-interaction in dwarf galaxies of  $\sim 10 \text{ cm}^2/\text{g}$



Zavala et al. (2013)

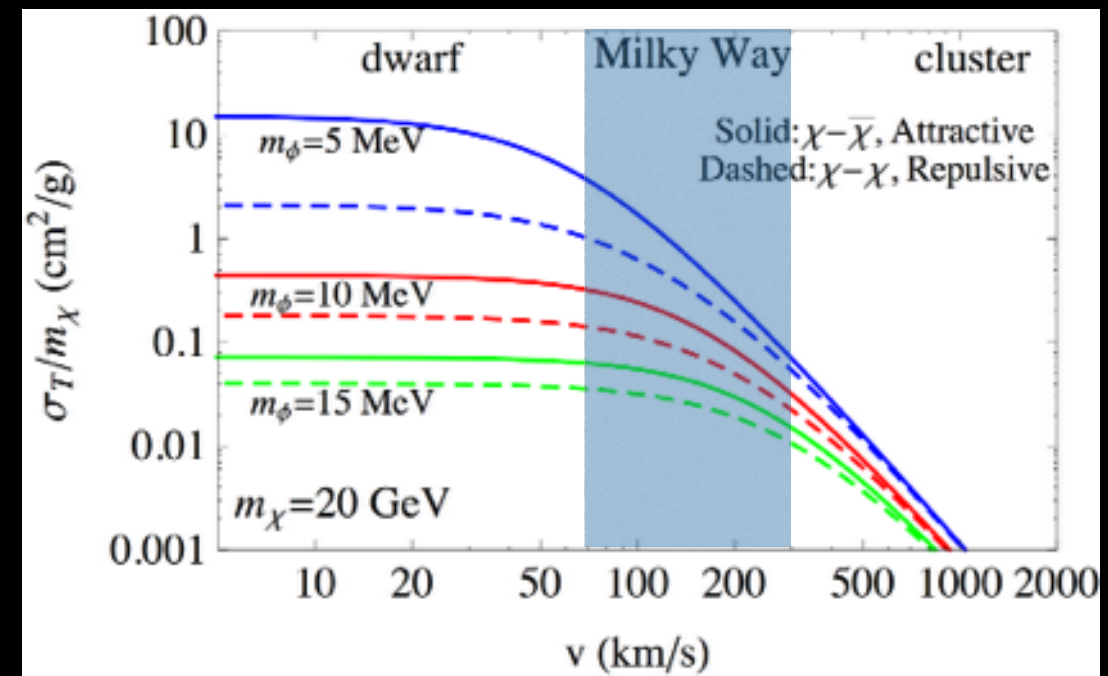


# The Milky Way Density Profile

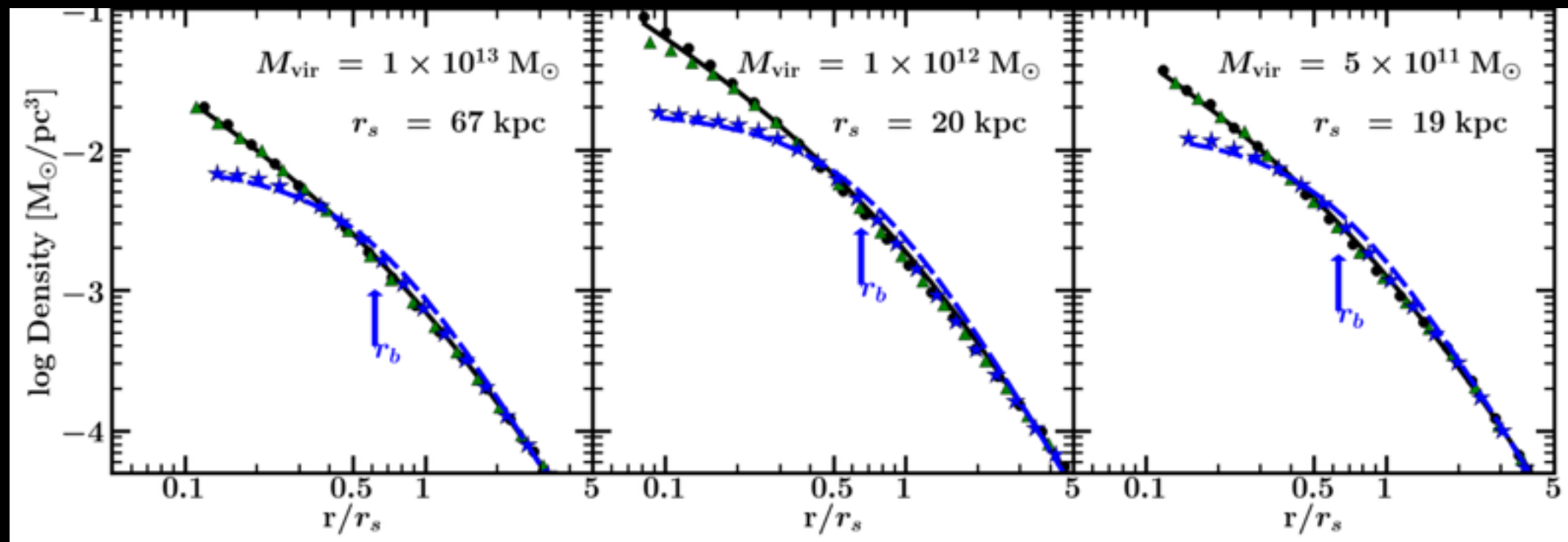
But what about the Milky Way Density Profile?

SIDM  $\sigma > 1$  in dwarf galaxies predicts  $\sigma \sim 1$  in the Milky Way

This produces a core with a radius  $\sim 10$  kpc

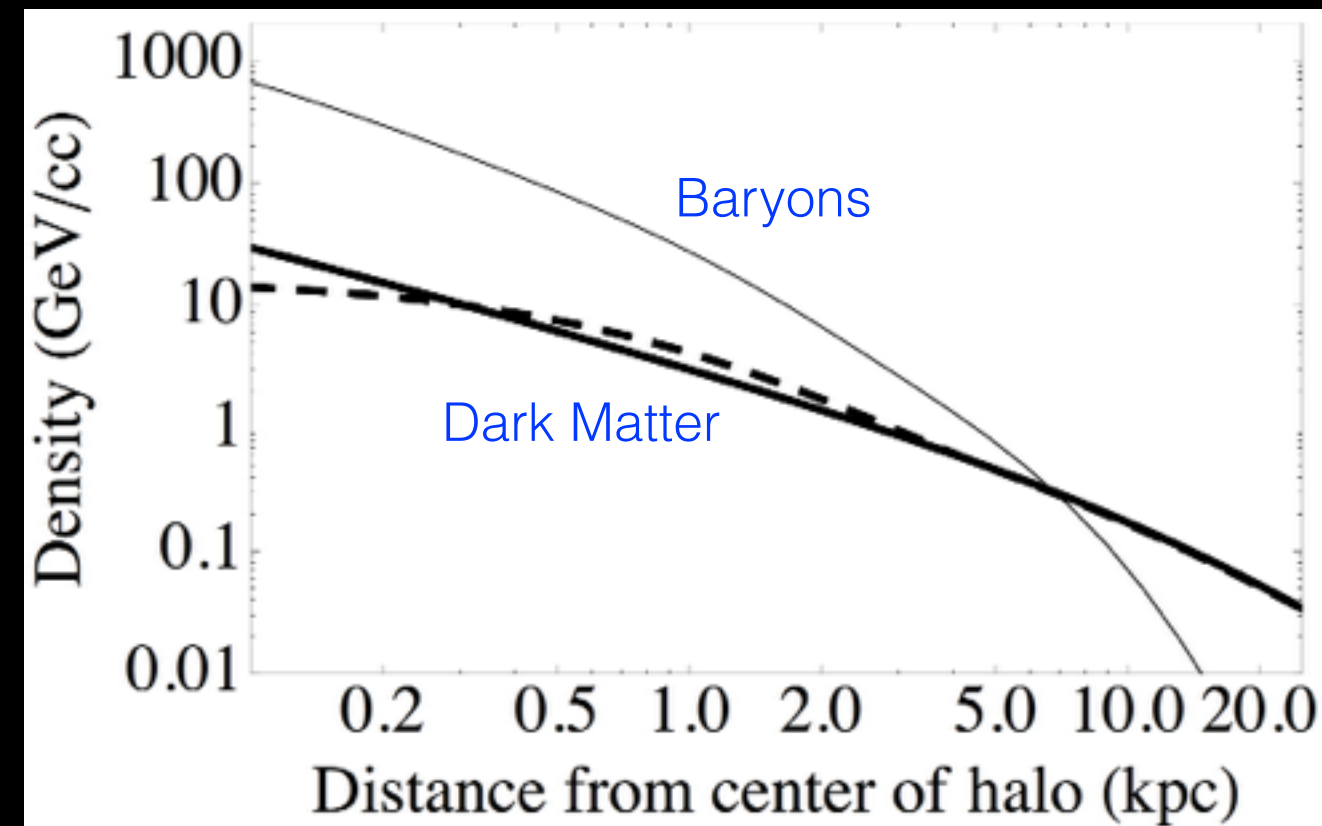


Rocha et al. (2012)

