

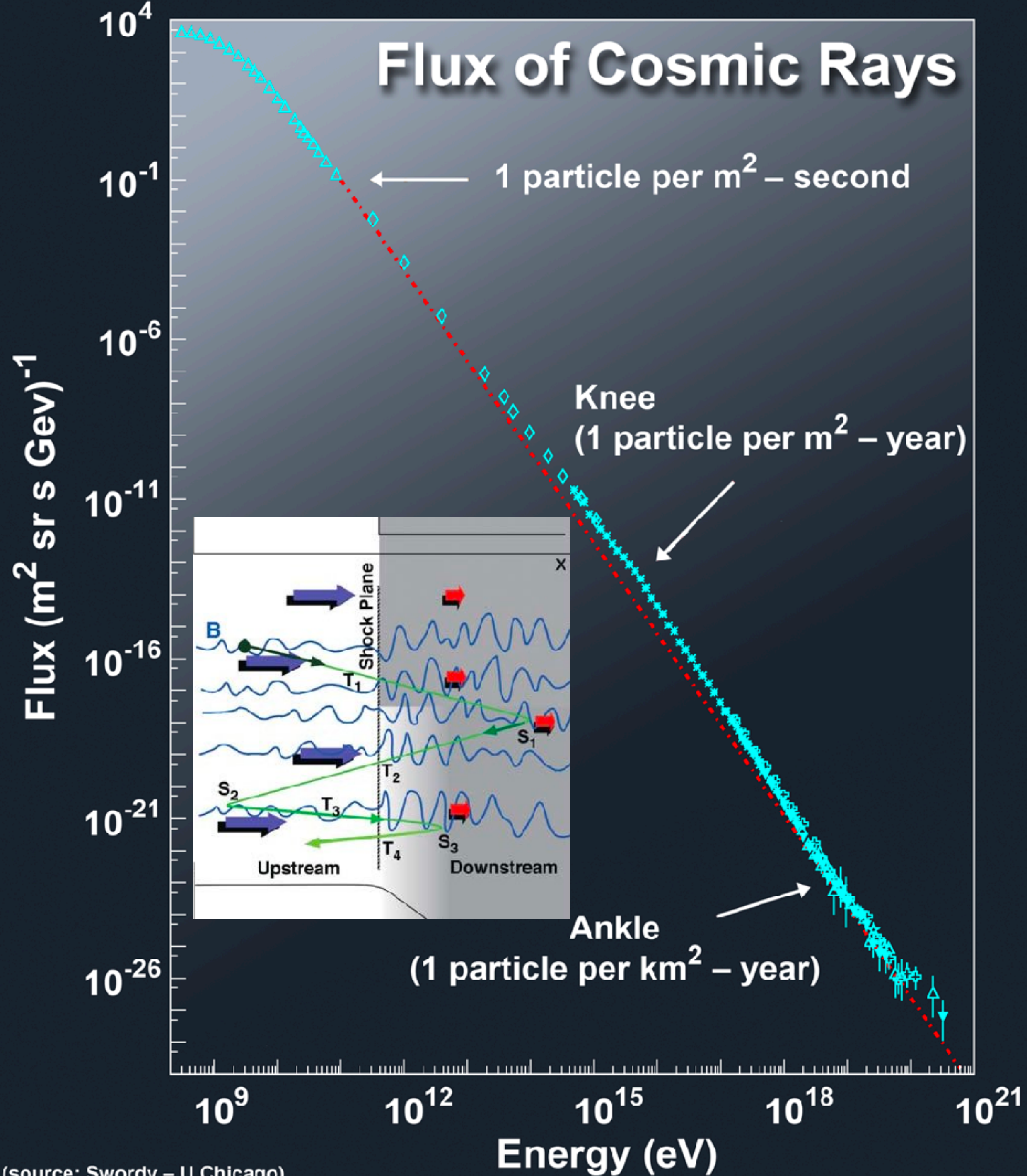


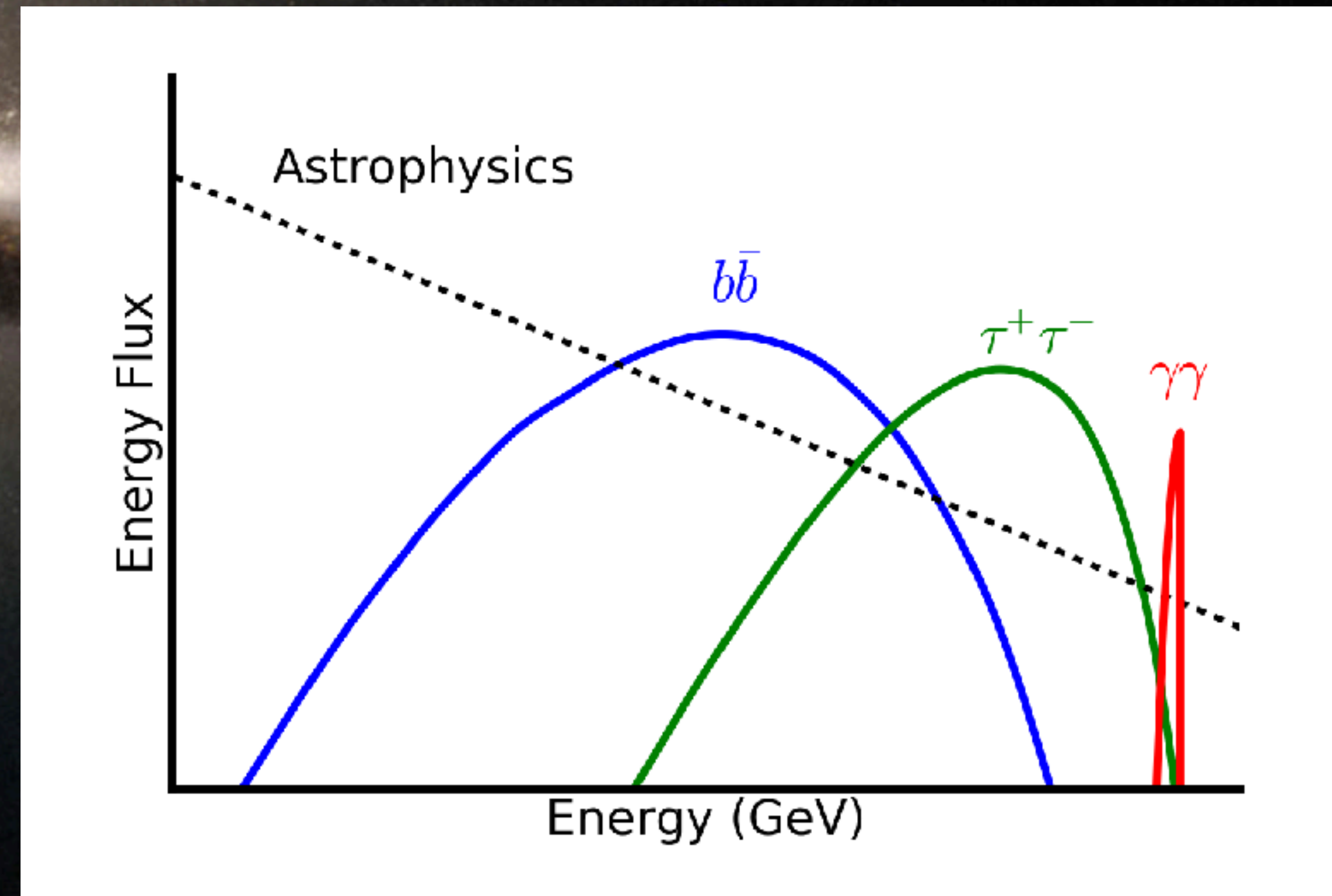
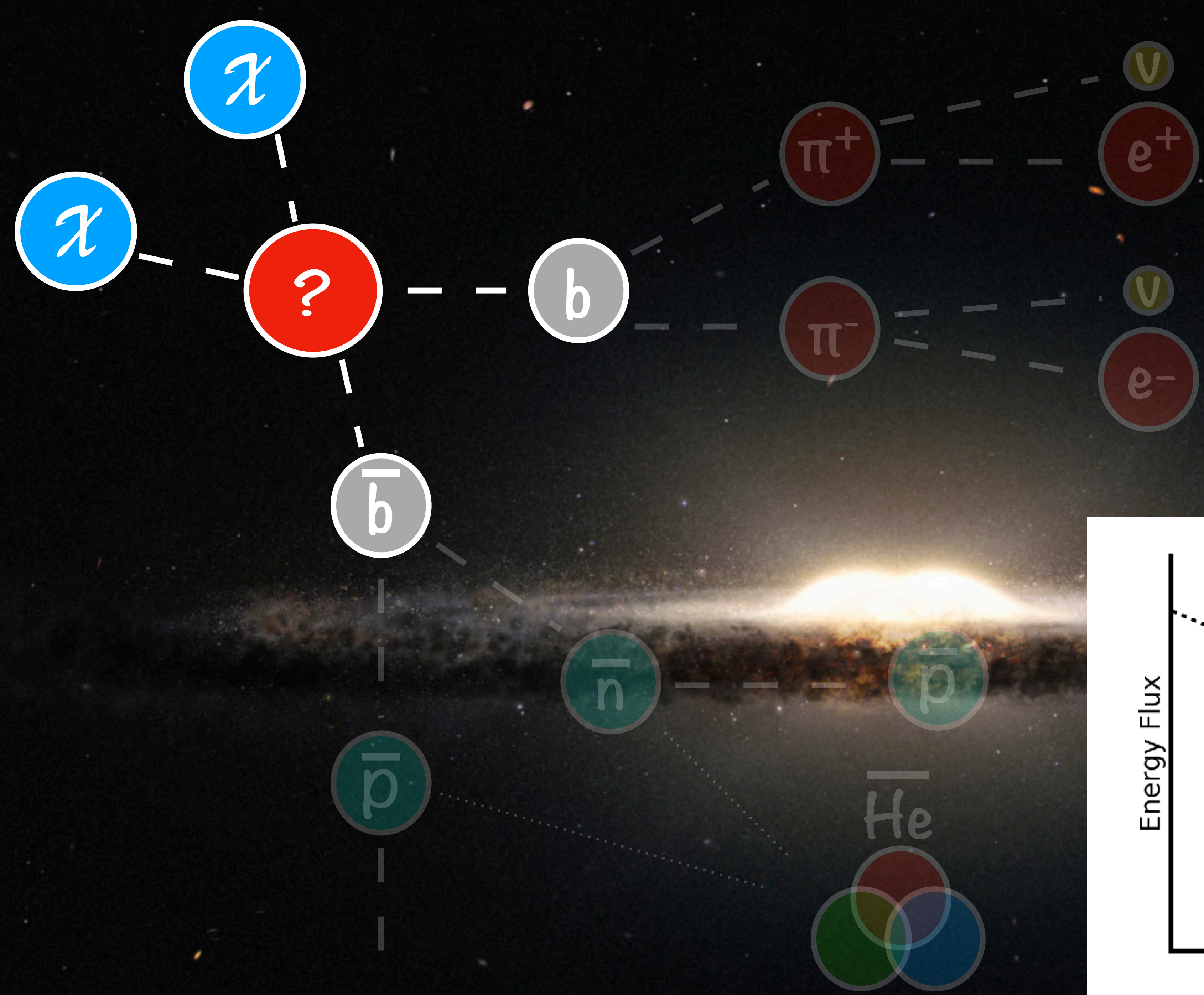
**Tim Linden**

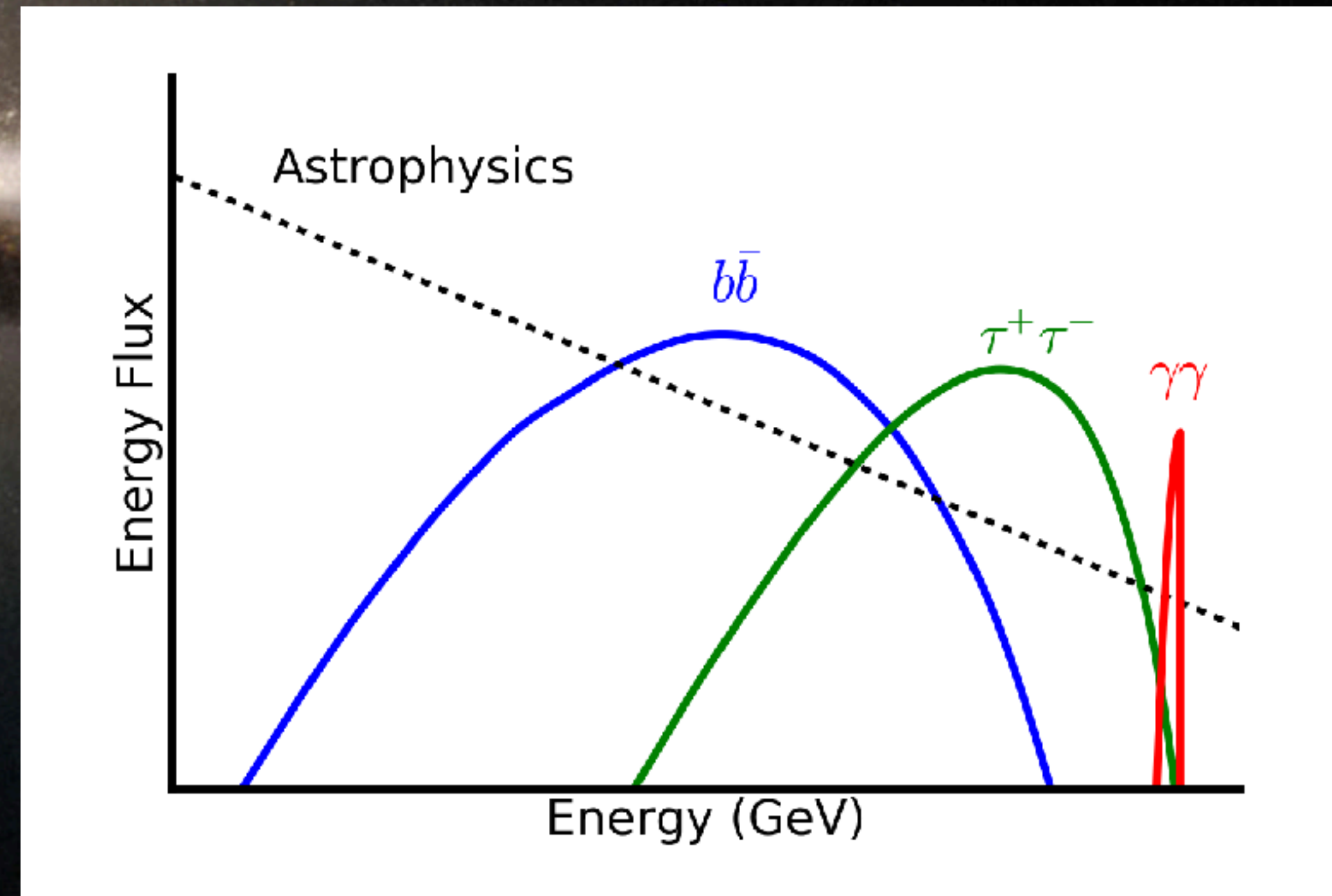
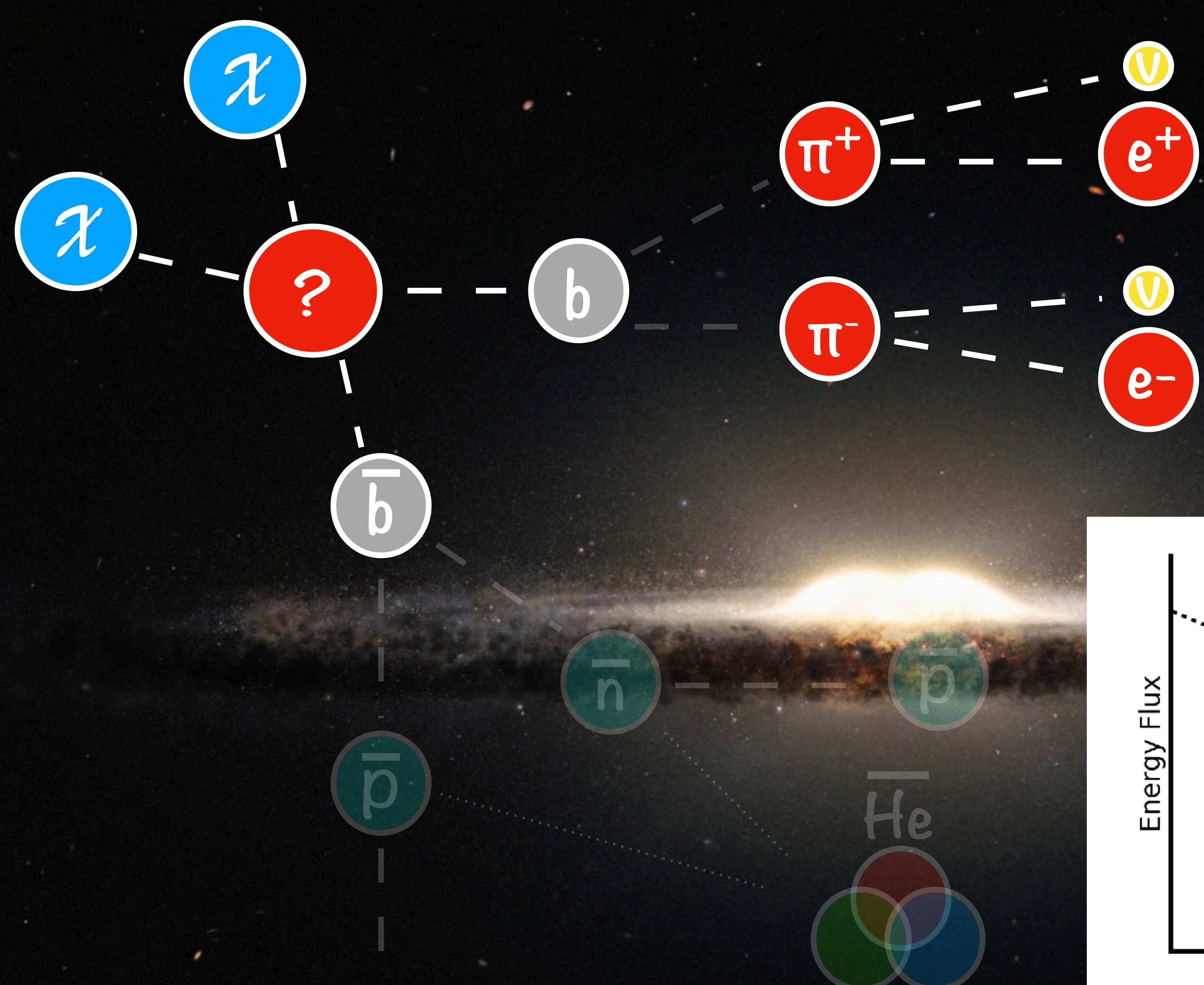


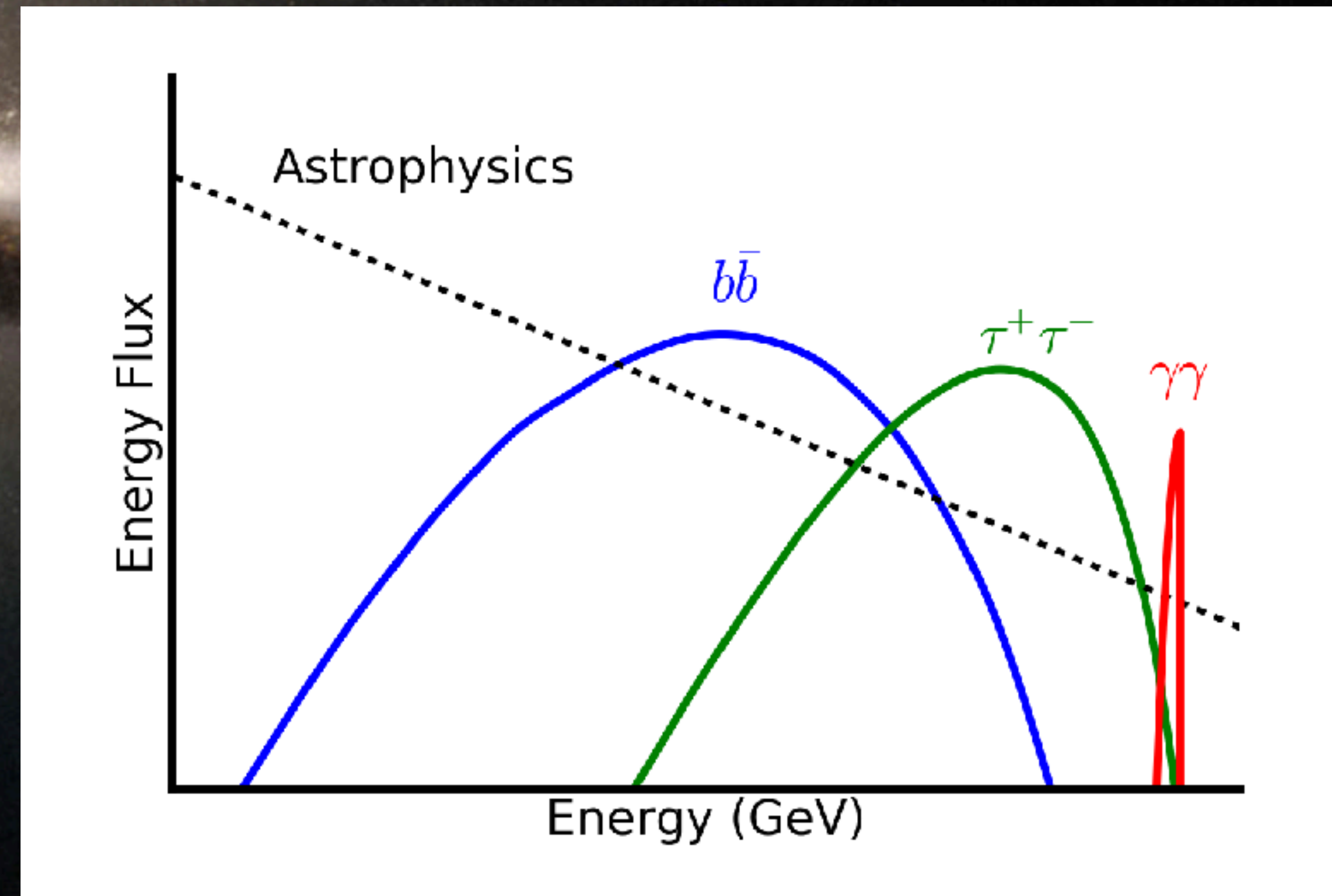
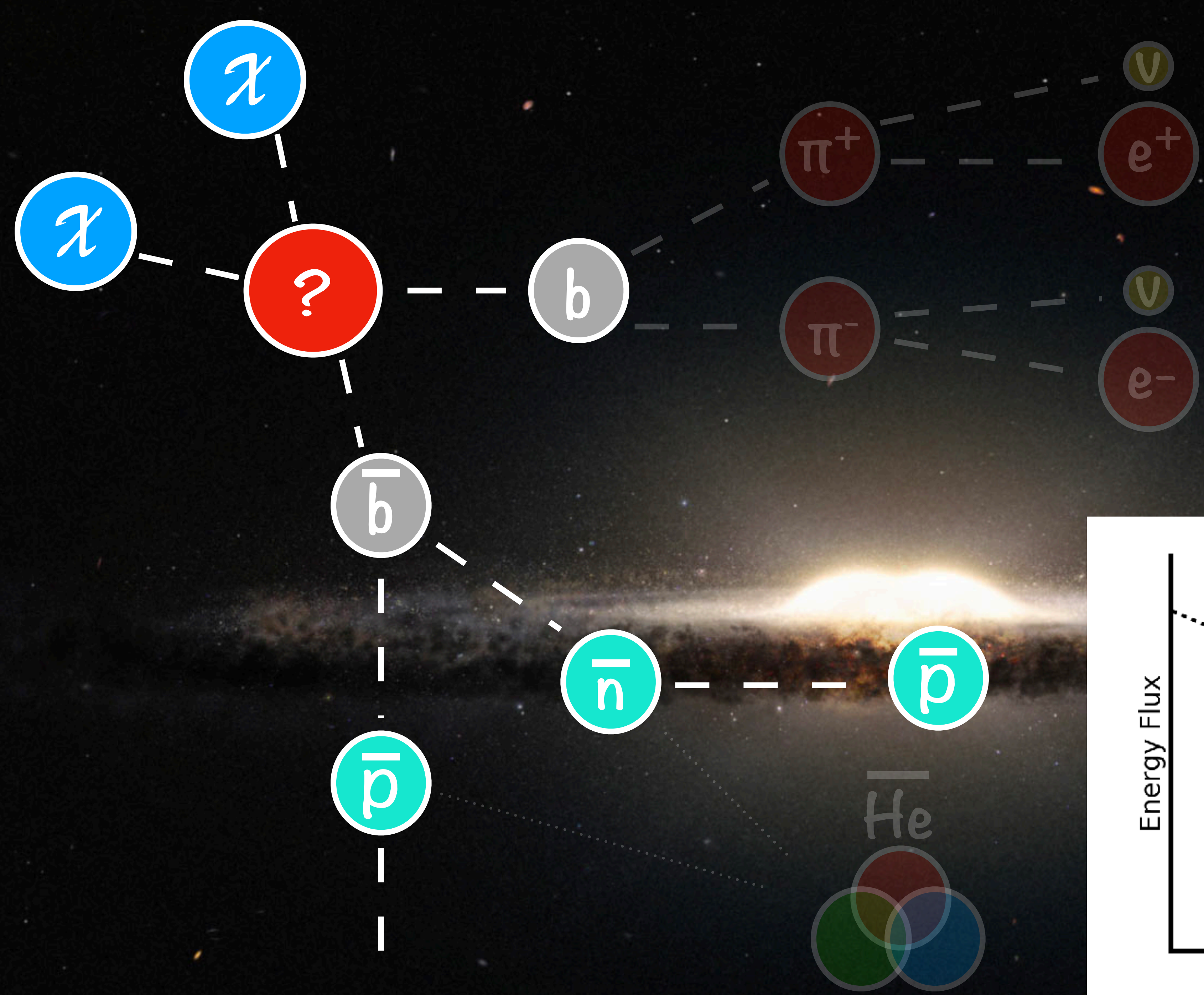
# **Cosmic-Ray Searches for Dark Matter: Yesterday, Today and Tomorrow**

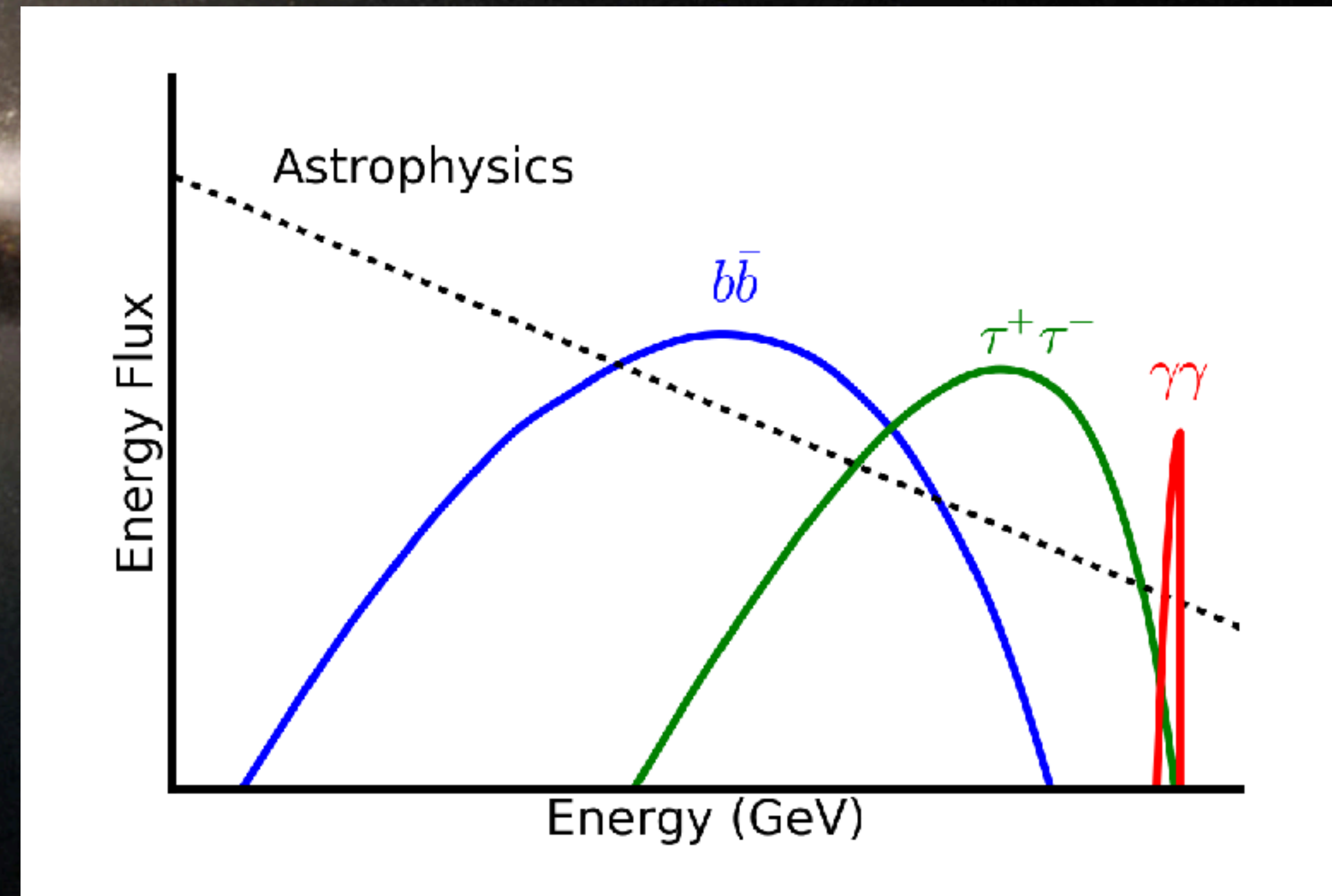
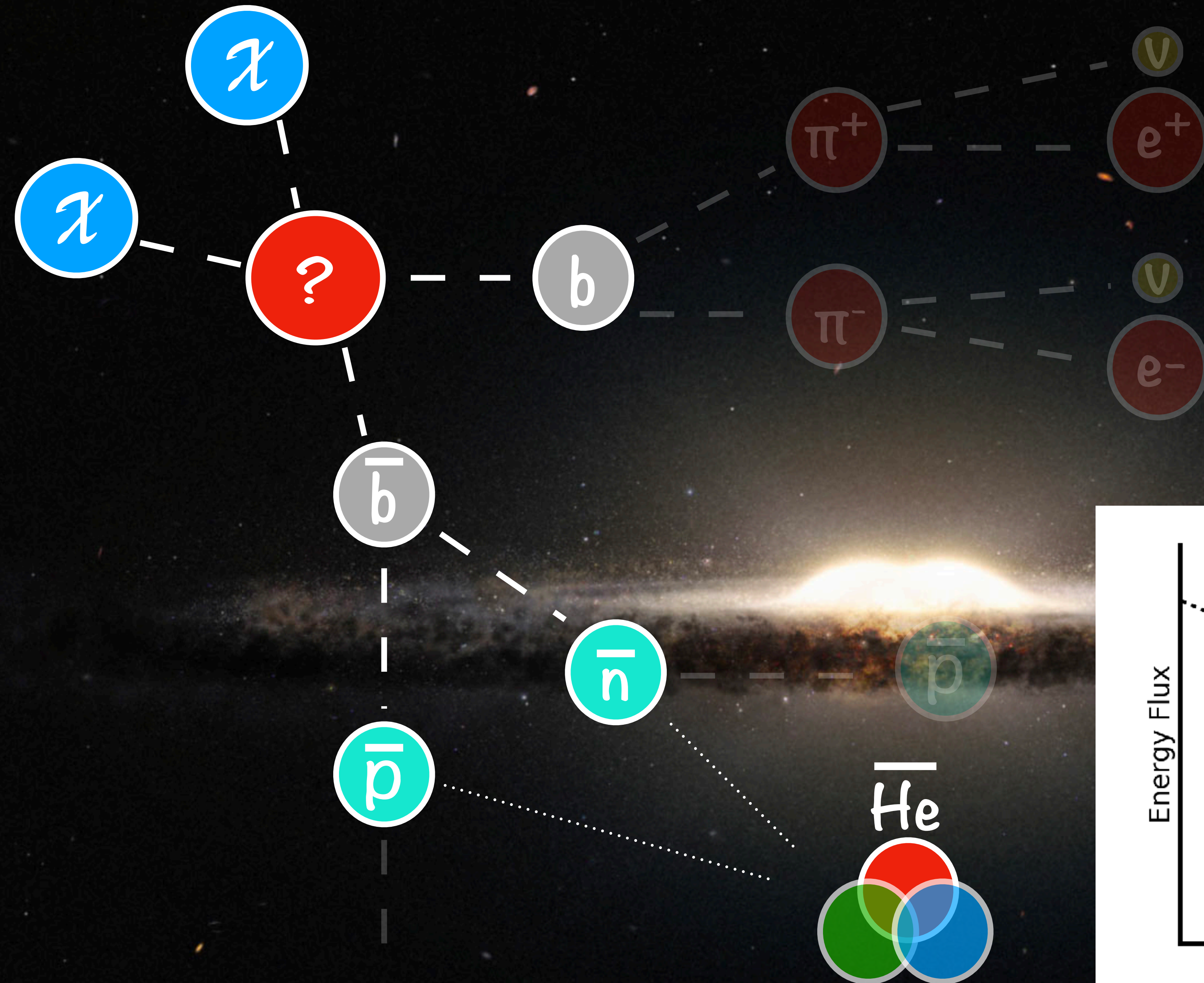
# Flux of Cosmic Rays



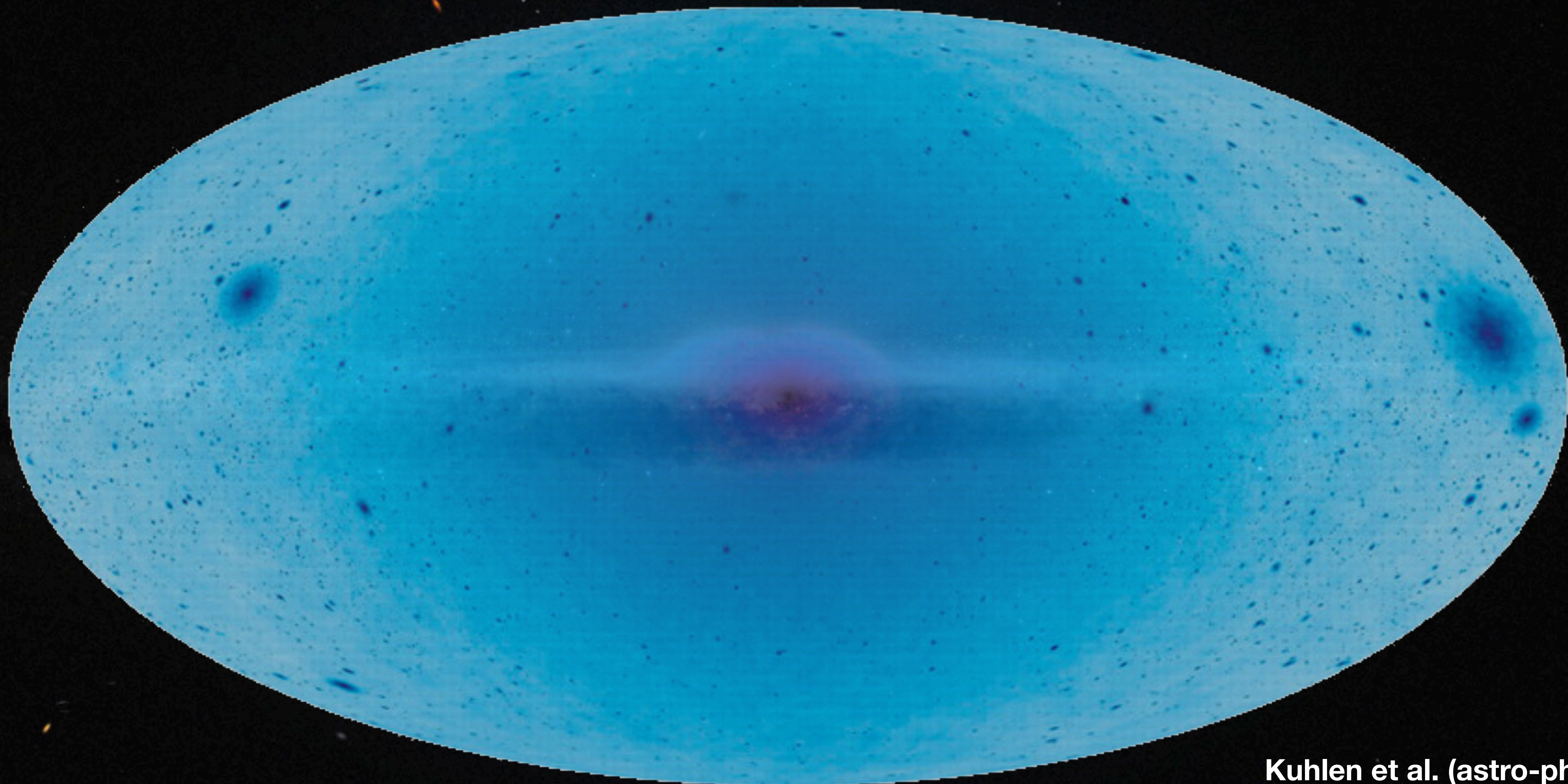












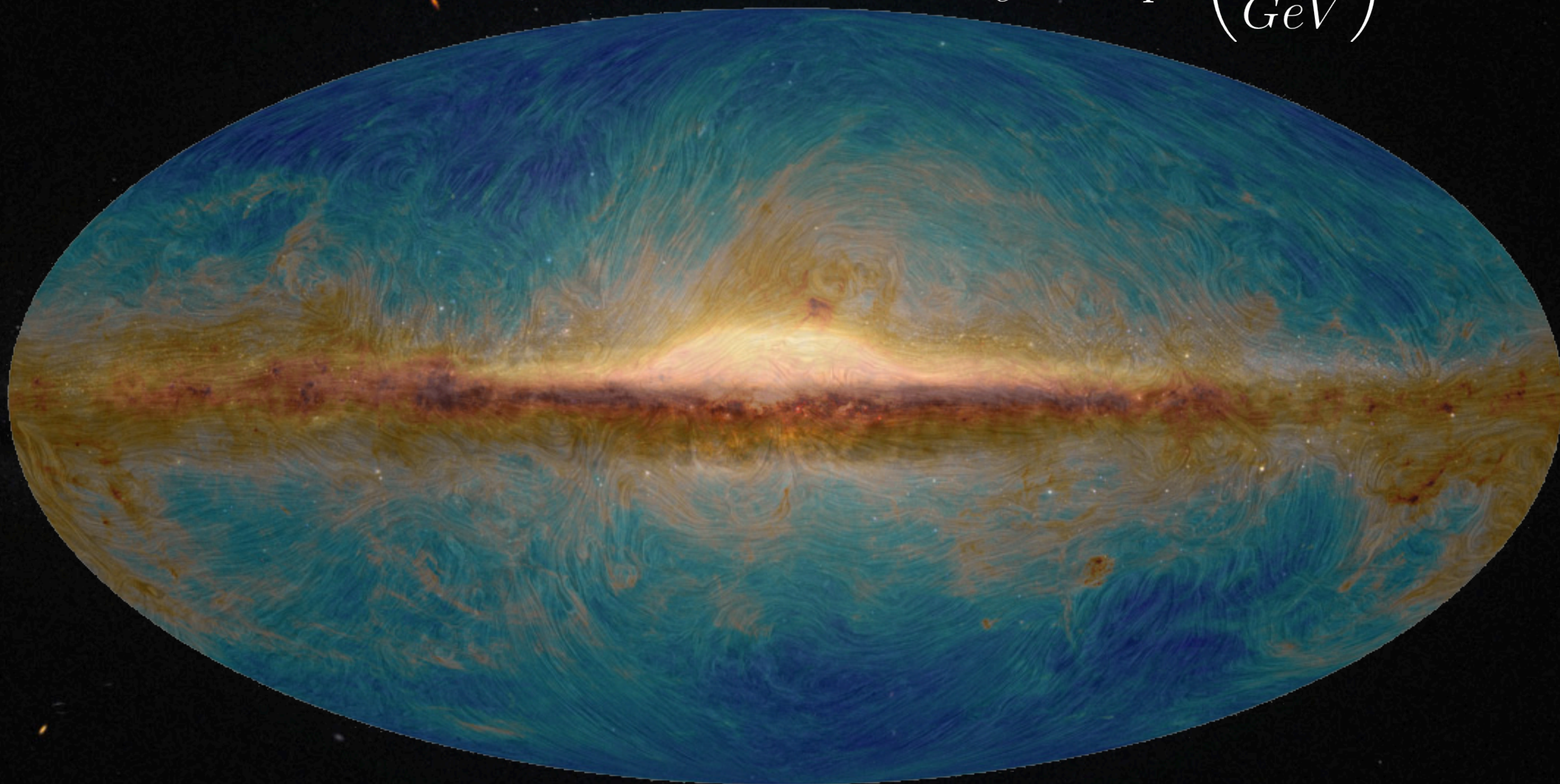




$$B \approx 5 \mu\text{G}$$

$$D_0 \approx 5 \times 10^{28} \text{ cm}^2 \text{ s}^{-1} \left( \frac{E}{\text{GeV}} \right)^{0.3}$$