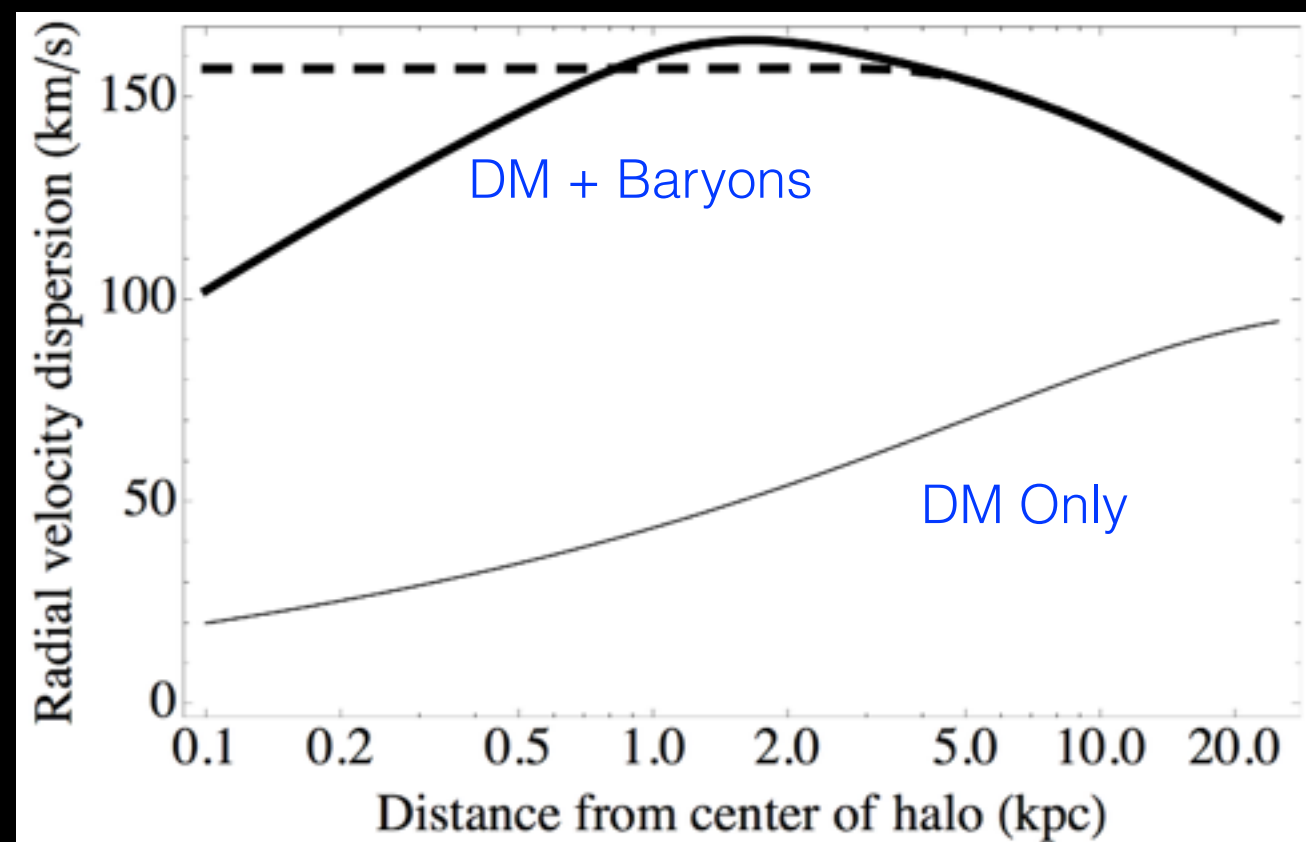


# Add a Dash of Baryons

But baryons dominate the potential within 10 kpc of the Galactic Center, they must affect the DM density profile



This means that at the SIDM critical point (1 interaction) the temperature is falling, and not rising





# Add a Dash of Baryons

$$\nabla_x^2 \left( h(\vec{x}) + \Phi_B(\vec{x})/\sigma_0^2 \right) + \frac{4\pi G_N \rho_0 r_0^2}{\sigma_0^2} \exp(h(\vec{x})) = 0$$

Use these conditions to solve the Jeans equation for the potential, which is dominated by baryons

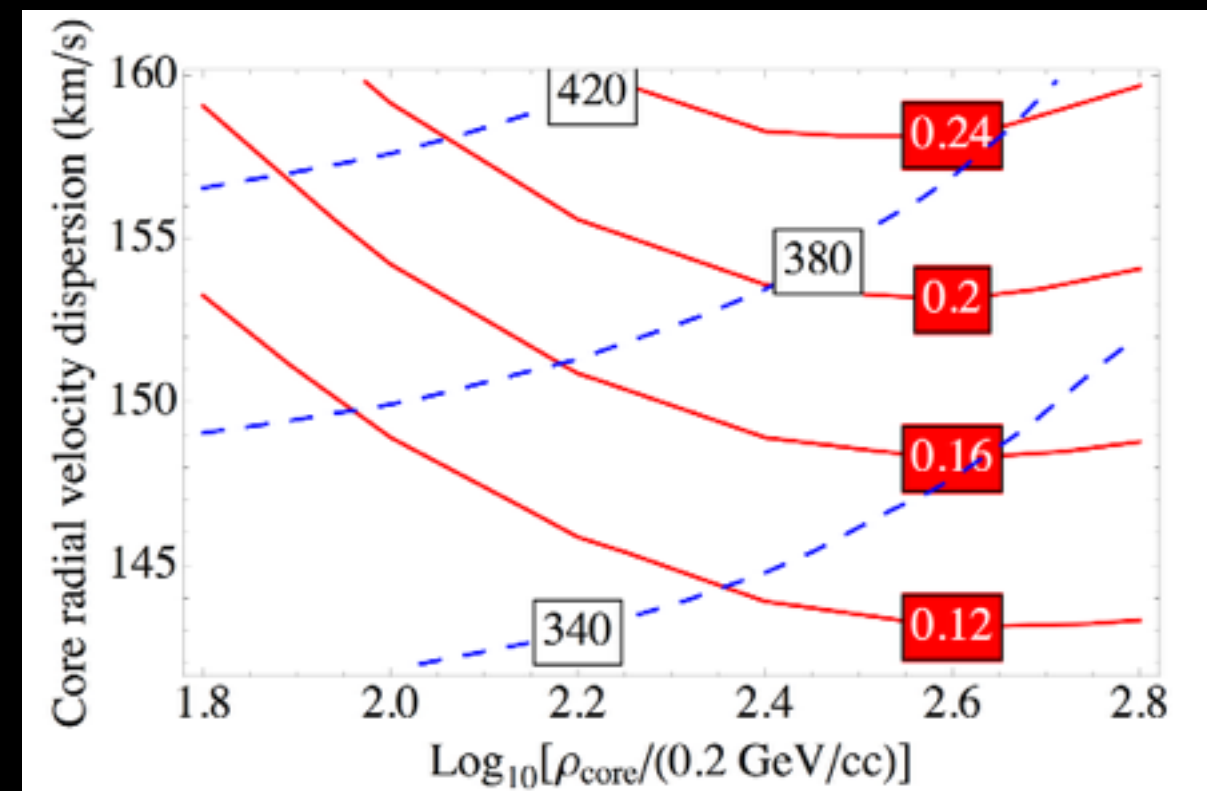
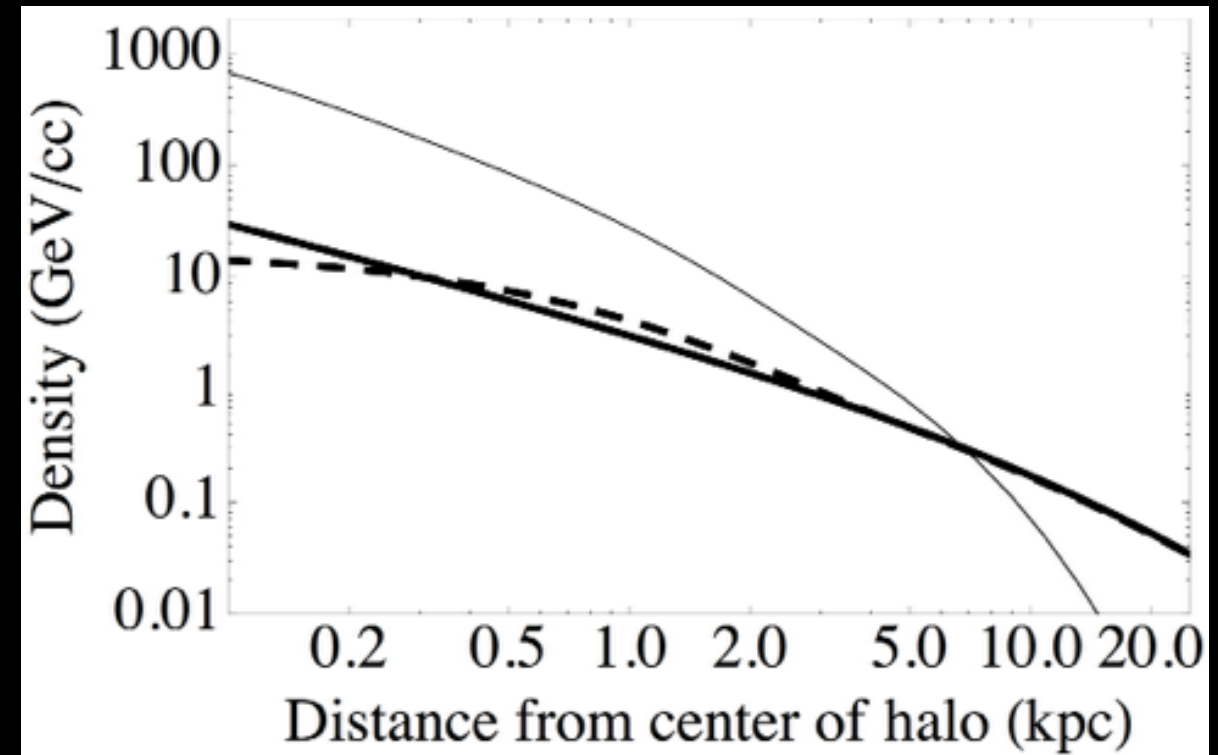
$$\frac{1}{y^2} \frac{d}{dy} \left[ y^2 \frac{d}{dy} w(y) \right] + \frac{2a_1}{y} + \frac{a_0}{(1-y)^4} \exp[w(y)] = 0$$

$$r_c \approx r_0 \frac{\sqrt{1 + (2/3) \ln(2) a_0 / a_1^2} - 1}{1 + a_0 / (3a_1) - \sqrt{1 + (2/3) \ln(2) a_0 / a_1^2}}$$

# Add a Dash of Baryons

Analytically solving the Jeans equation produces a dark matter density profile which increases until the inner few hundred pc