$$l_c \approx 0.1 \text{ pc} \left( \frac{E}{\text{GeV}} \right)^{0.3}$$

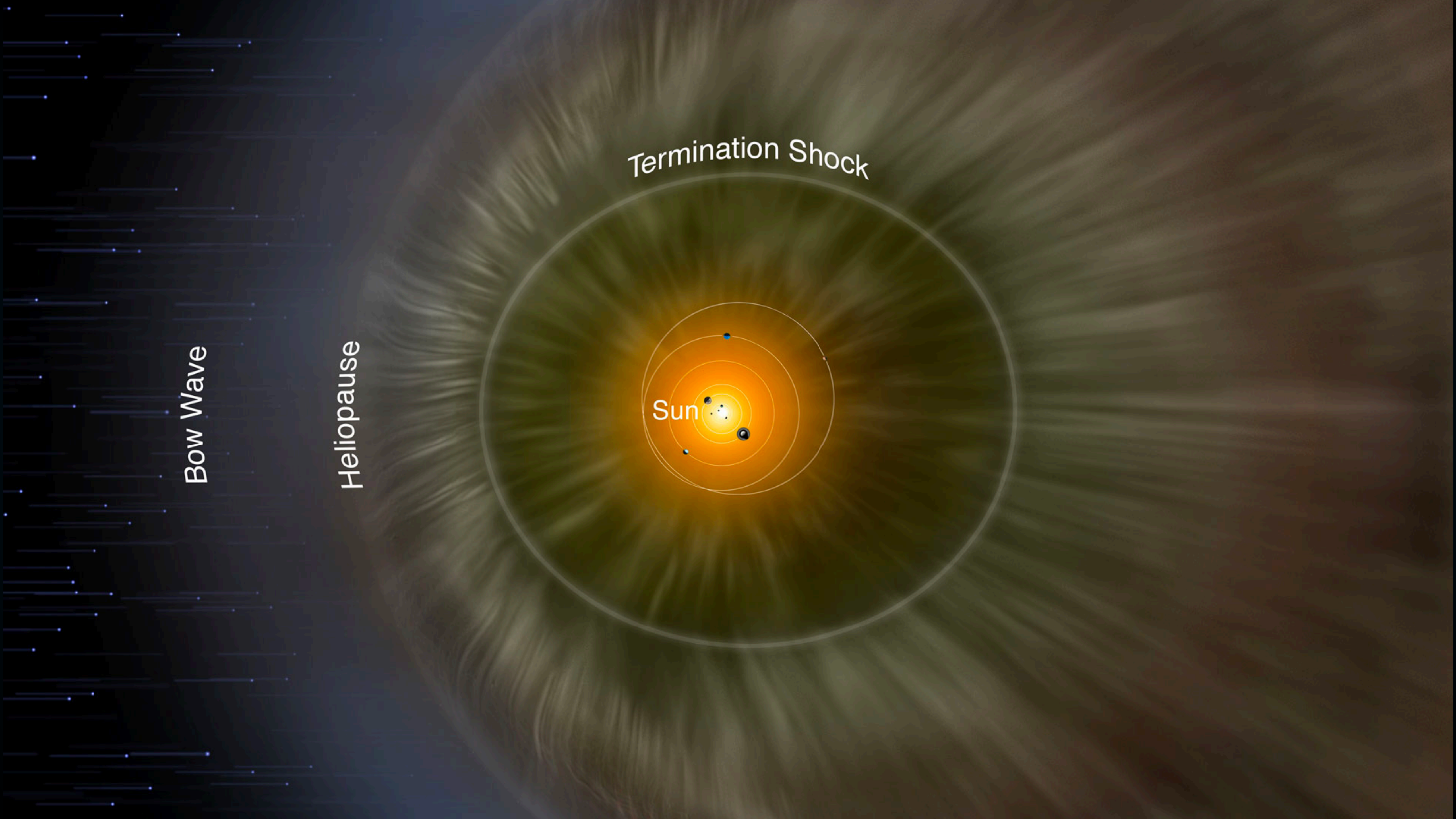


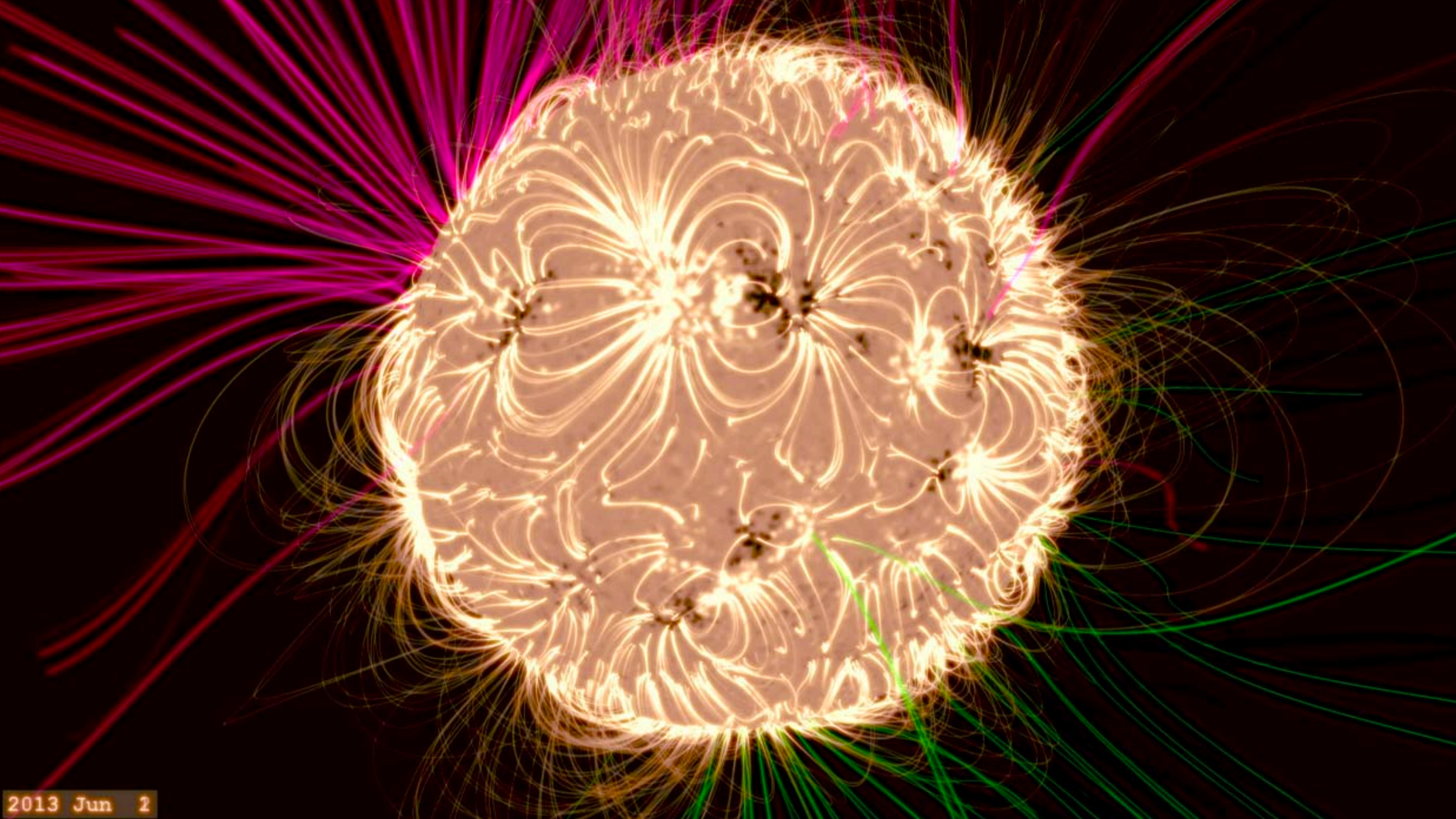
Bow Wave

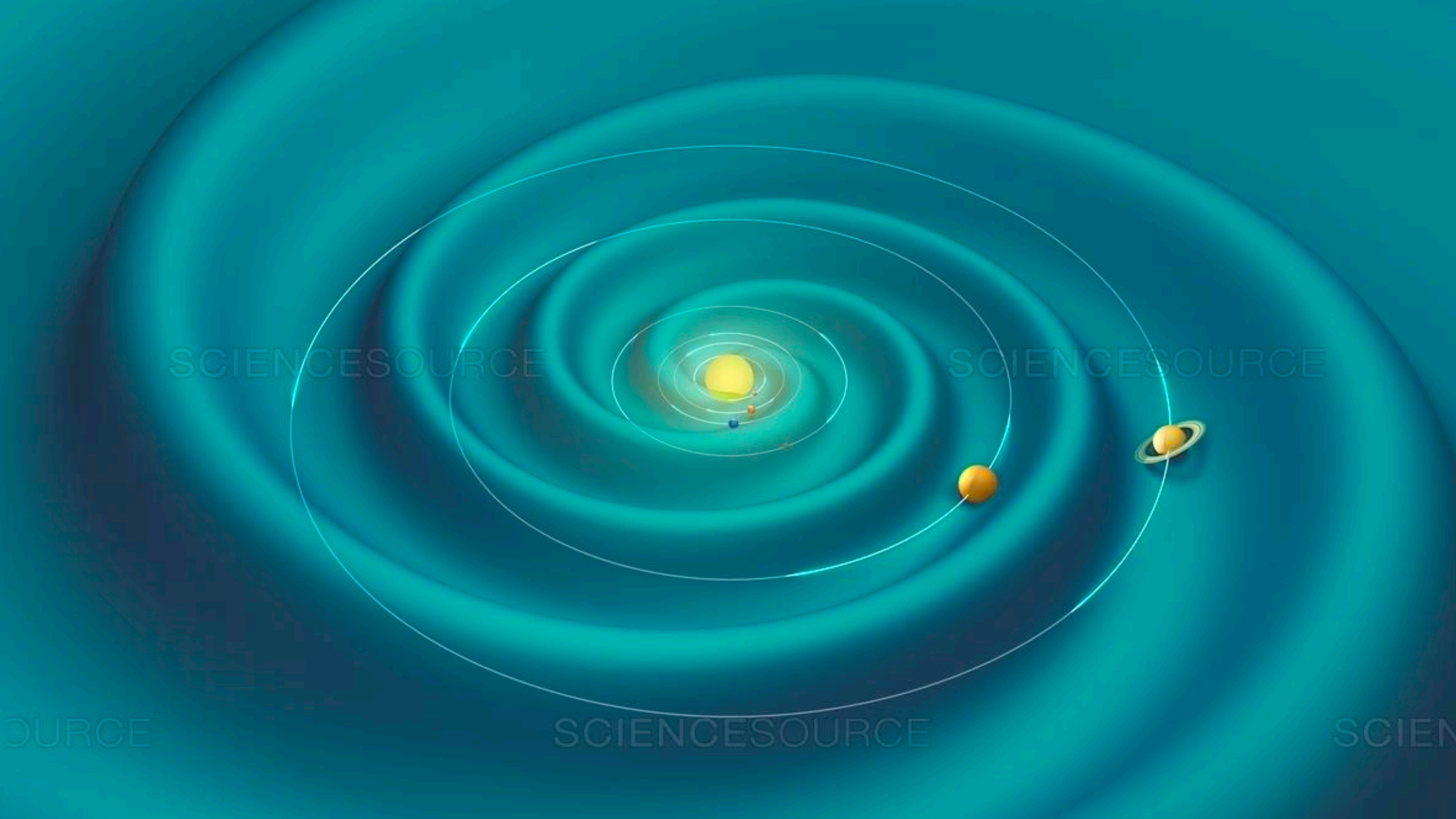
Heliopause

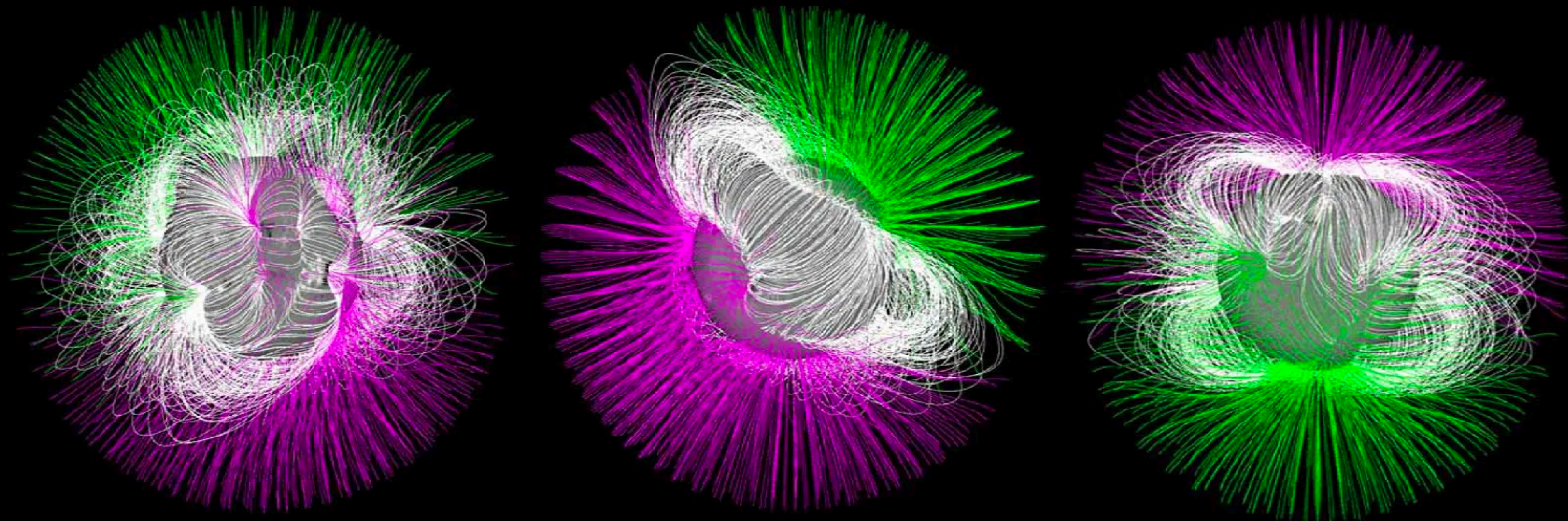
Termination Shock

Sun

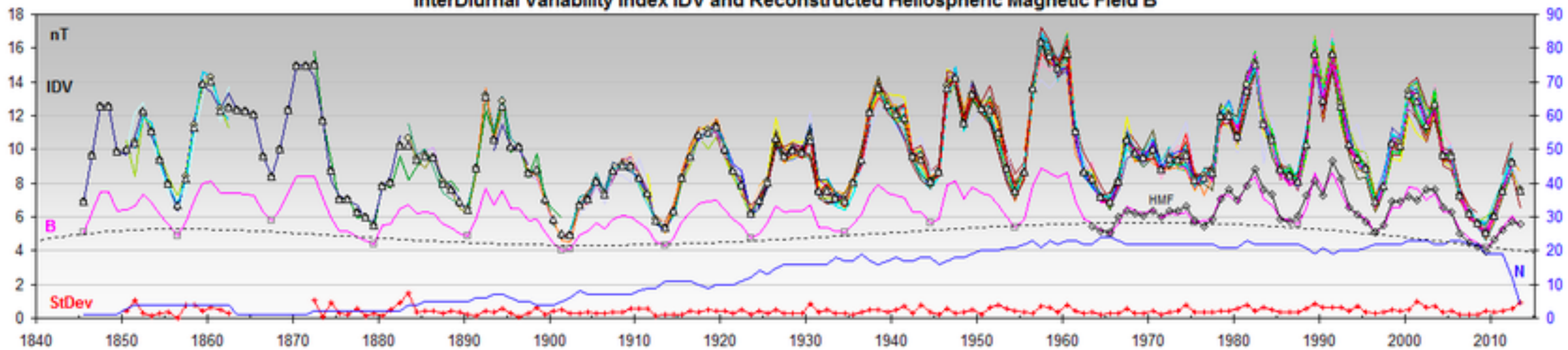




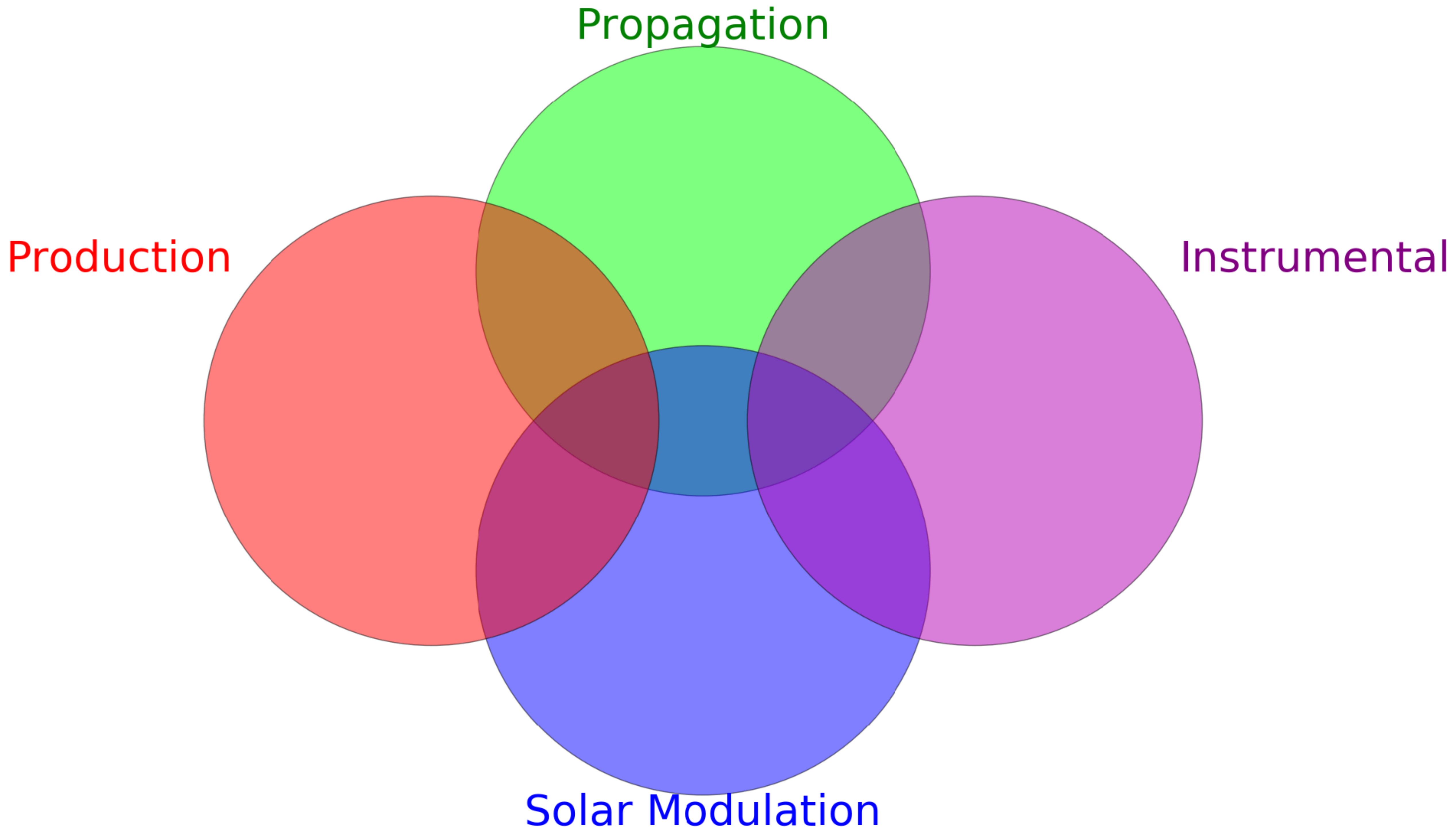




InterDiurnal Variability Index IDV and Reconstructed Heliospheric Magnetic Field B









Specificity (DM Flux / Astrophysics Flux)

Small Dark Matter Signal  
Small Astrophysical Background

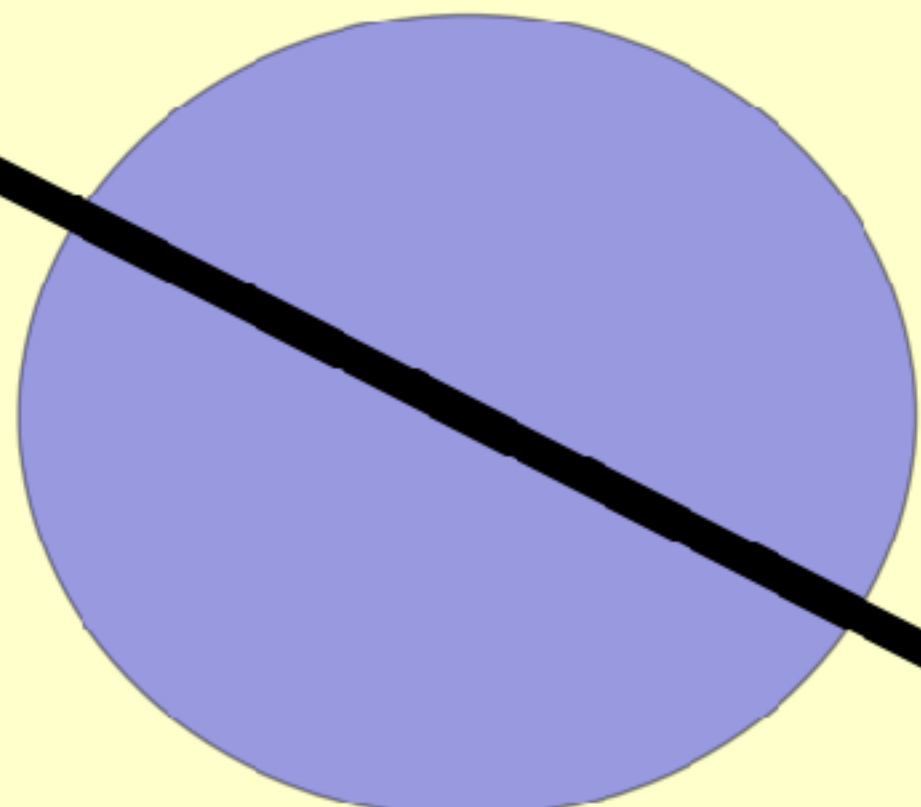
Large Dark Matter Signal  
Small Astrophysical Background

Small Dark Matter Signal  
Large Astrophysical Background

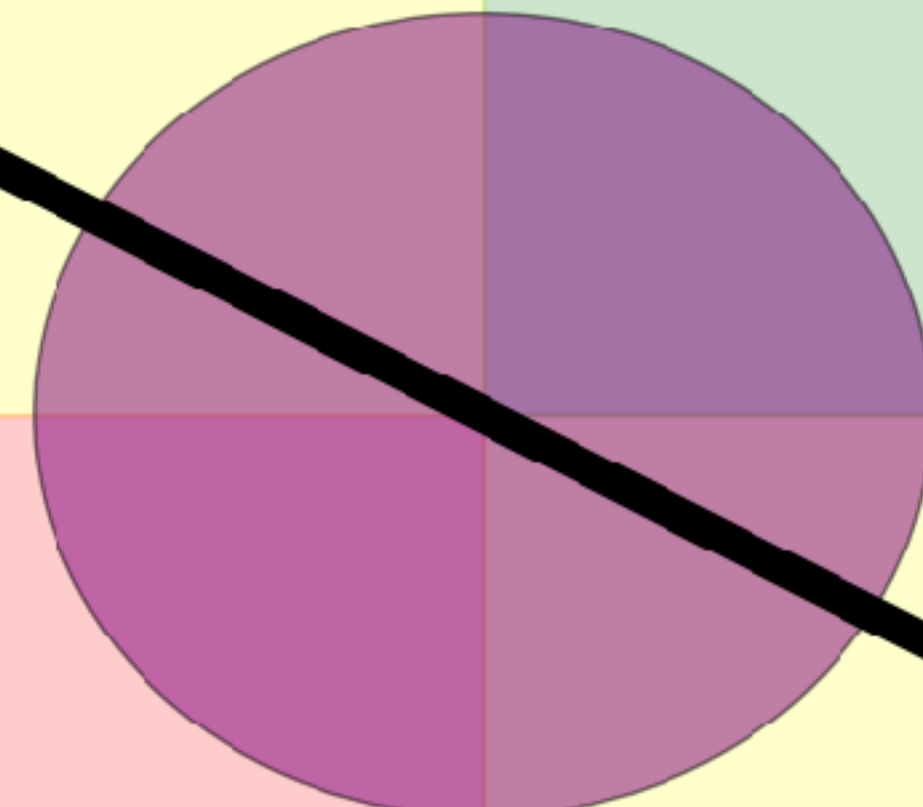
Large Dark Matter Signal  
Large Astrophysical Background

Fraction of Dark Matter Flux

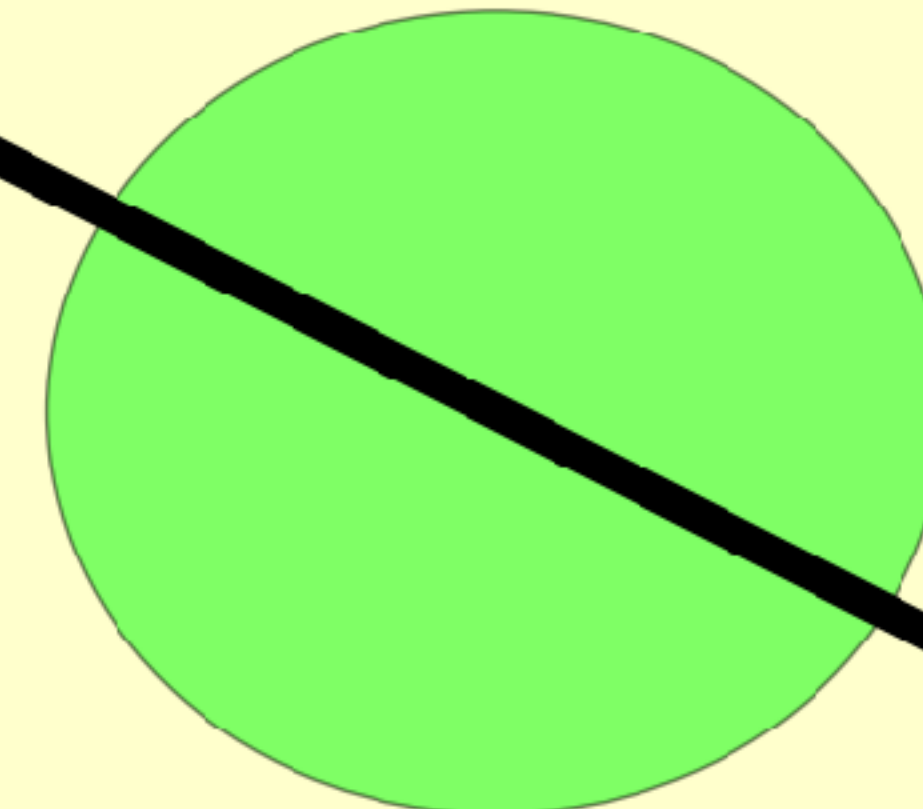
Specificity (DM Flux / Astrophysics Flux)



Anti-Nuclei

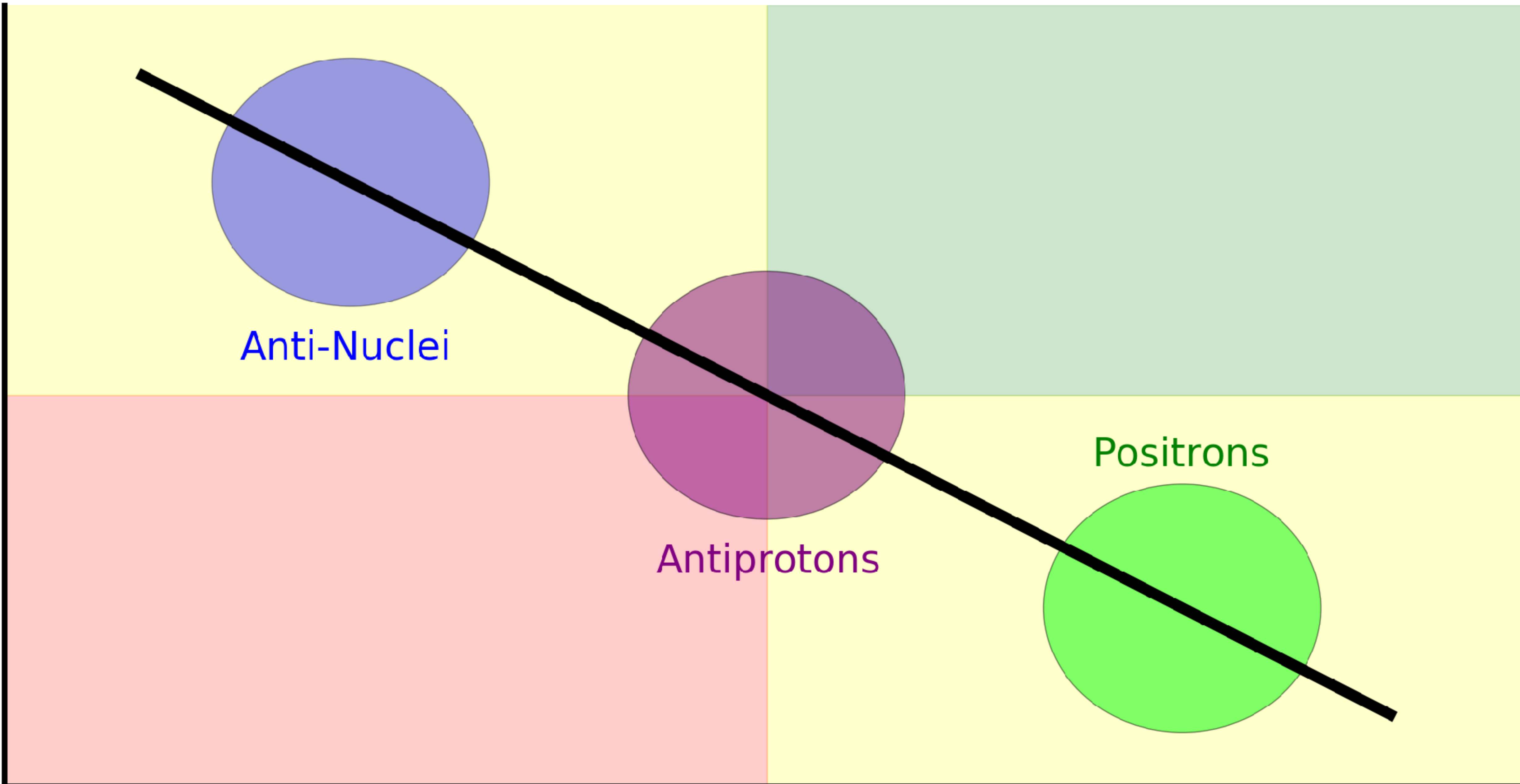
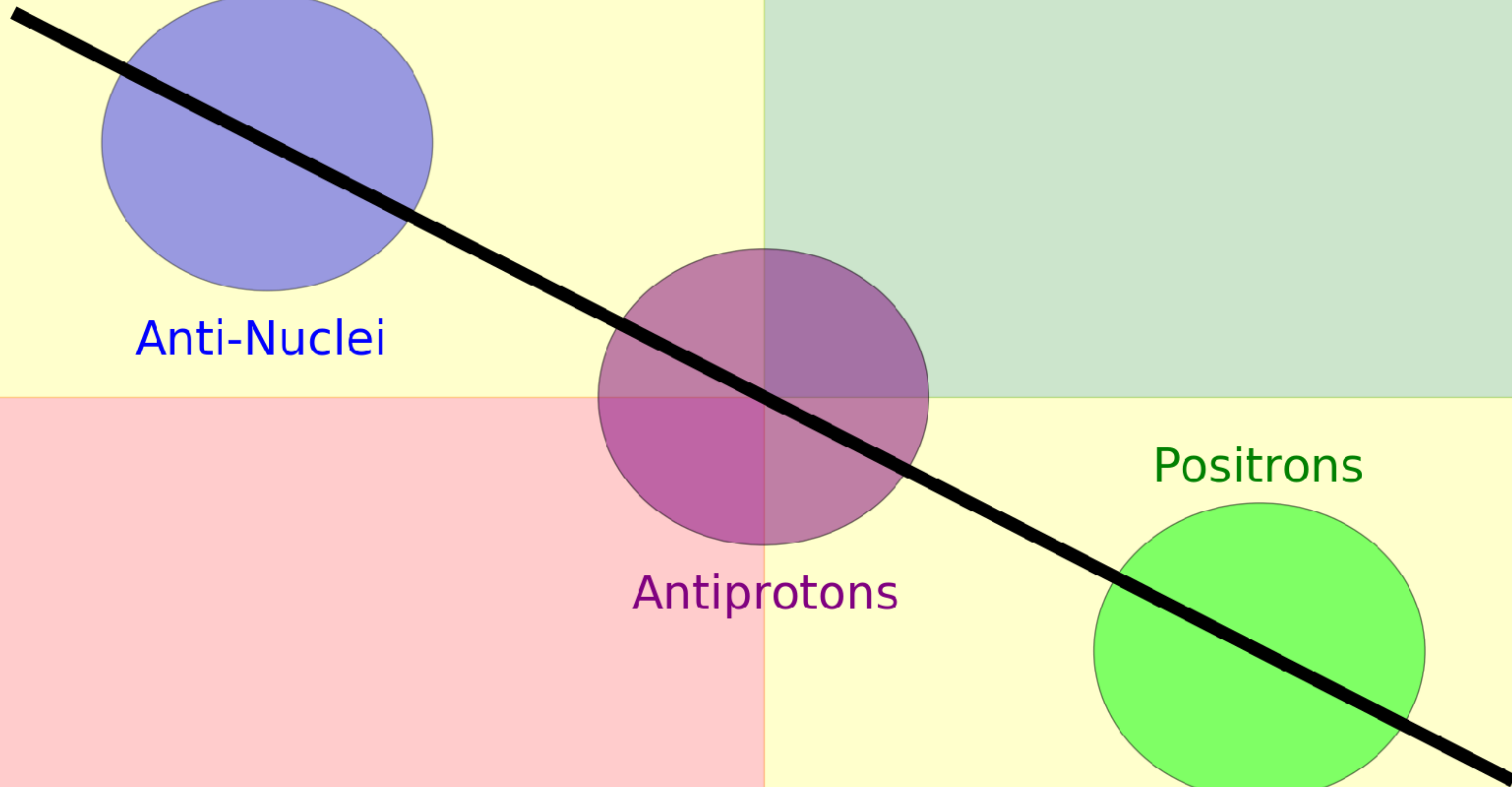


Antiprotons



Positrons

Fraction of Dark Matter Flux



# Dark Matter and Astrophysical Fluxes

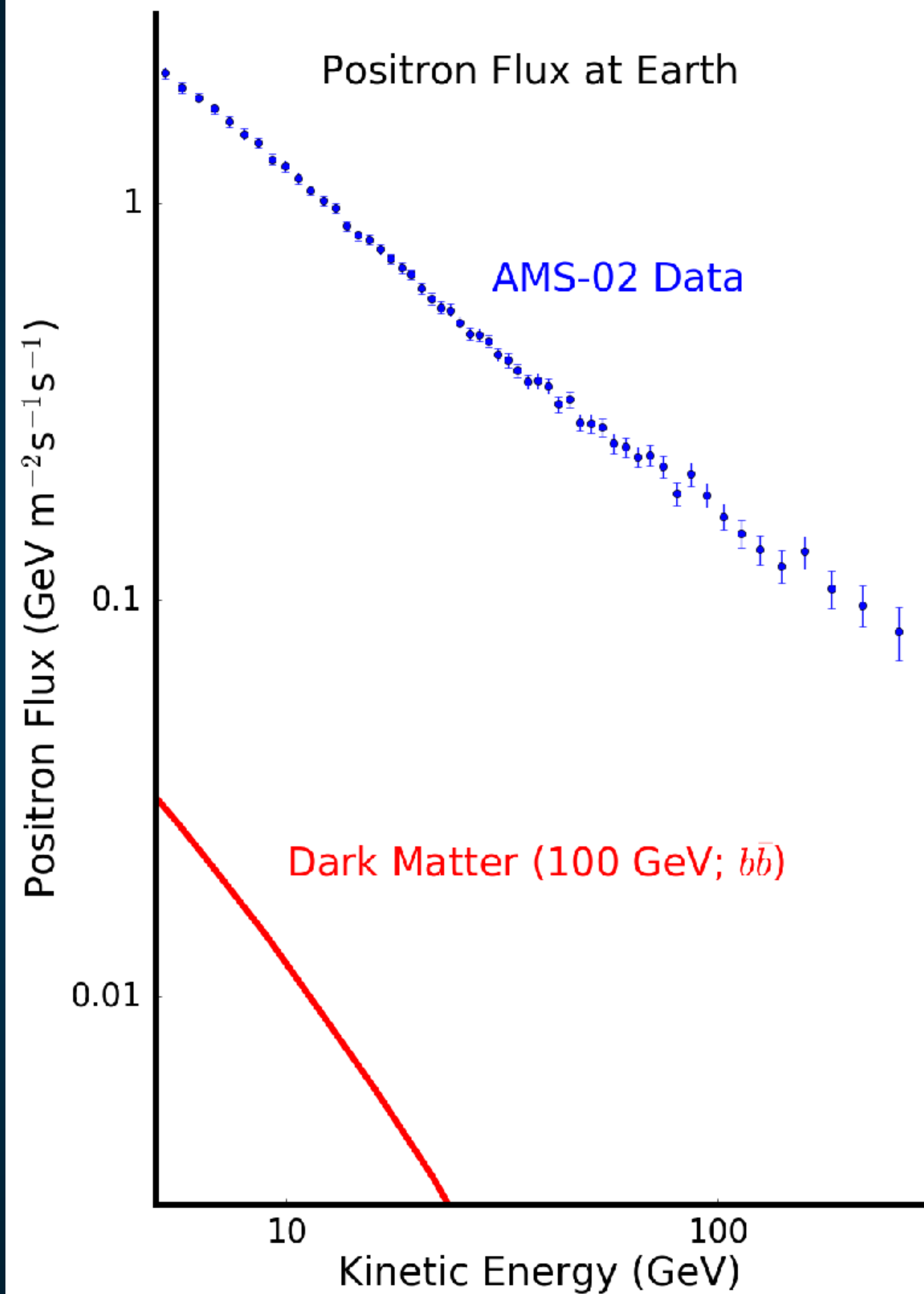
Local Dark Matter Density

Thermal Cross-Section (Early Universe)

Leptonic Component of Dark Matter Final State

Convection of Annihilation Products from GC (Winds?)

Astrophysical Flux ~ **100x larger**



# Dark Matter and Astrophysical Fluxes

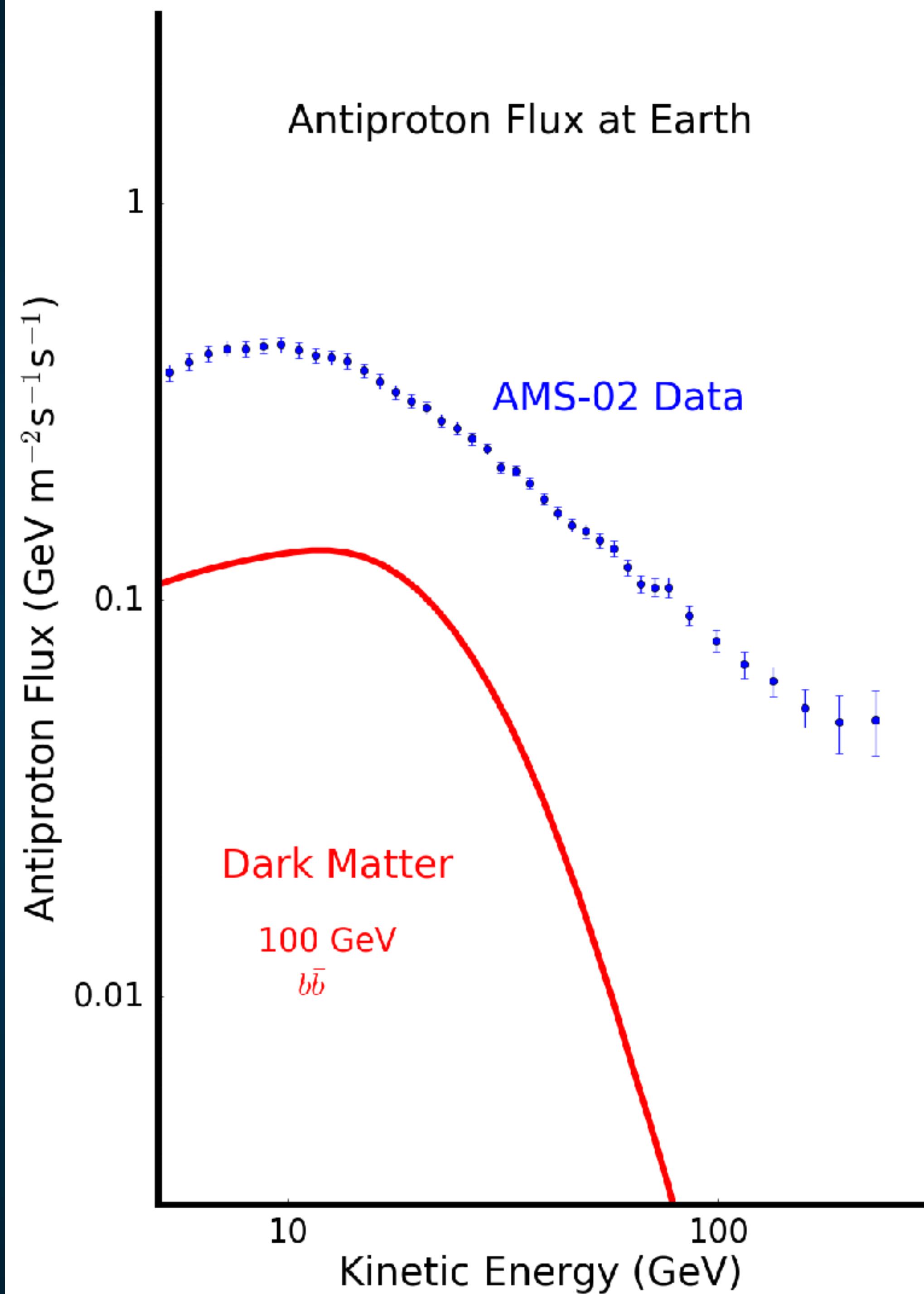
Local Dark Matter Density

Thermal Cross-Section (Early Universe)

Hadronic Component of Dark Matter Final State

Convection of Annihilation Products from GC (Winds?)

Astrophysical Flux ~ **10x larger**



# Dark Matter and Astrophysical Fluxes

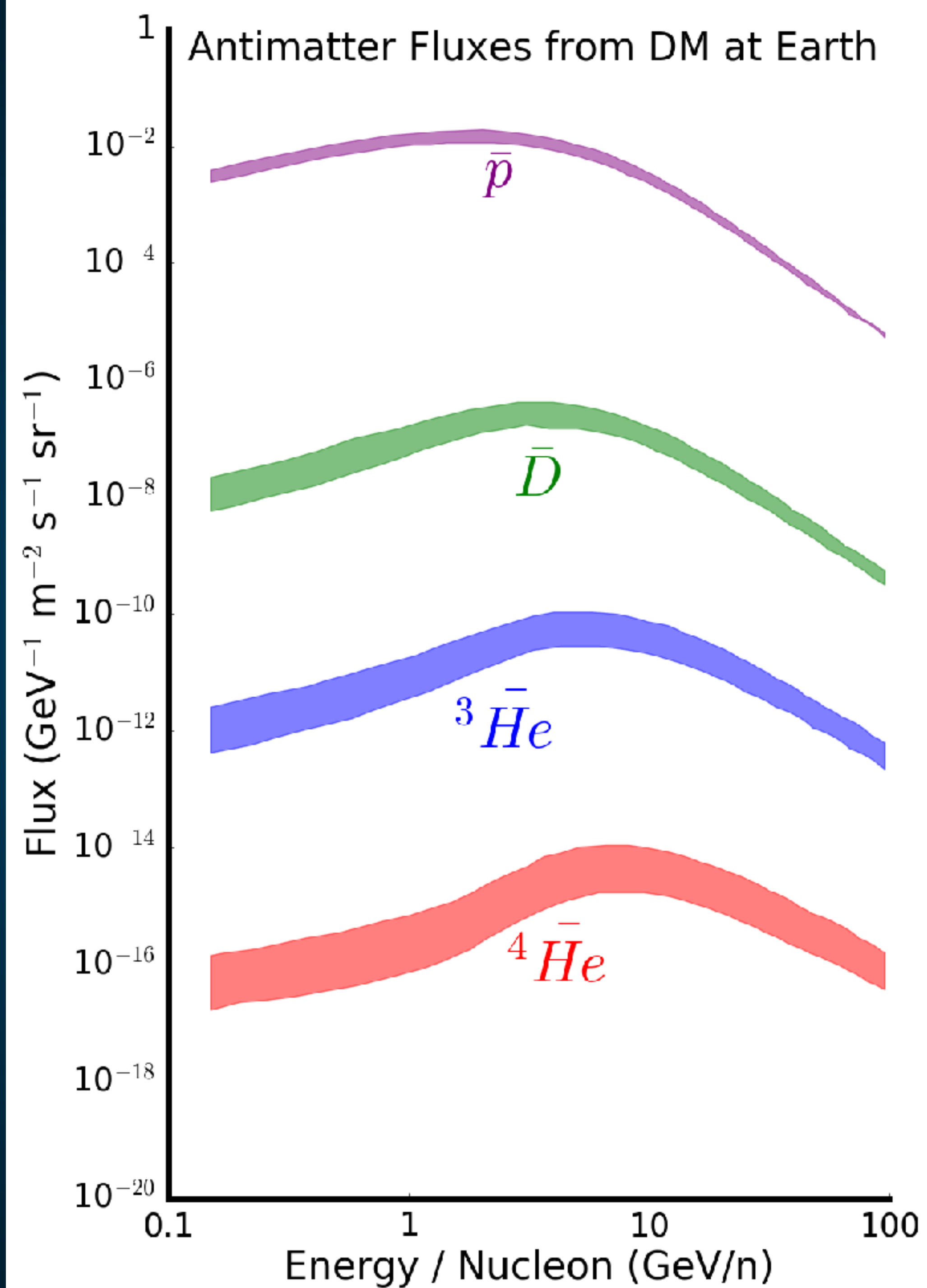
Local Dark Matter Density

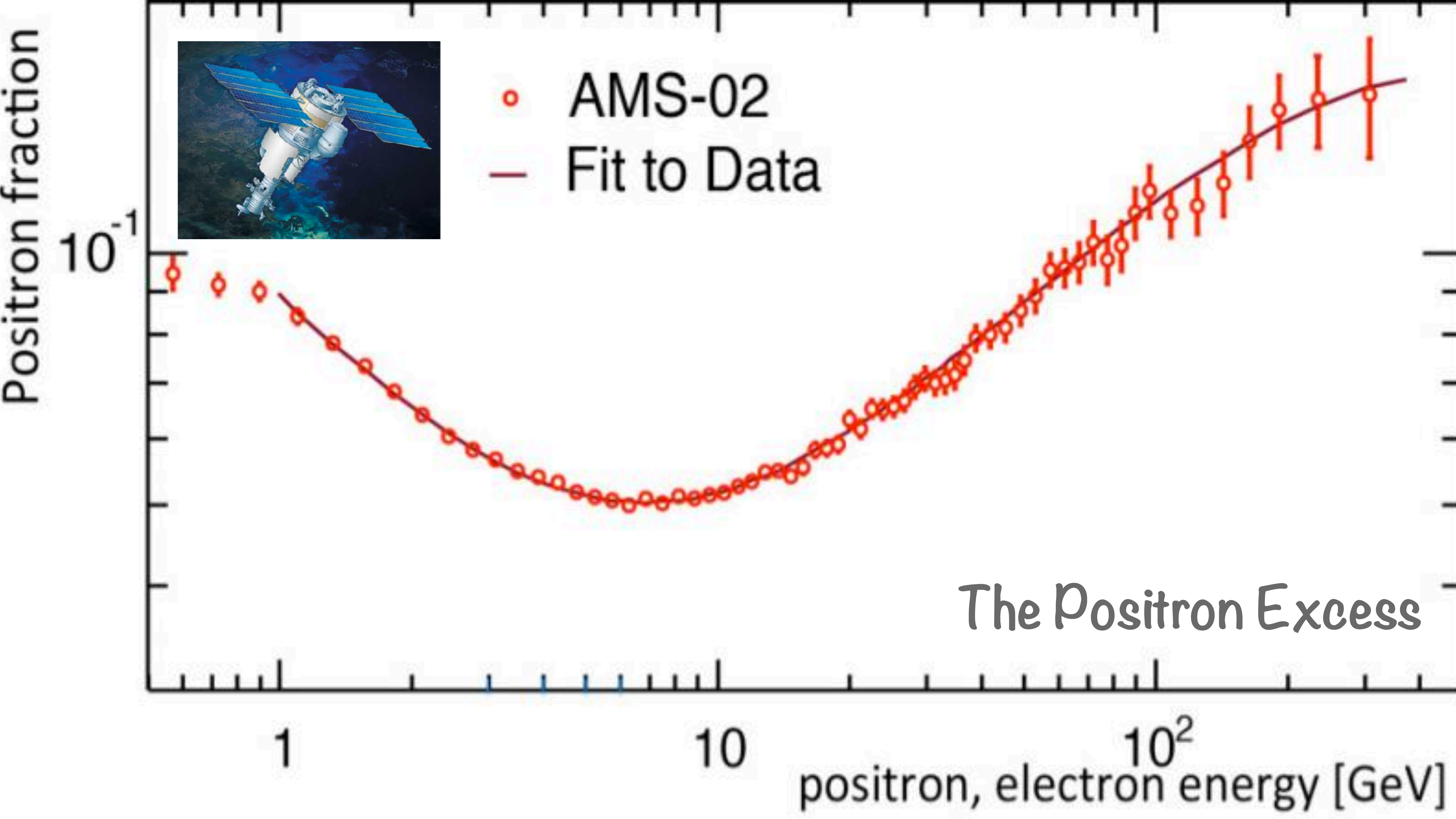
Thermal Cross-Section (Early Universe)

Coalescence of baryons into heavier nuclei

Convection of Annihilation Products from GC (Winds?)

Astrophysical Flux - **Undetected, likely much lower**

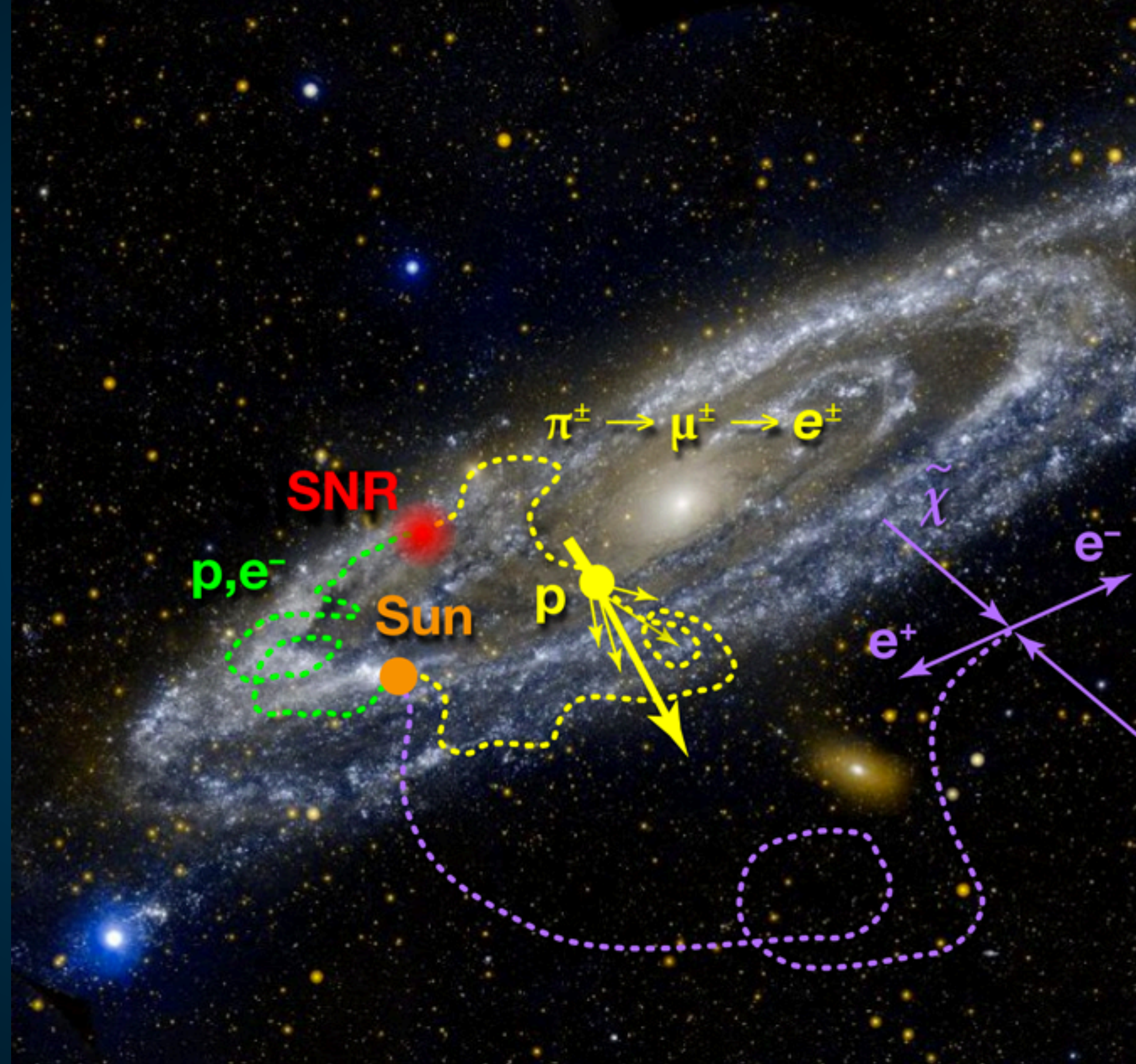
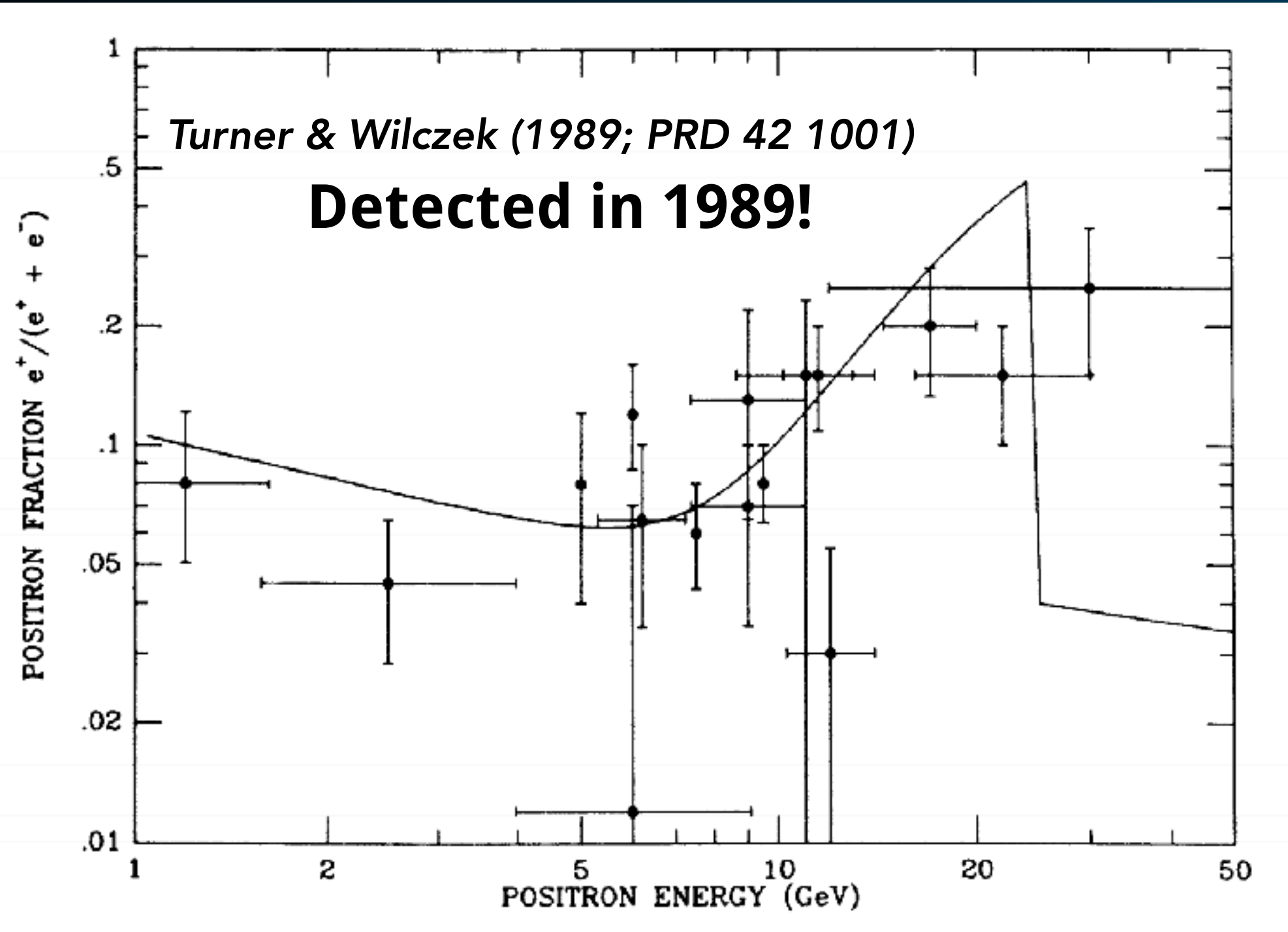




# The Positron Excess

Key Idea: Investigate the Positron Fraction!

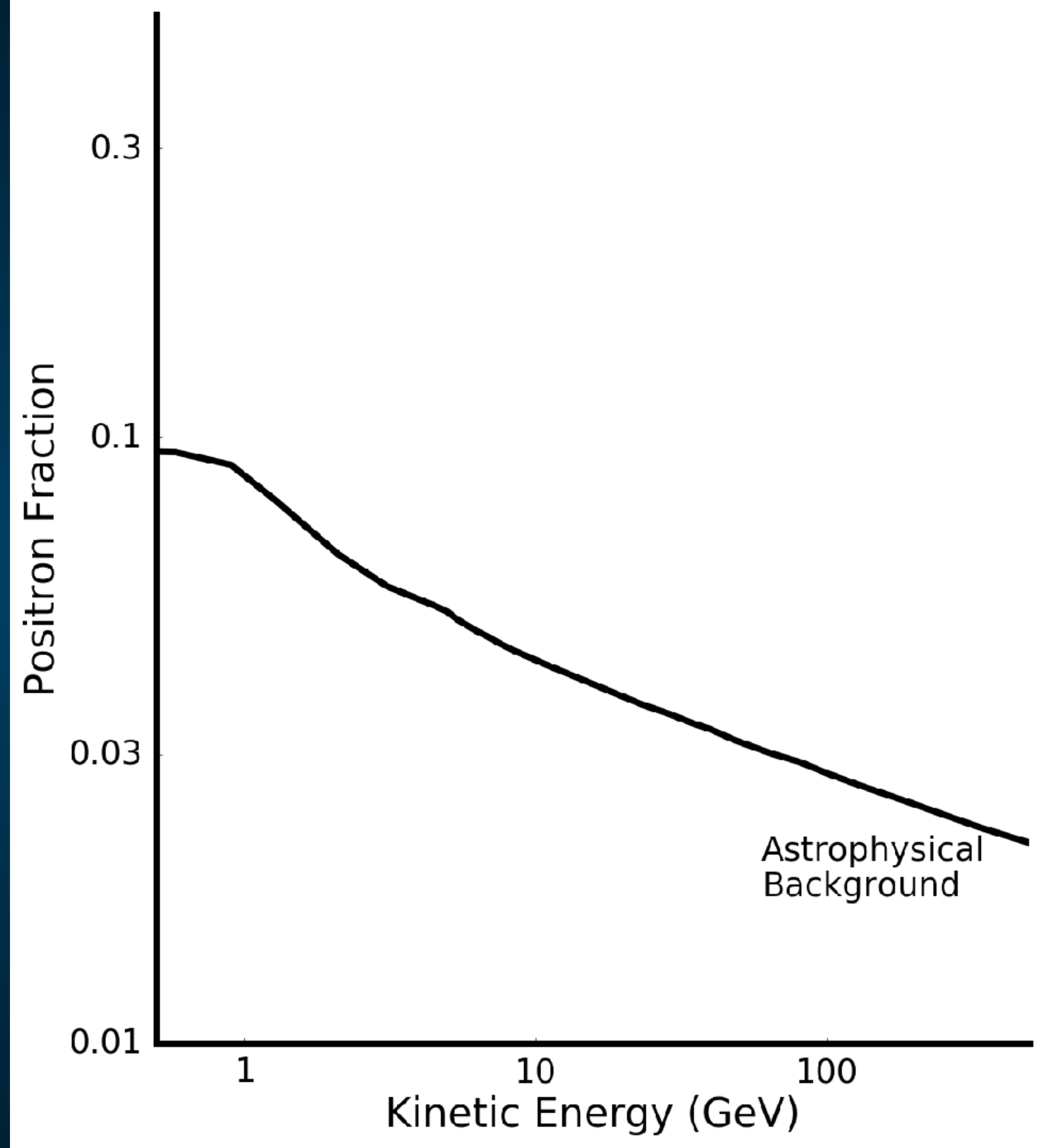
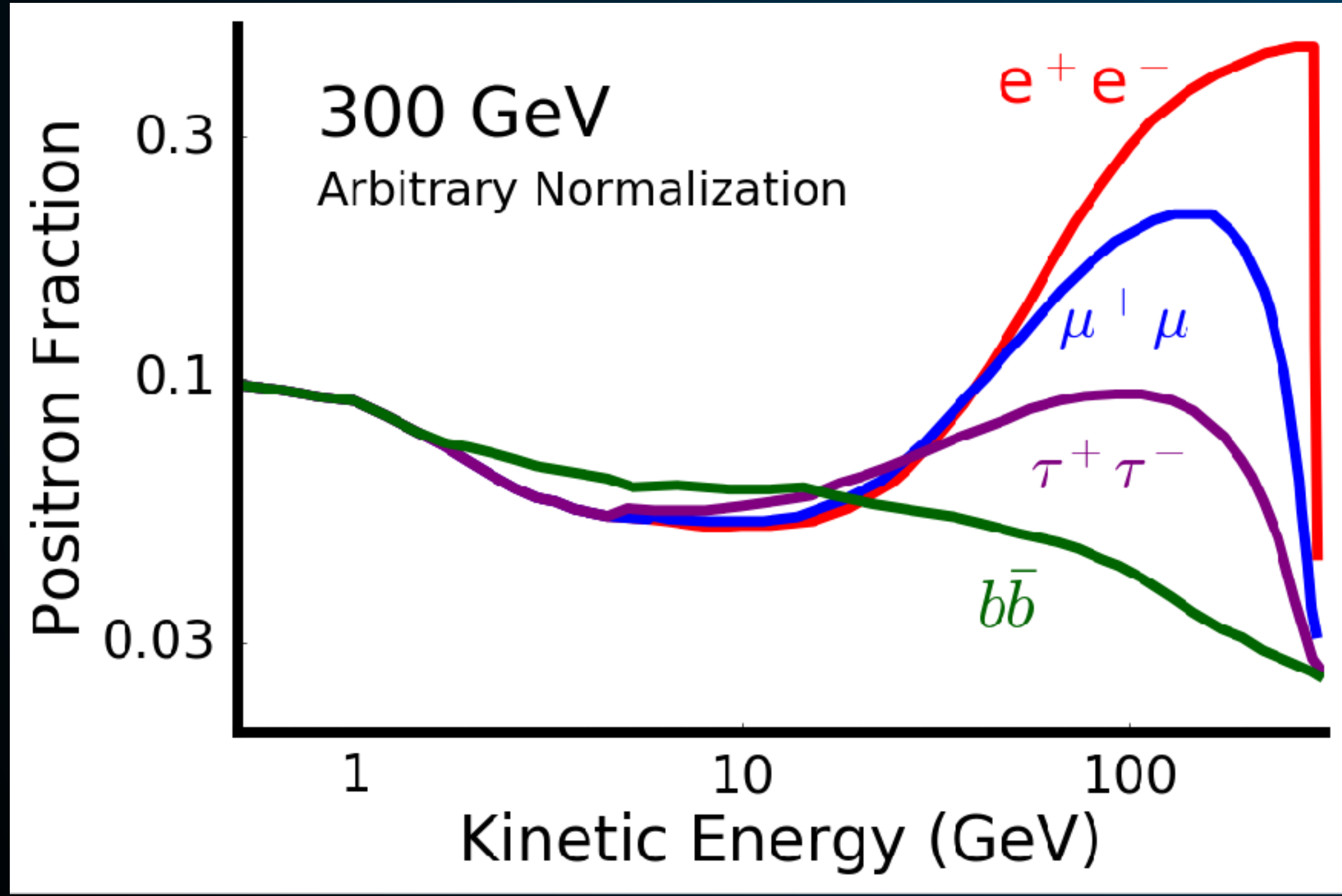
$$\frac{\phi_{e^+}}{\phi_{e^+} + \phi_{e^-}}$$



# The Positron Excess

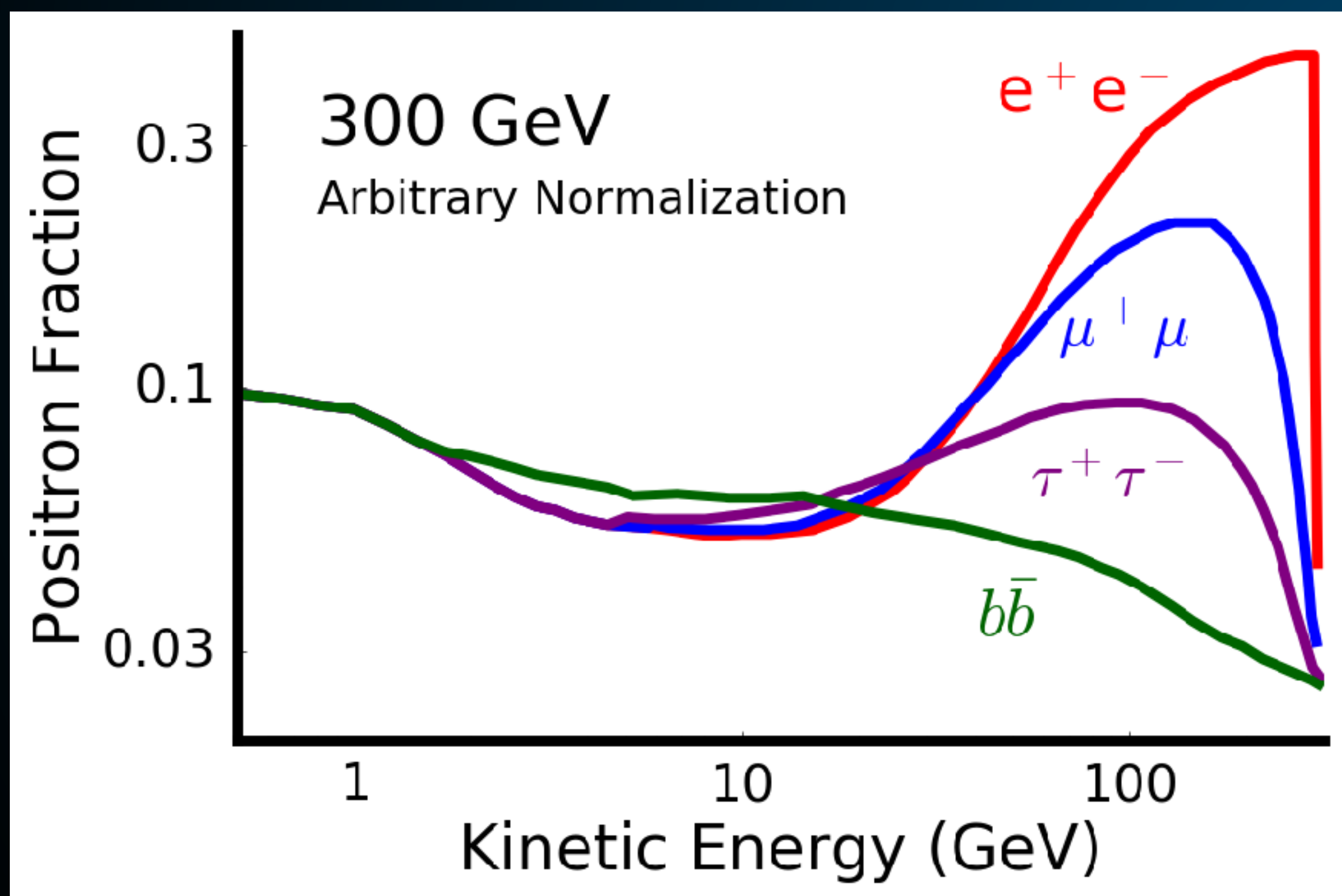
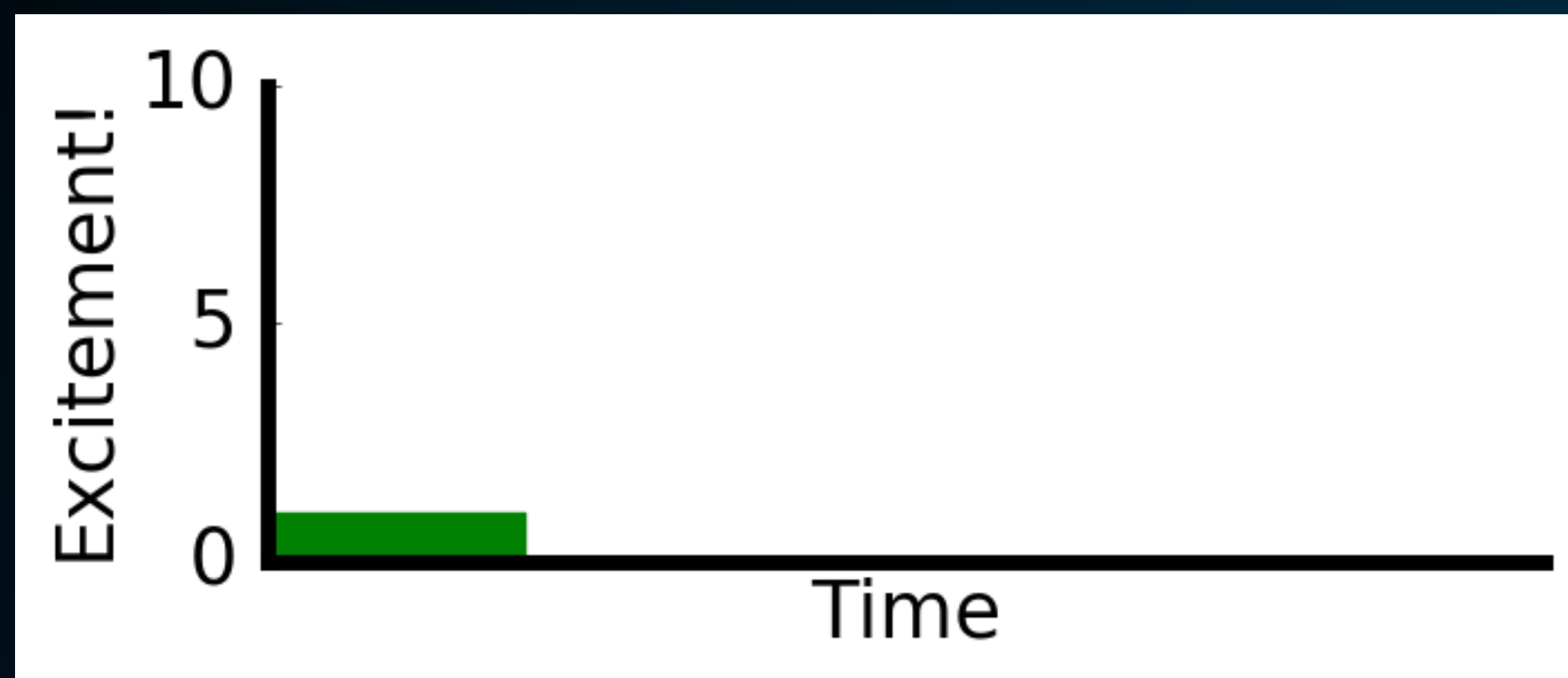
Astrophysics - Slowly Decreasing

Dark Matter - Sharp Bump!

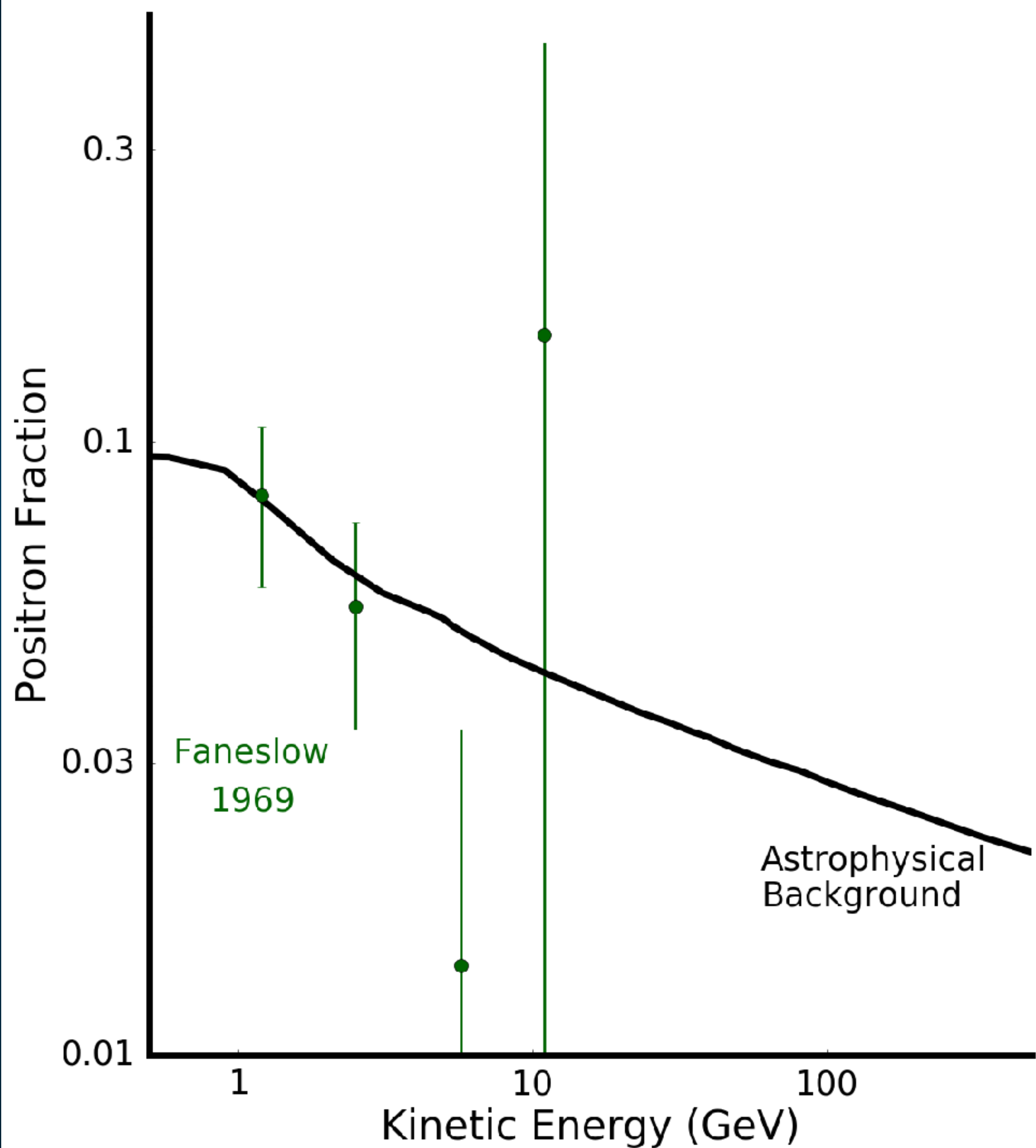




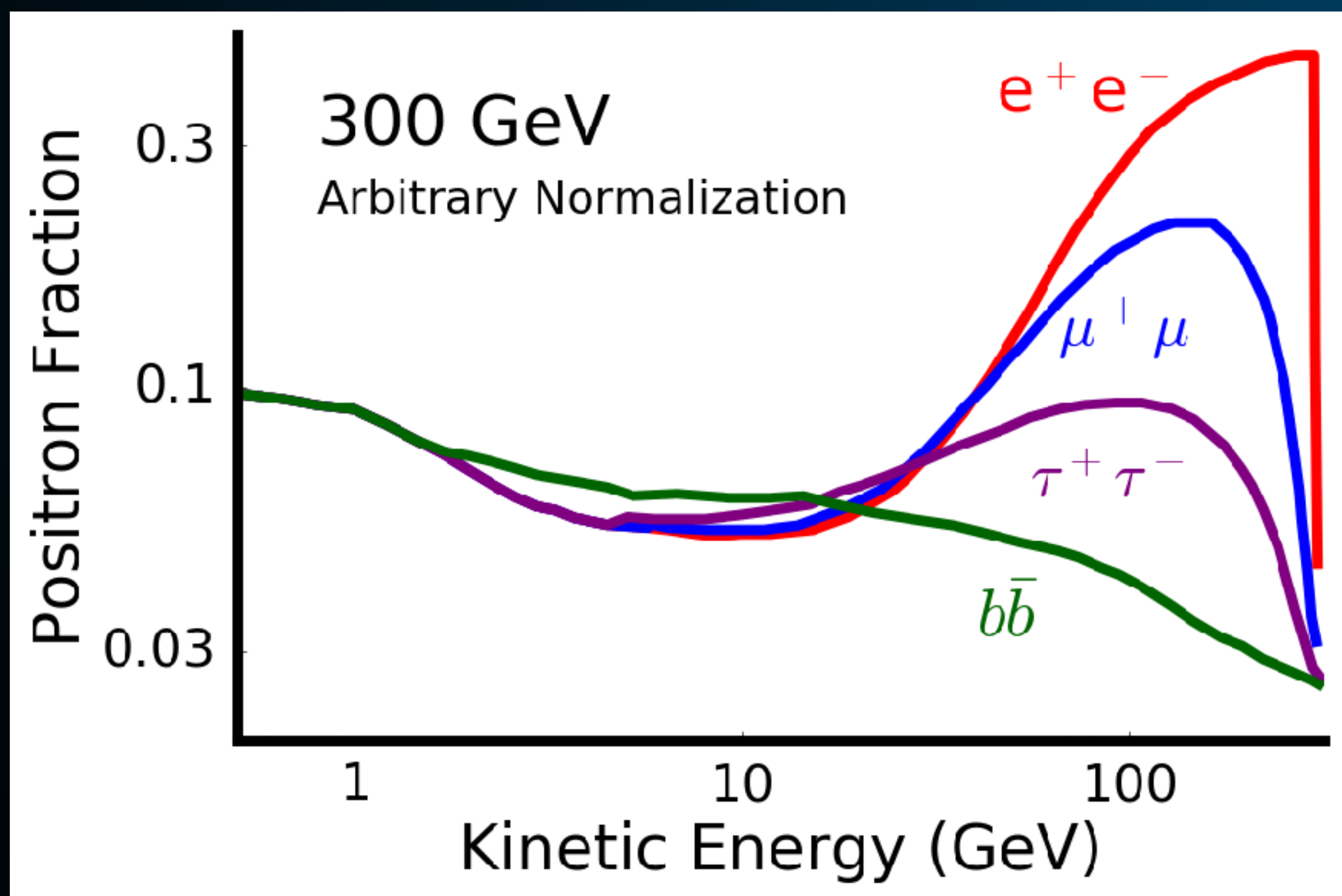
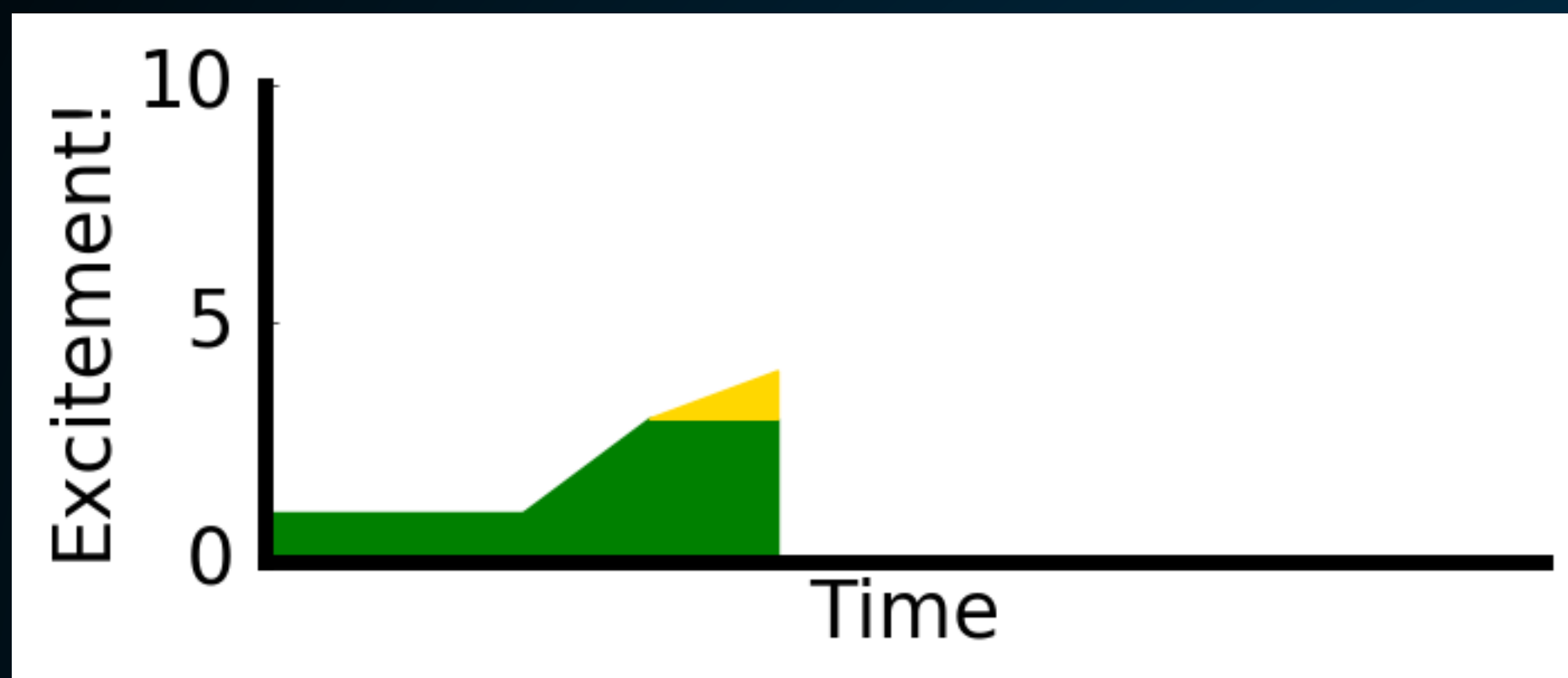
# The Positron Excess



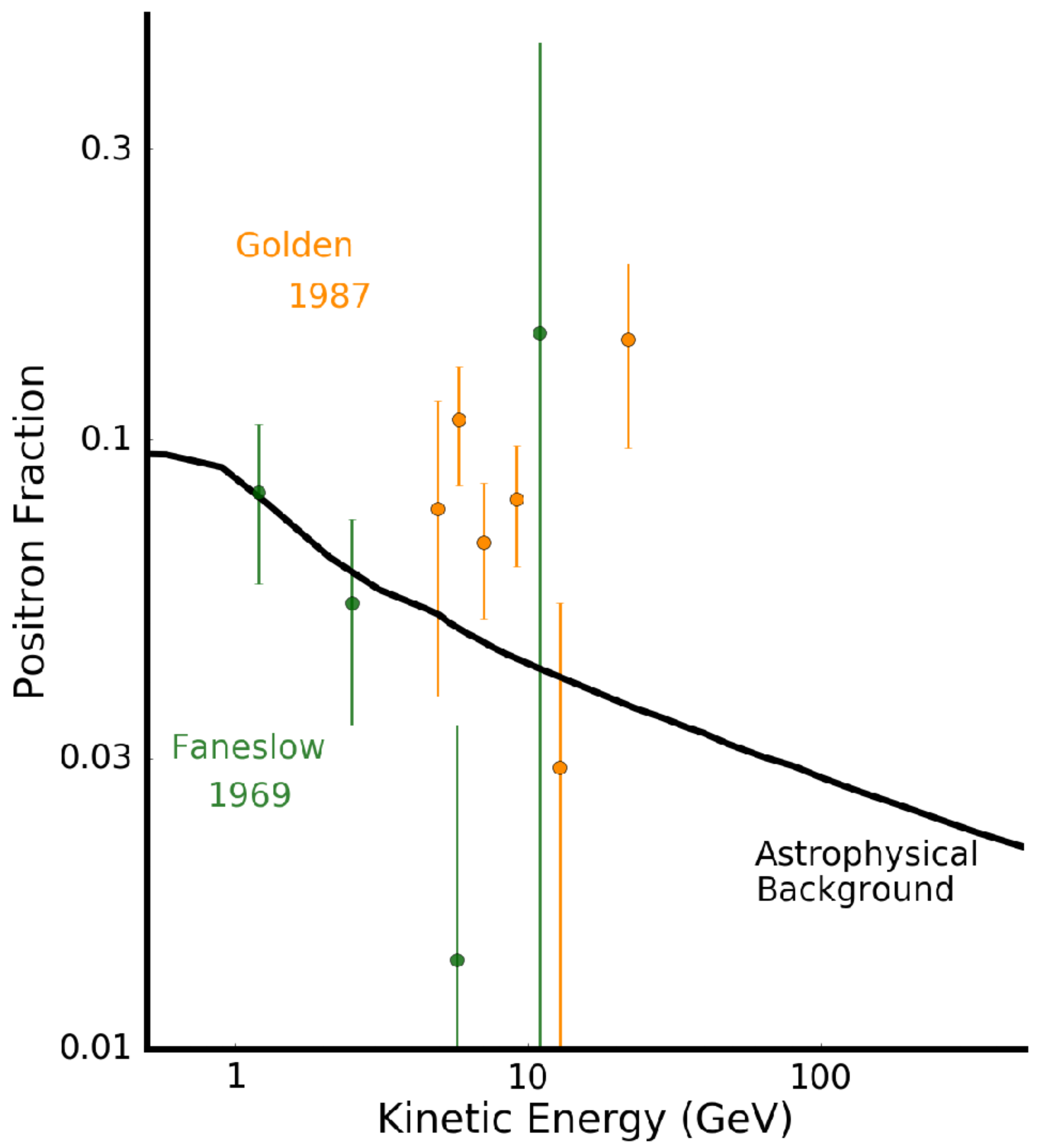
(Not an exhaustive list of observations)