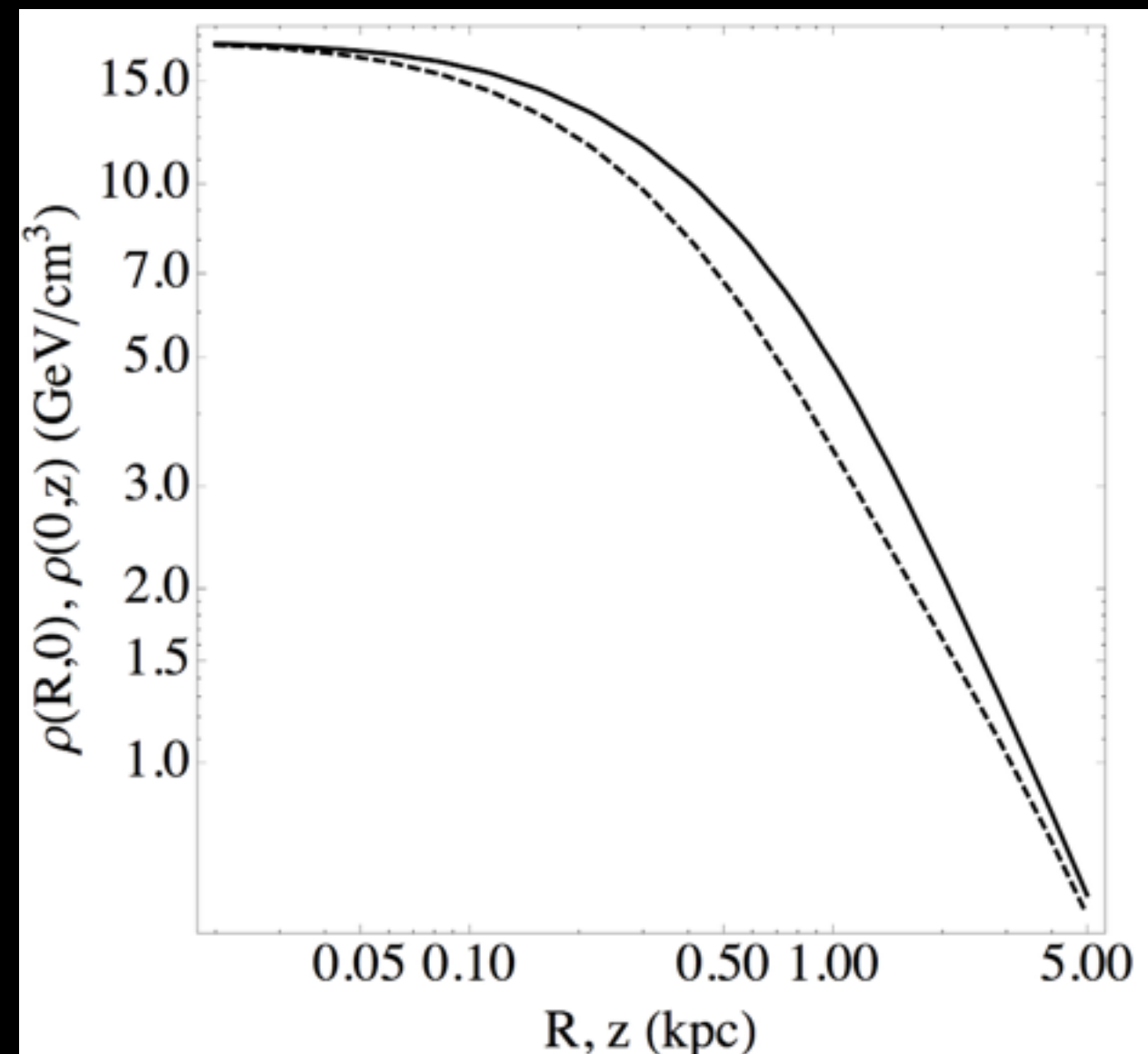
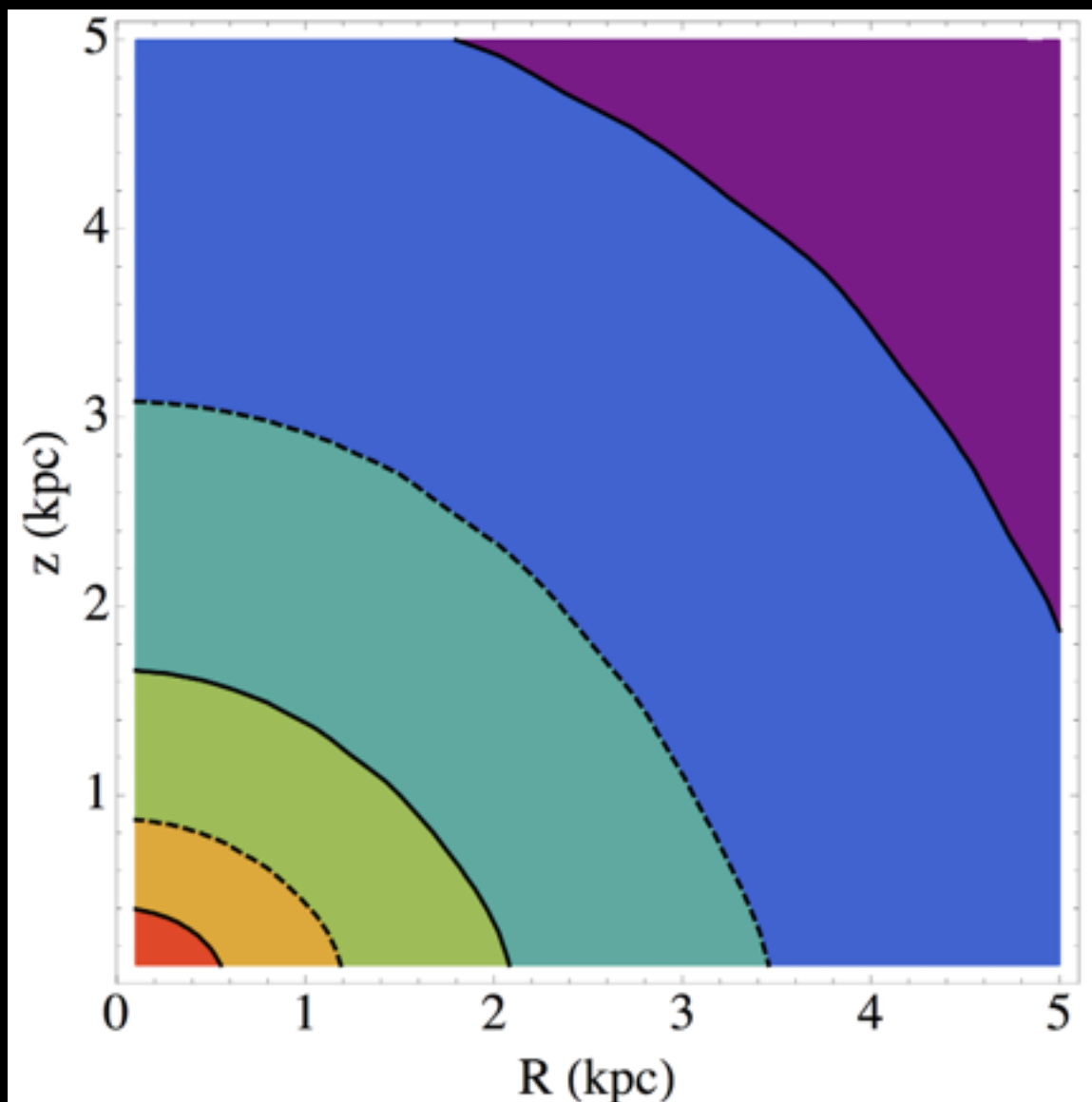
Density profile ends up being comparable to NFW - less steep within inner 200 pc, steeper from 200 - 4000 pc



# Add a Dash of Baryons

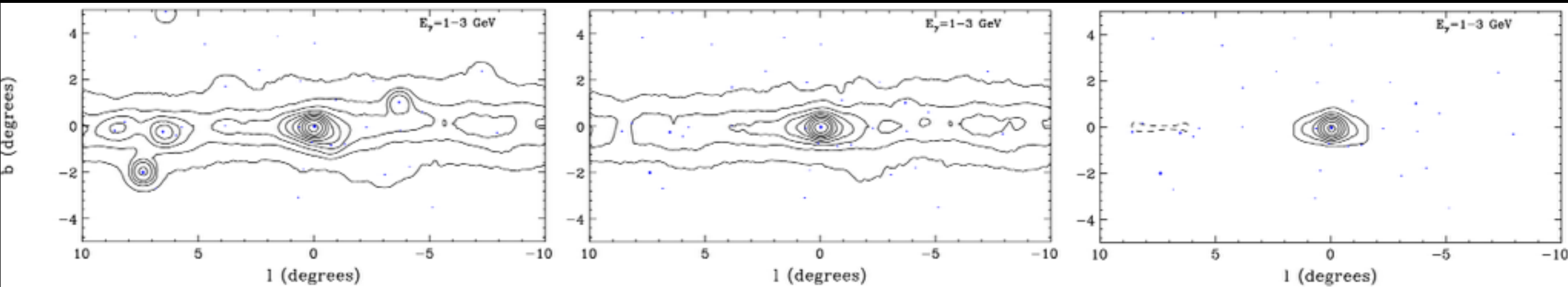
Can also solve in 2D!

$$\nabla_x^2 (h(\vec{x}) + \Phi_B(\vec{x})/\sigma_0^2) + \frac{4\pi G_N \rho_0 r_0^2}{\sigma_0^2} \exp(h(\vec{x})) = 0$$