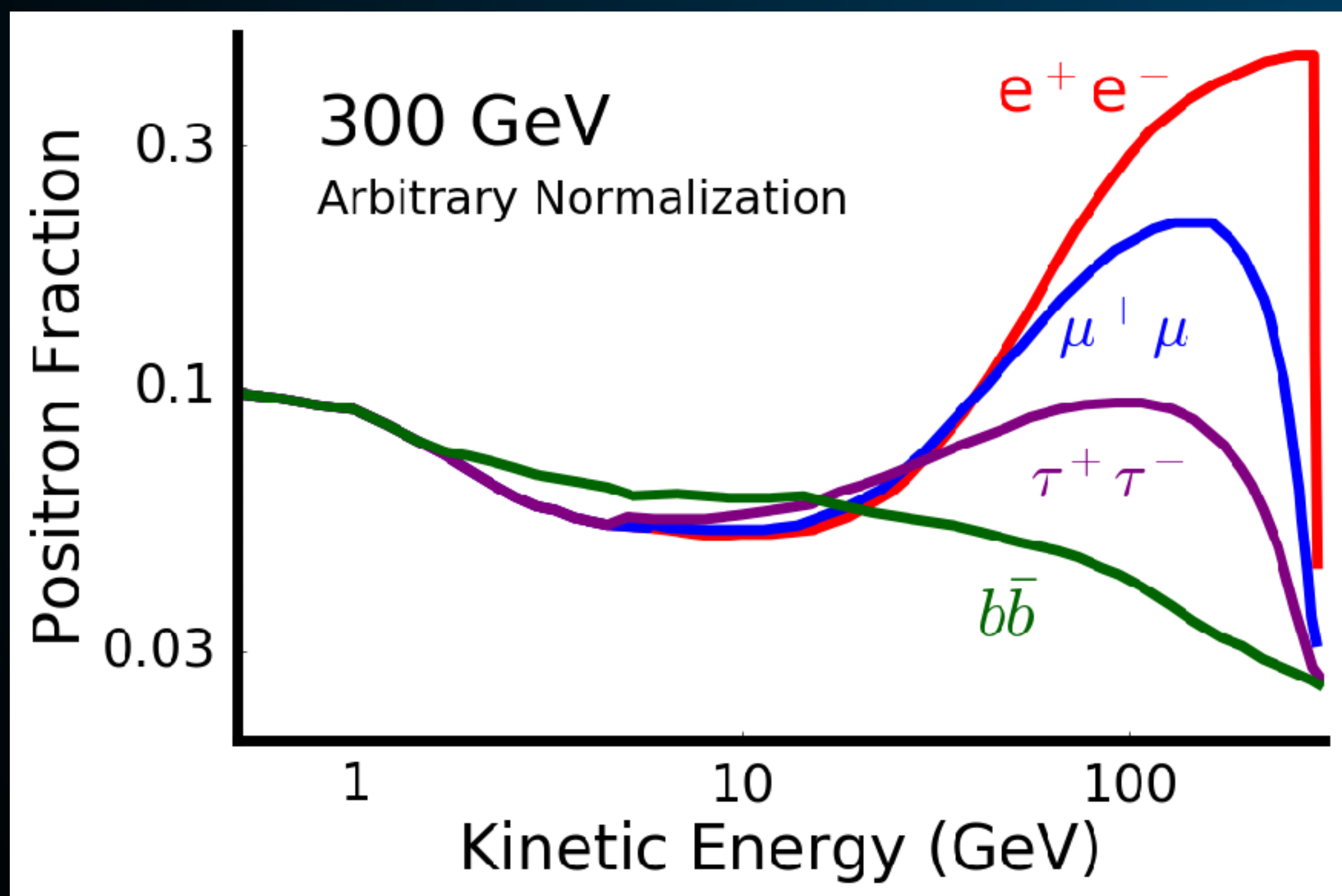
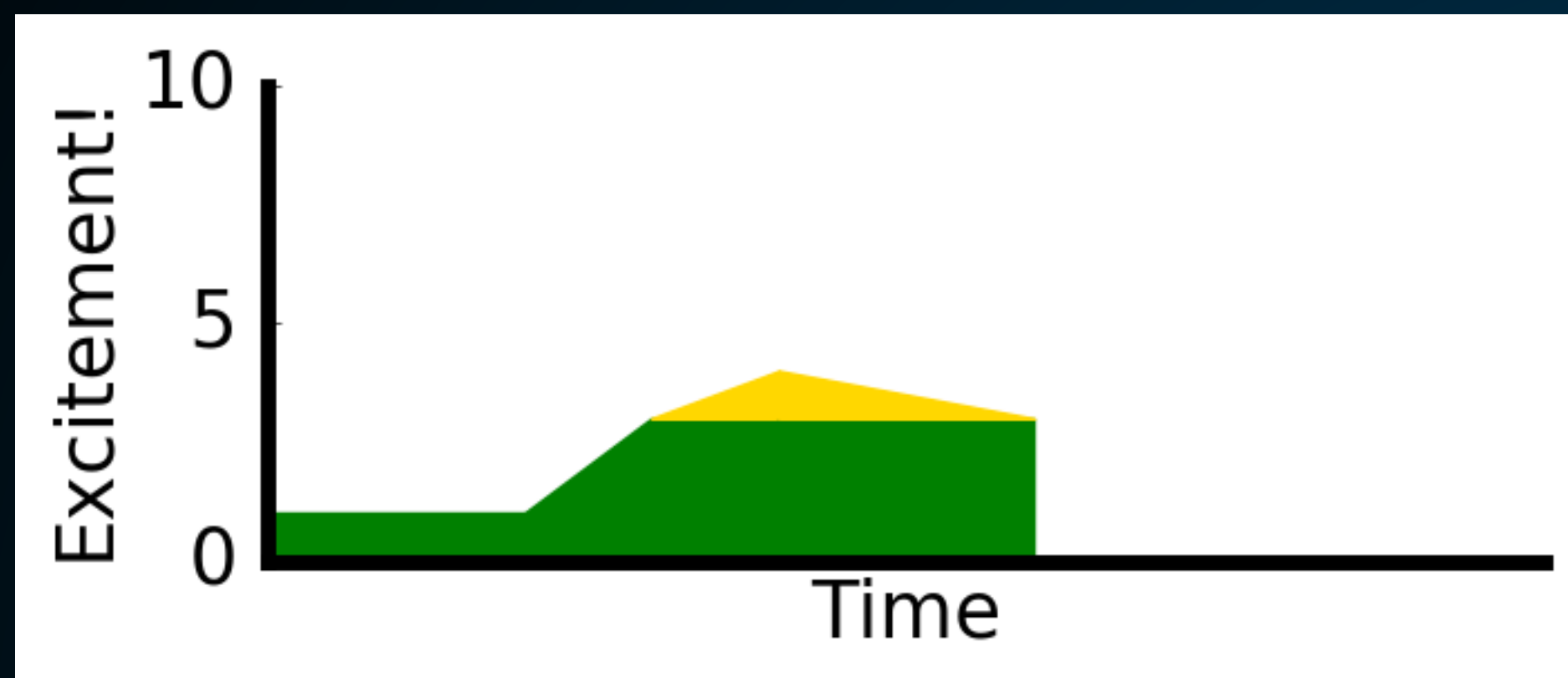
# The Positron Excess



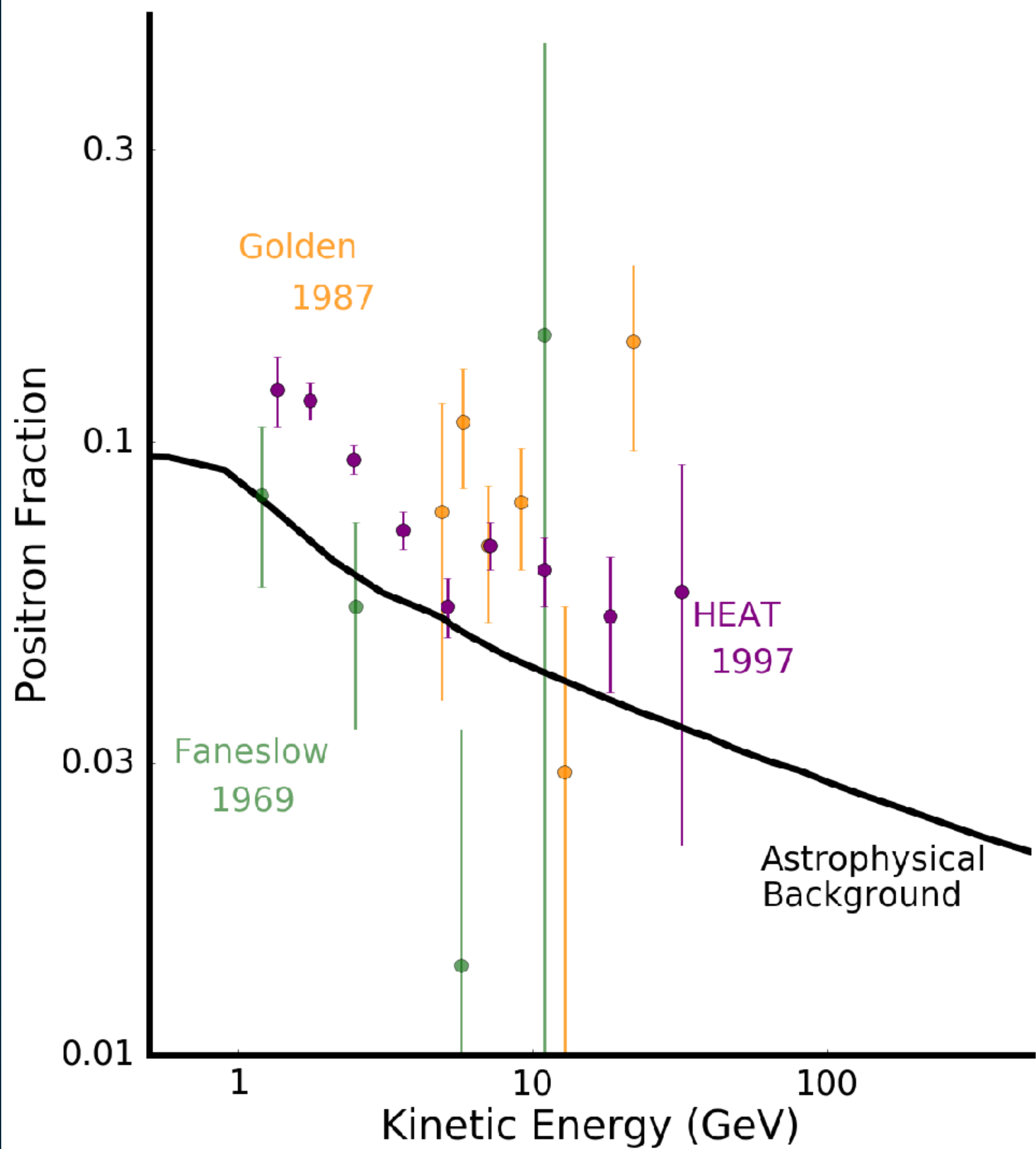
(Not an exhaustive list of observations)



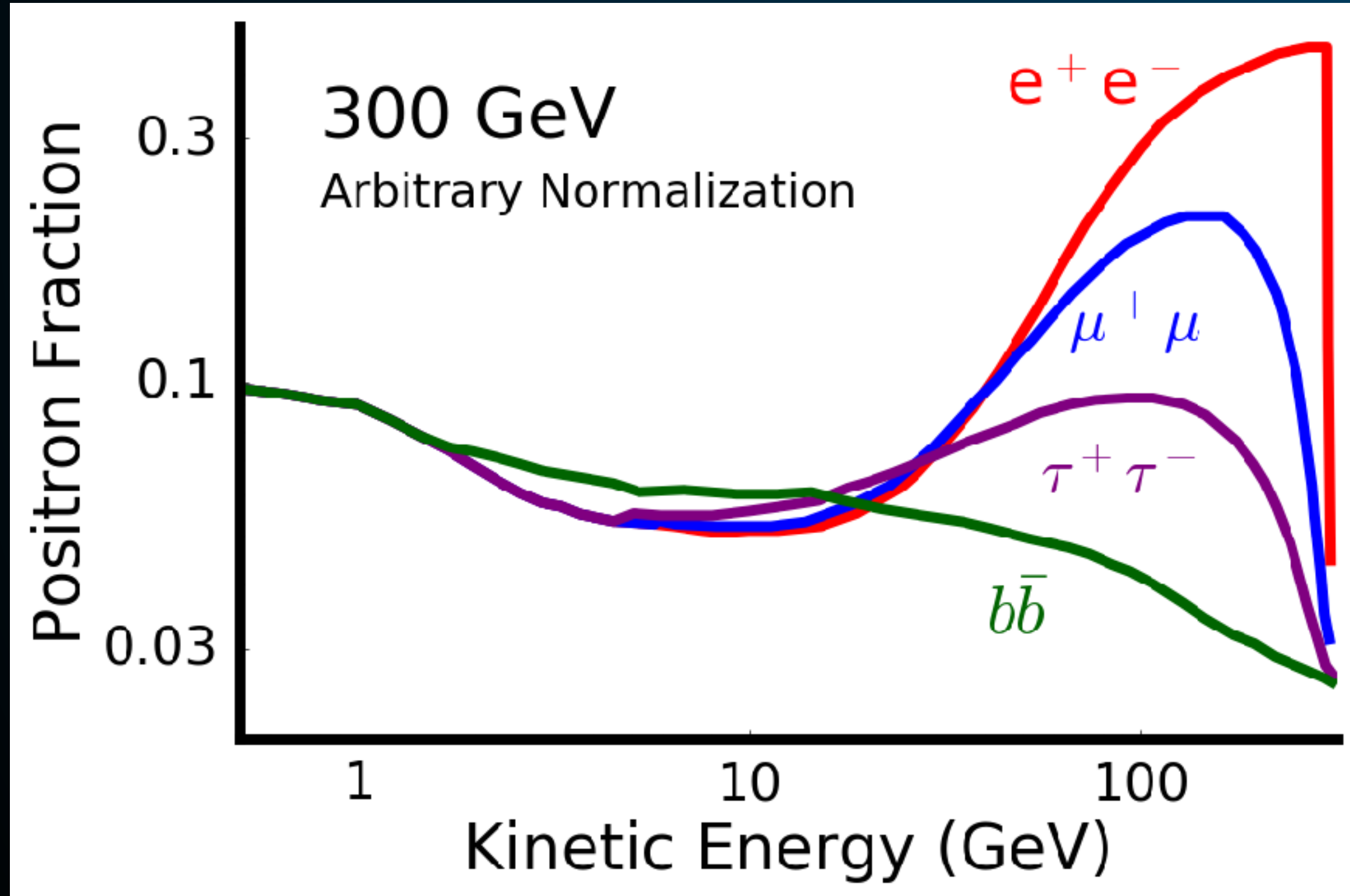
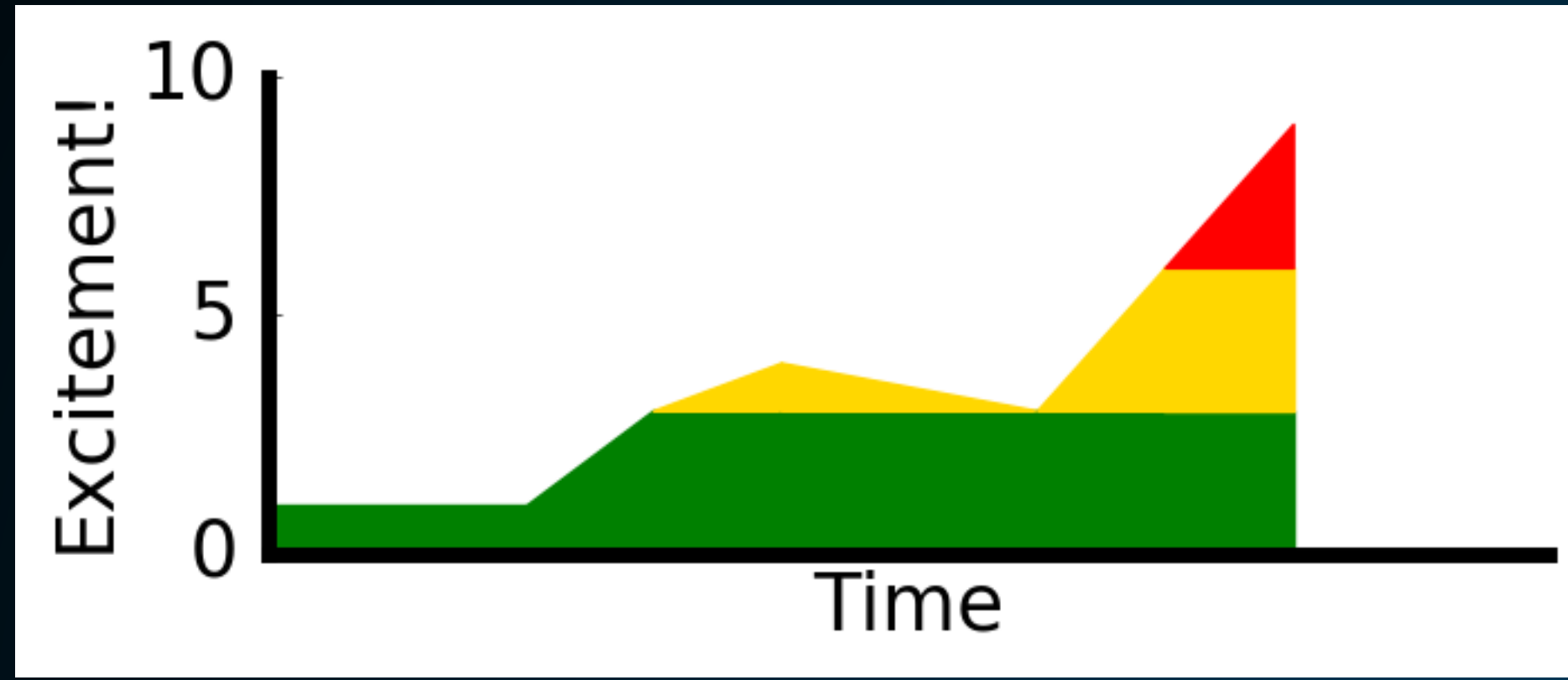
# The Positron Excess



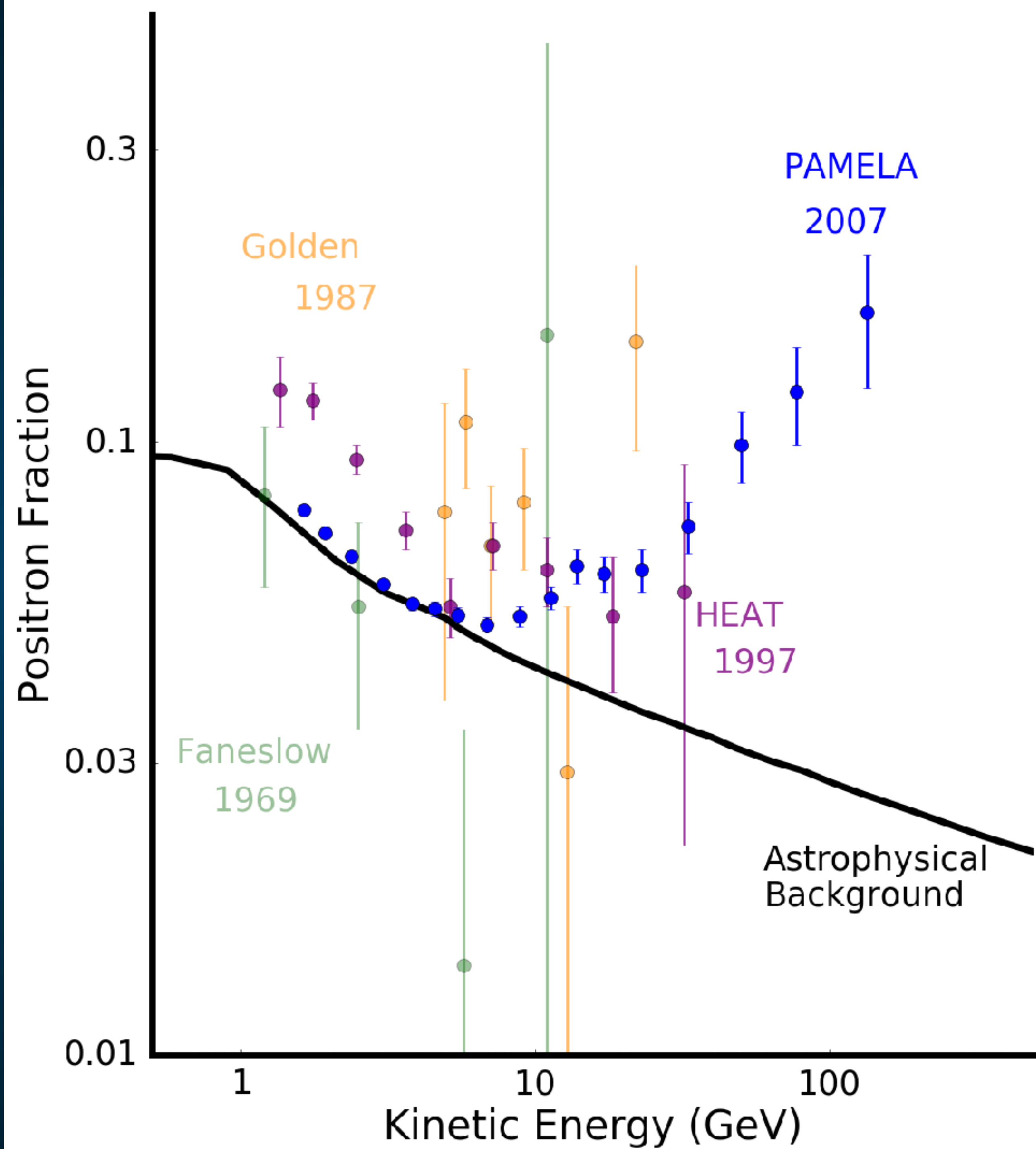
(Not an exhaustive list of observations)



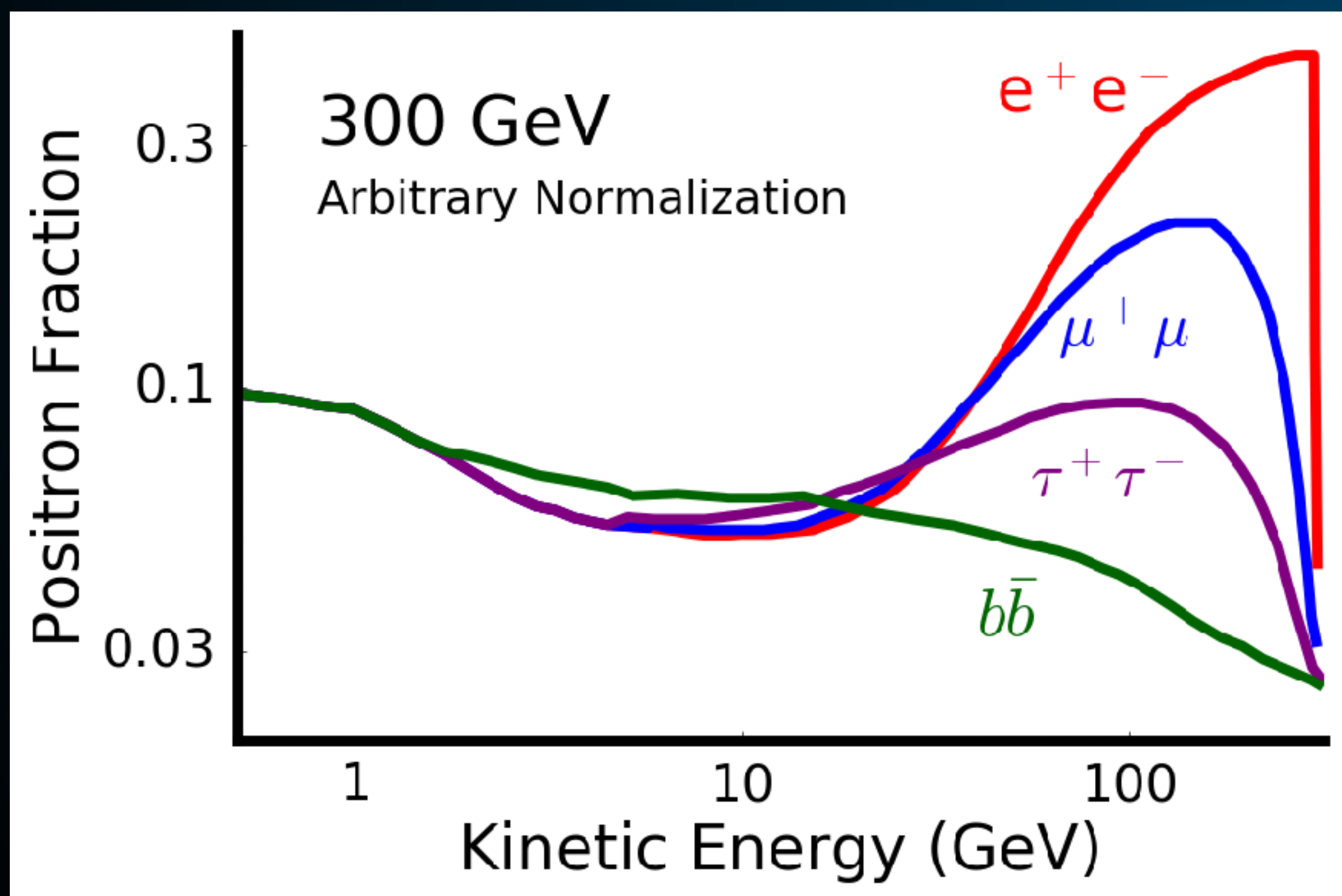
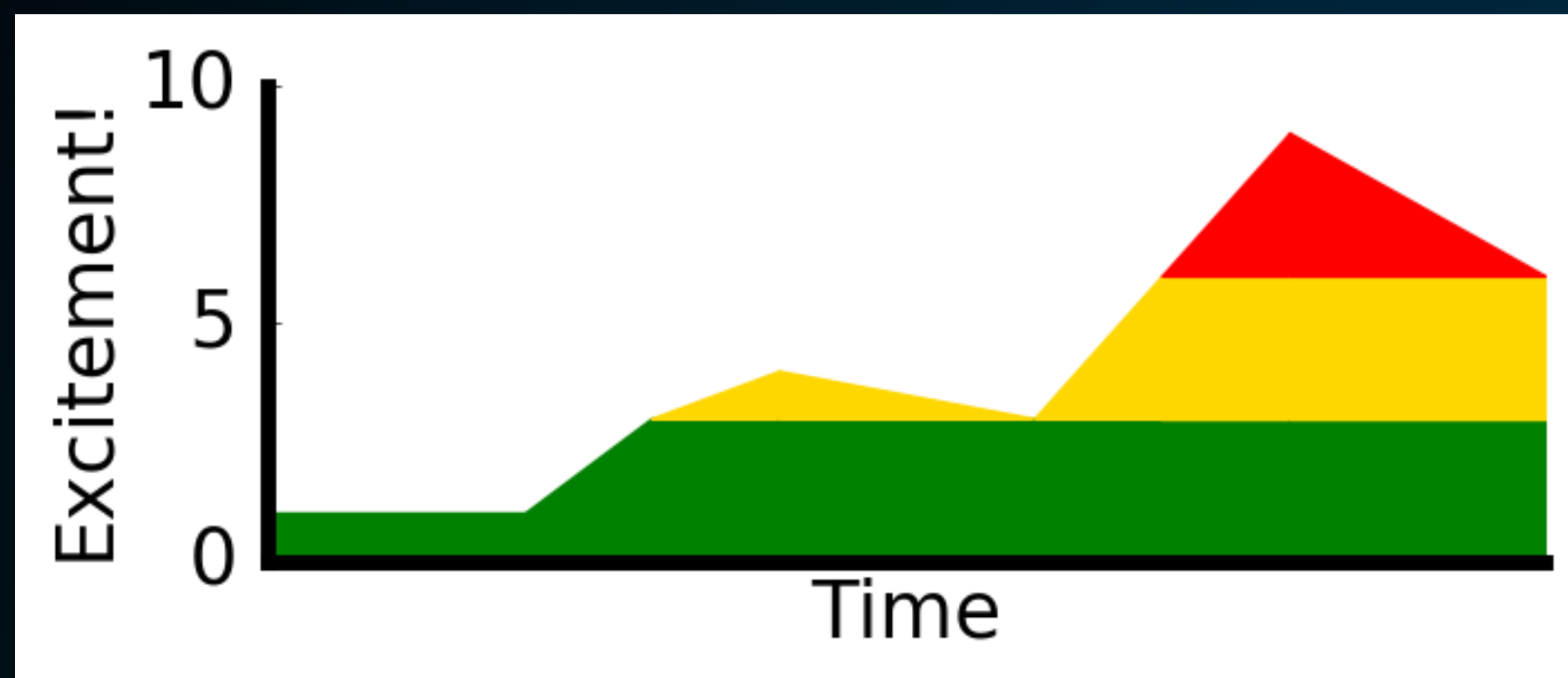
# The Positron Excess



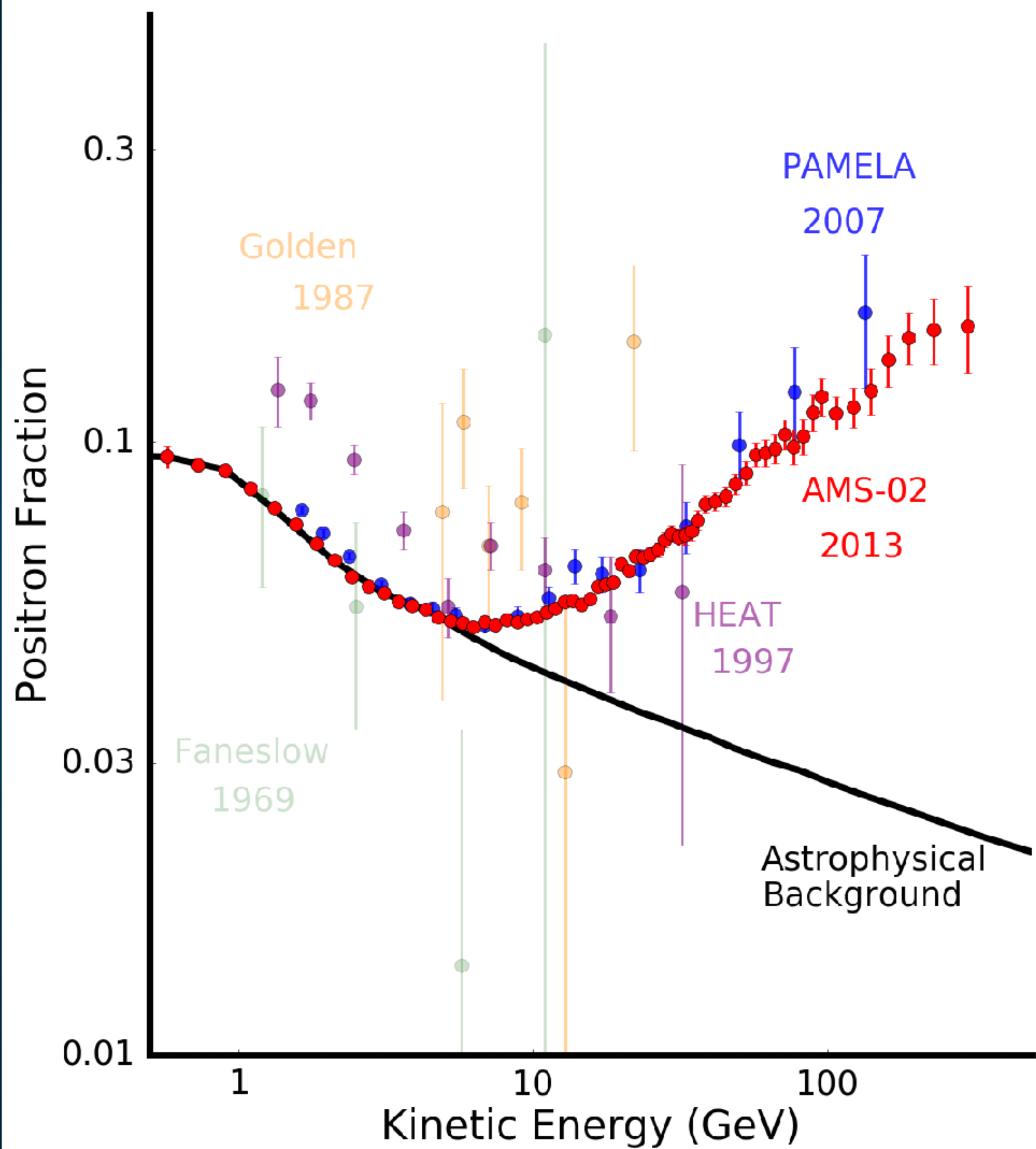
(Not an exhaustive list of observations)



# The Positron Excess



(Not an exhaustive list of observations)

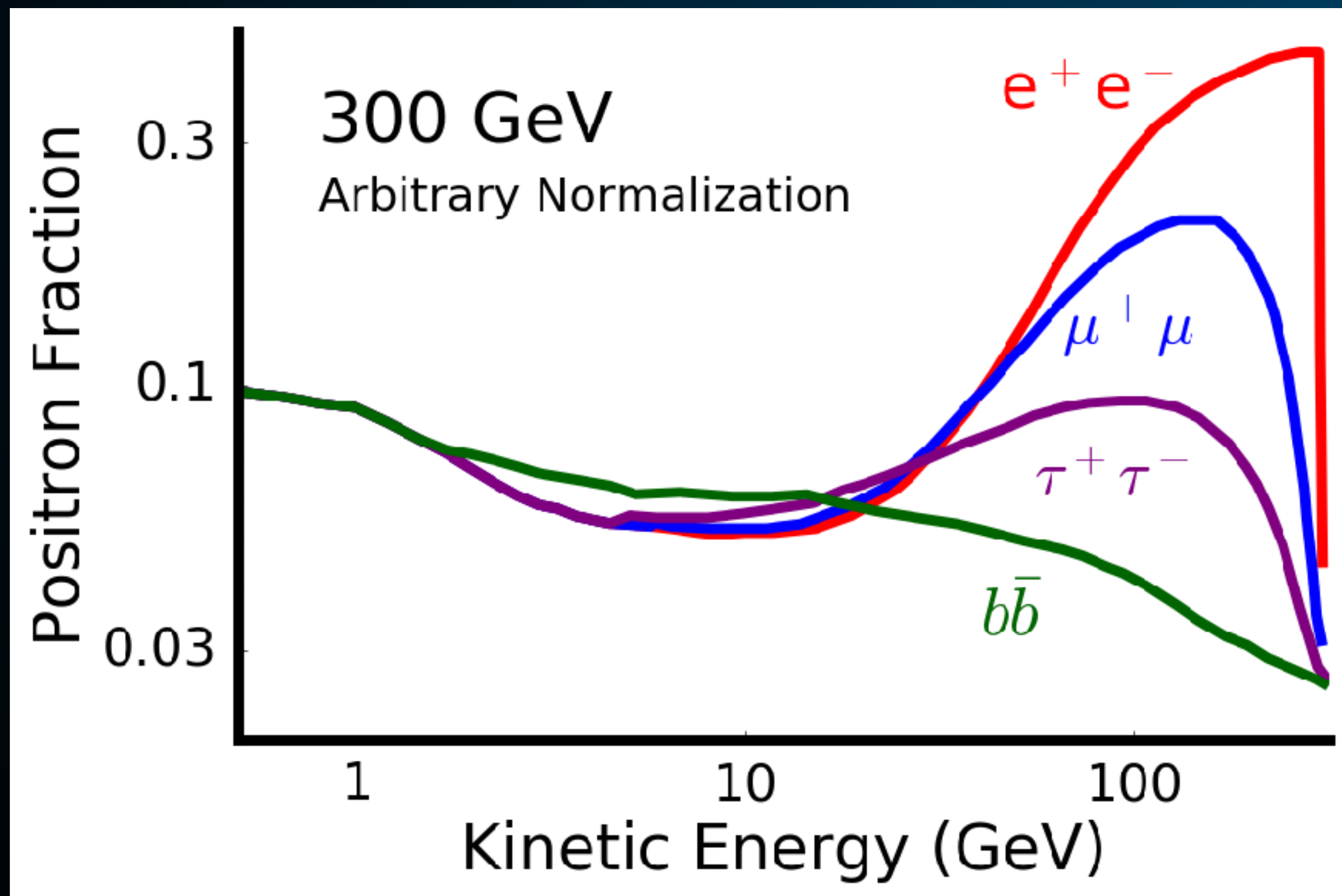


# The Positron Excess

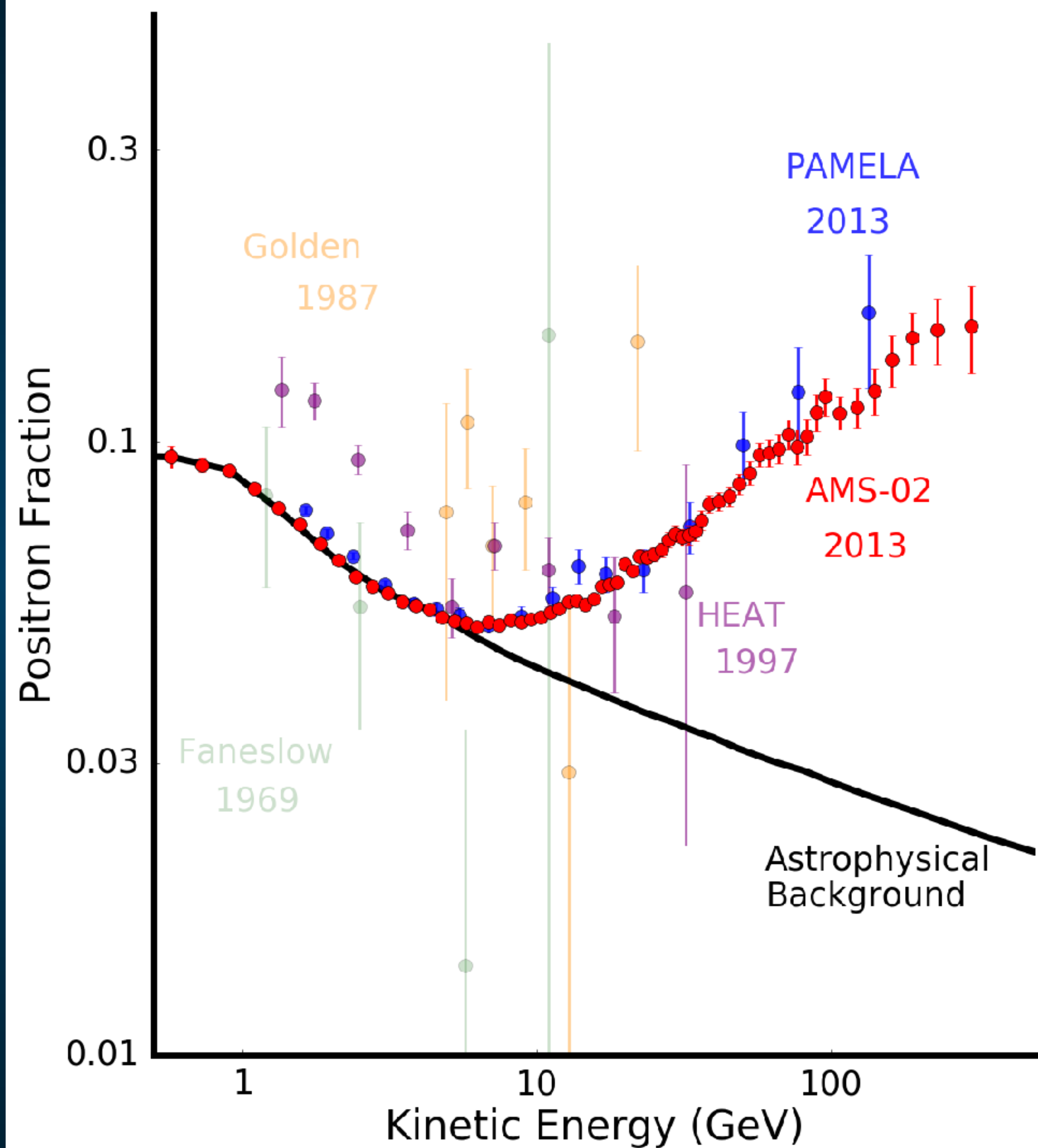
Why Less Excitement?

Continues to Higher Mass

Spectrum Relatively Smooth



(Not an exhaustive list of observations)

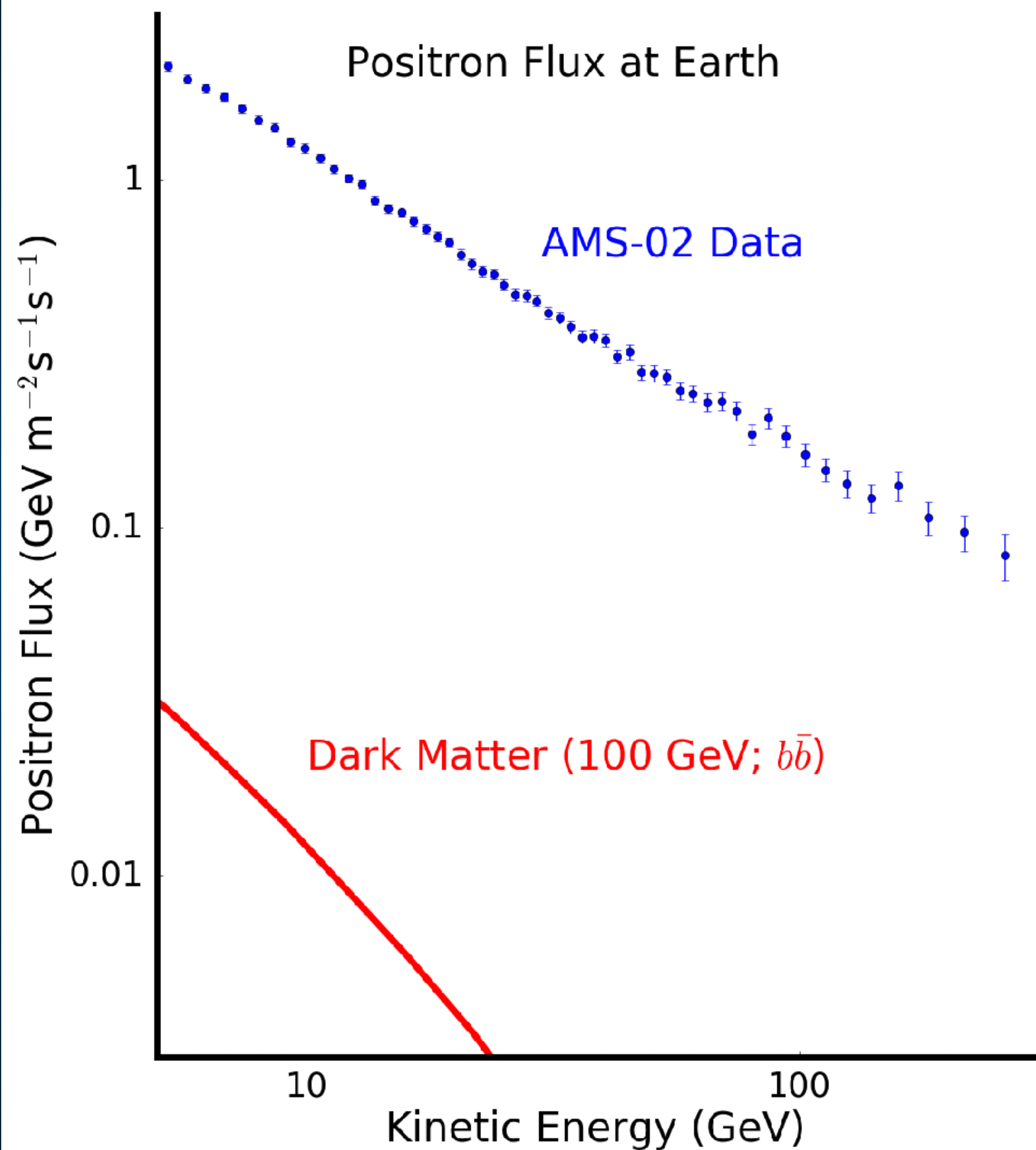
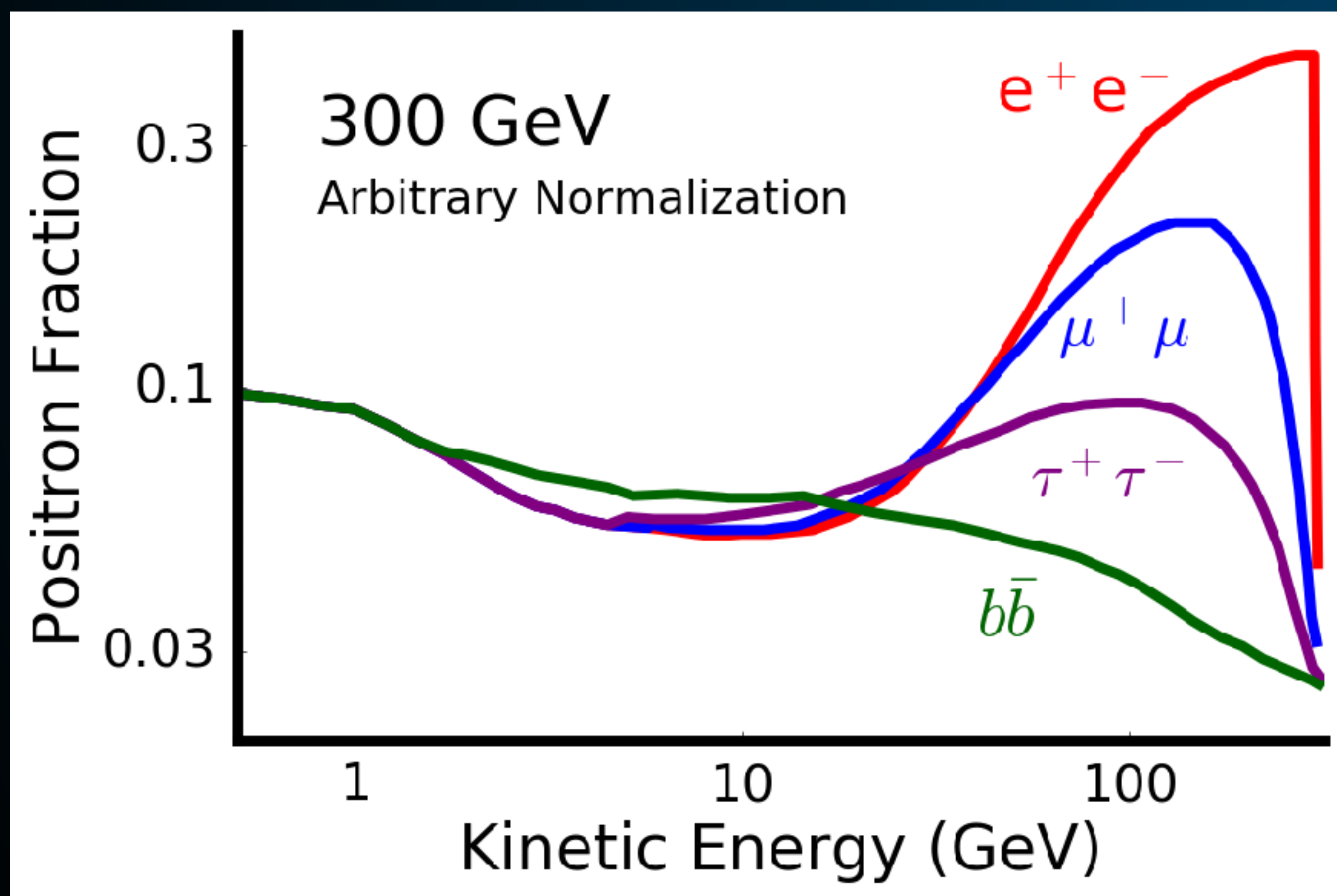


# The Positron Excess

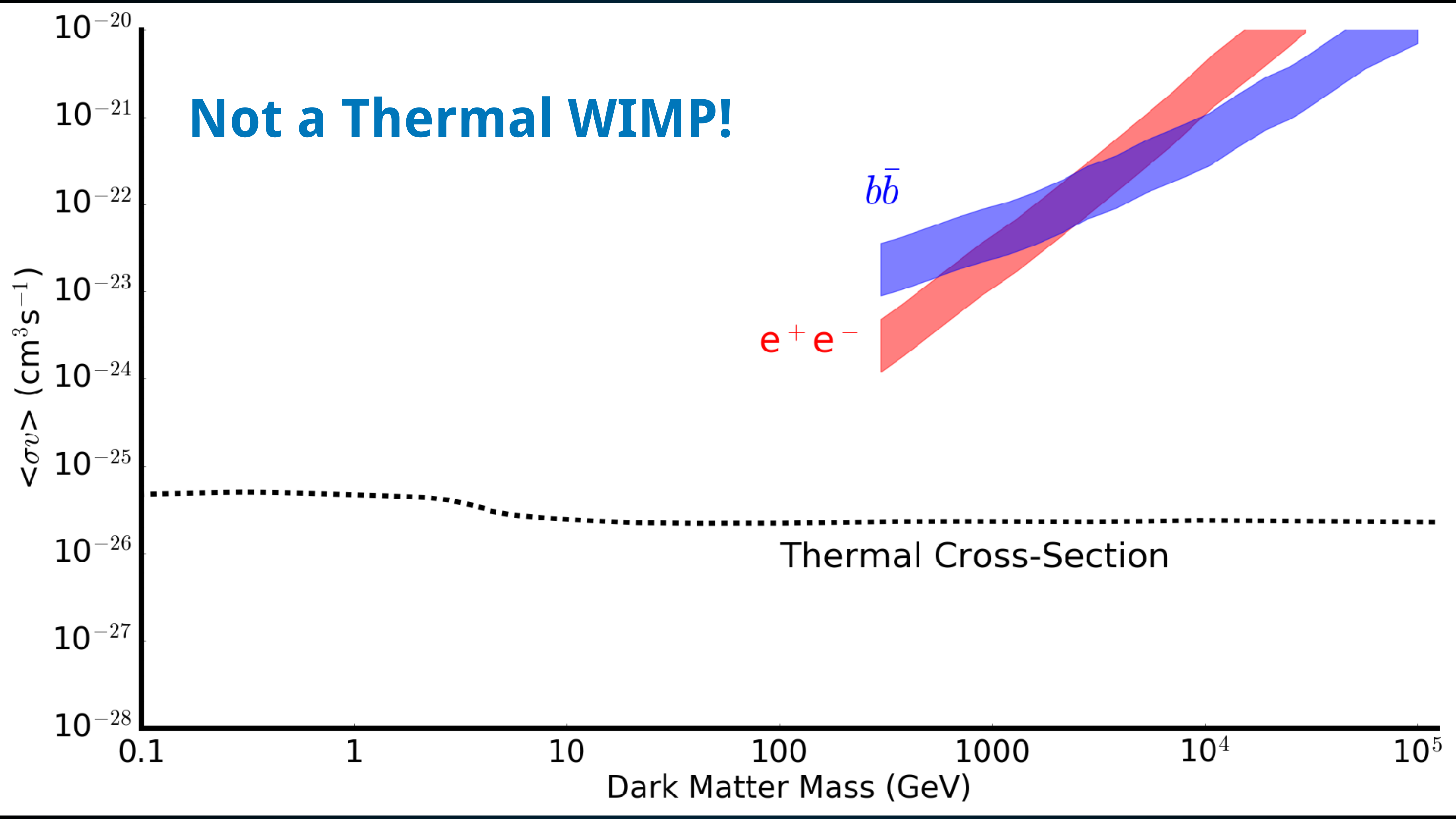
Why Less Excitement?

Continues to Higher Mass

Spectrum Relatively Smooth

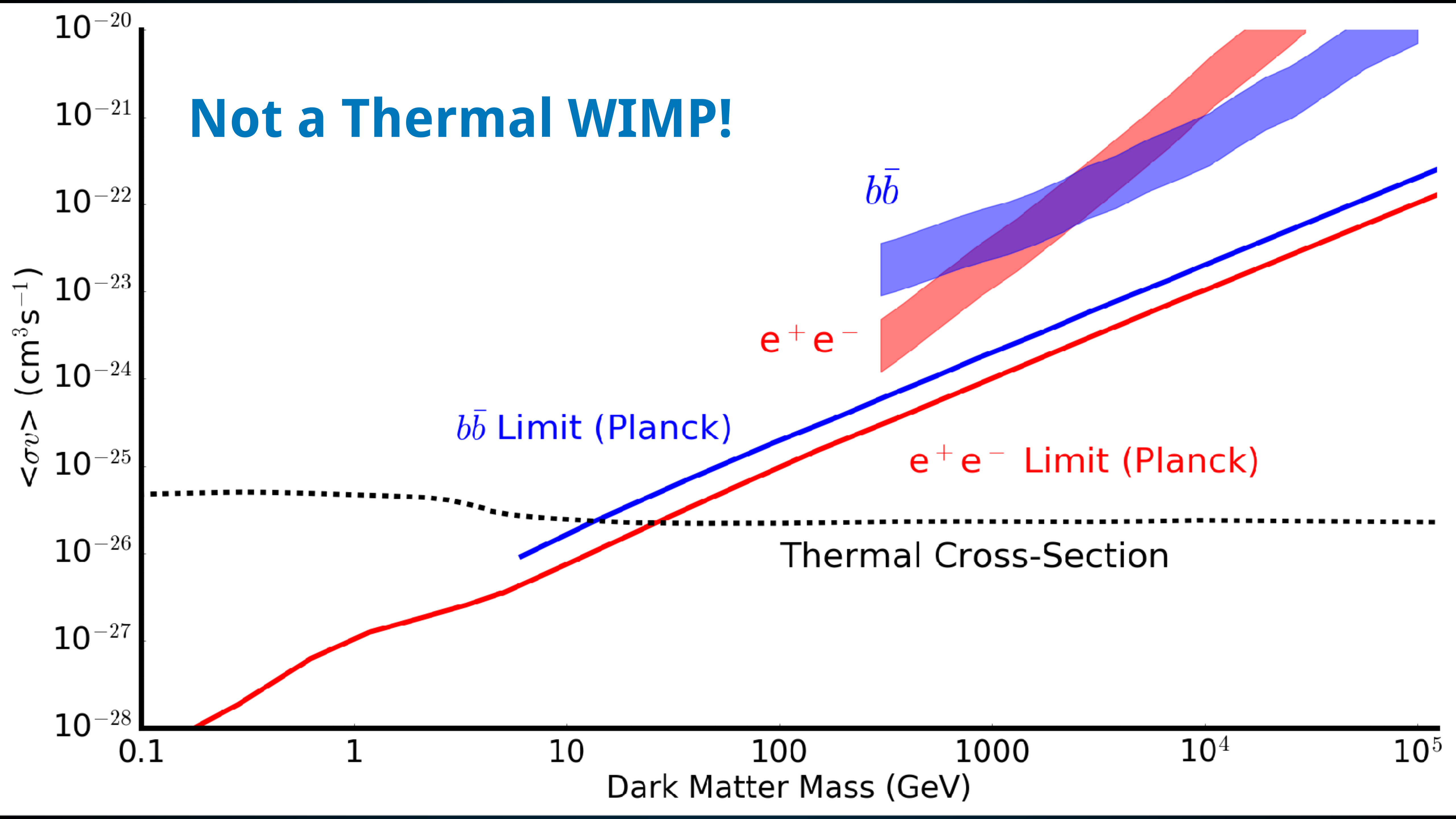


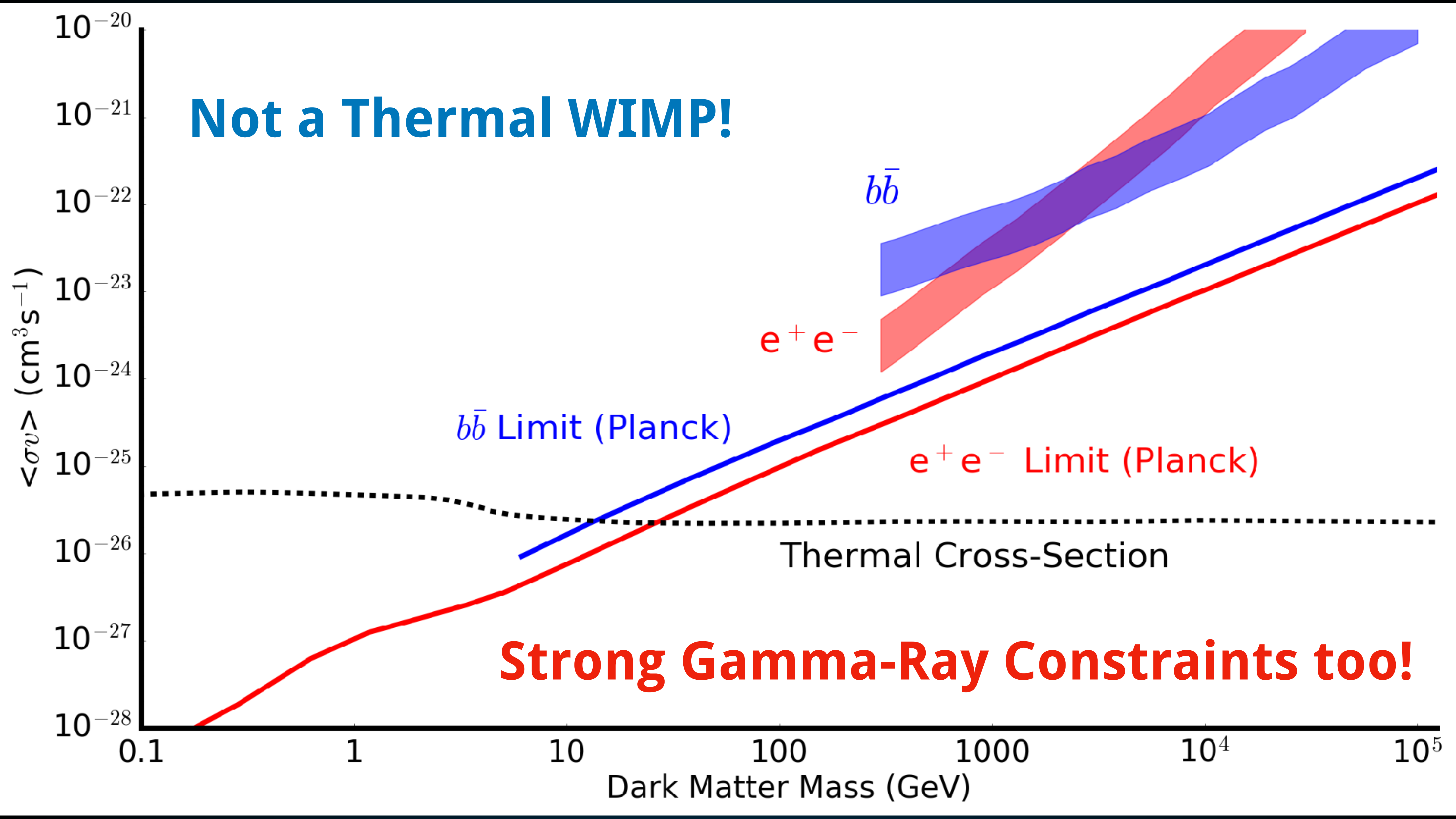
**Not a Thermal WIMP!**





**Not a Thermal WIMP!**



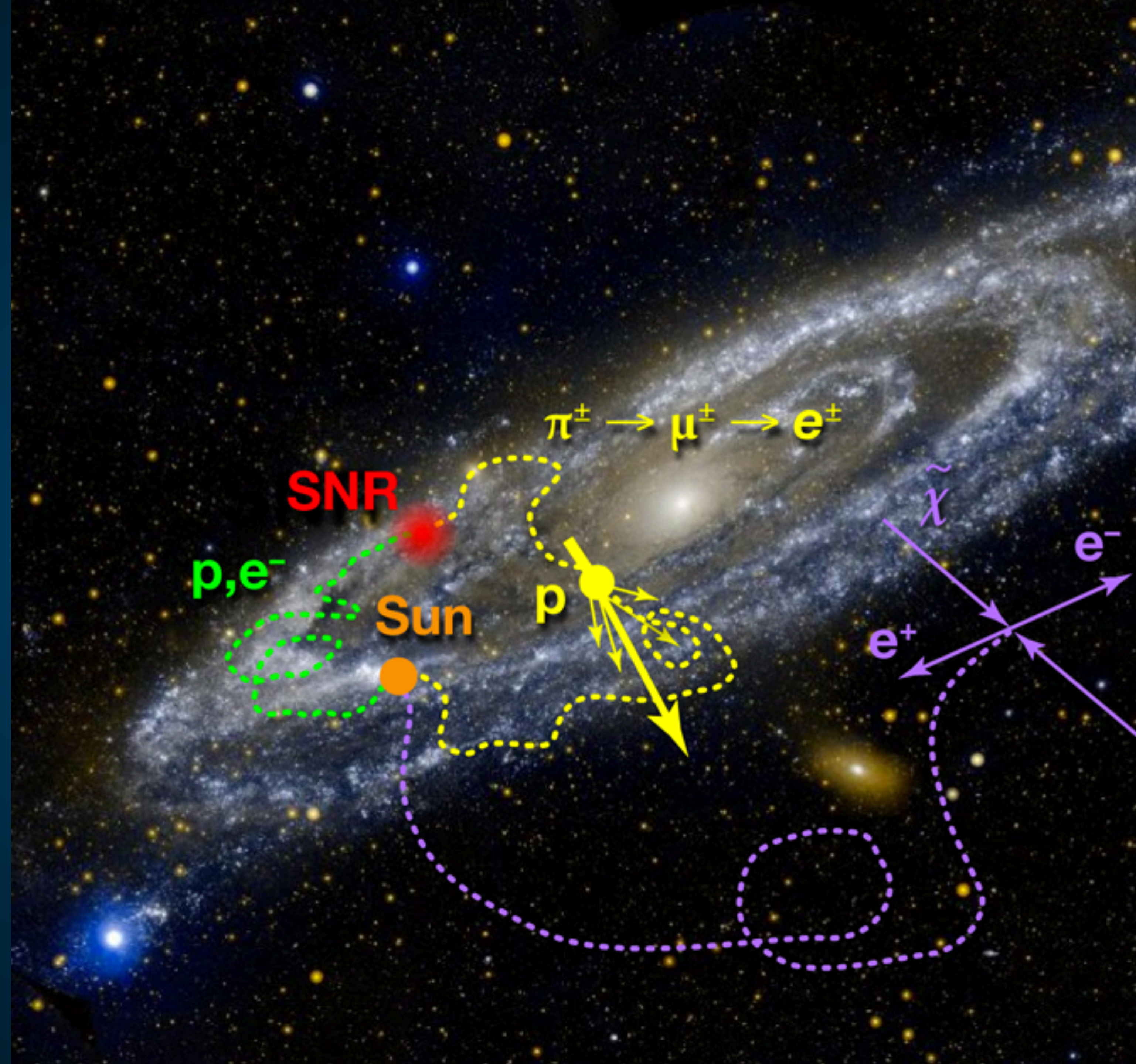
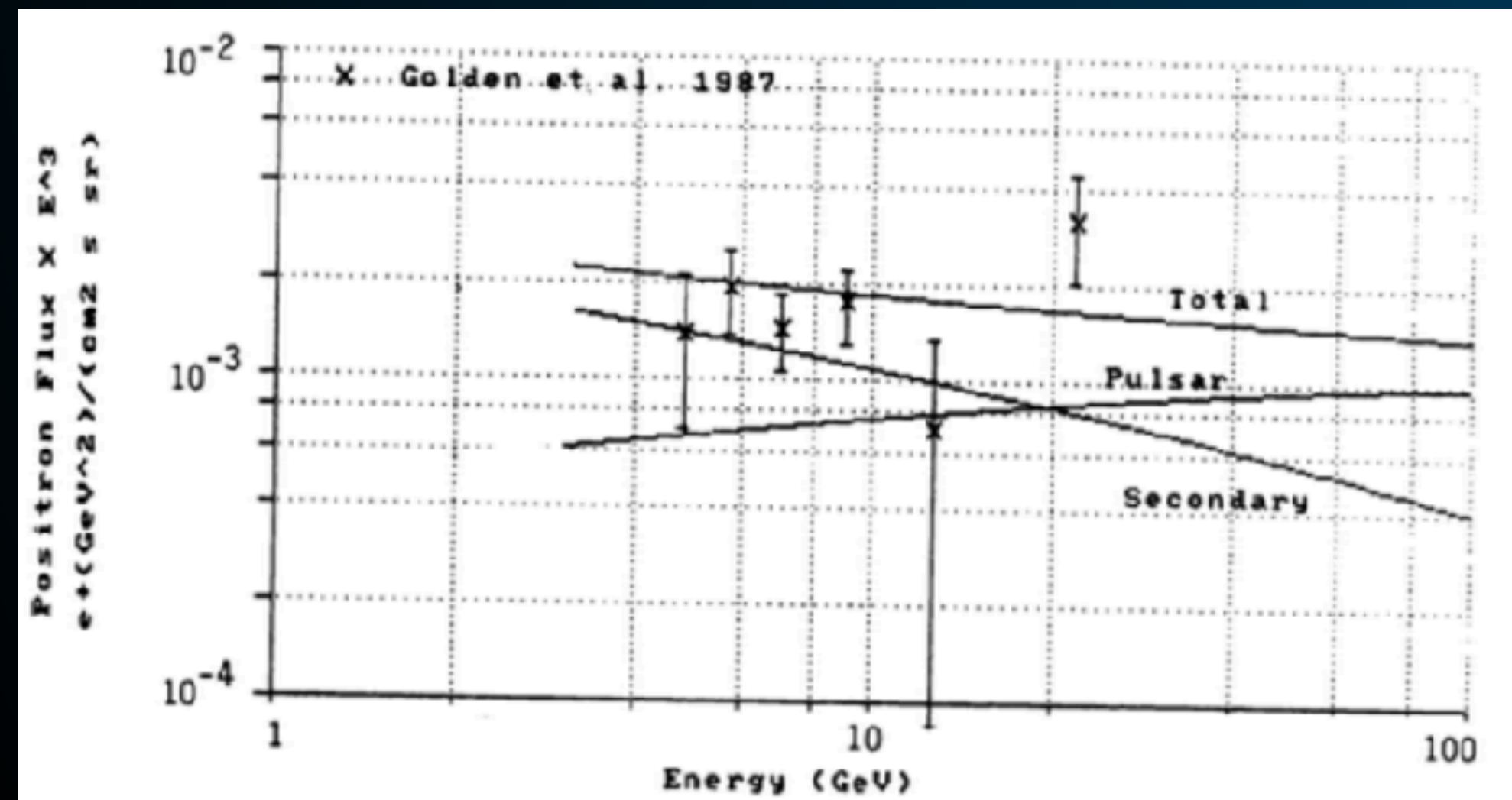


# The Positron Excess

Key Idea: Investigate the Positron Fraction!

$$\frac{\phi_{e^+}}{\phi_{e^+} + \phi_{e^-}}$$

Harding & Ramaty (ICRC! 1987)



## ► Uncertainties in pulsar models:

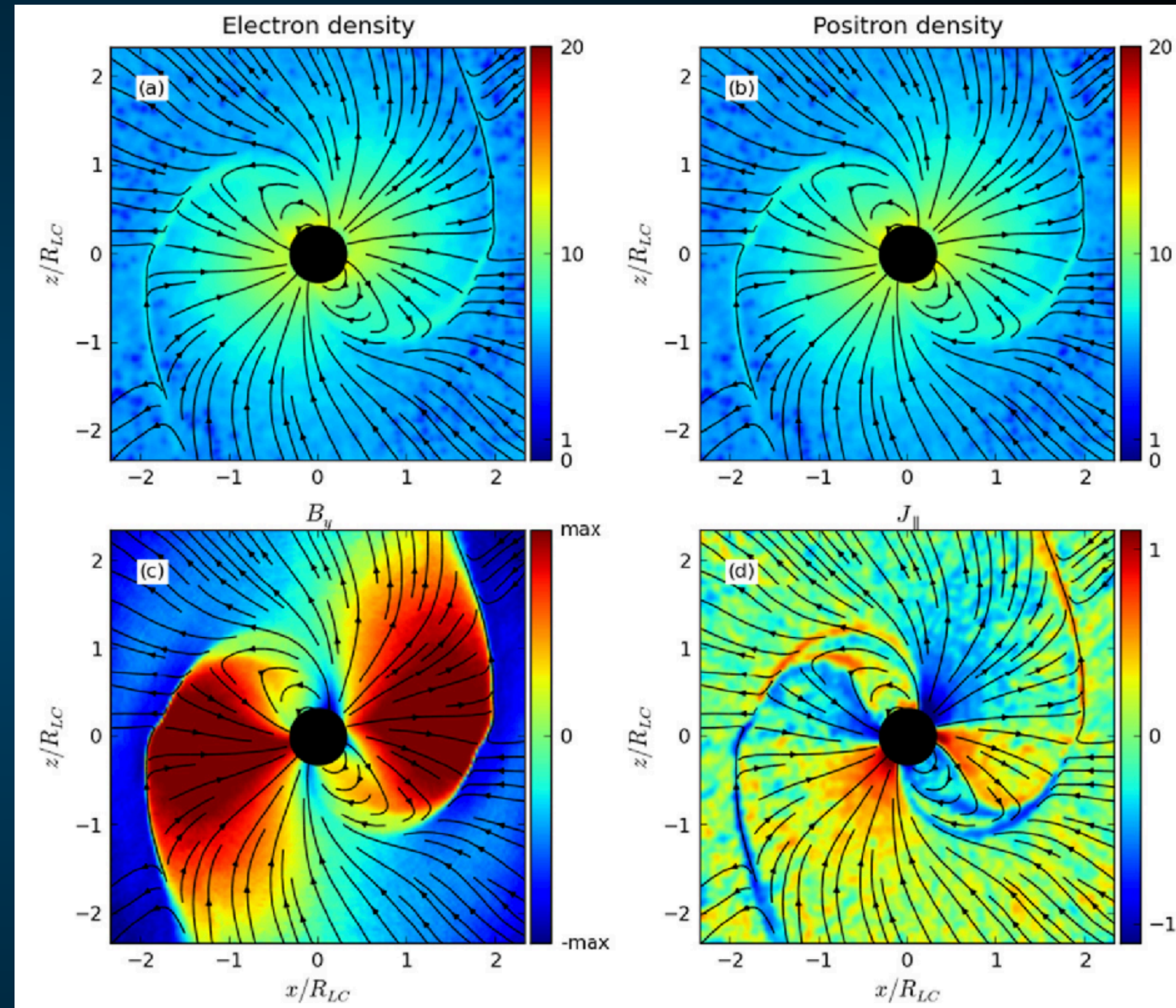
### ► I: The $e^+e^-$ production efficiency

Profumo (0812.4457); Malyshev et al. (0903.1310)

% . A quantitative discussion of plausible values for  $f_{e^\pm}$  was recently given in Ref. [38]. We shall not review their discussion here, but Ref. [38] argues (see in particular their very informative App. B and C) that in the context of a standard model for the pulsar wind nebulae, a reasonable range for  $f_{e^\pm}$  falls between 1% and 30%.

### ► II: The $e^+e^-$ spectrum.

### ► III: The propagation of $e^+e^-$ to Earth.



## ► Uncertainties in pulsar models:

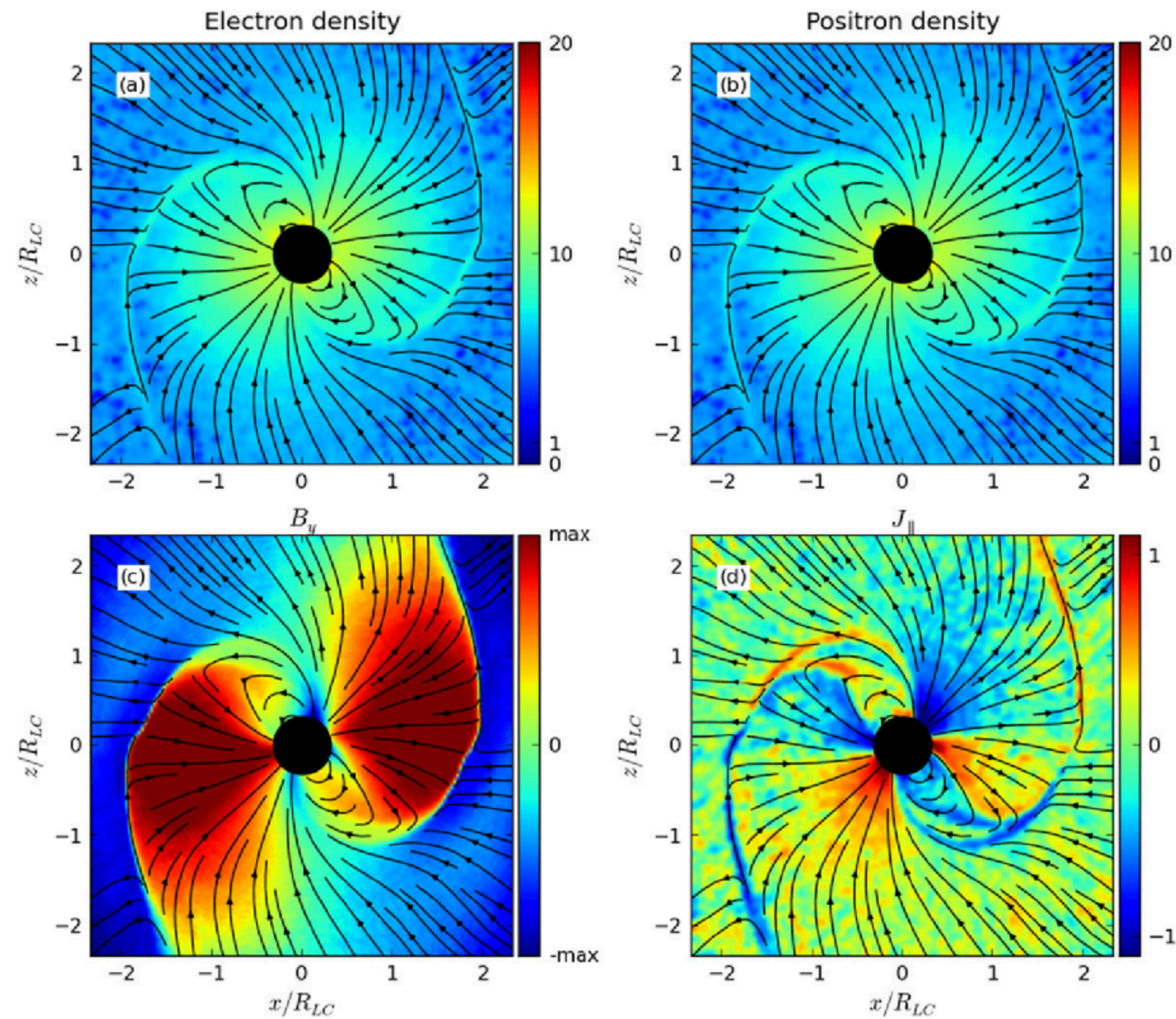
► I: The  $e^+e^-$  production efficiency

► II: The  $e^+e^-$  spectrum.

Hooper et al. (0810.1527)

part of their energy adiabatically because of the expansion of the wind. The energy spectrum injected by a single pulsar depends on the environmental parameters of the pulsar, but some attempts to calculate the average spectrum injected by a population of mature pulsars suggest that the spectrum may be relatively hard, having a slope of  $\sim 1.5-1.6$  [18]. This spectrum, however, results from a complex interplay of individual pulsar spectra, of the spatial and age distributions of pulsars in the Galaxy, and on the assumption that the chief channel for pulsar spin down is magnetic dipole radiation. Due to the related uncertainties, variations from this injection spectra cannot be ruled out. Typically, one concentrates the attention on pulsars of age  $\sim 10^5$  years because younger pulsars are likely to still

► III: The propagation of  $e^+e^-$  to Earth.



# Pulsar Fits to the Positron Excess

## ► Uncertainties in pulsar models:

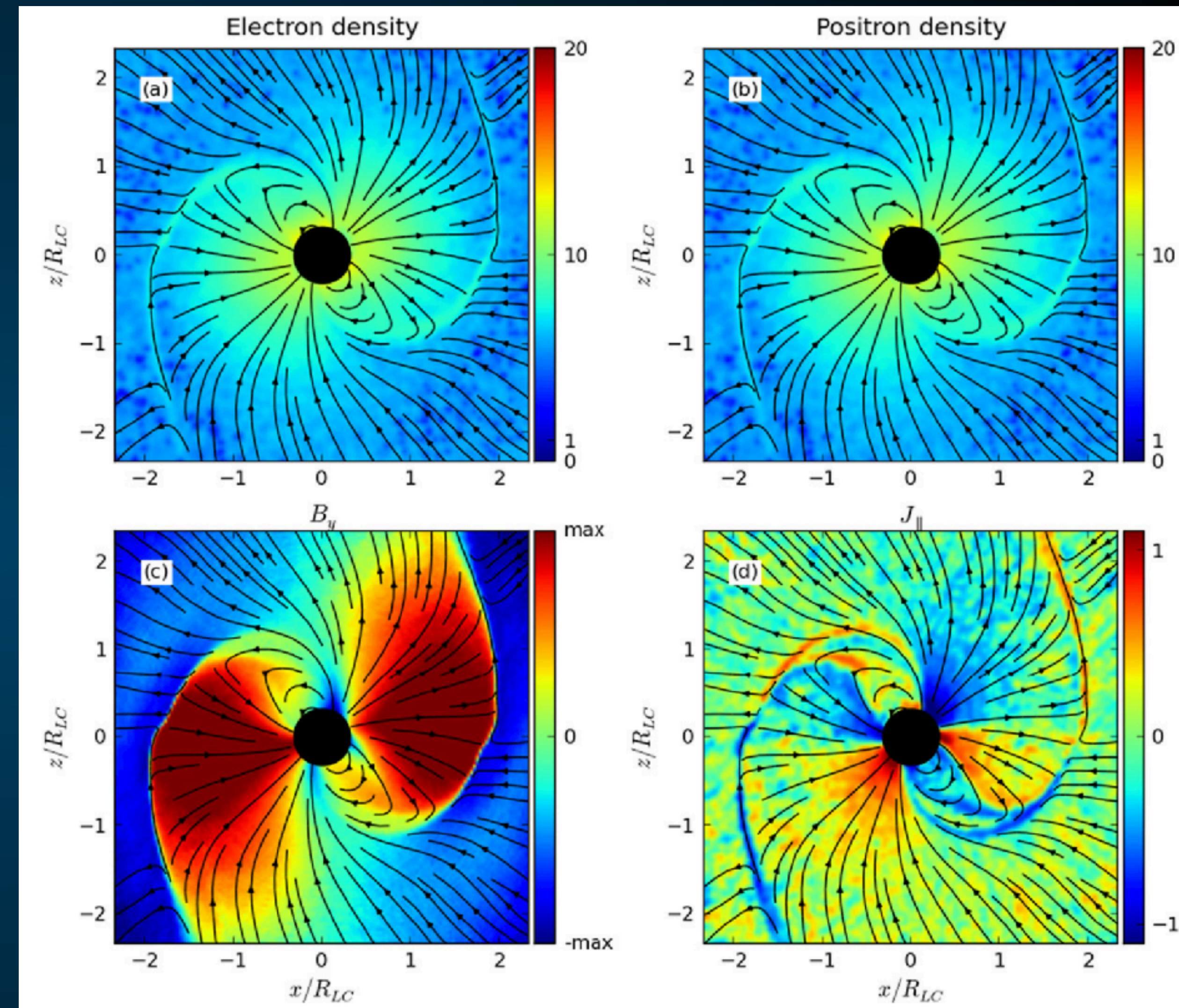
### ► I: The $e^+e^-$ production efficiency

### ► II: The $e^+e^-$ spectrum.

### ► III: The propagation of $e^+e^-$ to Earth.

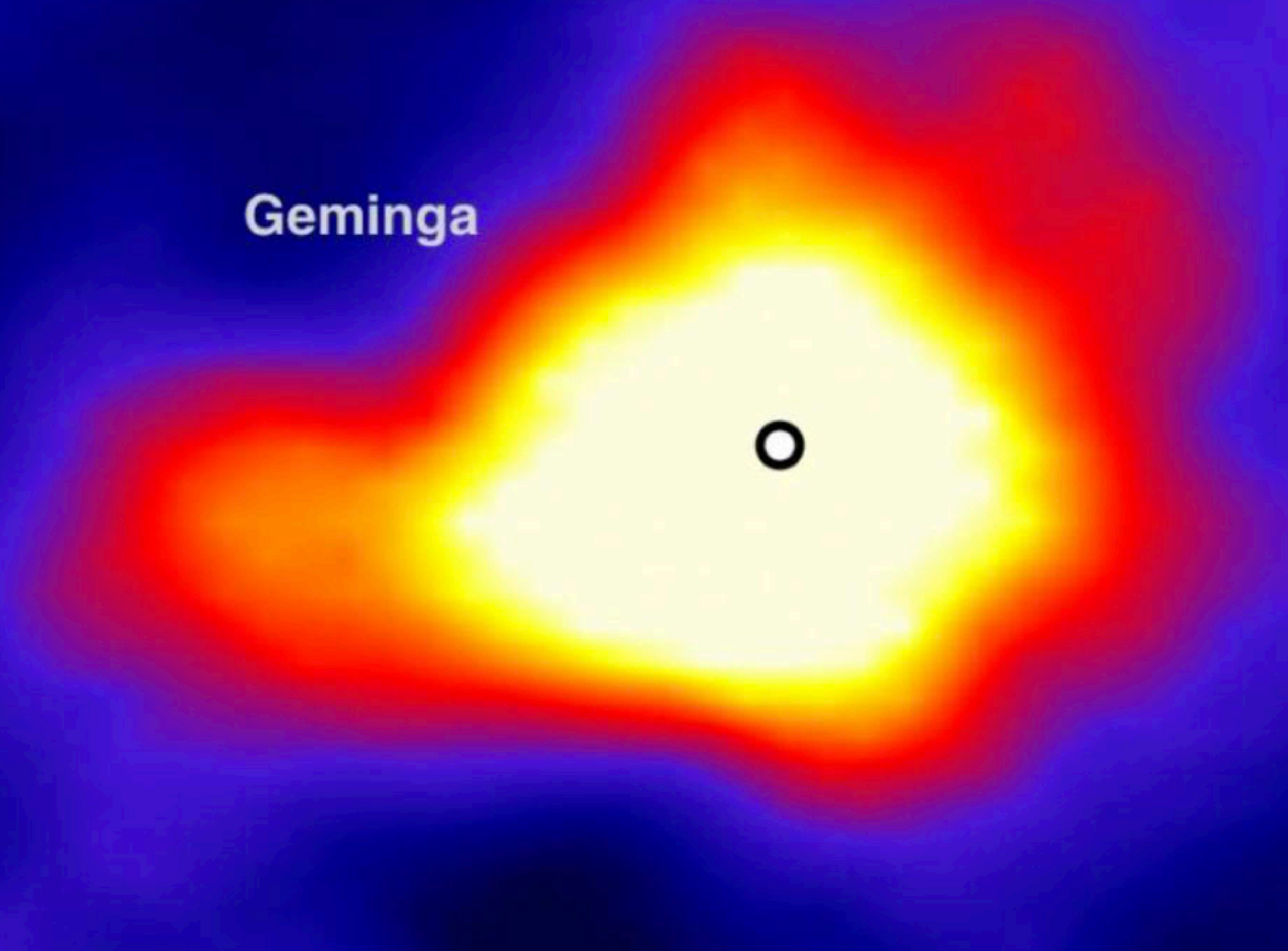
Malyshev et al. (0903.1310)

The observed spectrum on Earth of electrons and positrons injected by pulsars is also strongly dependent on propagation effects. In particular, the observed cutoff in the flux of electrons from a pulsar can be much smaller than the injection cutoff due to energy losses (“cooling”) during propagation. We define the cooling break,  $E_{br}(t)$ , as the maximal energy electrons can have after propagating for time  $t$ . Since – as stated above – the typical