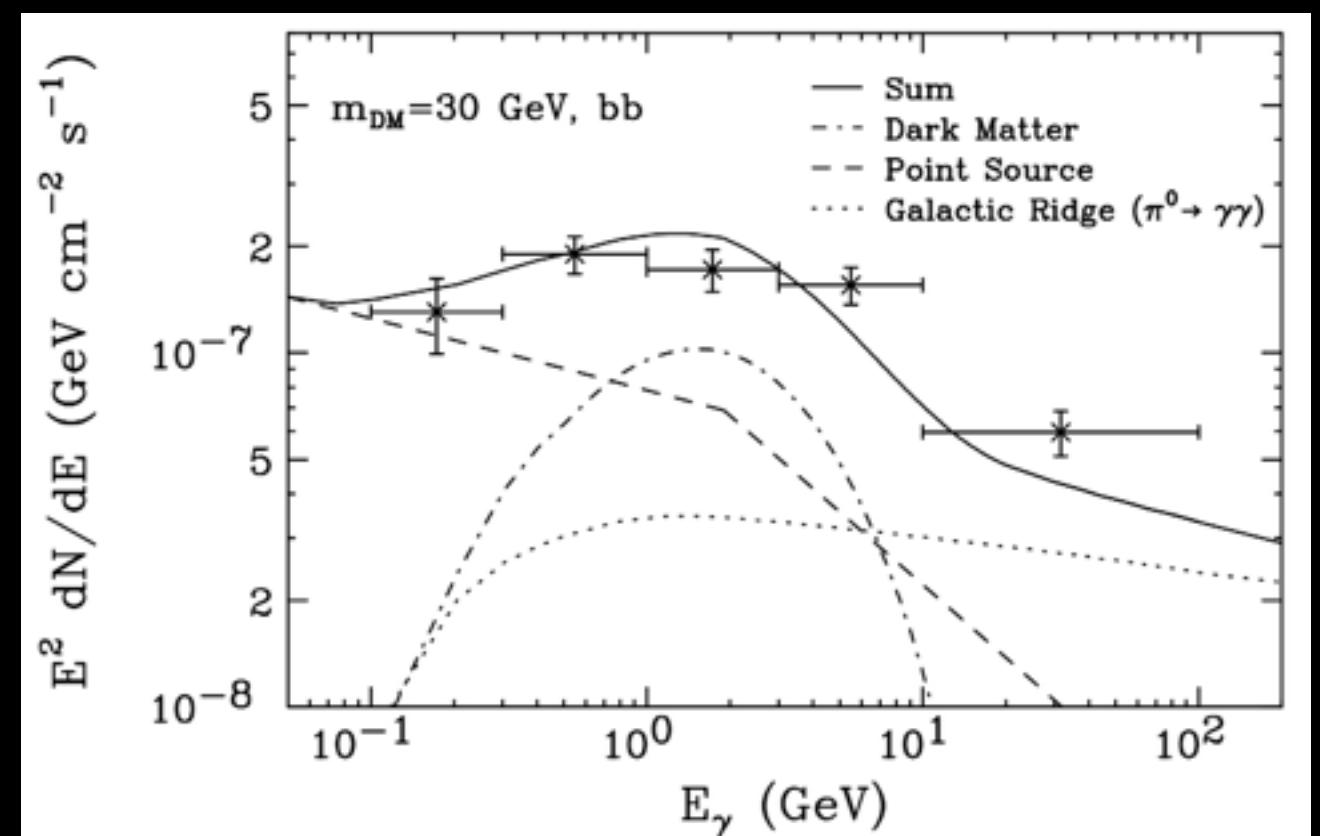
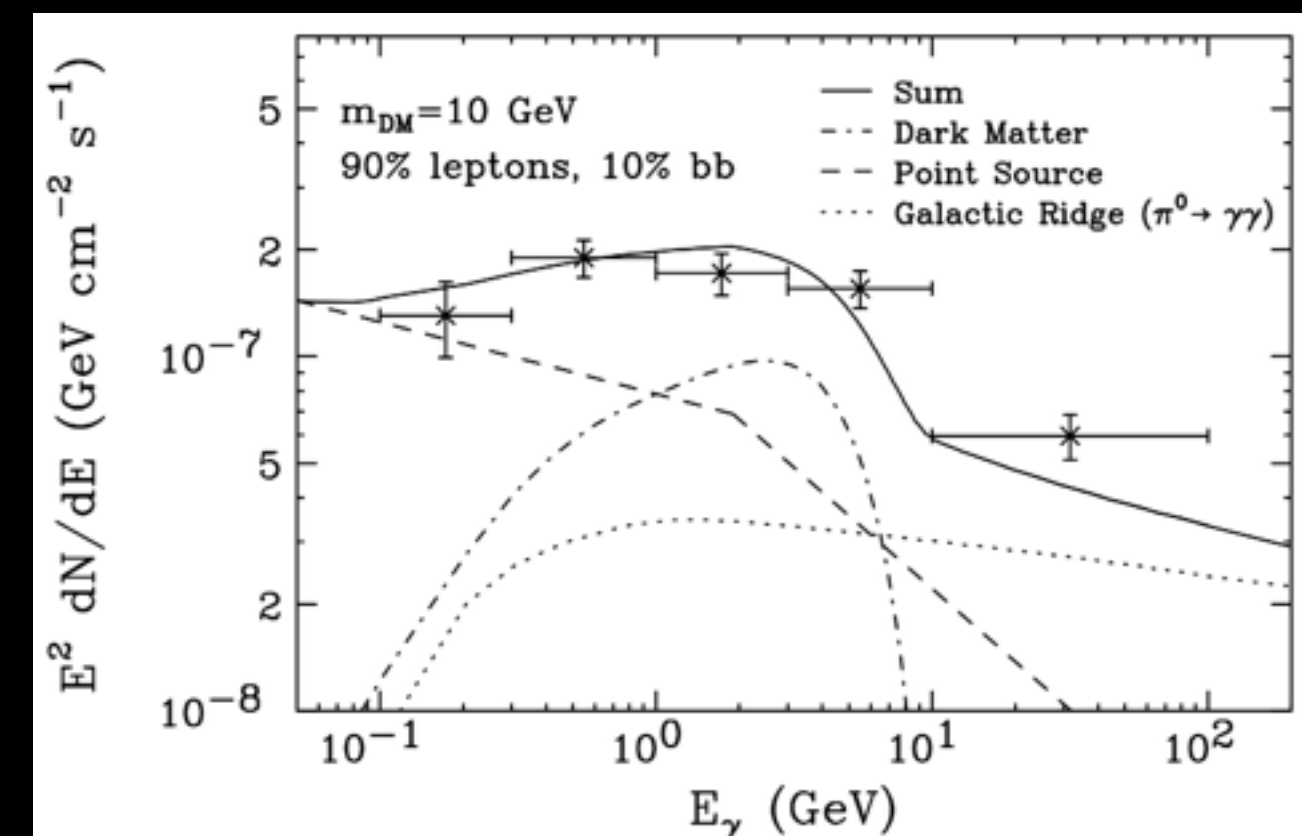




# Gamma-Rays from The Galactic Center



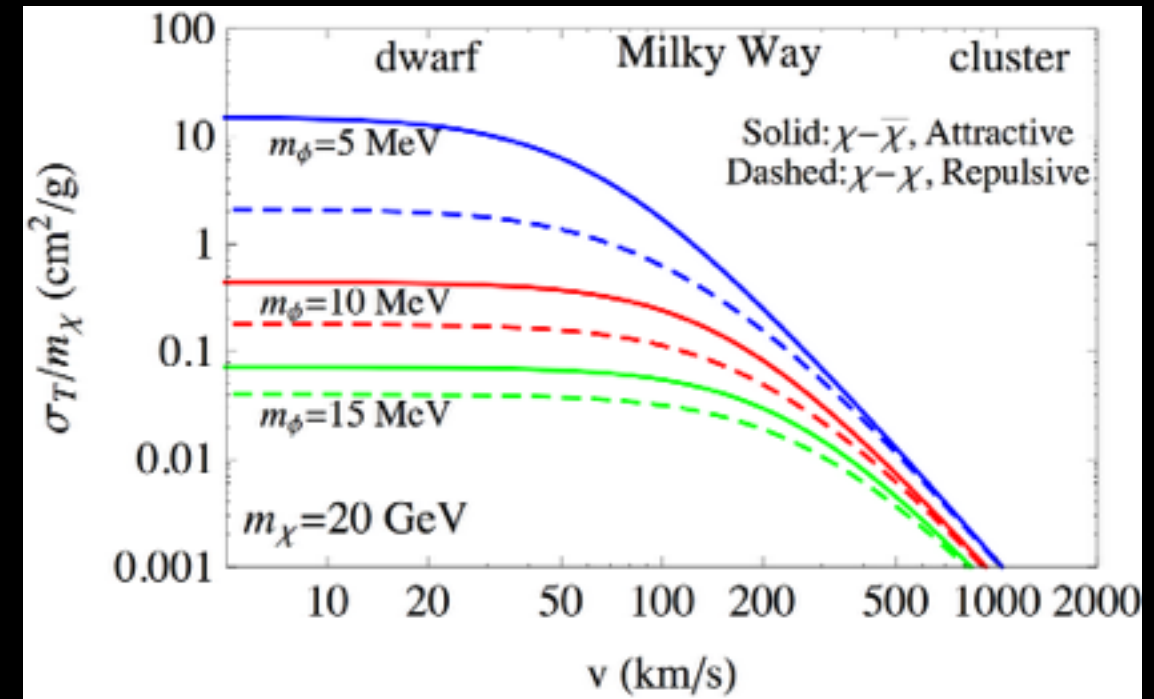
Hooper & Linden (2011)



# Particle Physics Model

Remember SIDM mediator  
mass is  $< 20$  MeV

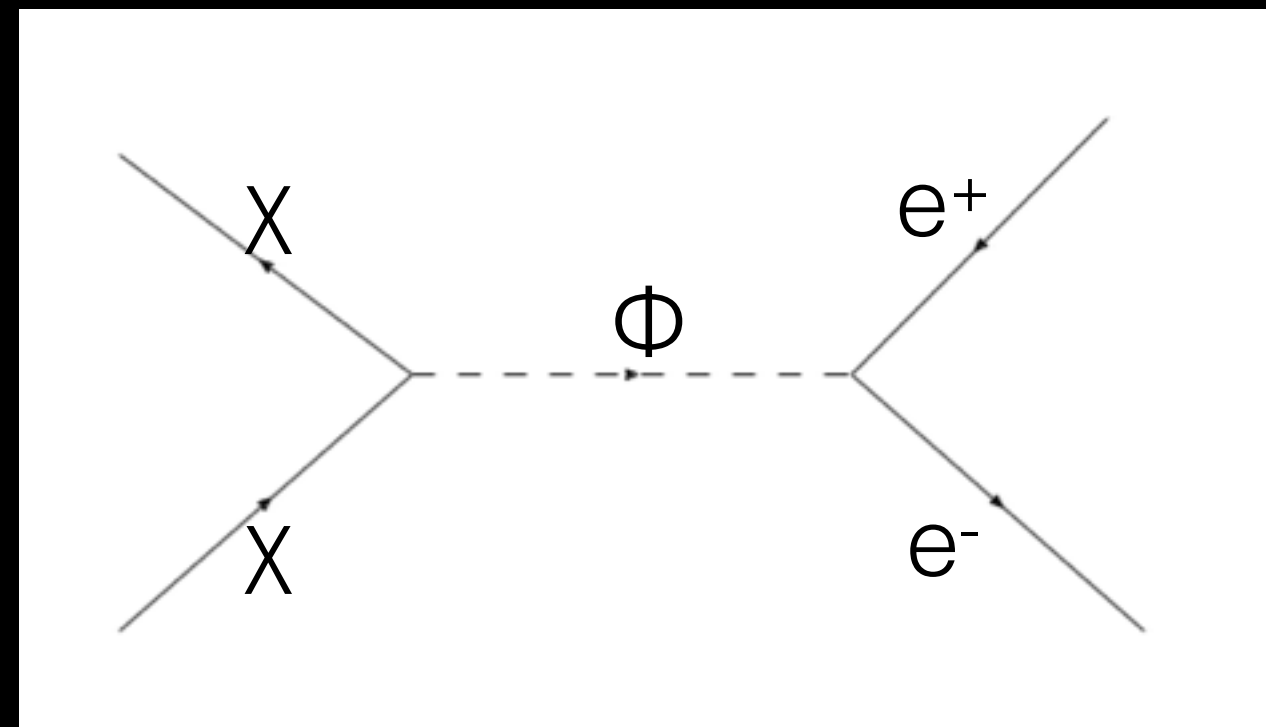
Annihilations through mediator  
can produce only  $e^+e^-$



$$\mathcal{L}_{\text{int}} = g_\chi \bar{\chi} \gamma^\mu \chi \phi_\mu + \frac{\epsilon_Y}{2} \phi^{\mu\nu} B_{\mu\nu}$$

$$\Omega_\chi h^2 \simeq 0.11 \xi_f \left[ \frac{4.4 \times 10^{-26} \text{ cm}^3/\text{s}}{\langle \sigma v \rangle_{\text{ann}}} \right]$$

$$\rho_{DM}(r) = \begin{cases} 0.2 \frac{\text{GeV}}{\text{cm}^3} \left( \frac{r_\odot^2 + r_c^2}{r^2 + r_c^2} \right), & r \leq r_1 \simeq 15 \text{ kpc.} \\ 0.065 \frac{\text{GeV}}{\text{cm}^3} \left( \frac{r_1(r_1 + r_s)^2}{r(r + r_s)^2} \right), & r > r_1 \text{ kpc.} \end{cases}$$



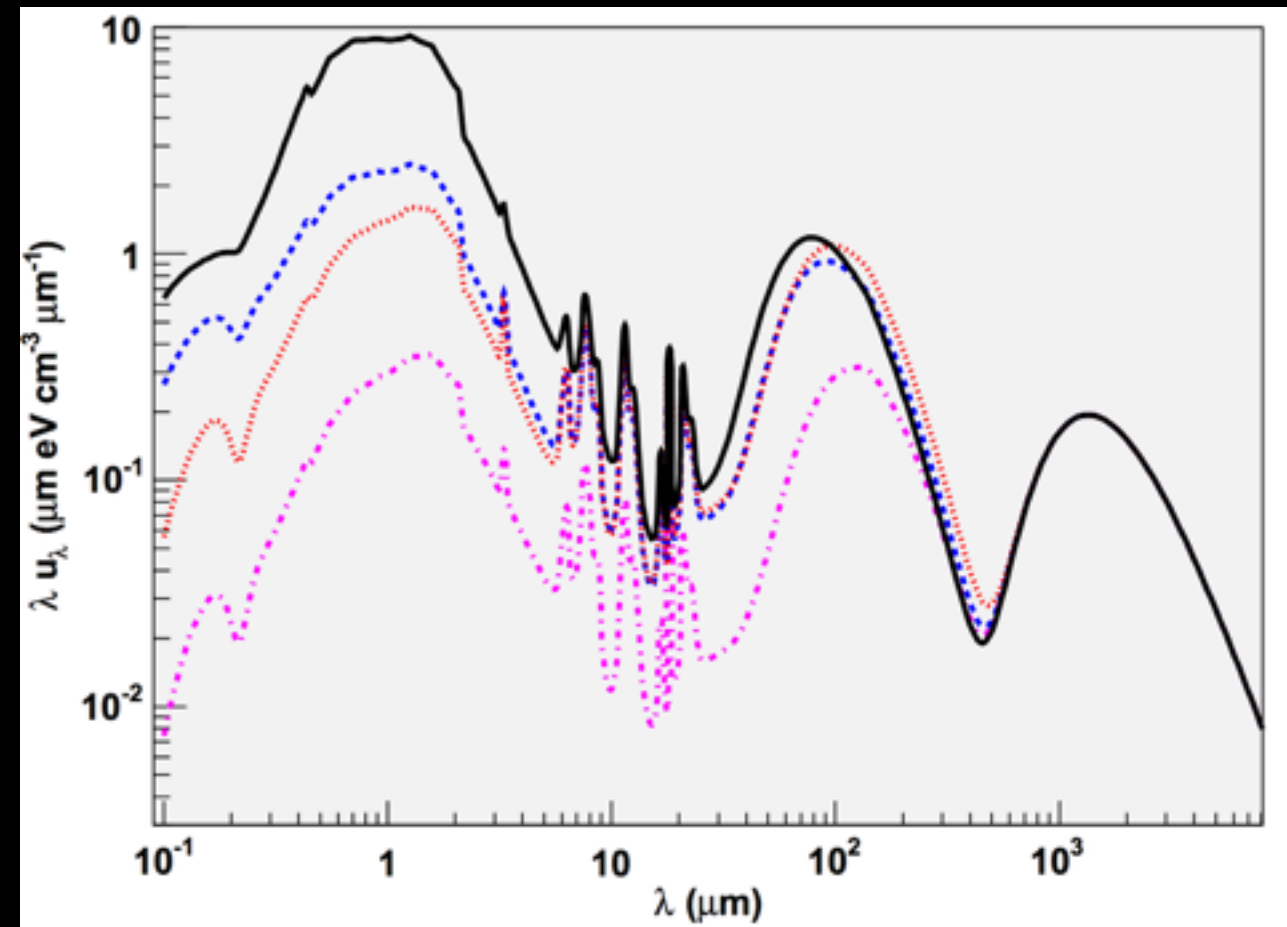
We choose  $M_{DM} = 20$  GeV

# Gamma-Rays from The Galactic Center

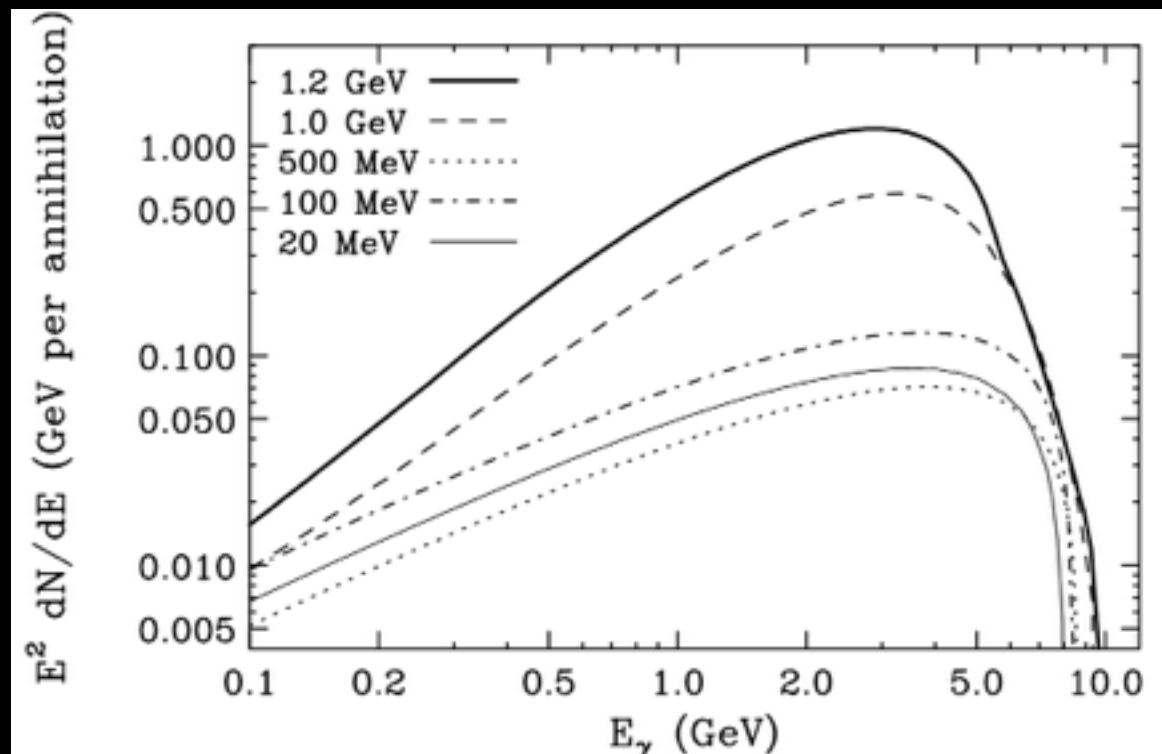
Now that SIDM has a large dark matter density around the GC, can you explain the GC excess with SIDM?

Annihilation to  $e^+e^-$  only

Need some mechanism to produce gamma-rays



Porter et al. (2006)



What about Inverse Compton Scattering?

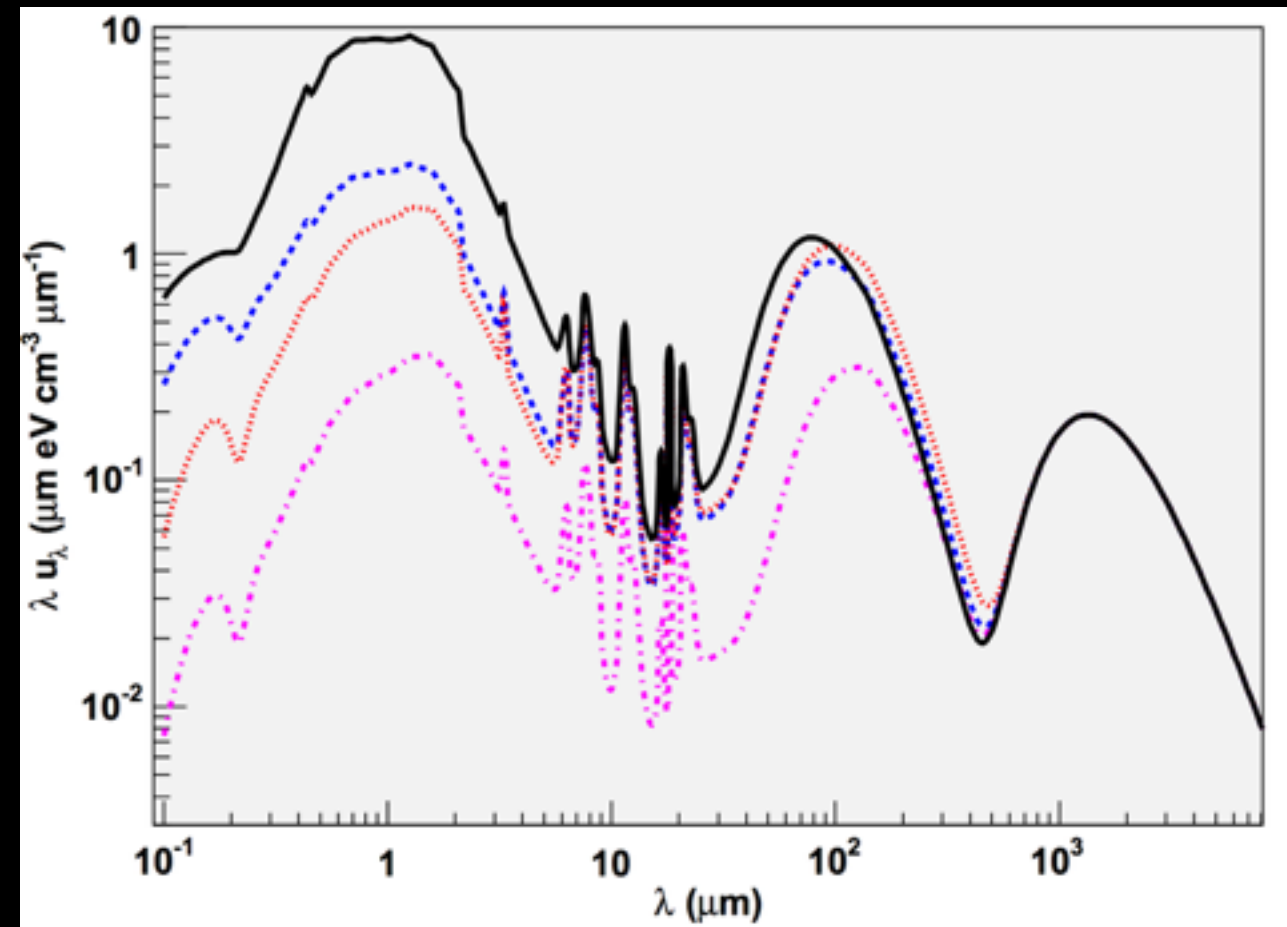
FSR intensity is small

Hooper et al. (2012)