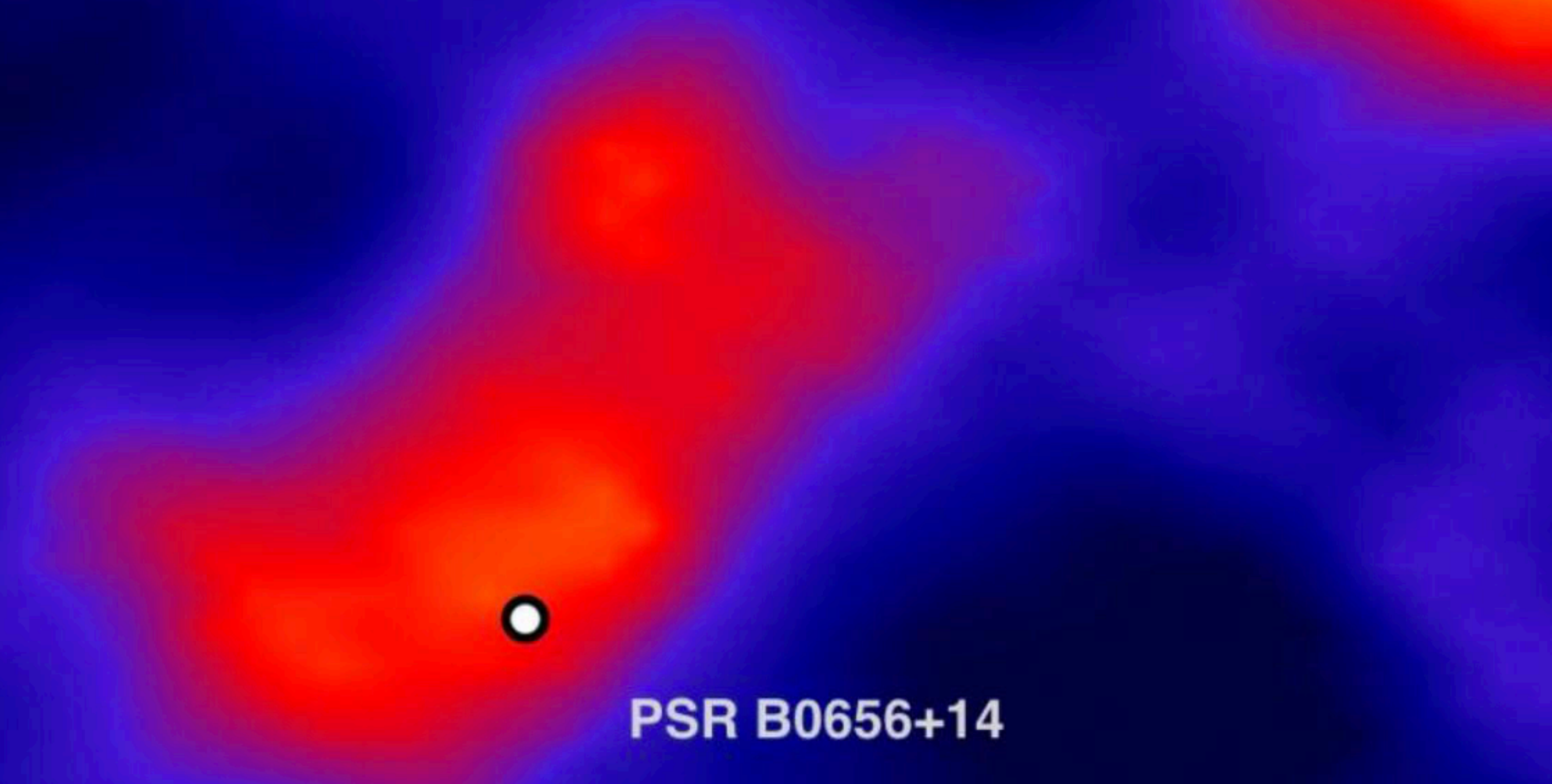




Geminga



PSR B0656+14

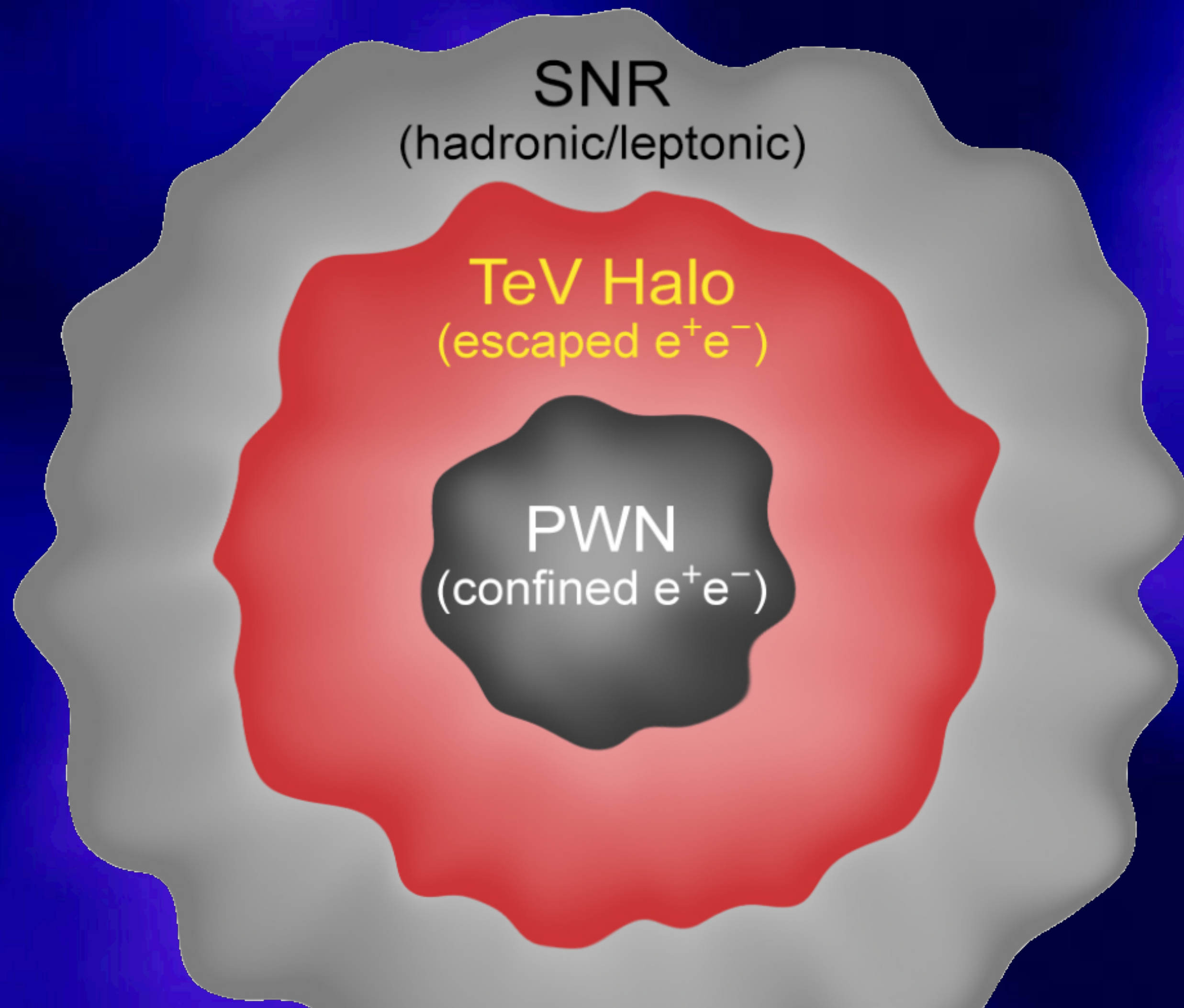




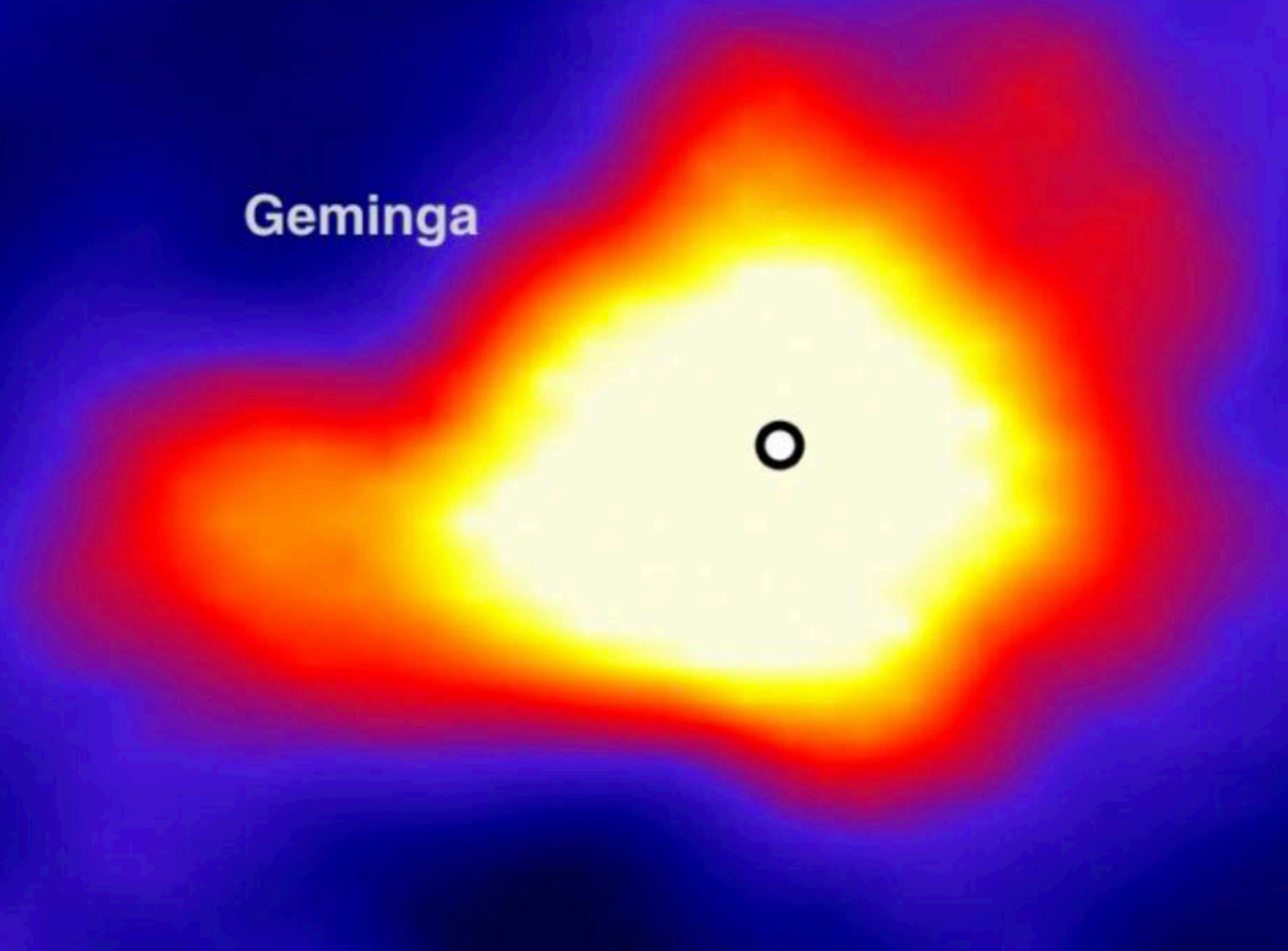
Moon (To Scale)

Linden et al. (2017; 1703.09704)

Sudoh, TL, Beacom (2019; 1902.08203)



Geminga



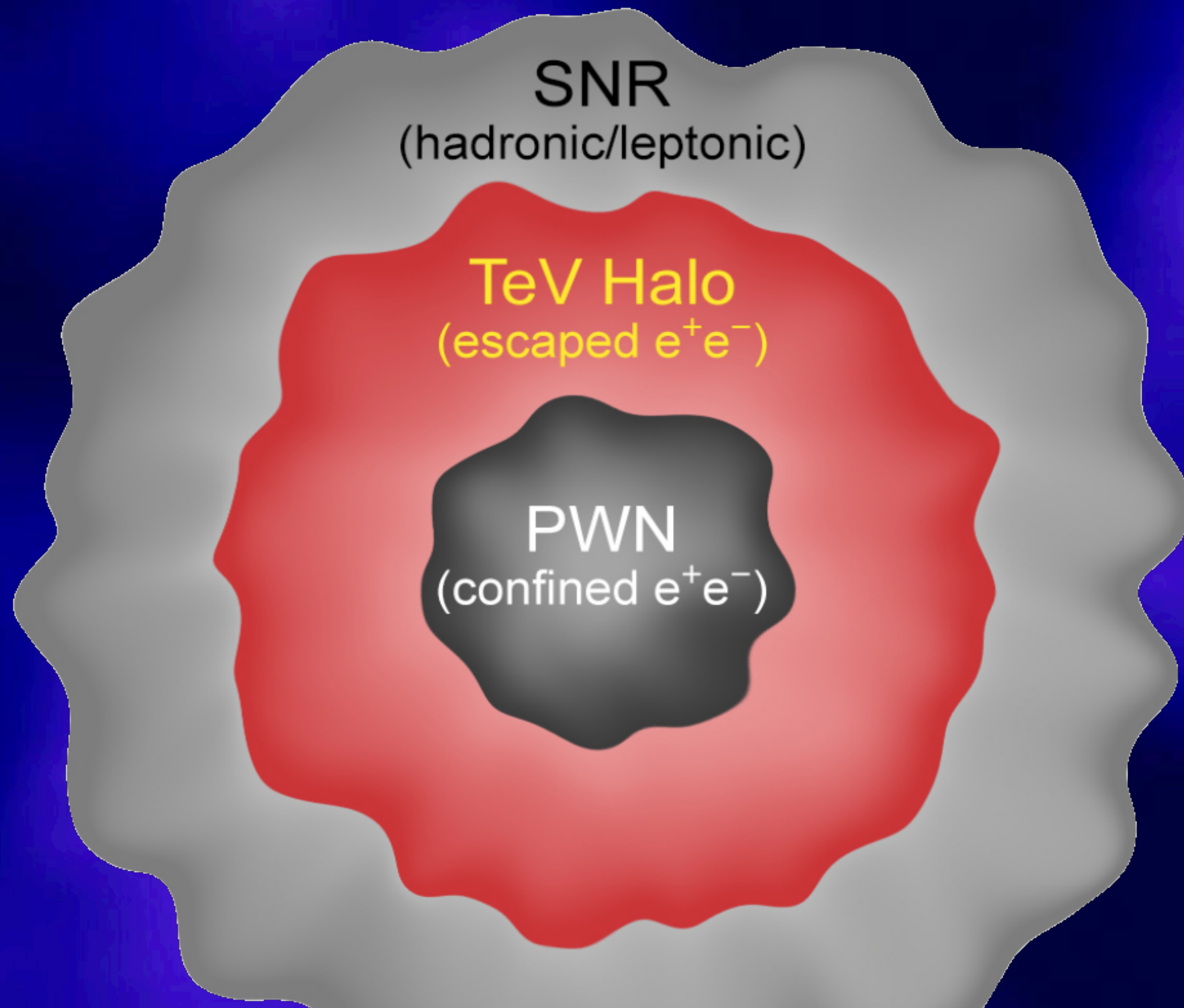




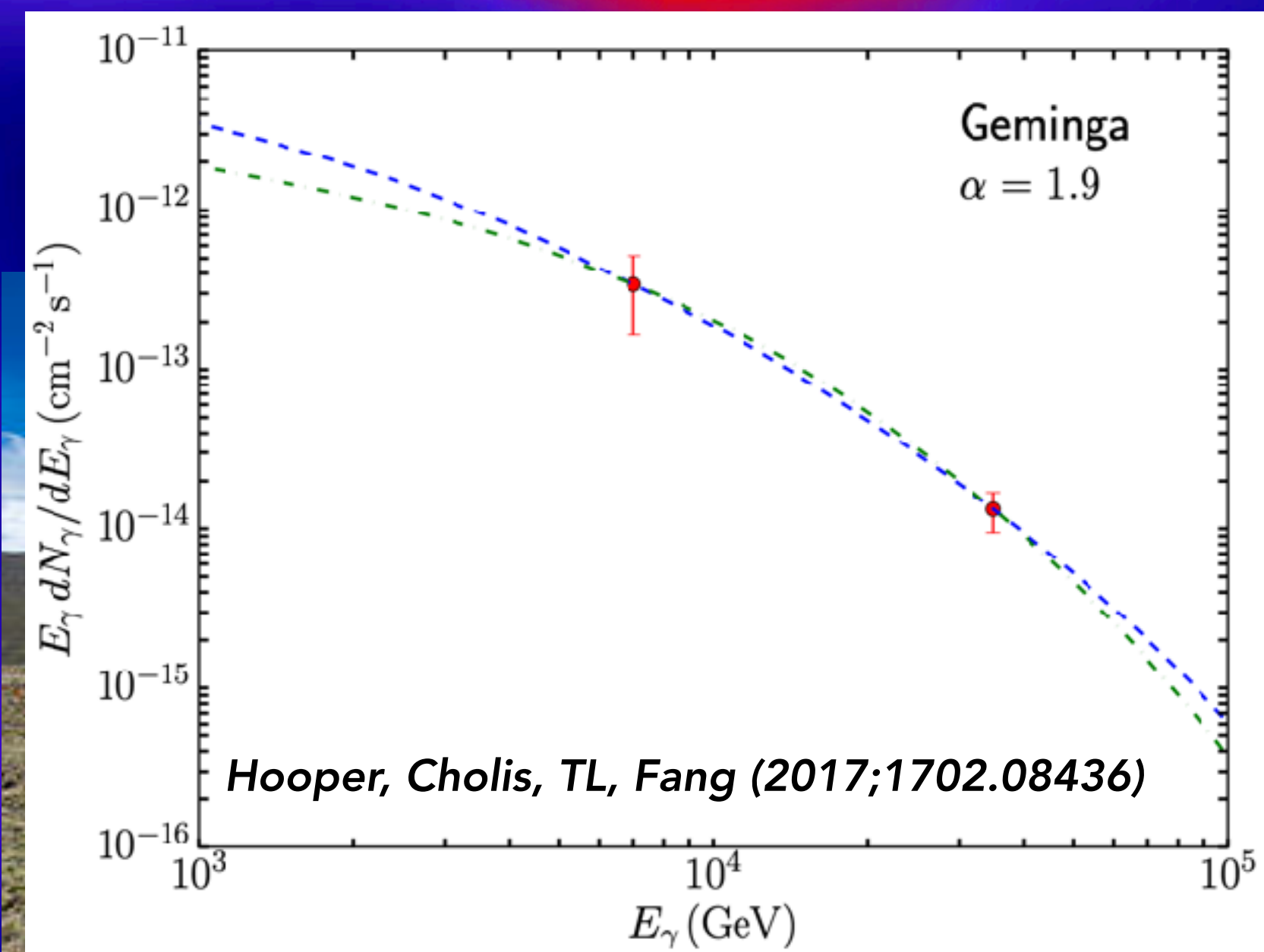
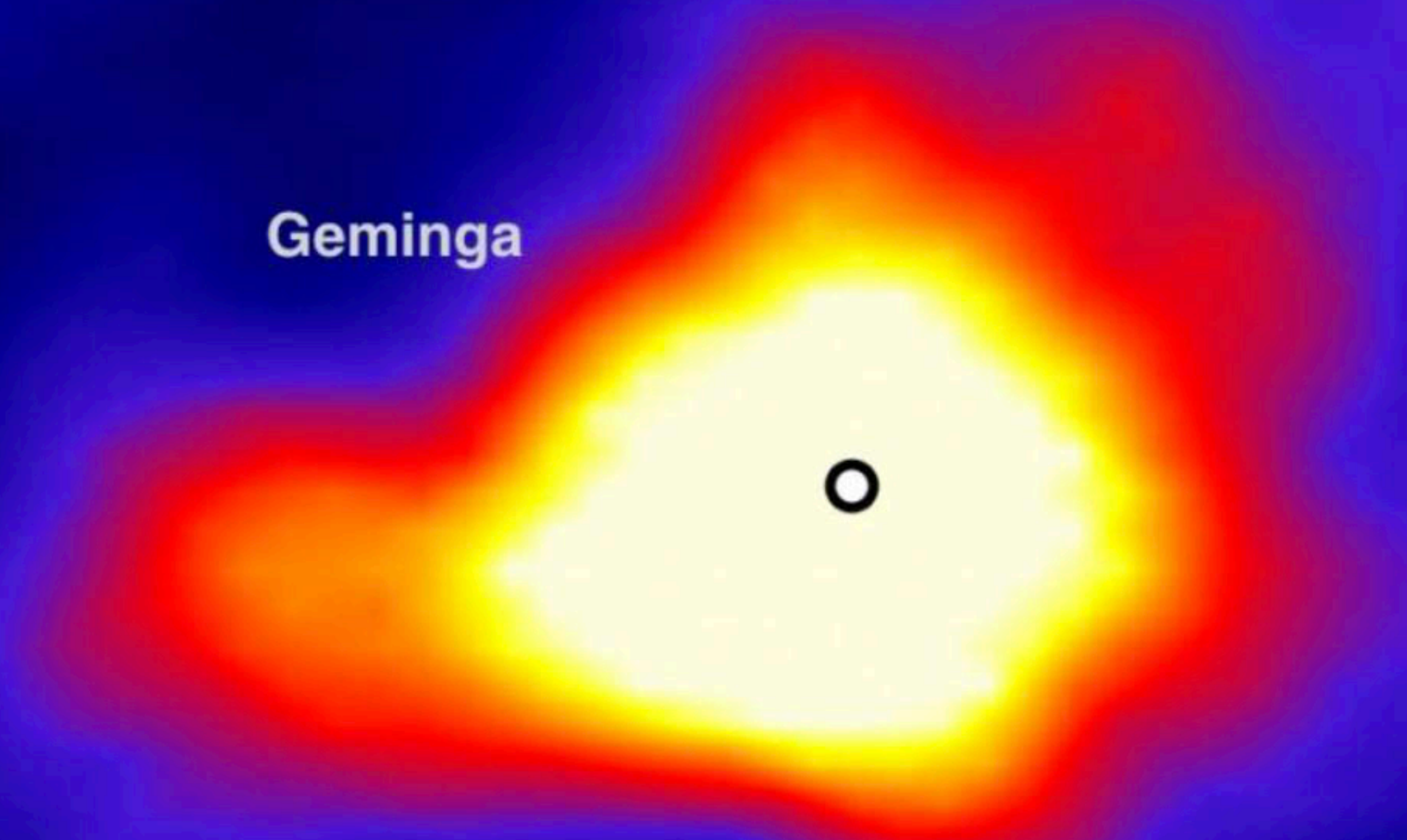
Moon (To Scale)

Linden et al. (2017; 1703.09704)

Sudoh, TL, Beacom (2019; 1902.08203)



Geminga

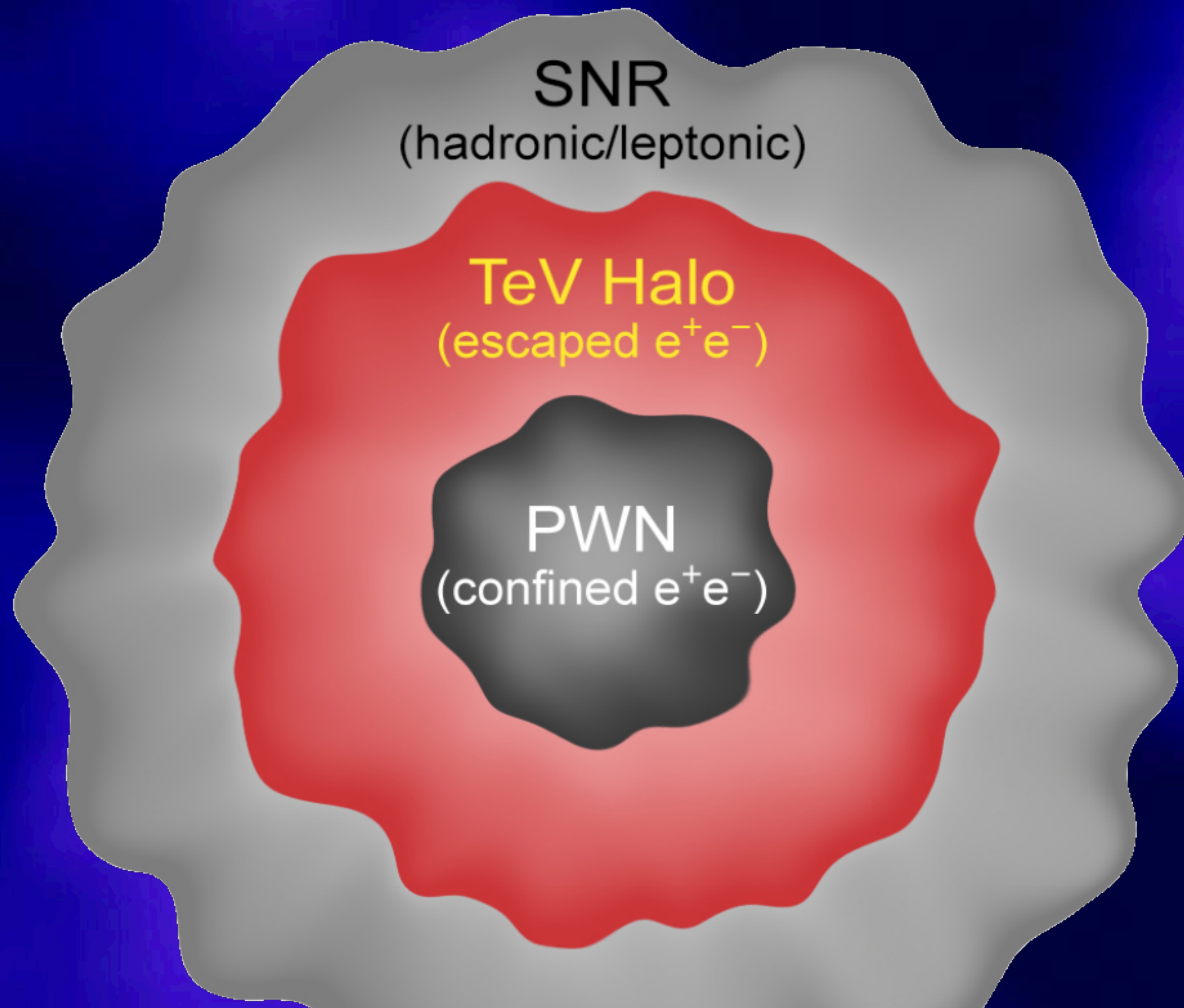




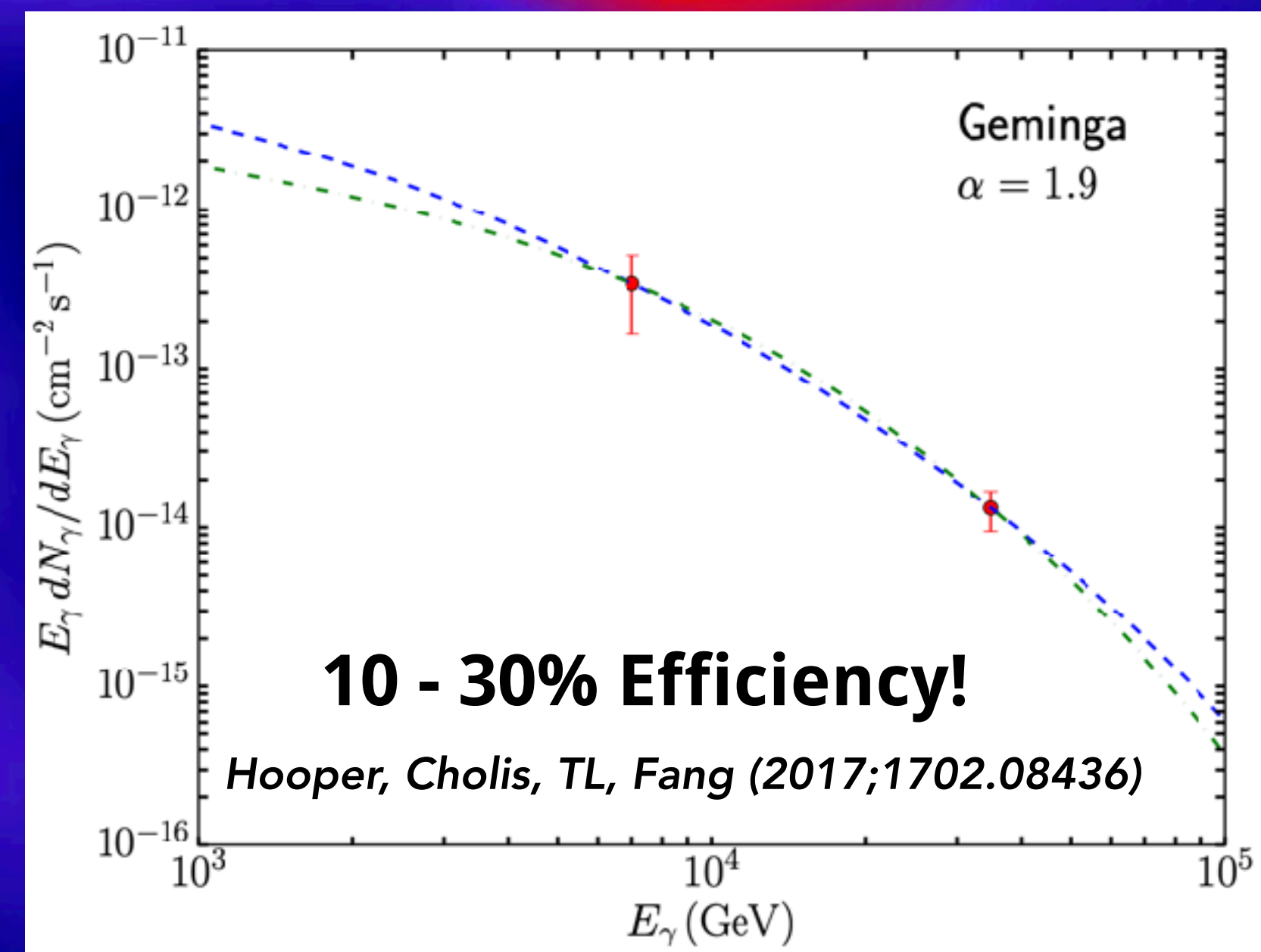
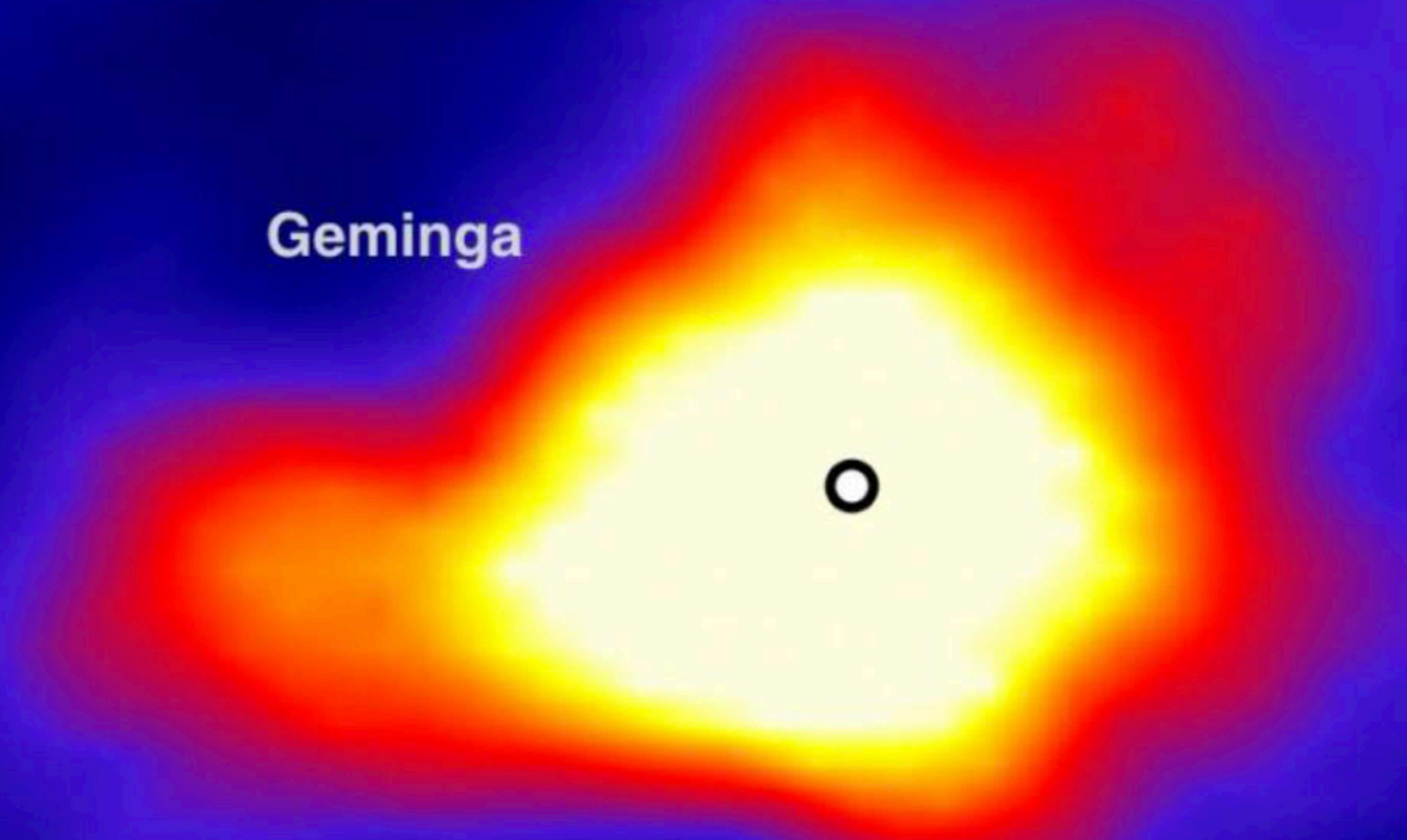
Moon (To Scale)

Linden et al. (2017; 1703.09704)

Sudoh, TL, Beacom (2019; 1902.08203)



Geminga

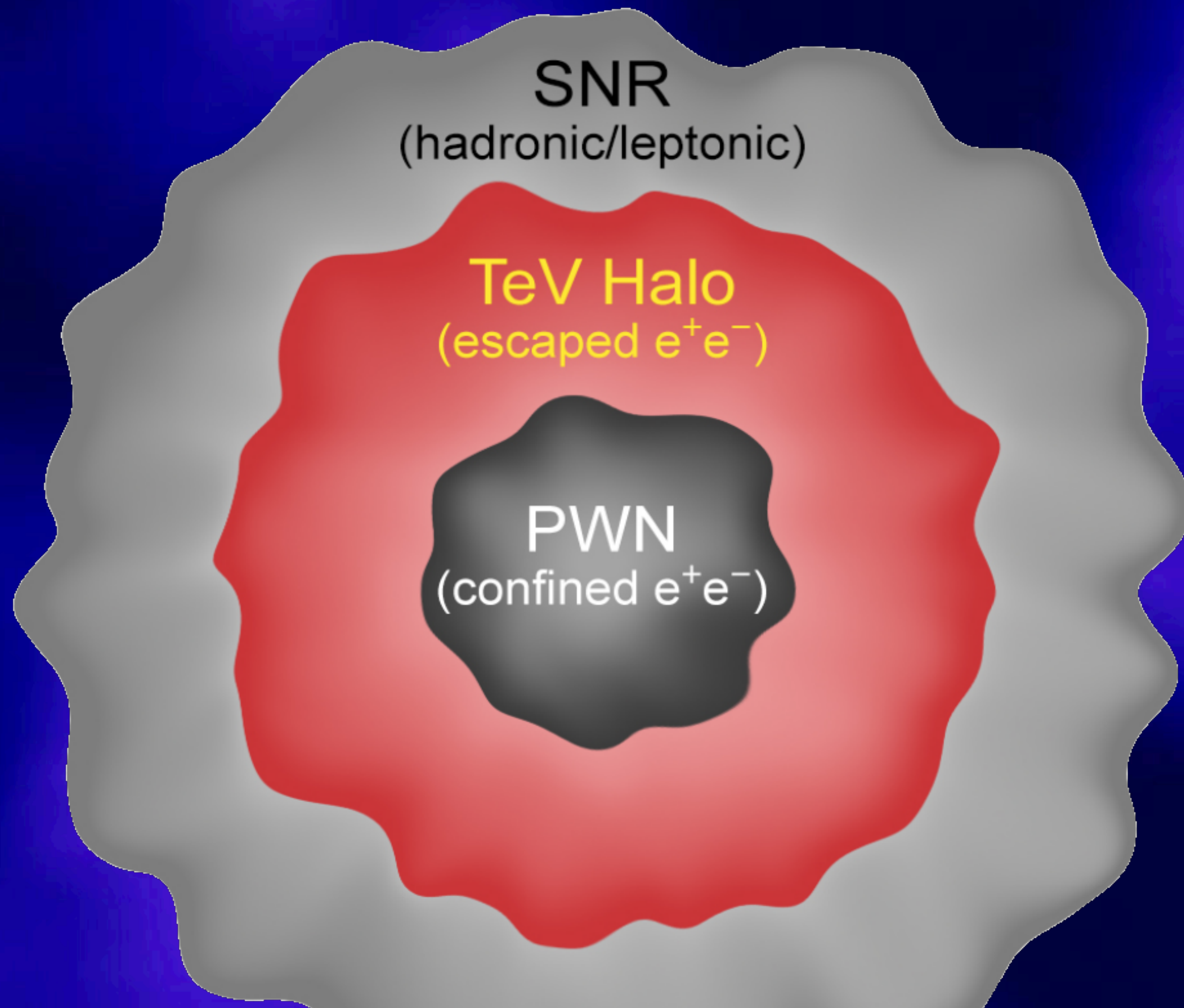




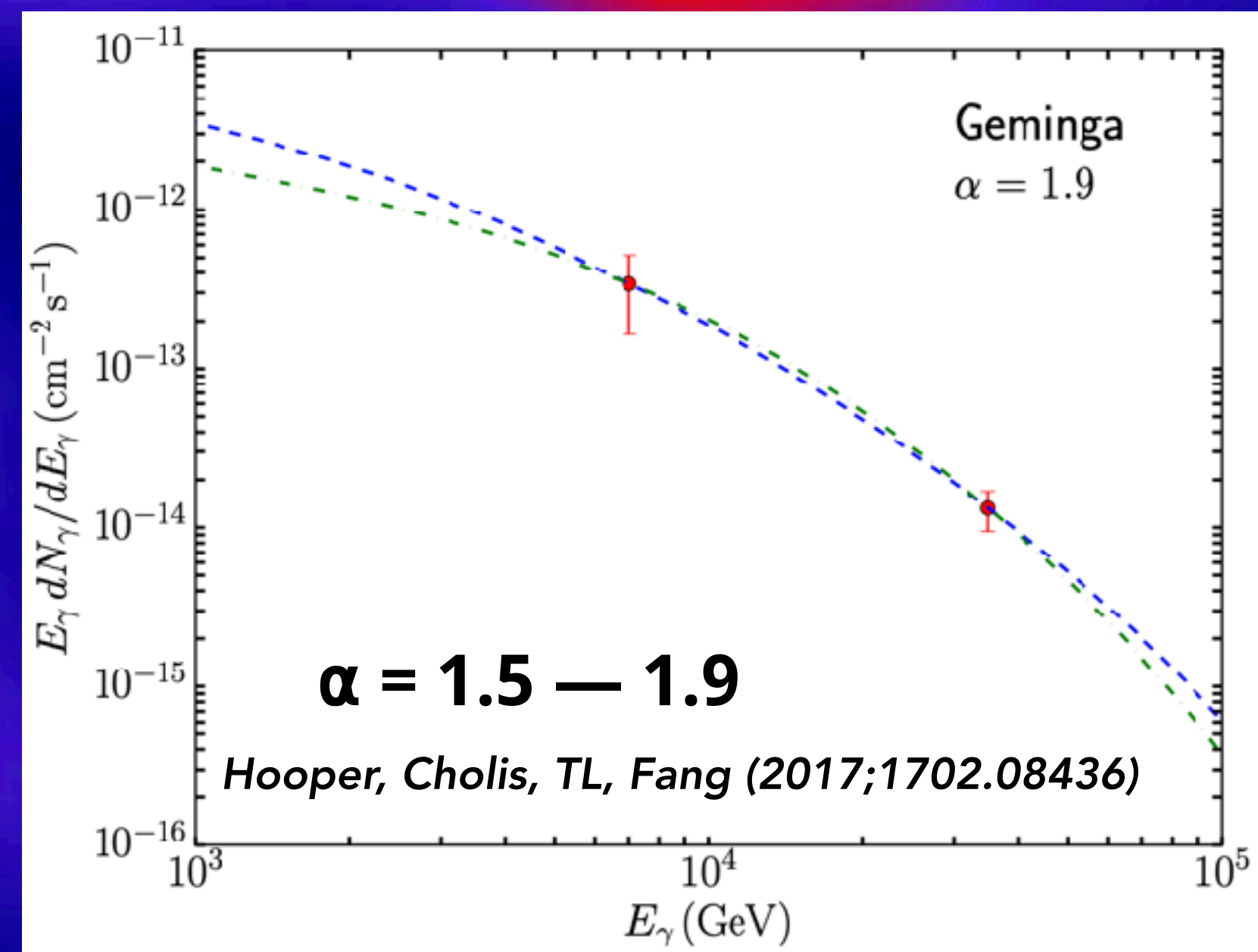
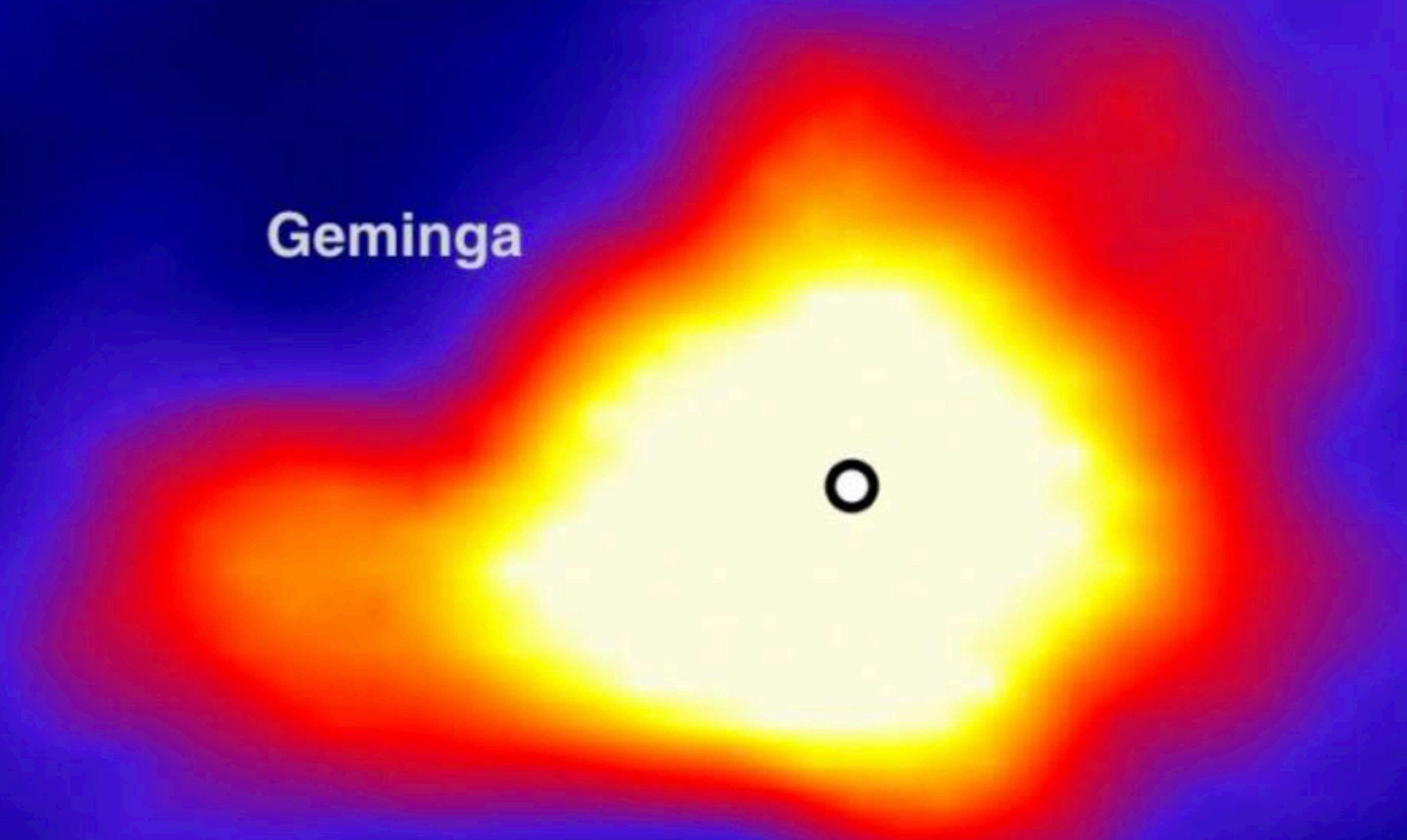
Moon (To Scale)