# Gamma-Rays from The Galactic Center

In fact, the ISRF near the GC is much more intense than this, and dominates the magnetic field density within 10 pc of the GC



Porter et al. (2006)

## PHYSICAL CONDITIONS IN PHOTODISSOCIATION REGIONS: APPLICATION TO GALACTIC NUCLEI

MARK G. WOLFIRE,<sup>1,2</sup> A. G. G. M. TIELENS,<sup>2</sup> AND DAVID HOLLENBACH<sup>2</sup>

*Received 1989 June 6; accepted 1990 January 23*

Within 5 pc of the Galactic center we find  $\sim 100$  clouds of size  $r \approx 0.4$  pc and density  $n \approx 10^5 \text{ cm}^{-3}$ . A far-ultraviolet radiation field, most likely from a central source with  $L \approx 2\text{--}3 \times 10^7 L_\odot$ , illuminates the clouds with an intensity  $\sim 10^5$  times greater than the local Galactic field and heats gas in the surface layers to  $\sim 700$  K. The gas phase Si abundance is  $\lesssim 4.7 \times 10^{-5}$  in the atomic layers of these clouds. For the case of M82, we

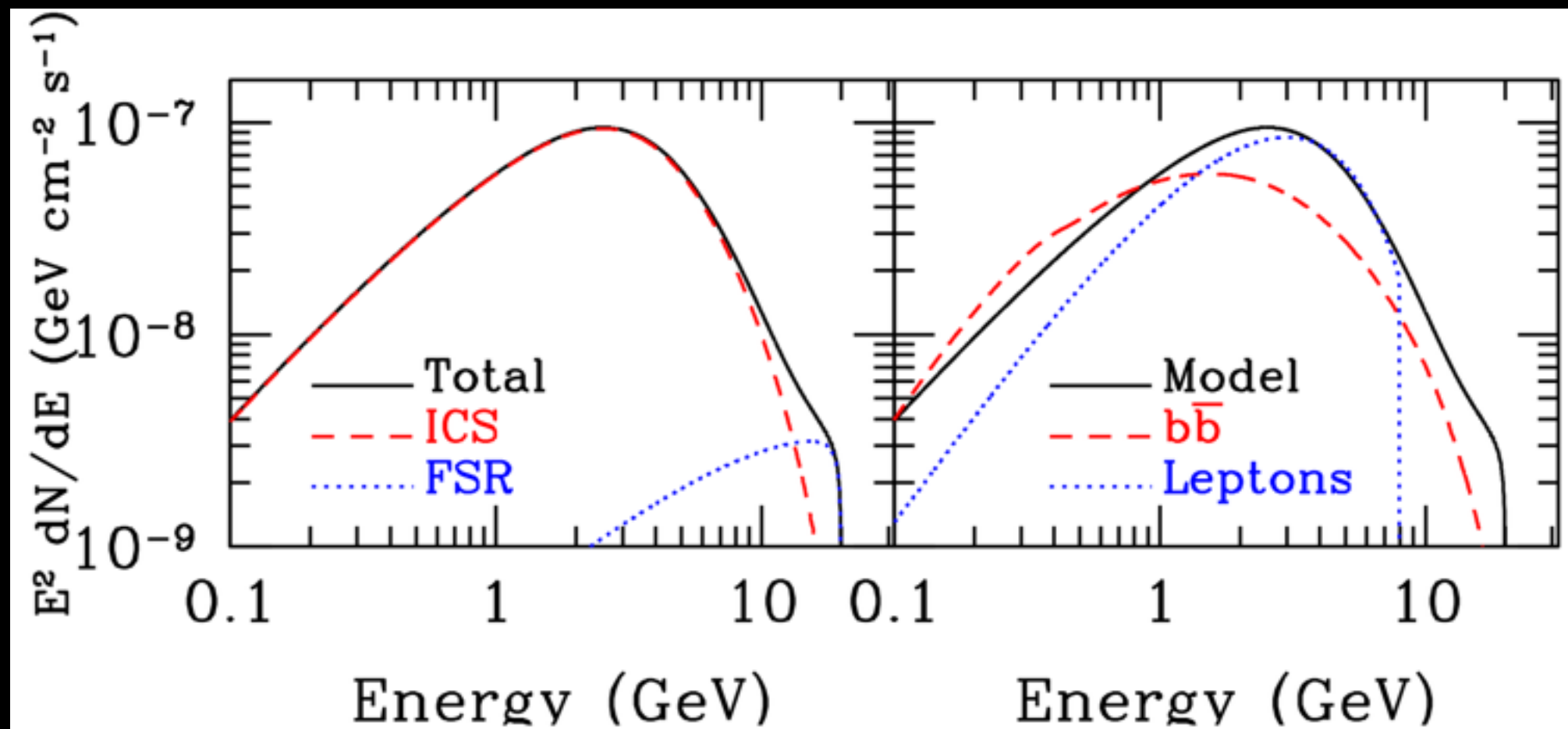
# The Gamma-Ray Spectrum

Multiple Uncertainties:

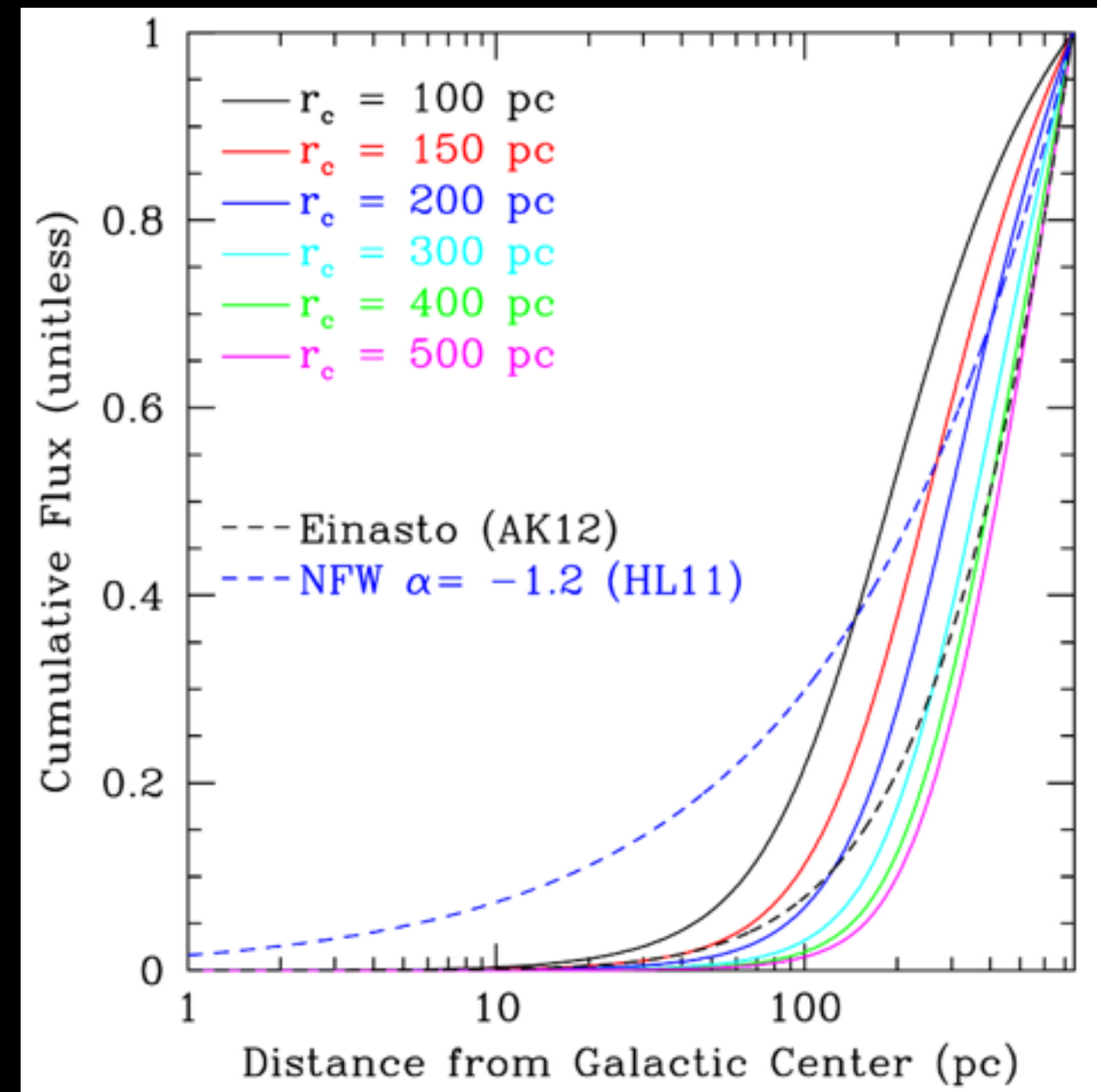
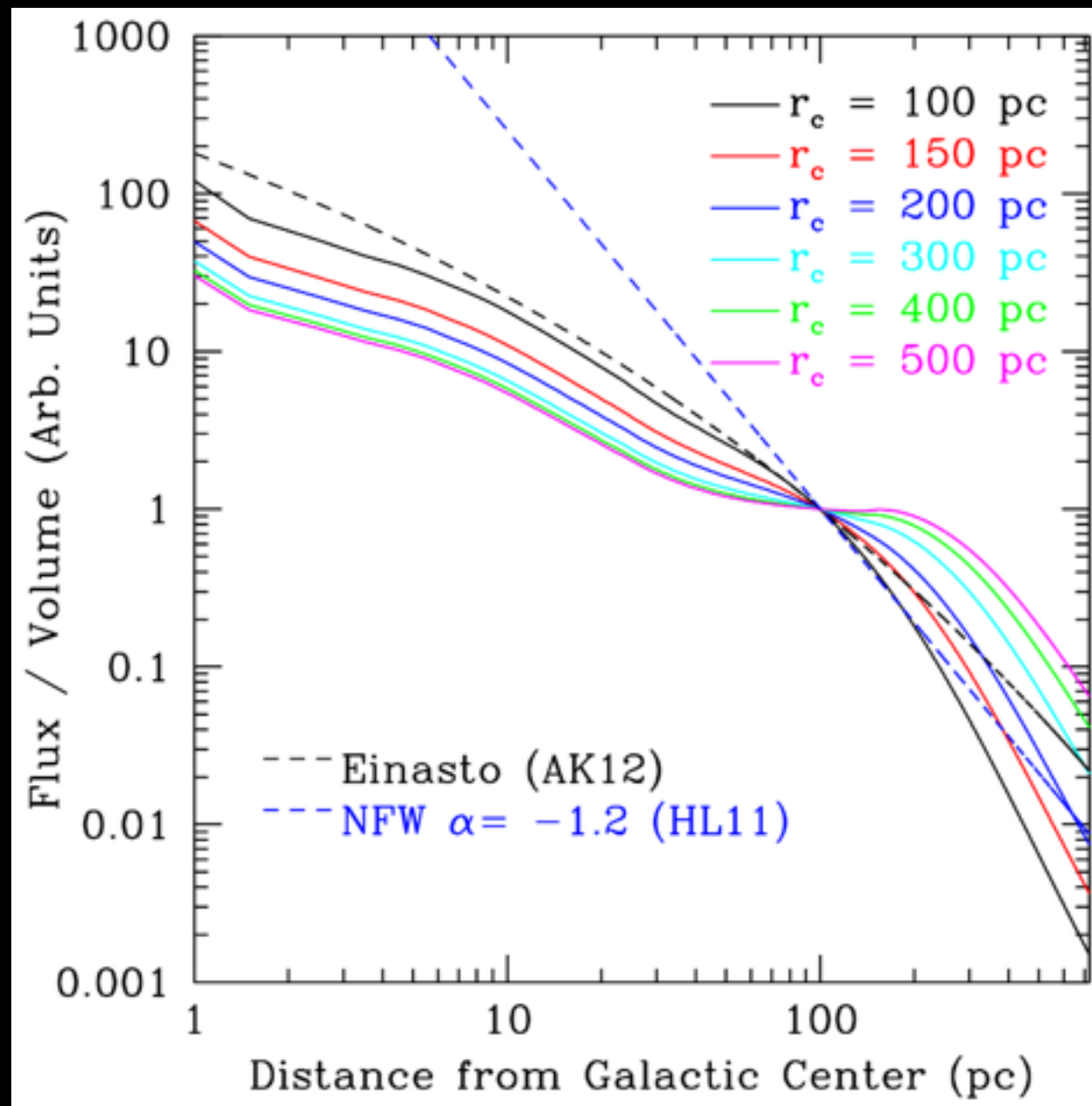
Electron Diffusion Length and Energy Loss