Linden et al. (2017; 1703.09704)

Sudoh, TL, Beacom (2019; 1902.08203)

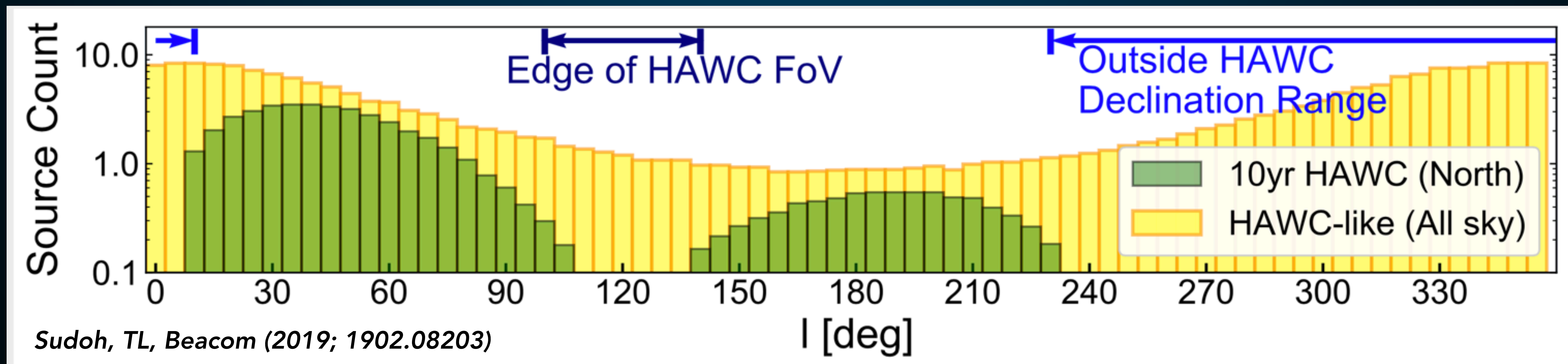
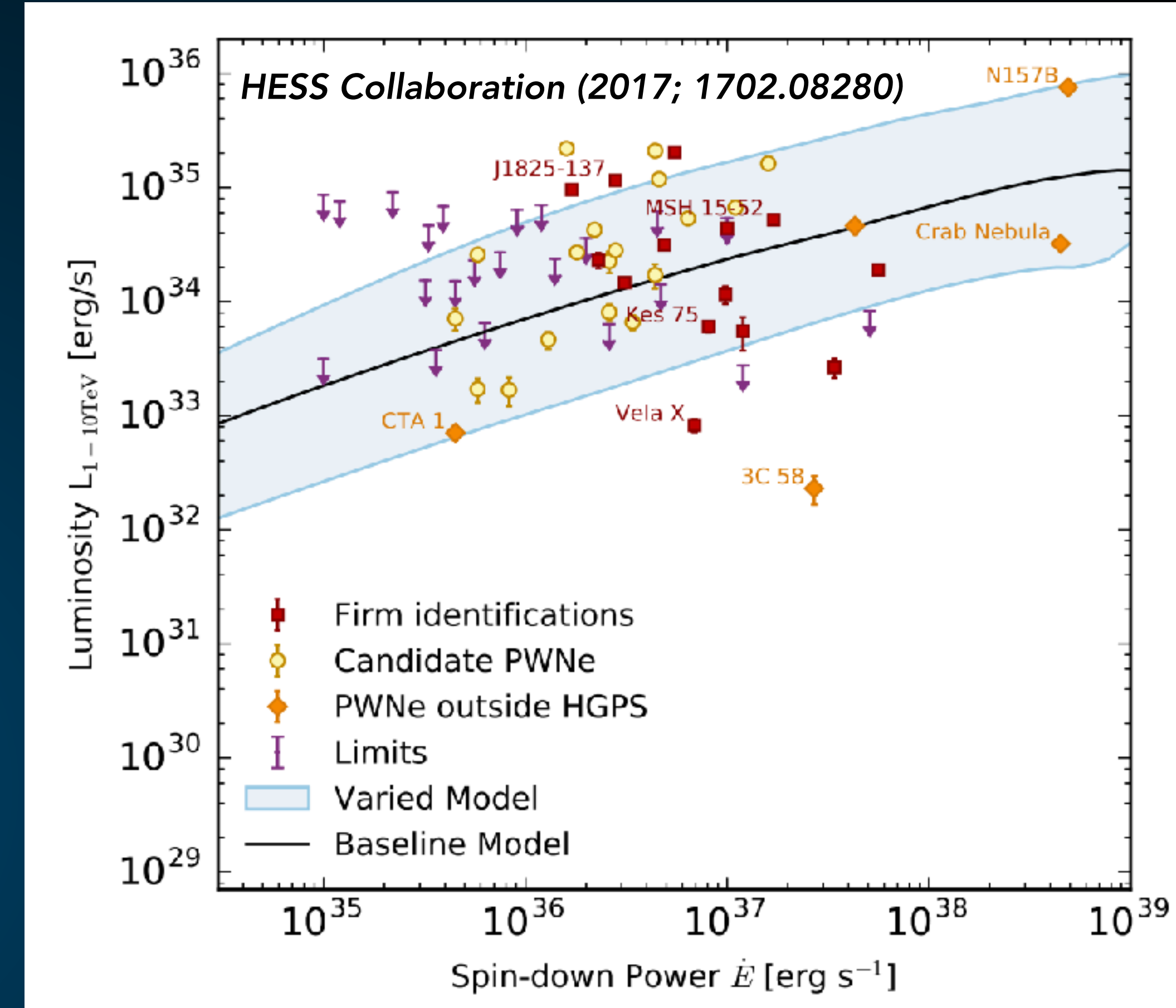


Geminga



# The Positron Excess

- ▶ **Geminga and Monogem are not unique!**
- ▶ **Nearly every bright pulsar observation is consistent with  $\sim 10\%$   $e^+e^-$  efficiency!**
- ▶ **Upcoming observations will detect  $>100$  systems**



# Pulsar Fits to the Positron Excess

## ► Uncertainties in pulsar models:

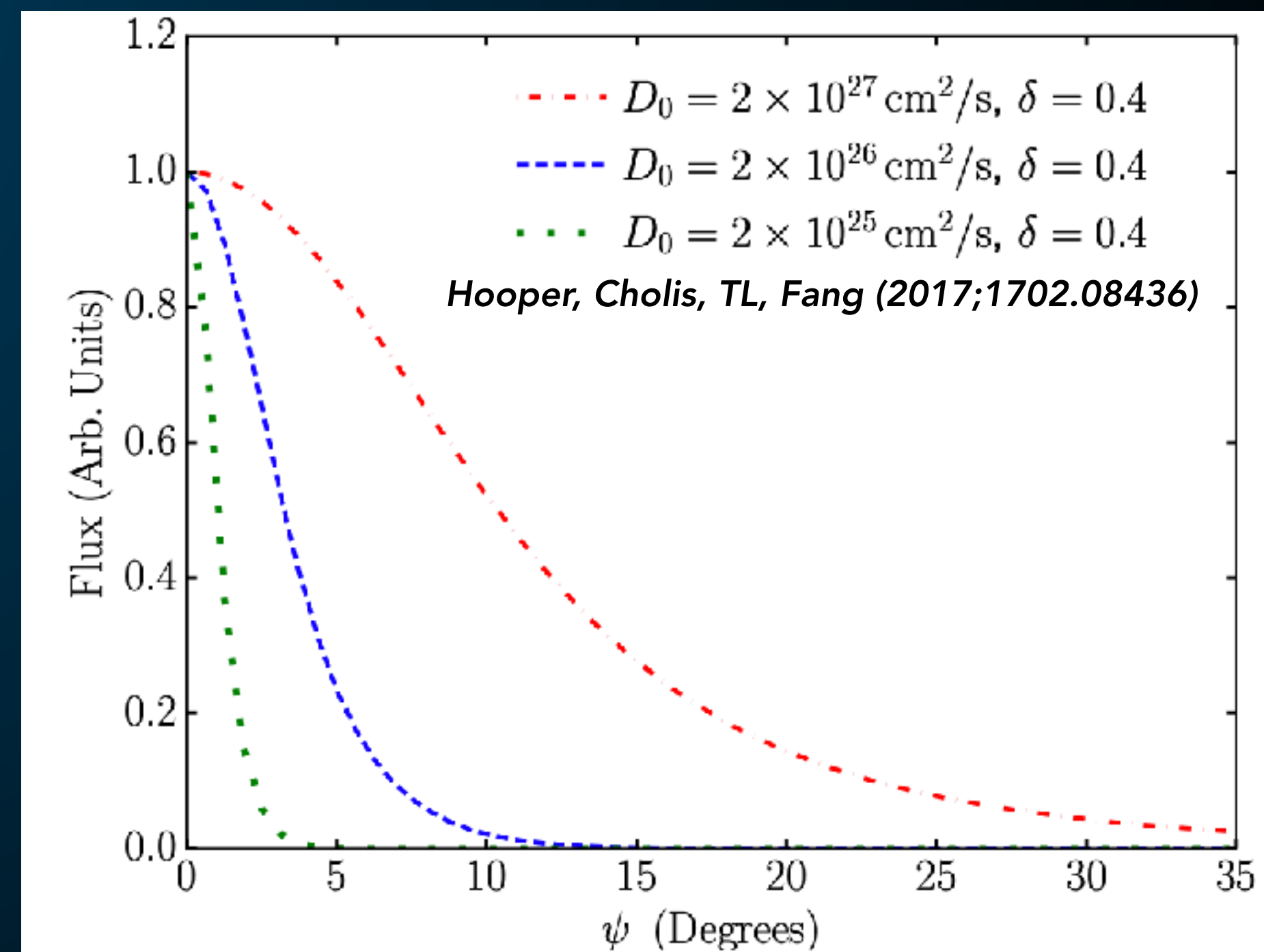
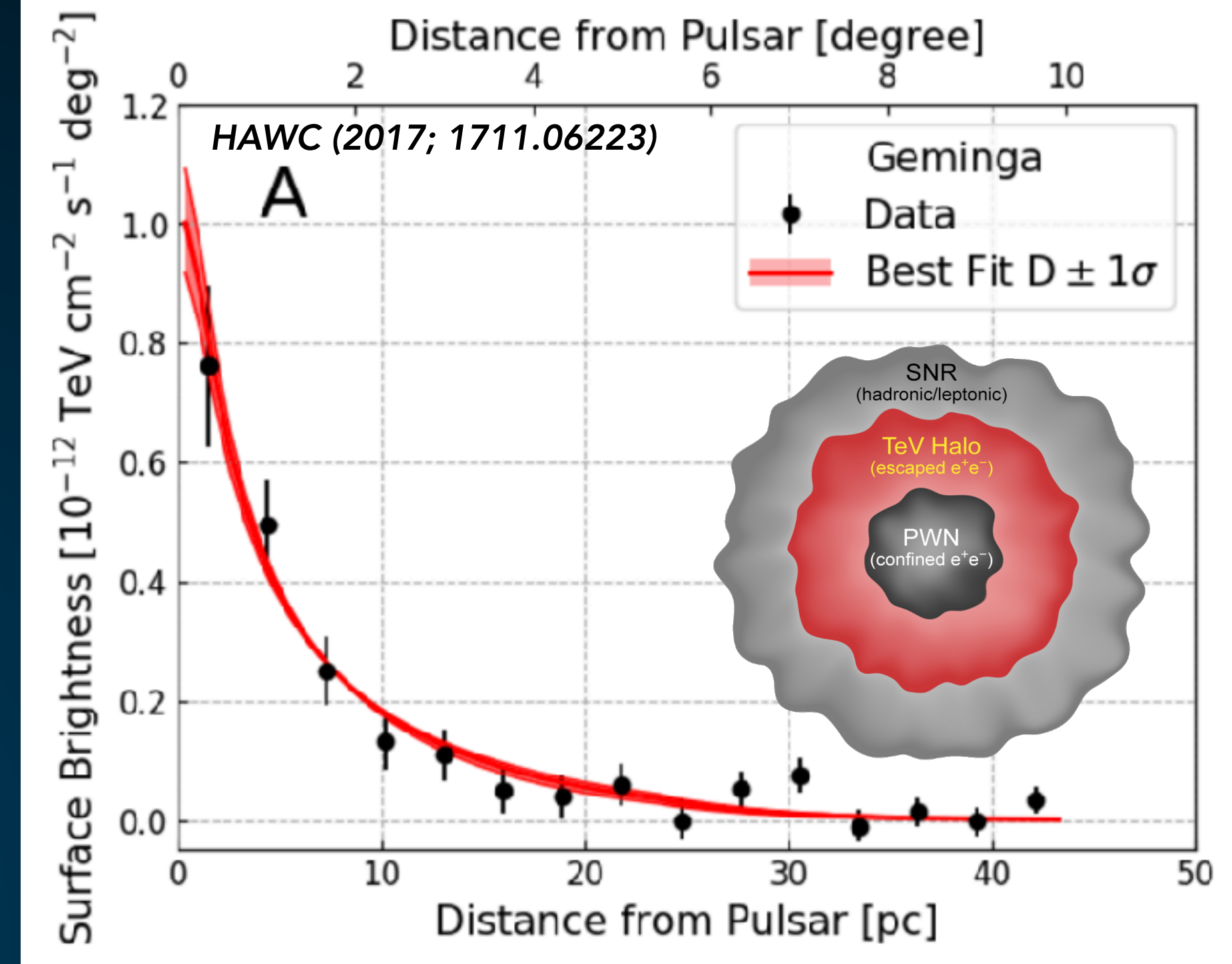
► I: The  $e^+e^-$  production efficiency

► II: The  $e^+e^-$  spectrum.

► III: The propagation of  $e^+e^-$  to Earth.

Malyshev et al. (0903.1310)

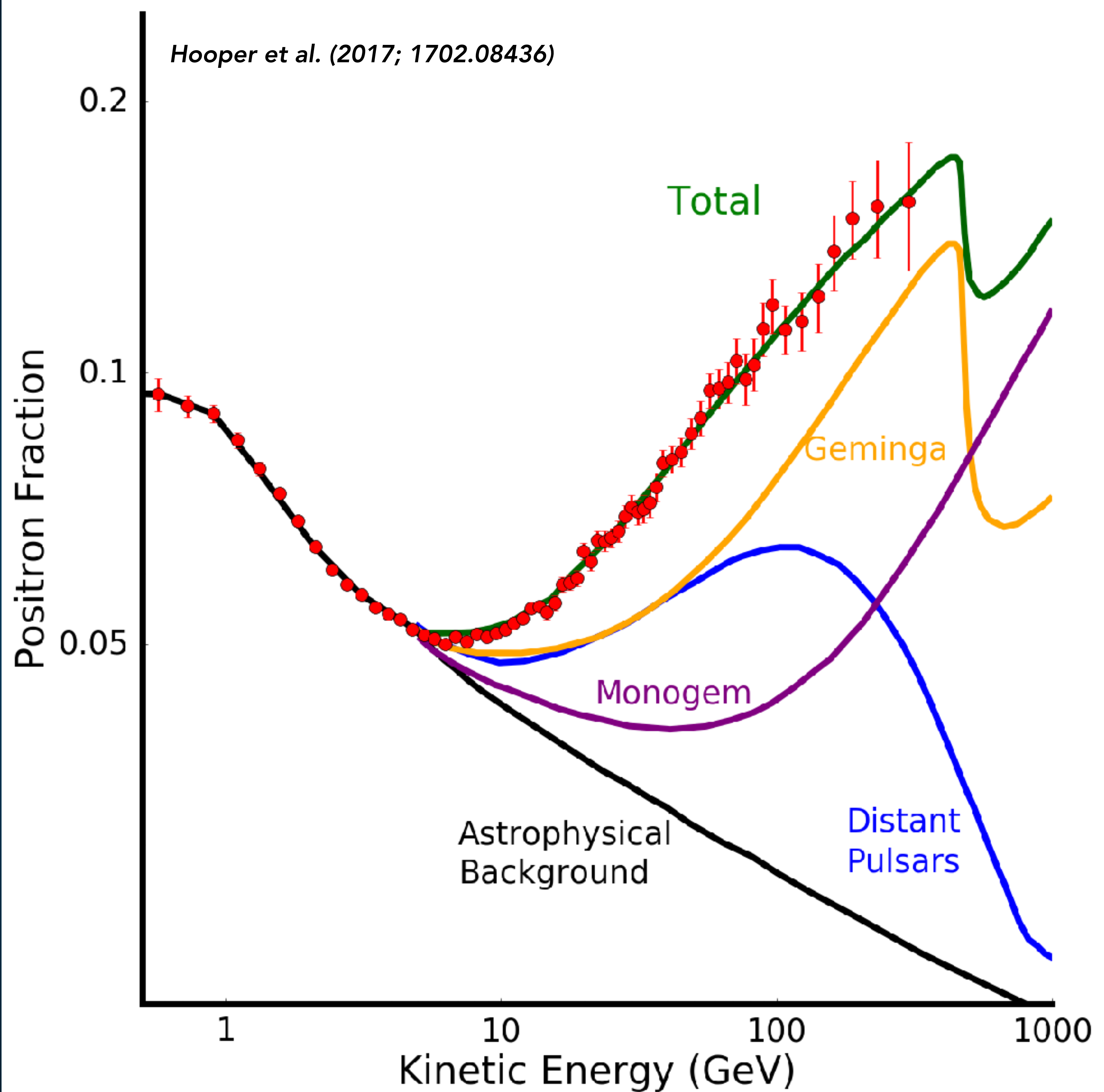
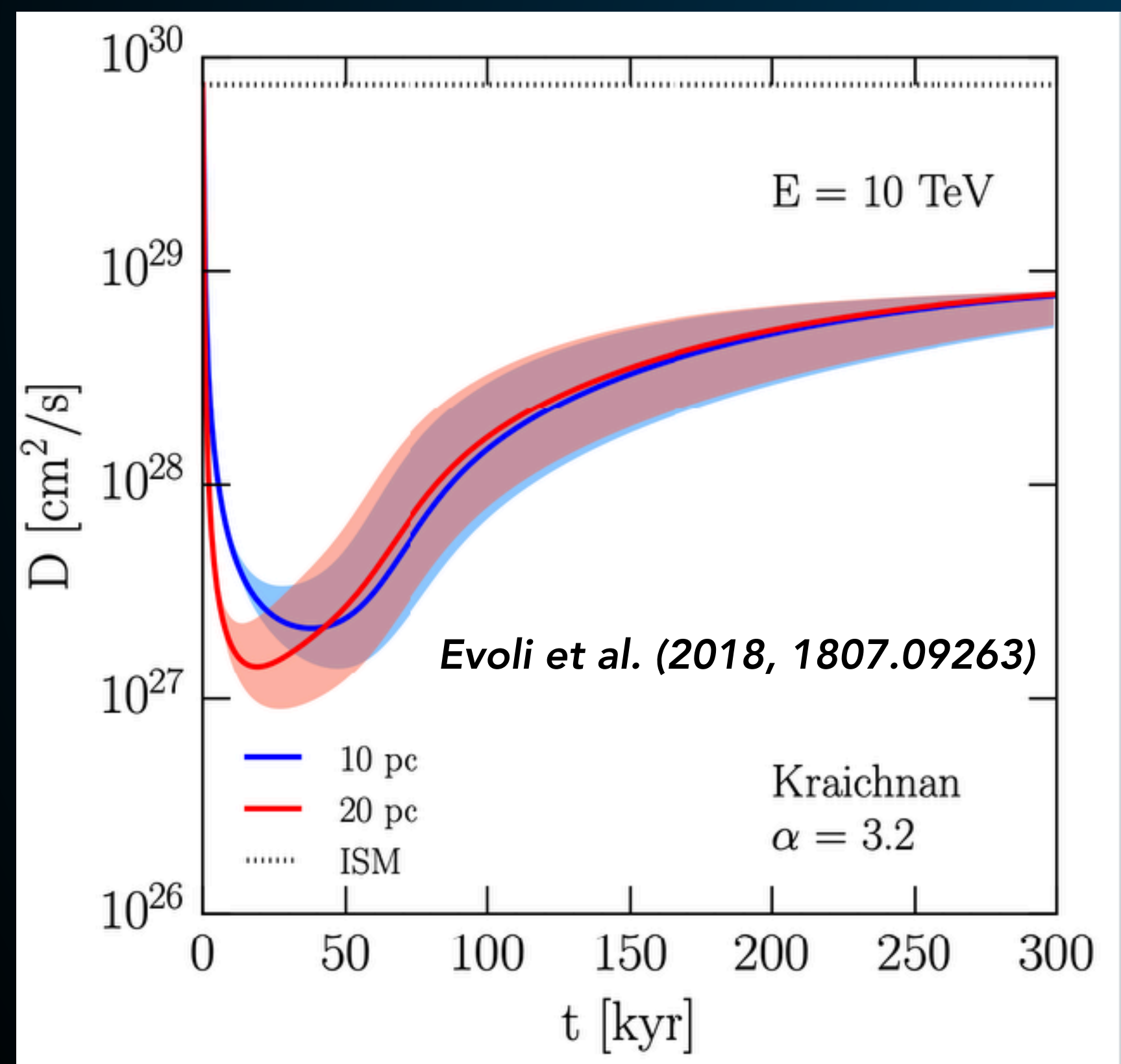
The observed spectrum on Earth of electrons and positrons injected by pulsars is also strongly dependent on propagation effects. In particular, the observed cutoff in the flux of electrons from a pulsar can be much smaller than the injection cutoff due to energy losses (“cooling”) during propagation. We define the cooling break,  $E_{br}(t)$ , as the maximal energy electrons can have after propagating for time  $t$ . Since – as stated above – the typical



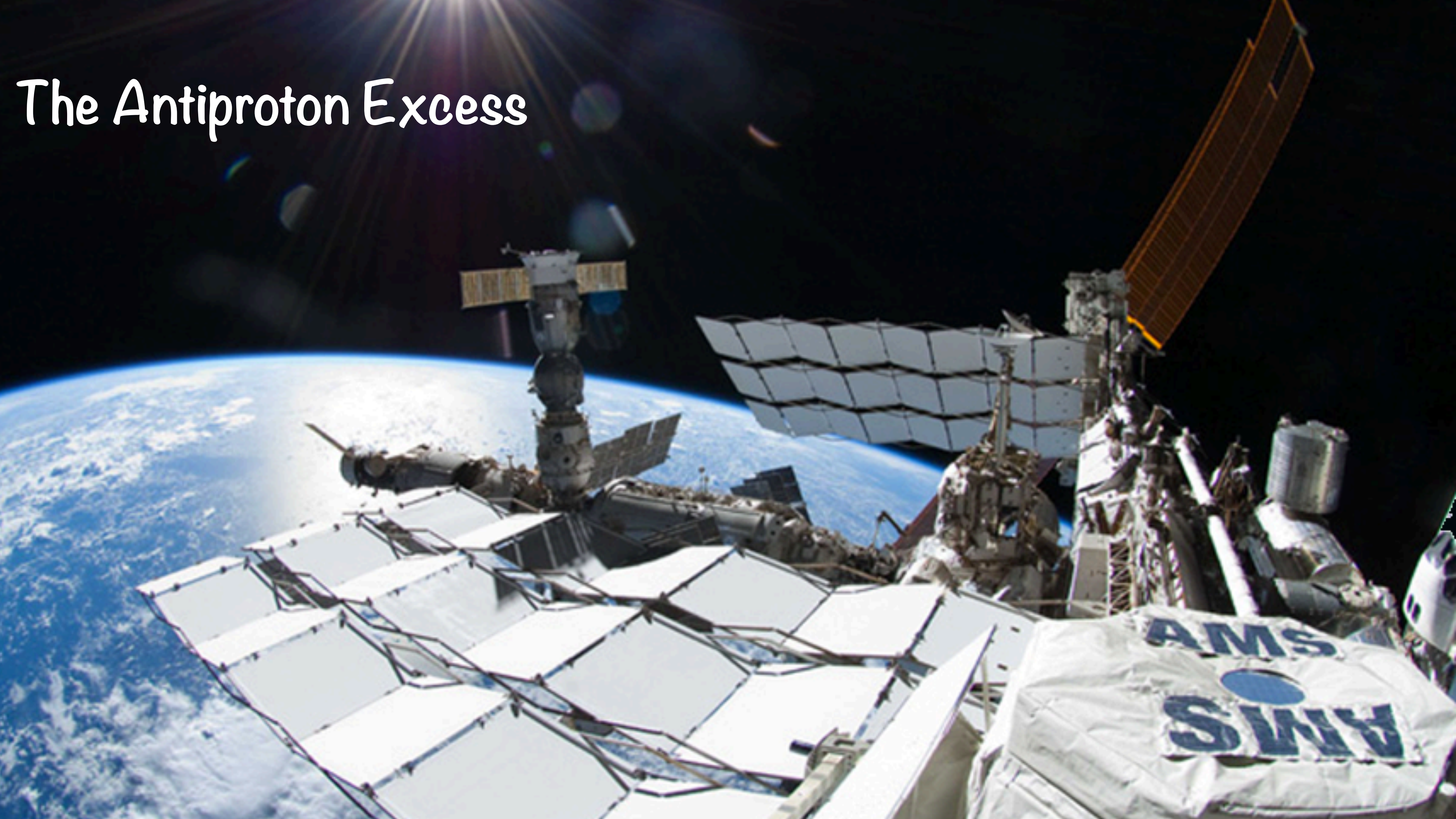
# The Positron Excess

This can easily match the positron fraction

Some transport issues are possible - but easy to solve in models with inhomogeneous diffusion.



# The Antiproton Excess



# The Antiproton Excess

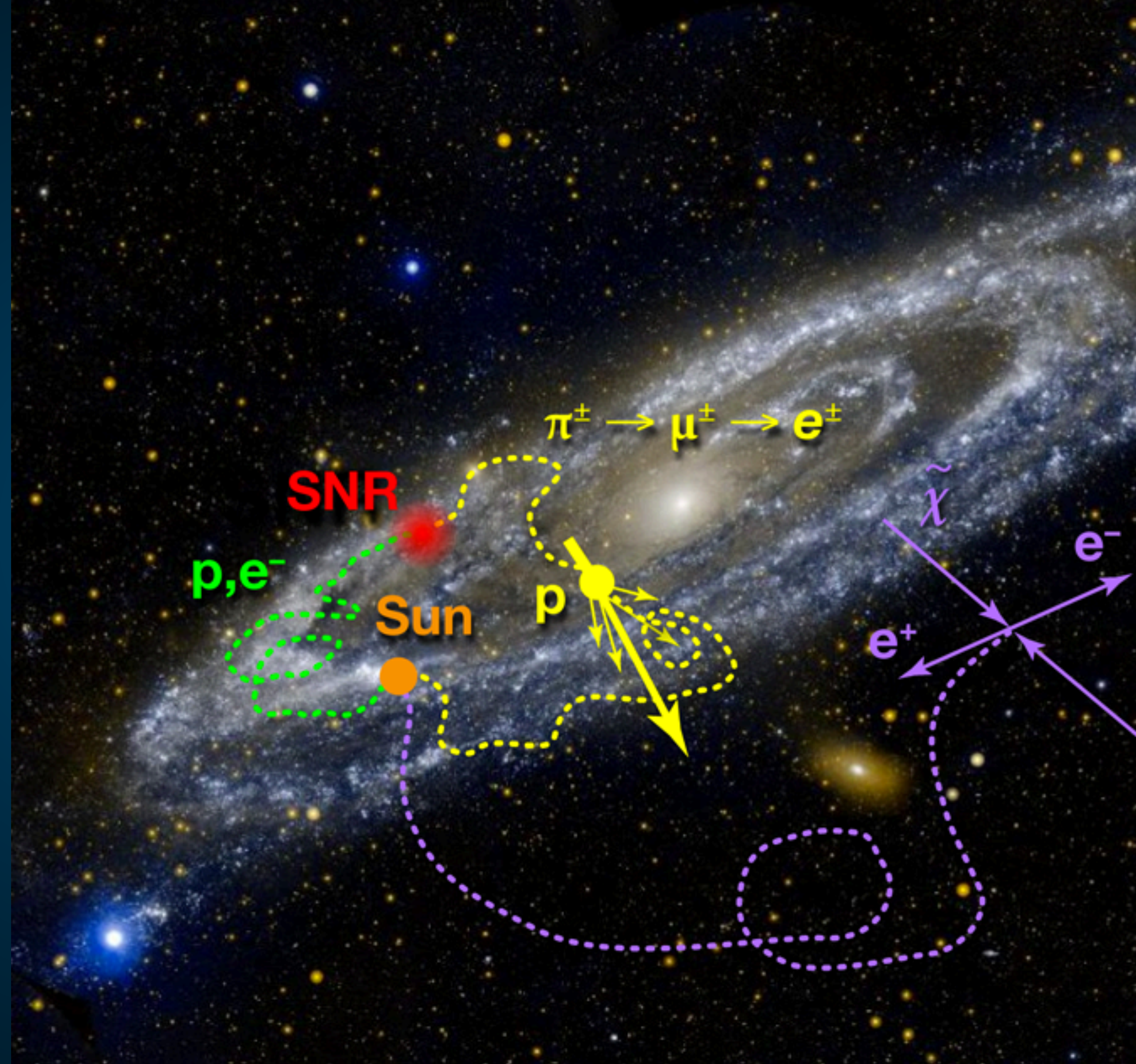
Investigate the Antiproton Fraction!

$$\frac{\phi_{\bar{p}}}{\phi_p}$$

Two Changes:

Ratio is much smaller (don't need to add antiprotons into denominator).

Hadronic Energy losses are slower (sensitive to antiproton production throughout the Galaxy)

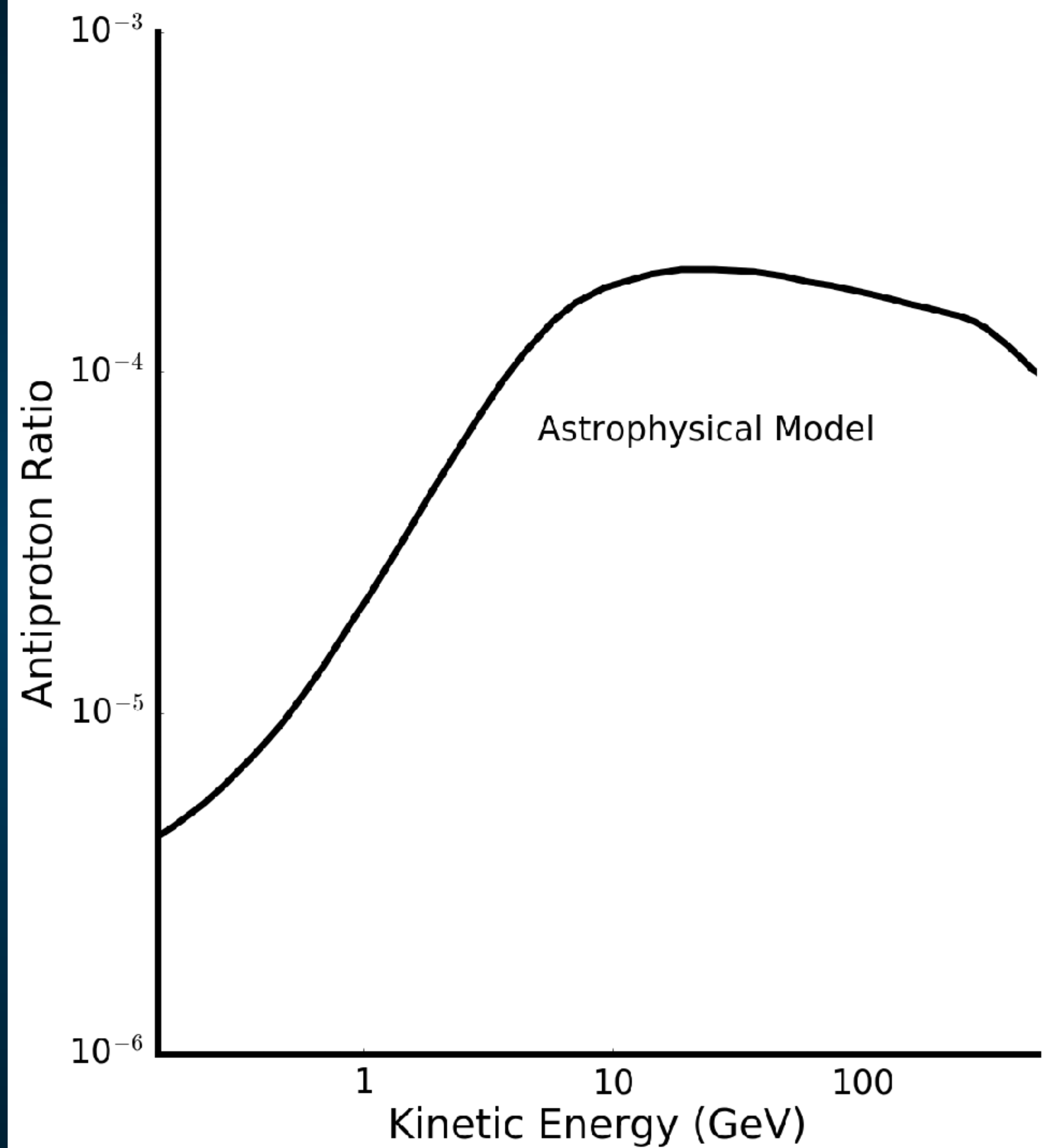
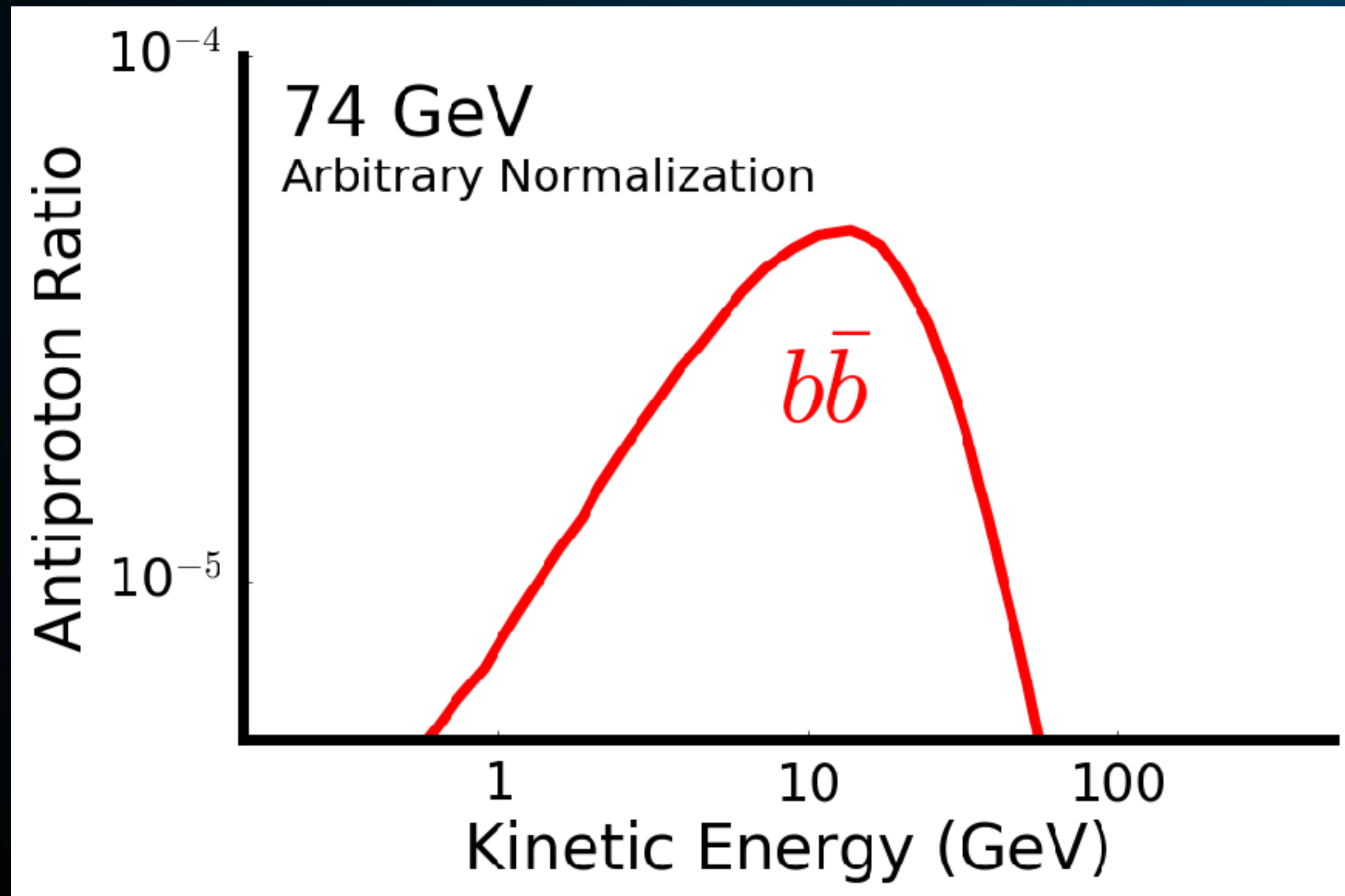




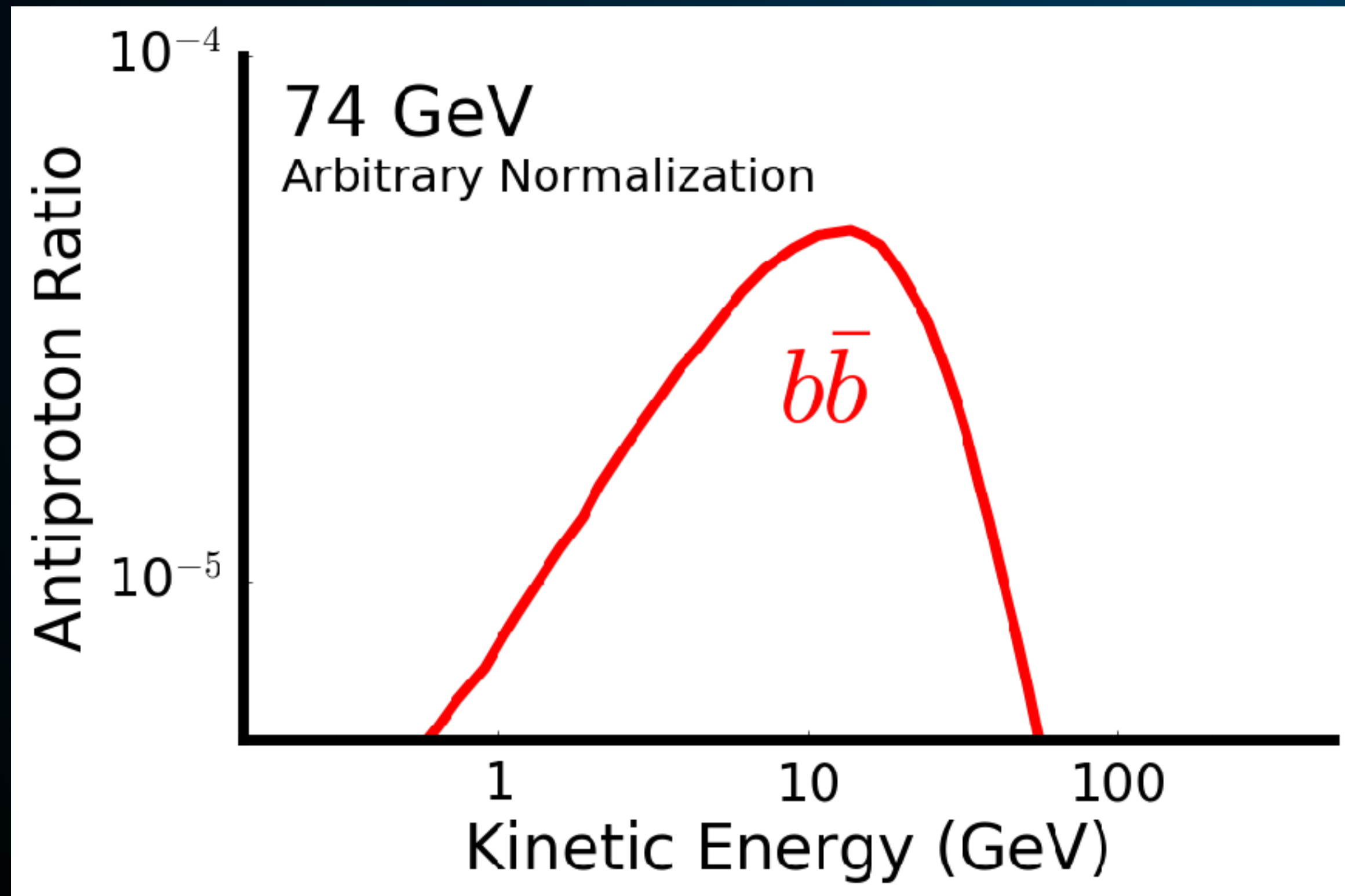
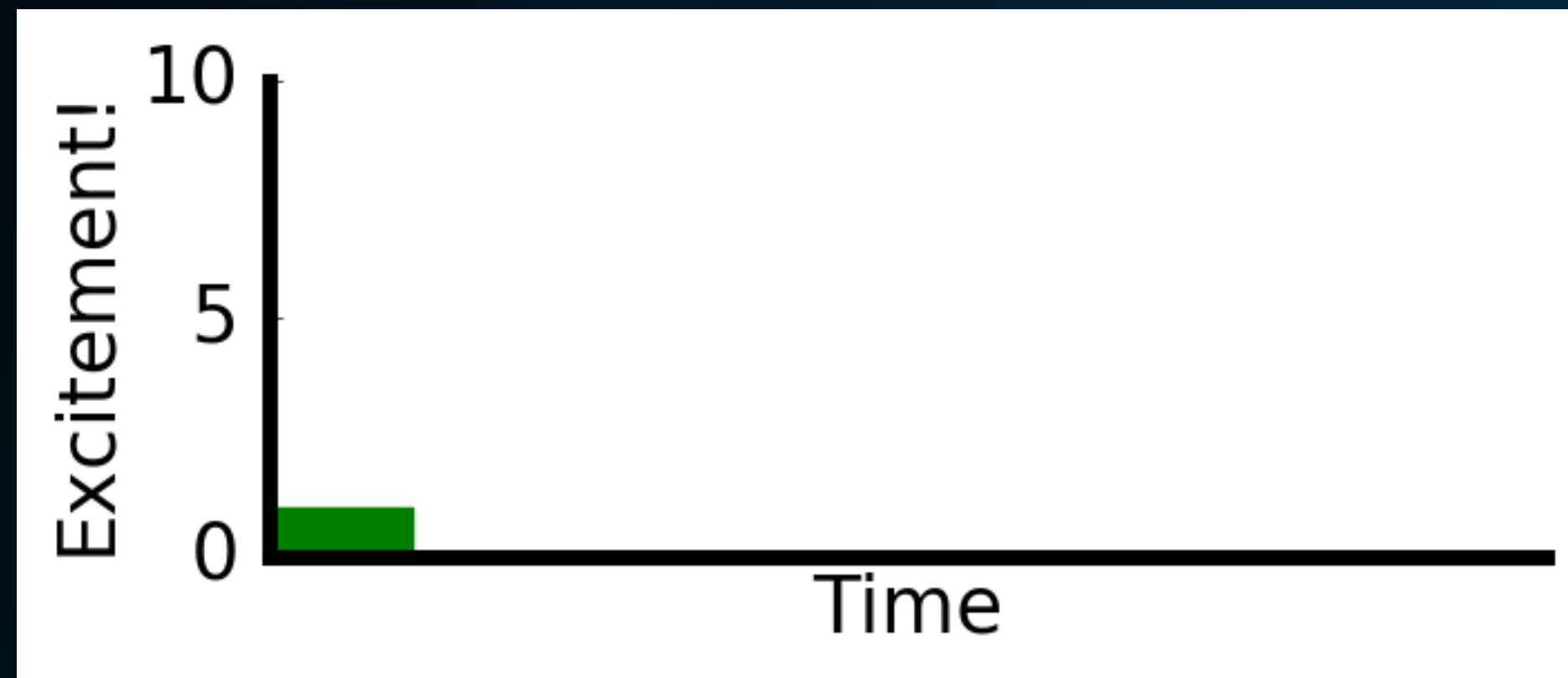
# The Antiproton Excess

Astrophysics - Smooth Profile

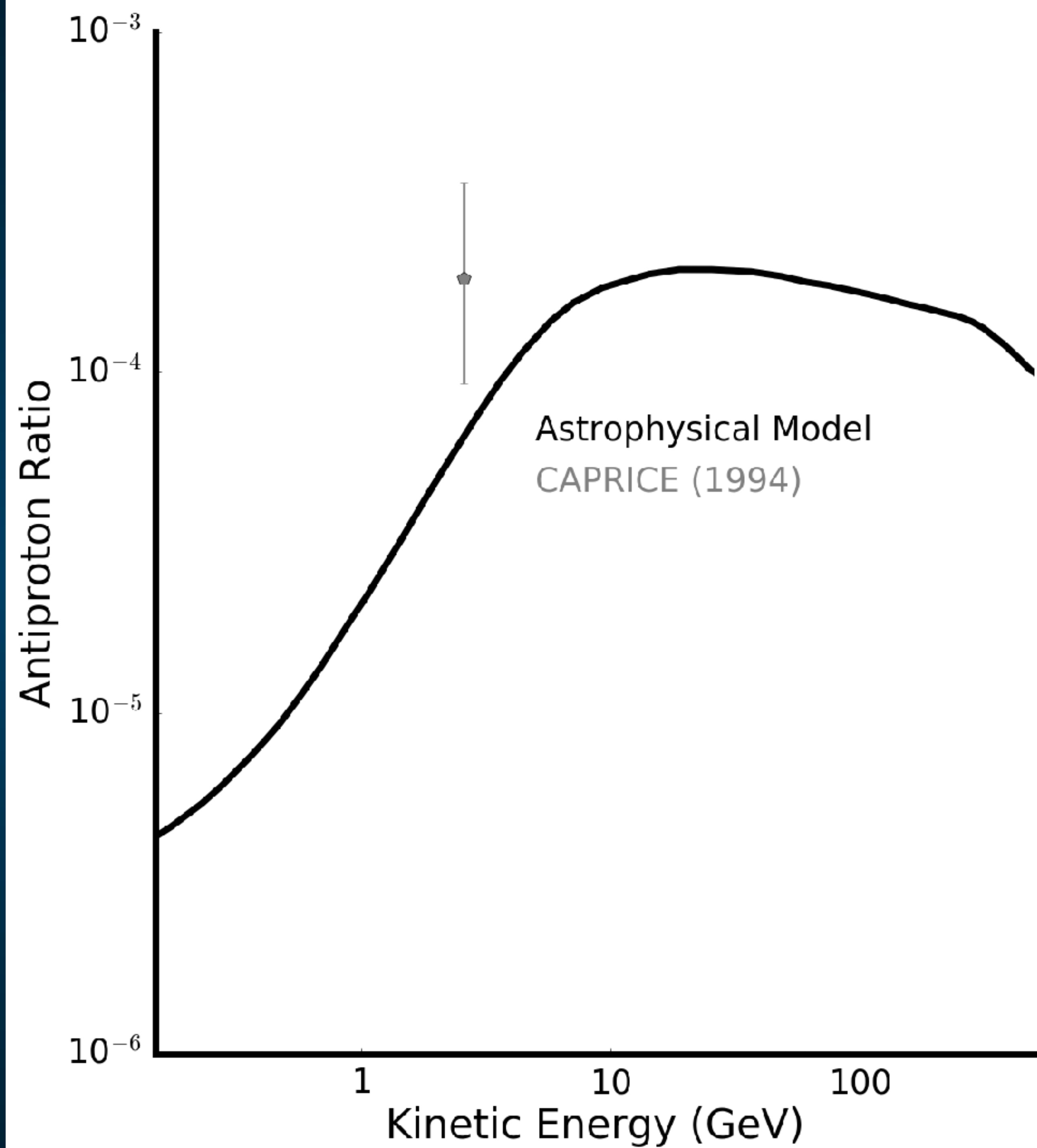
Dark Matter - Sharp Bump!