Magnetic Field Strength and ISRF Energy Density

Plus Normal Uncertainties (Cross-Section, DM Density, Mass etc.)



# The Gamma-Ray Morphology





# Local Positron Excess

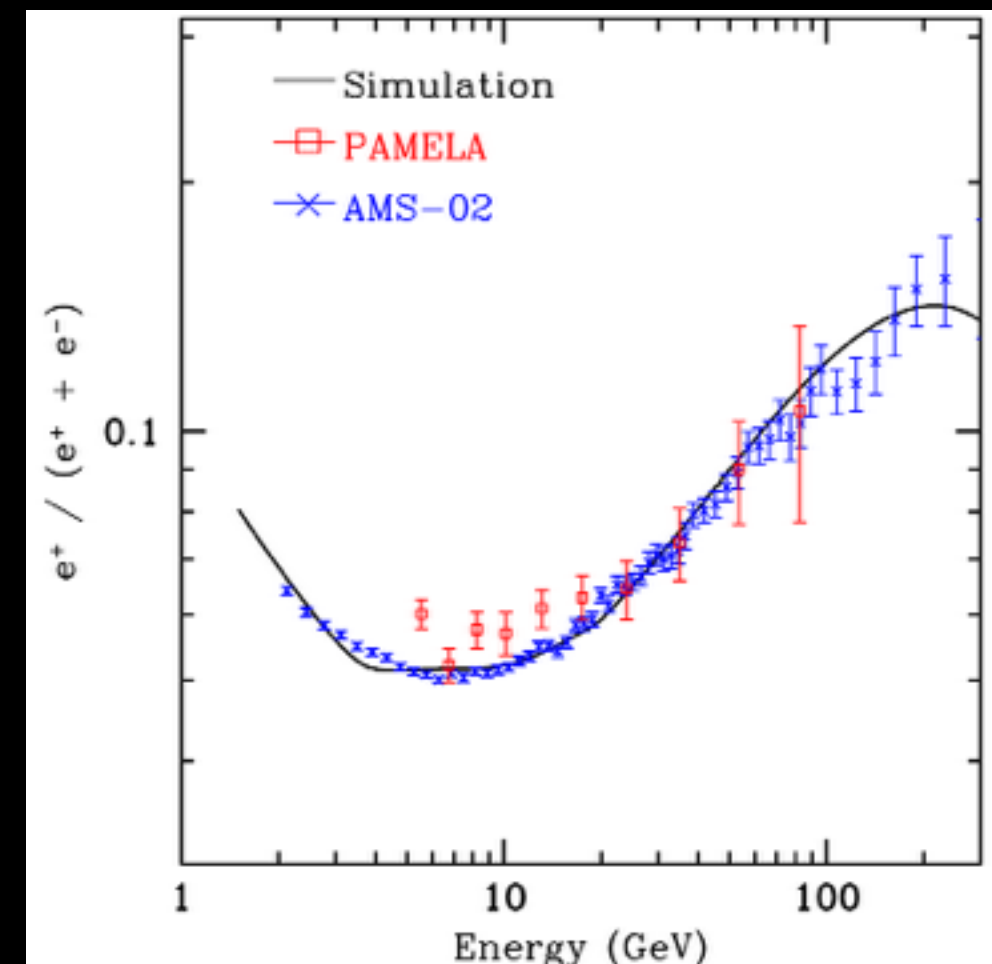
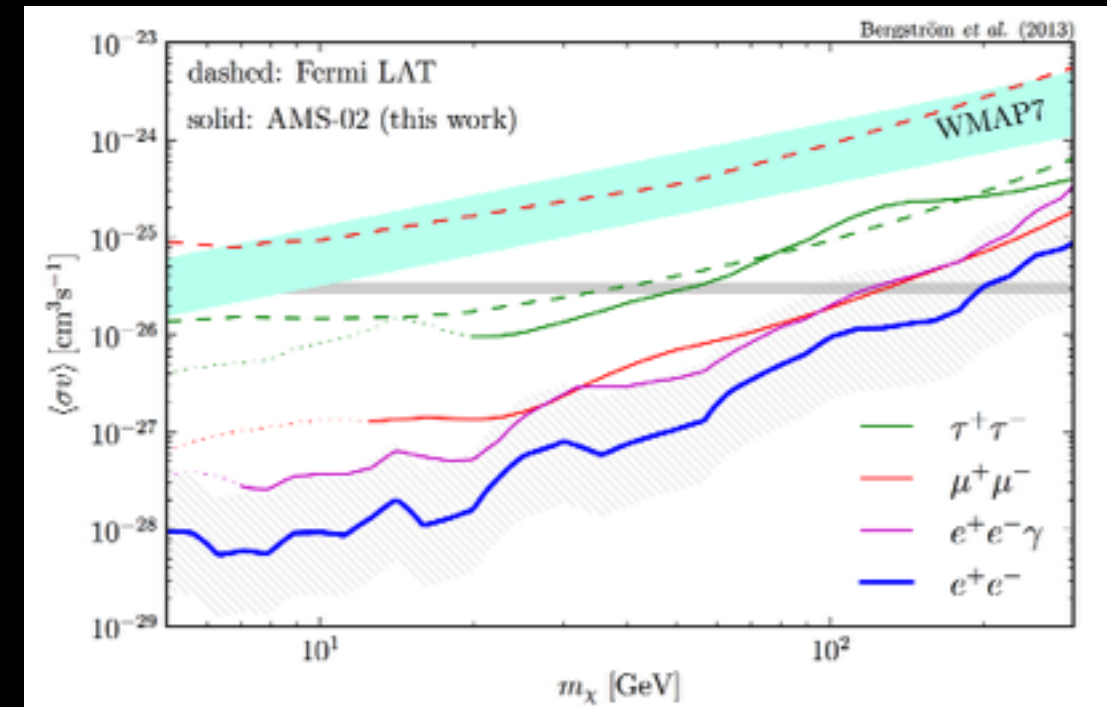
This SIDM model should also be highly detectable by AMS-02, as a bump in the positron excess

This is currently avoidable so long as:

DM Annihilation is sub thermal  
( $< 6 \times 10^{-27} \text{ cm}^3\text{s}^{-1}$ )

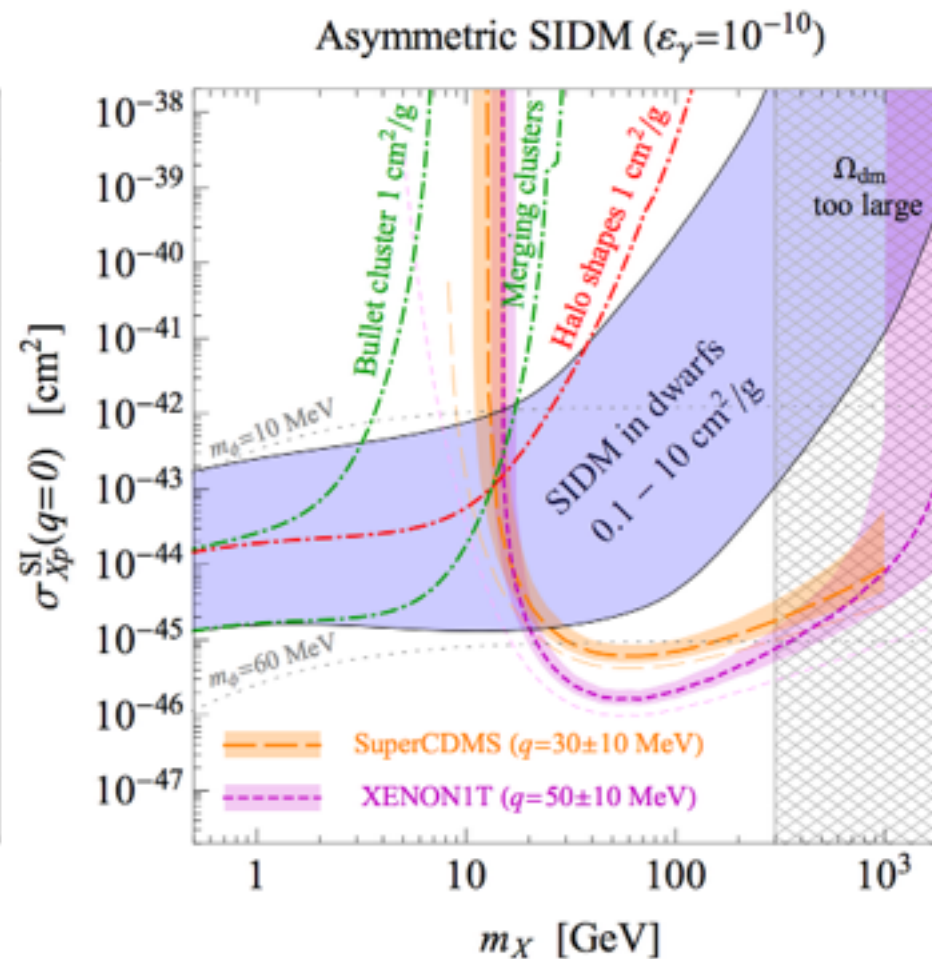
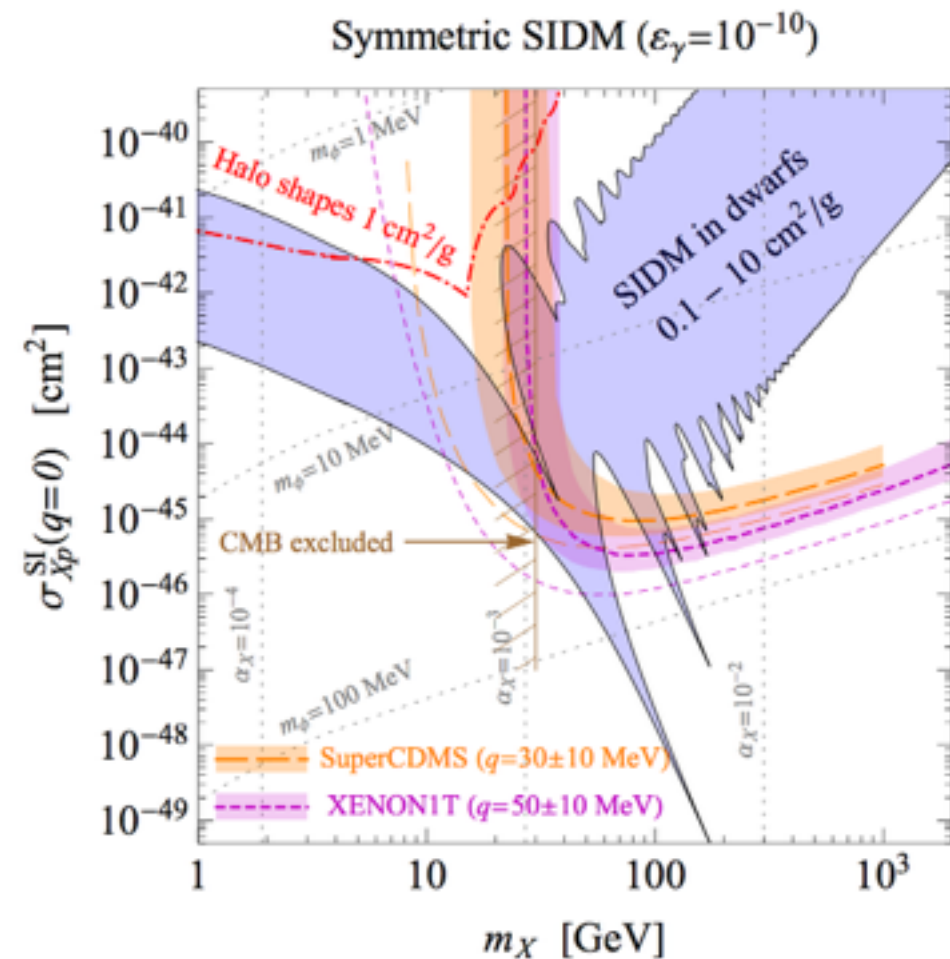
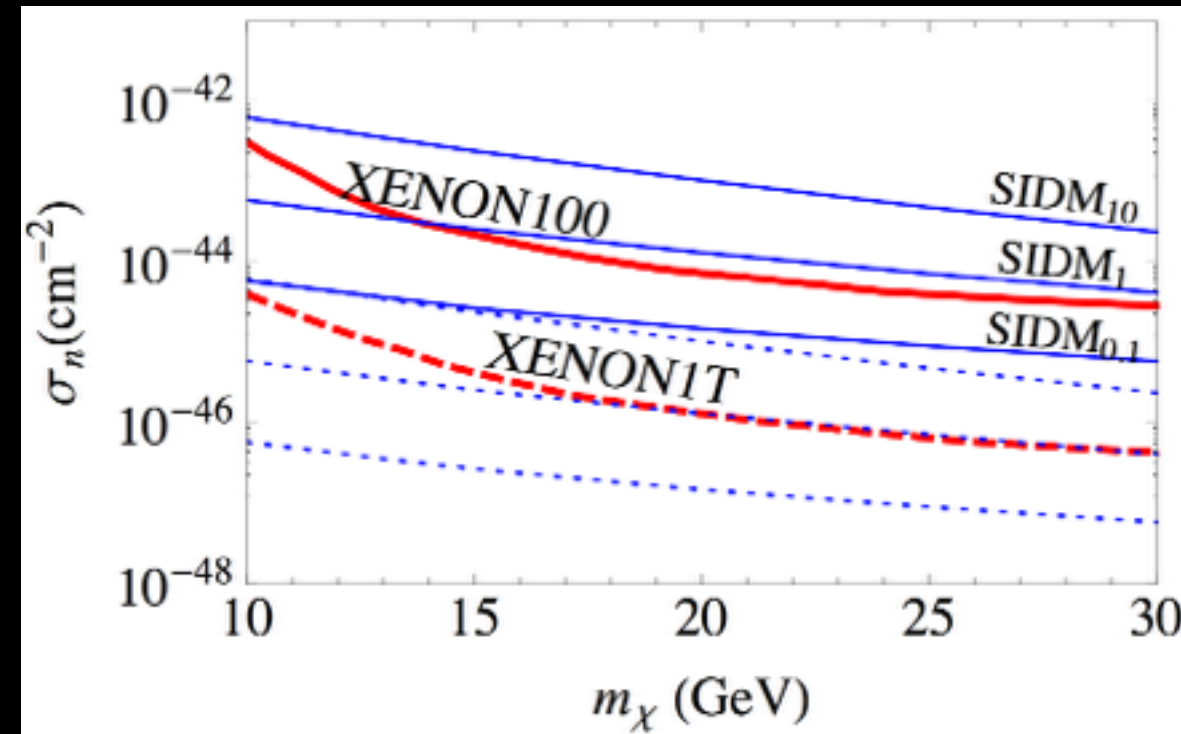
The local density is slightly low  
( $\sim 0.2 \text{ GeV cm}^{-3}$ )

Should strongly test models in next few years



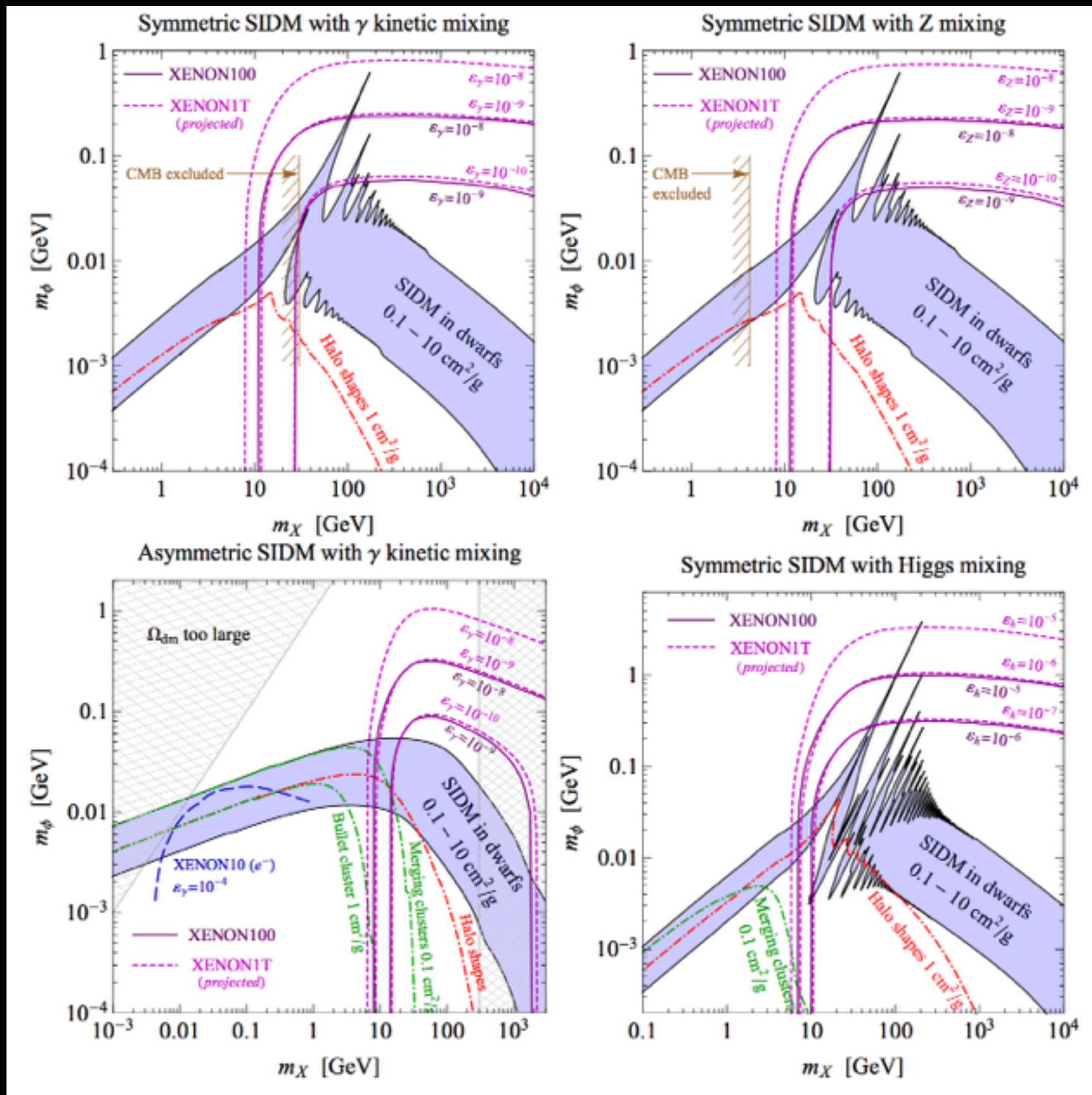
# Dark Matter Direct Detection

Future ton-scale direct detection experiments can also rule out much of the SIDM parameter space





# Dark Matter Direct Detection



# Conclusions

SIDM is an interesting empirical dark matter model, it solves many problems with small scale structure

Until recently, it was thought that the Indirect Detection of SIDM was difficult, since high densities would not be observed - but this is not true in systems dominated by a baryonic potential

Interestingly, SIDM produces very good fits to the observed GC excess