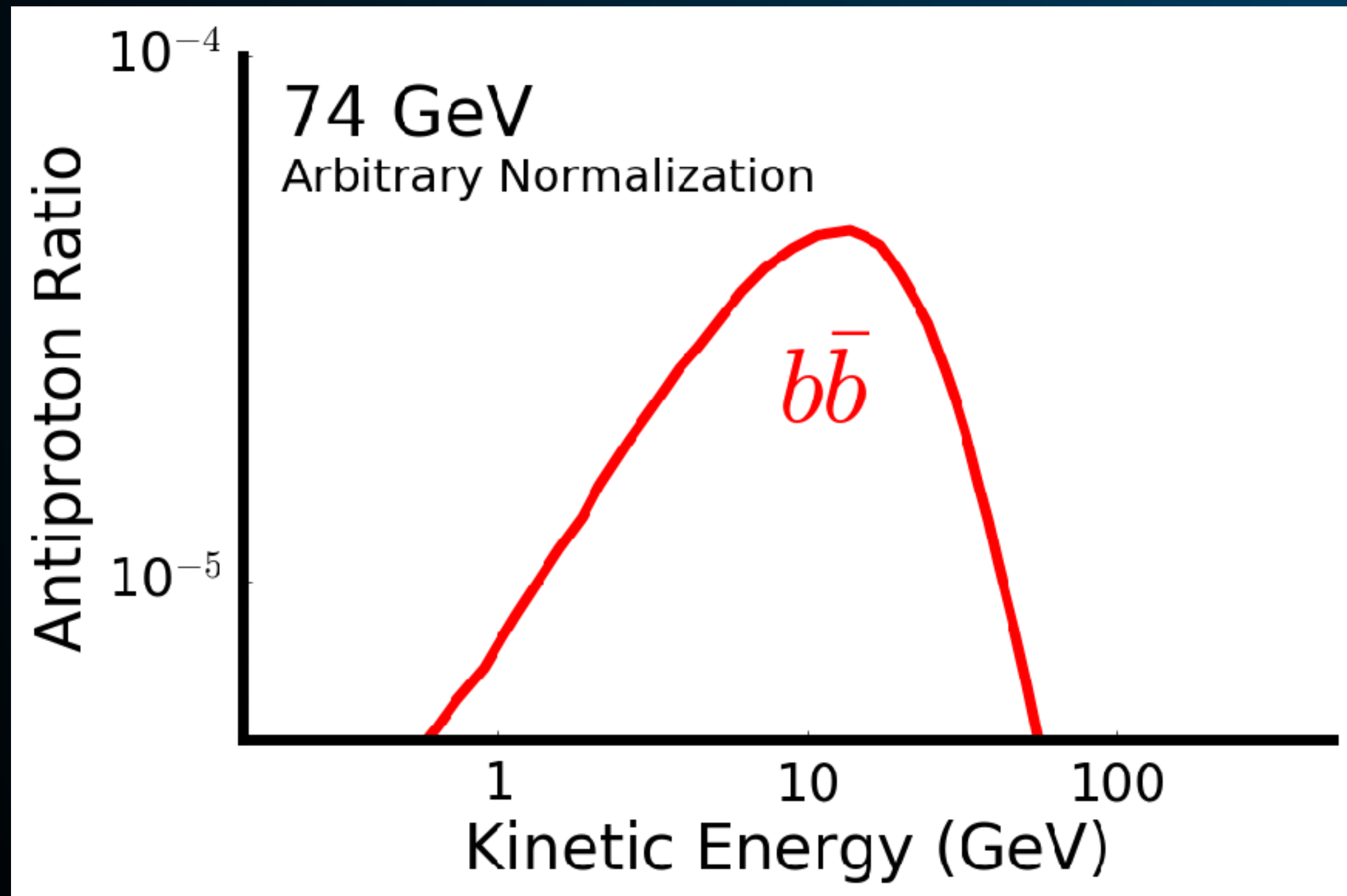
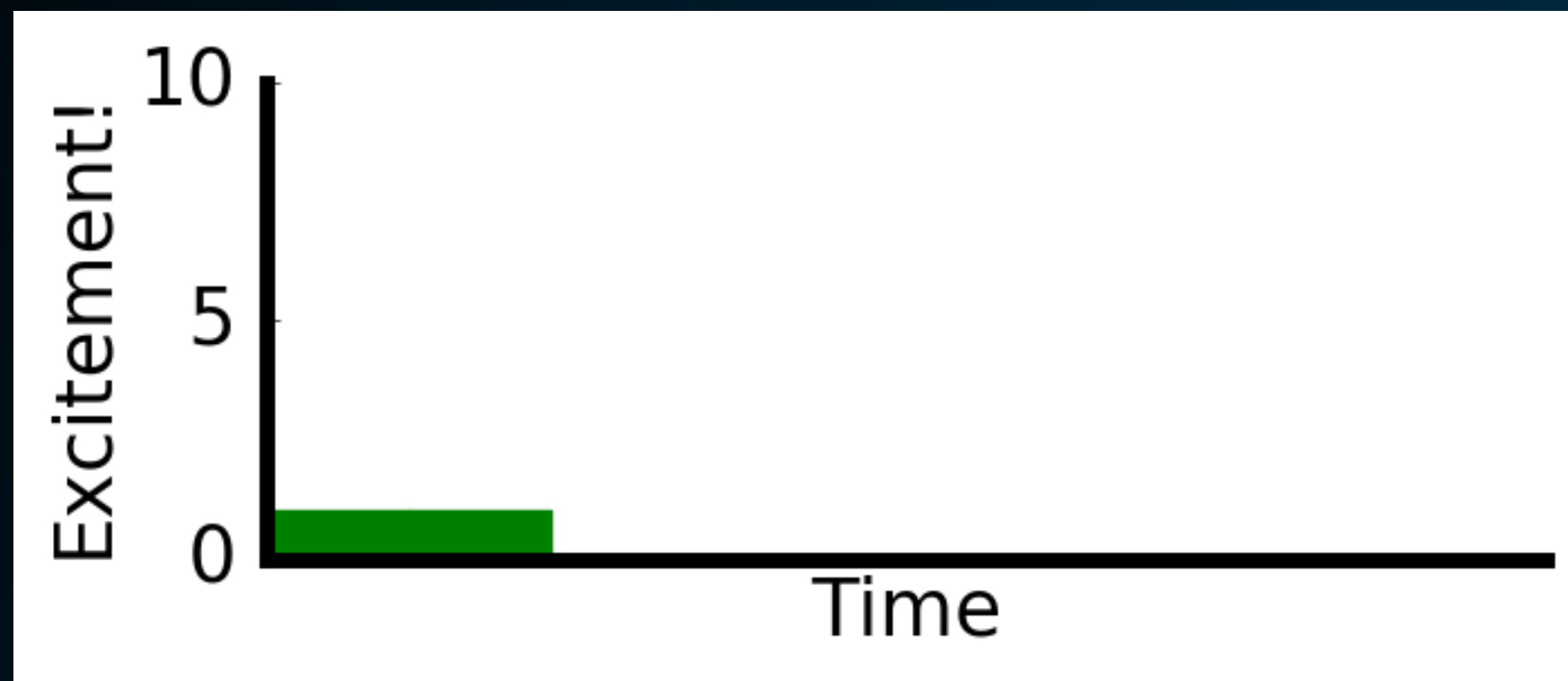
# The Antiproton Excess



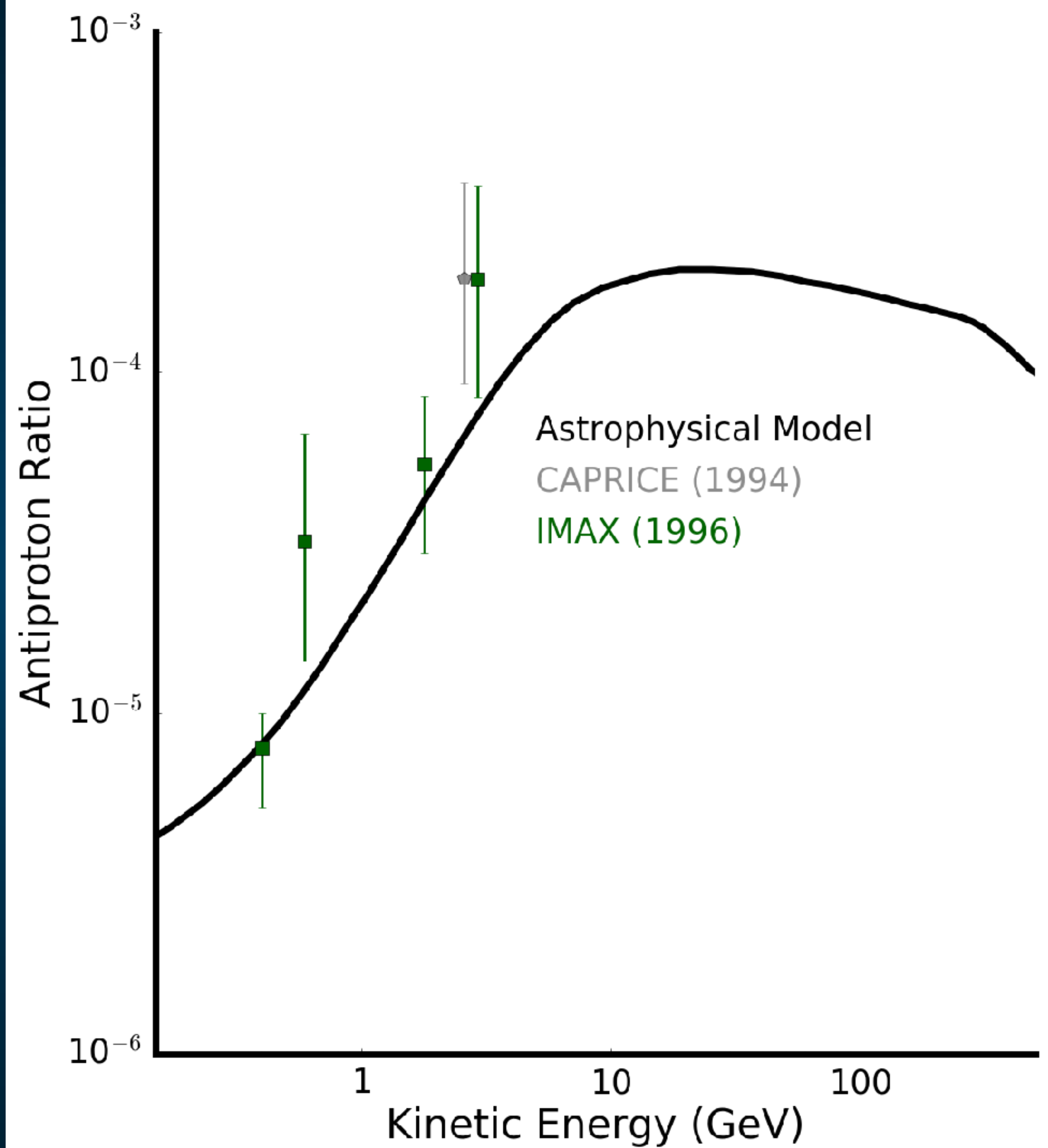
(Not an exhaustive list of observations)



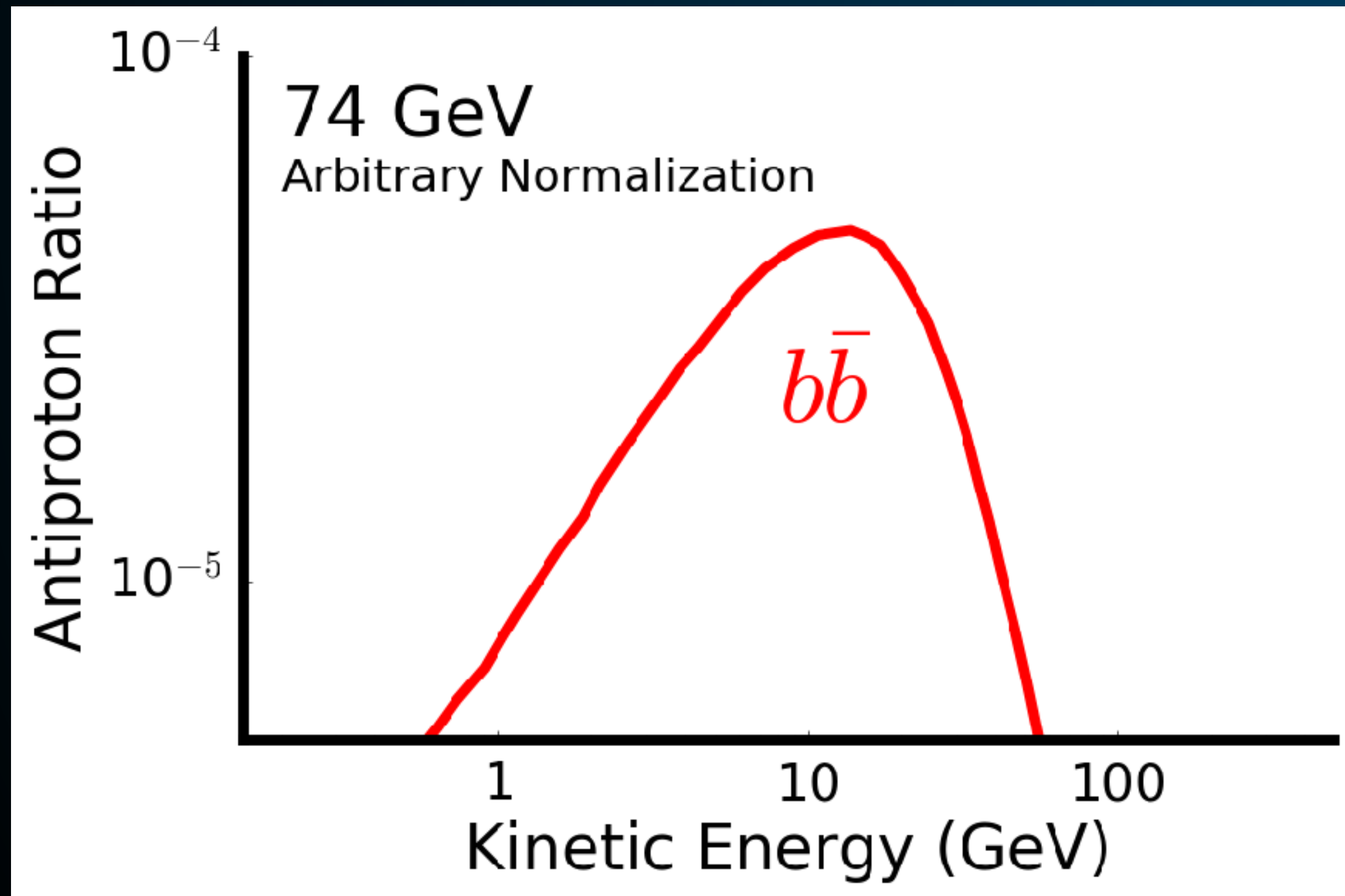
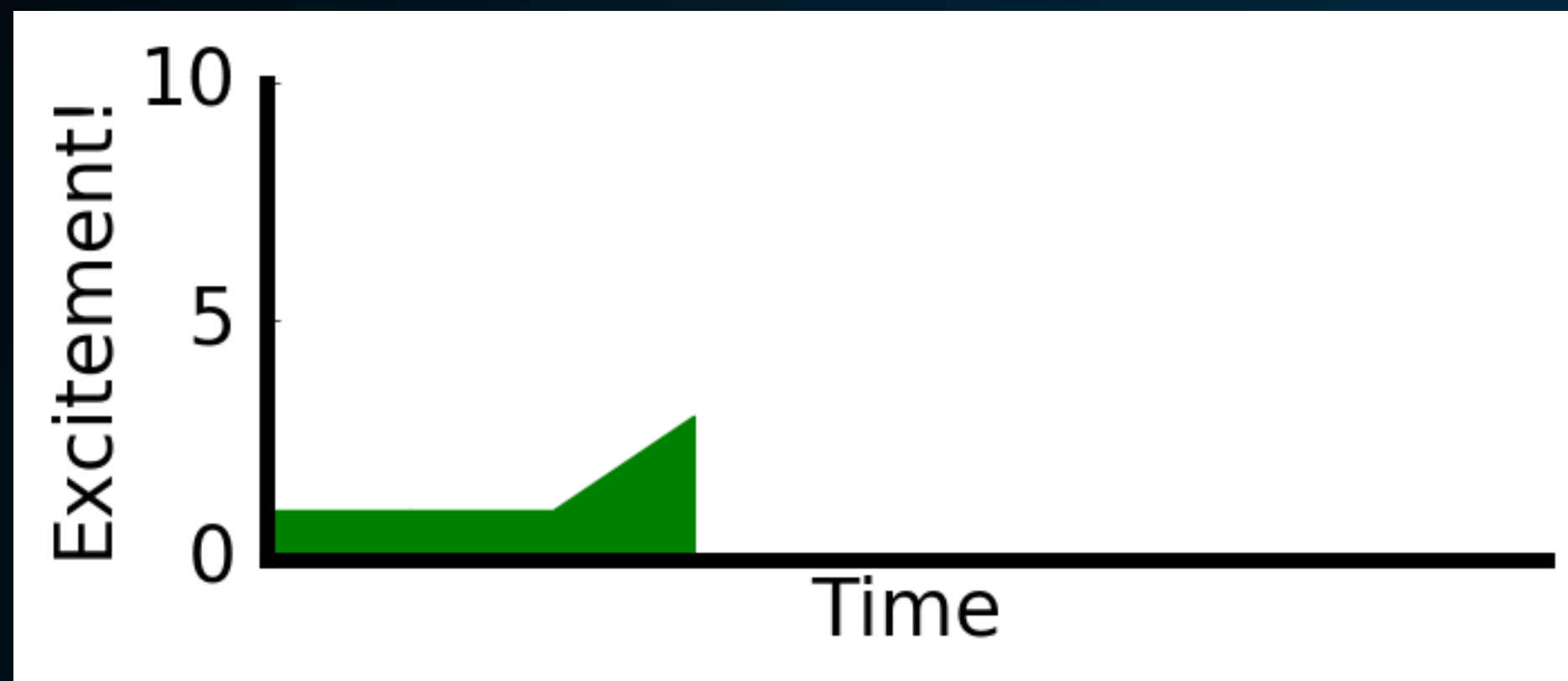
# The Antiproton Excess



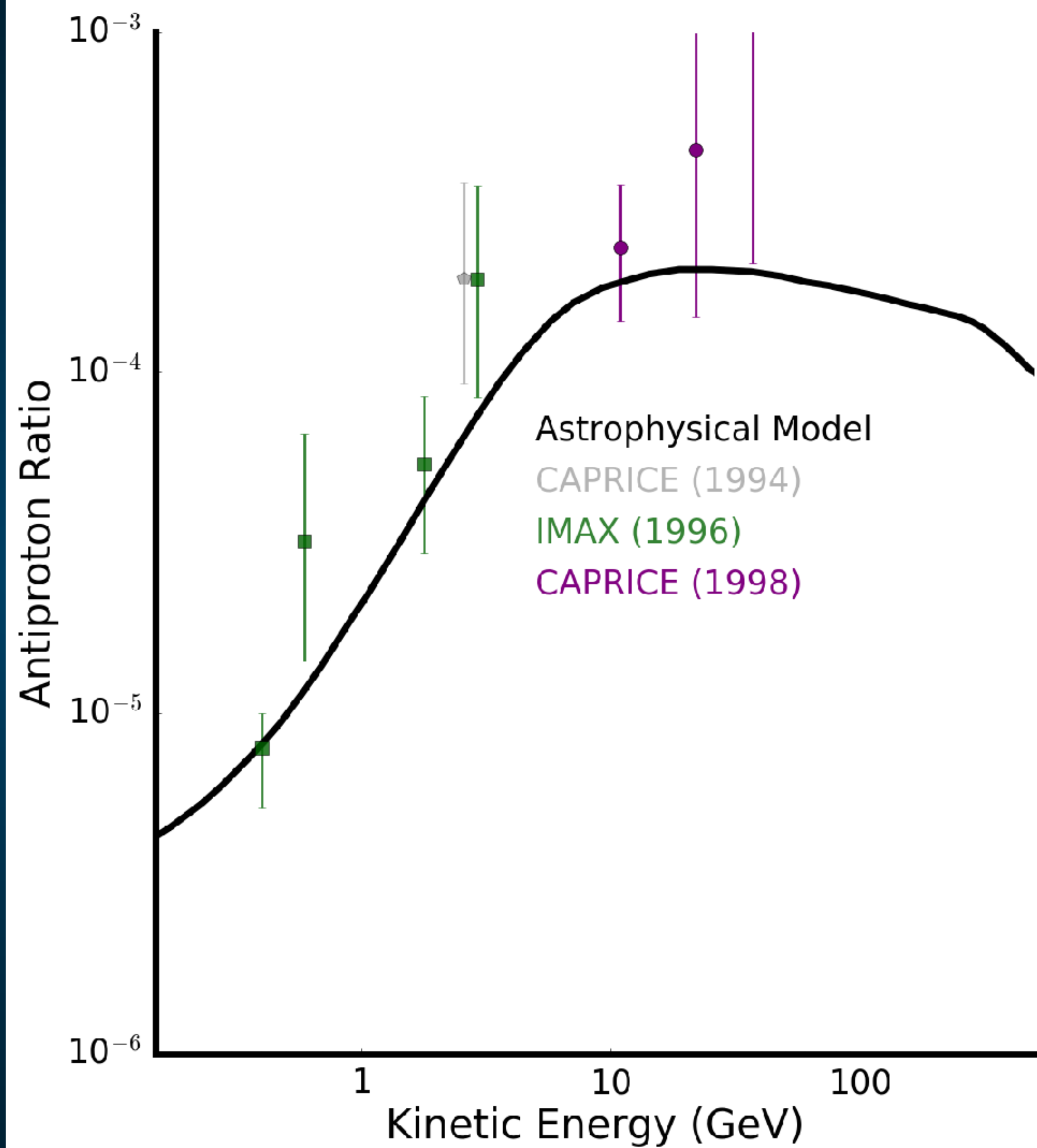
(Not an exhaustive list of observations)



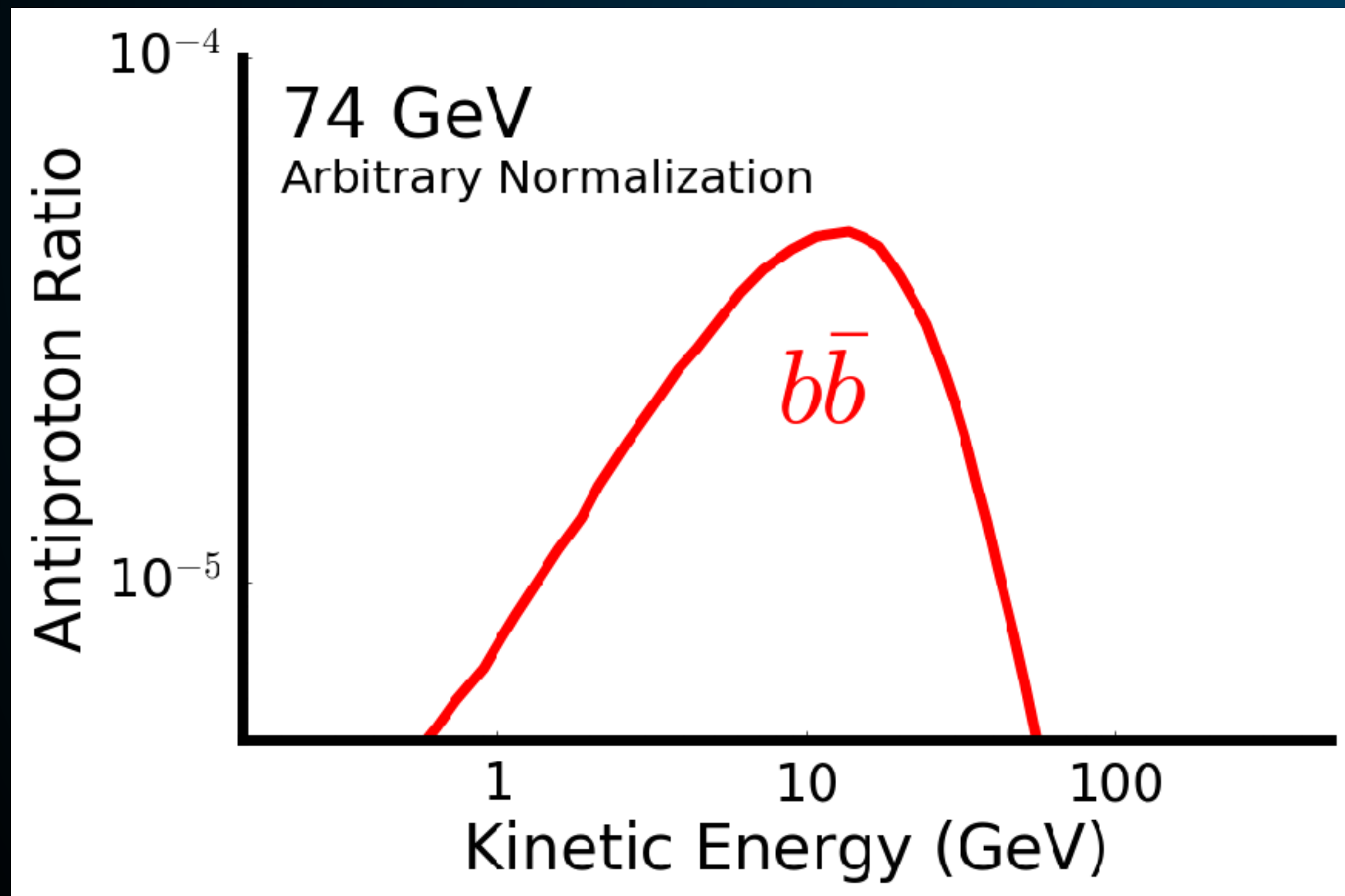
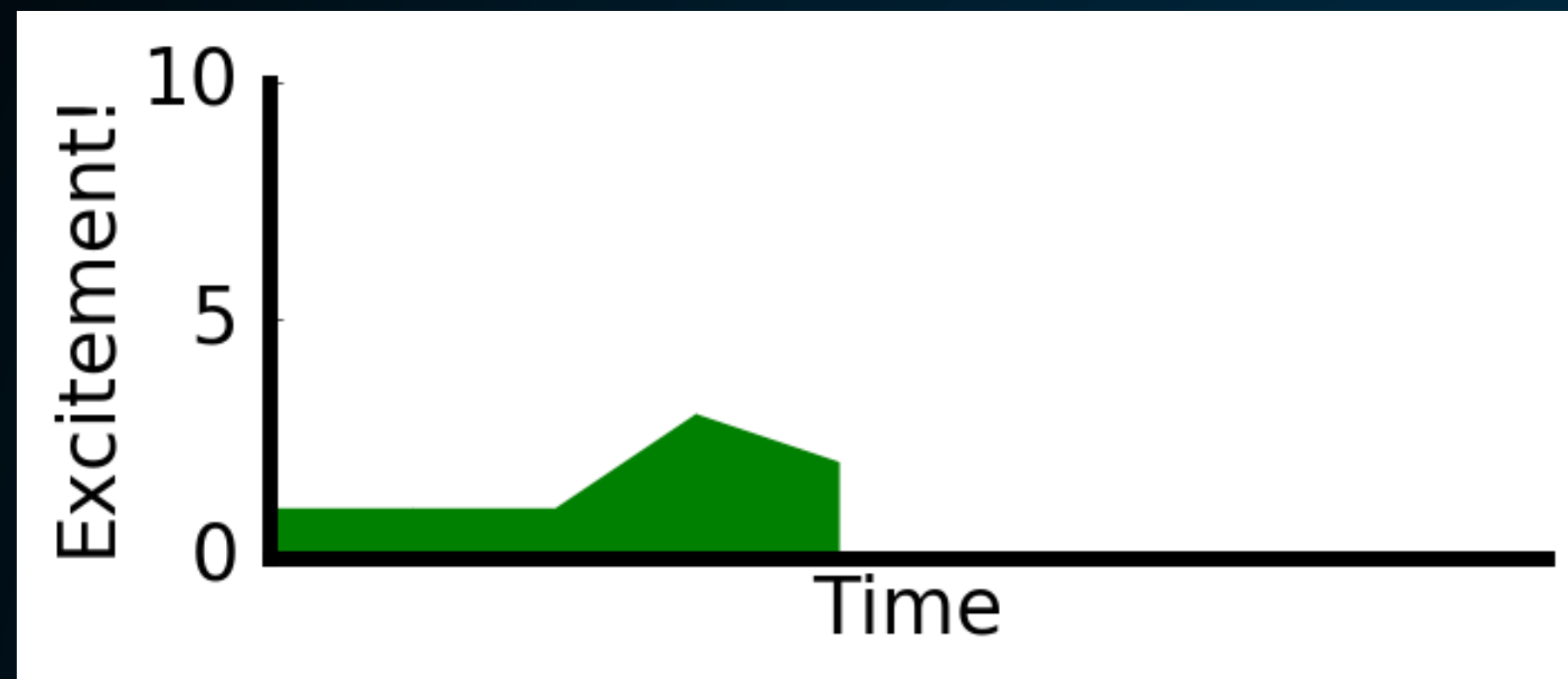
# The Antiproton Excess



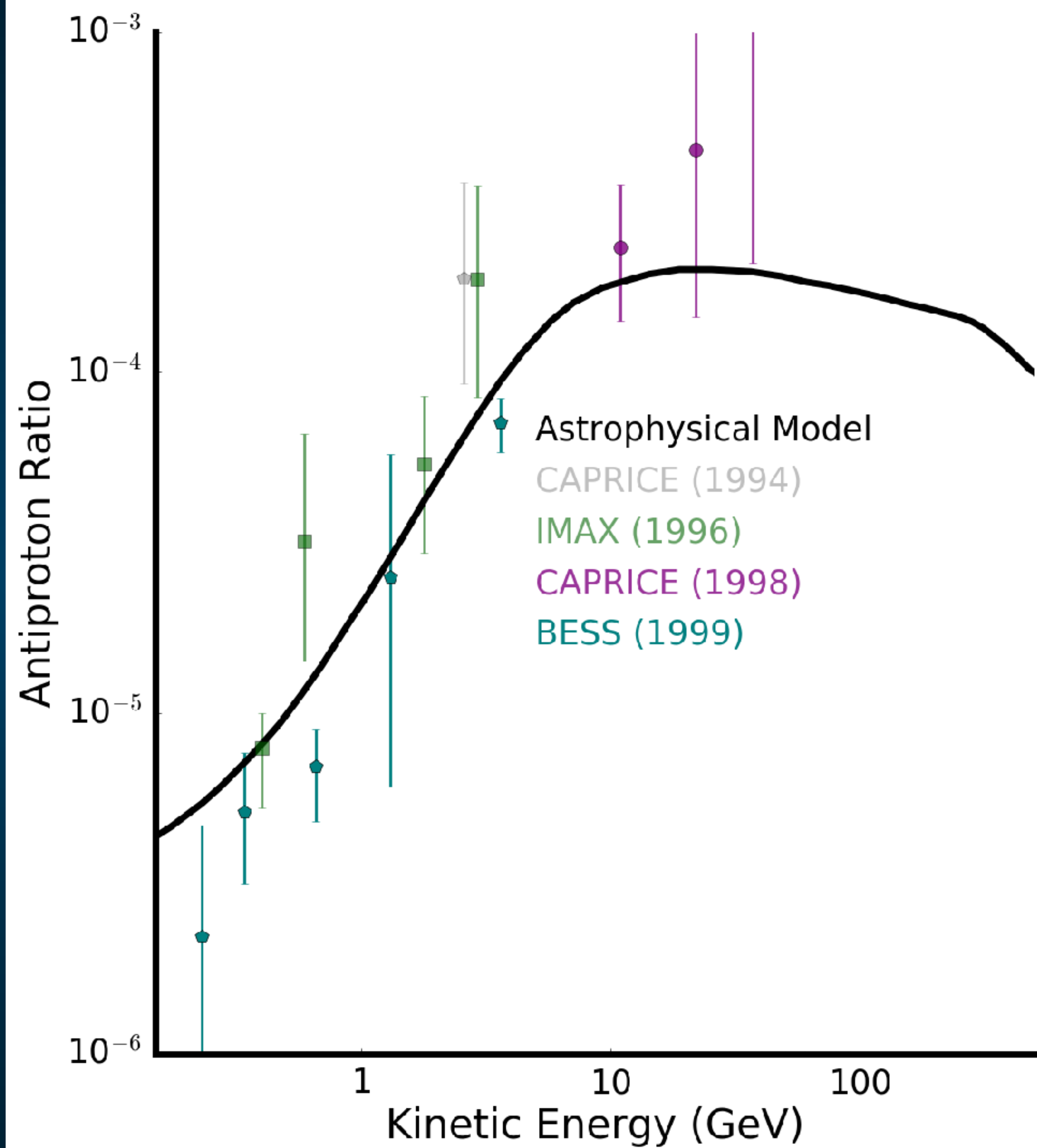
(Not an exhaustive list of observations)



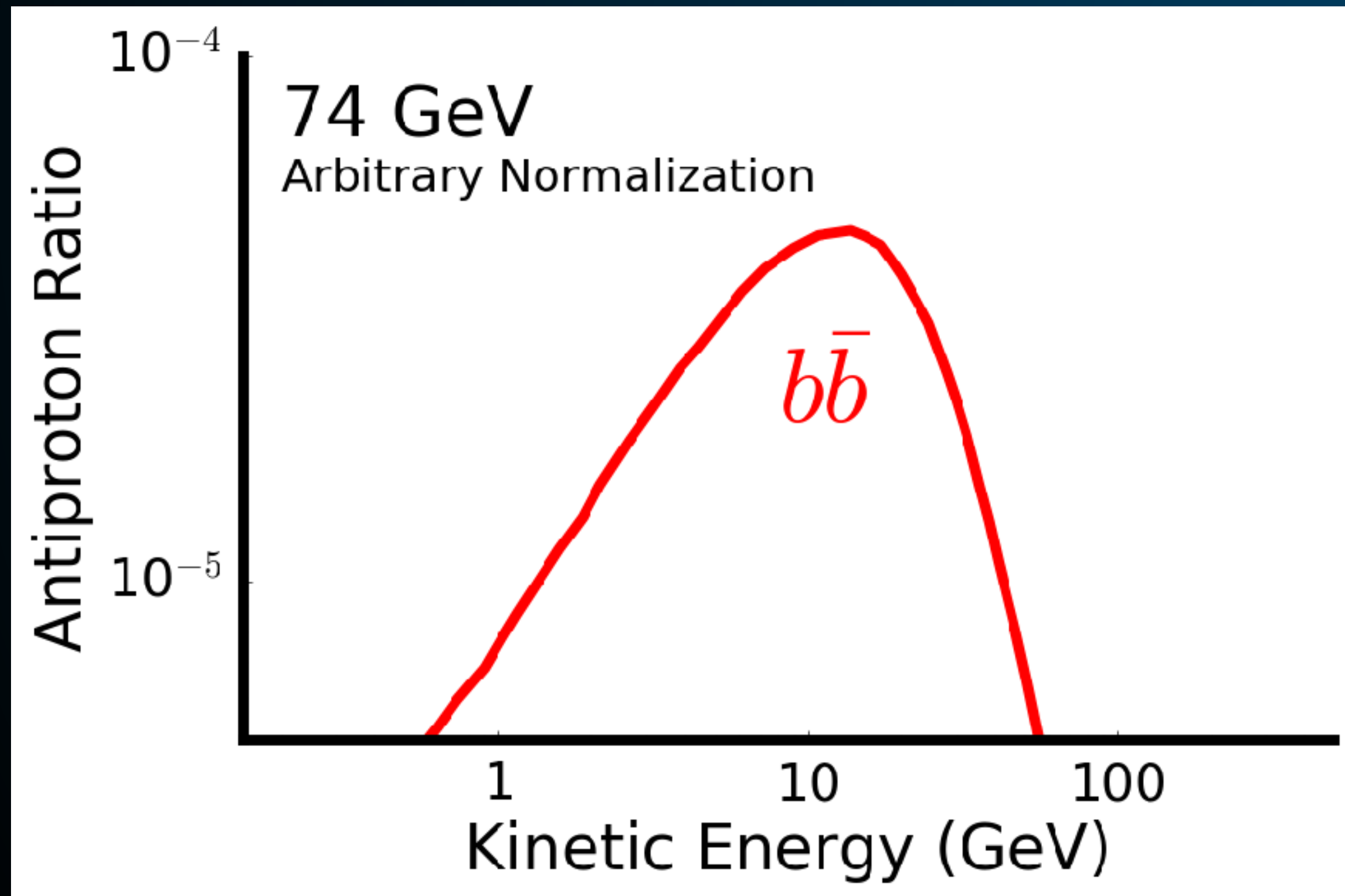
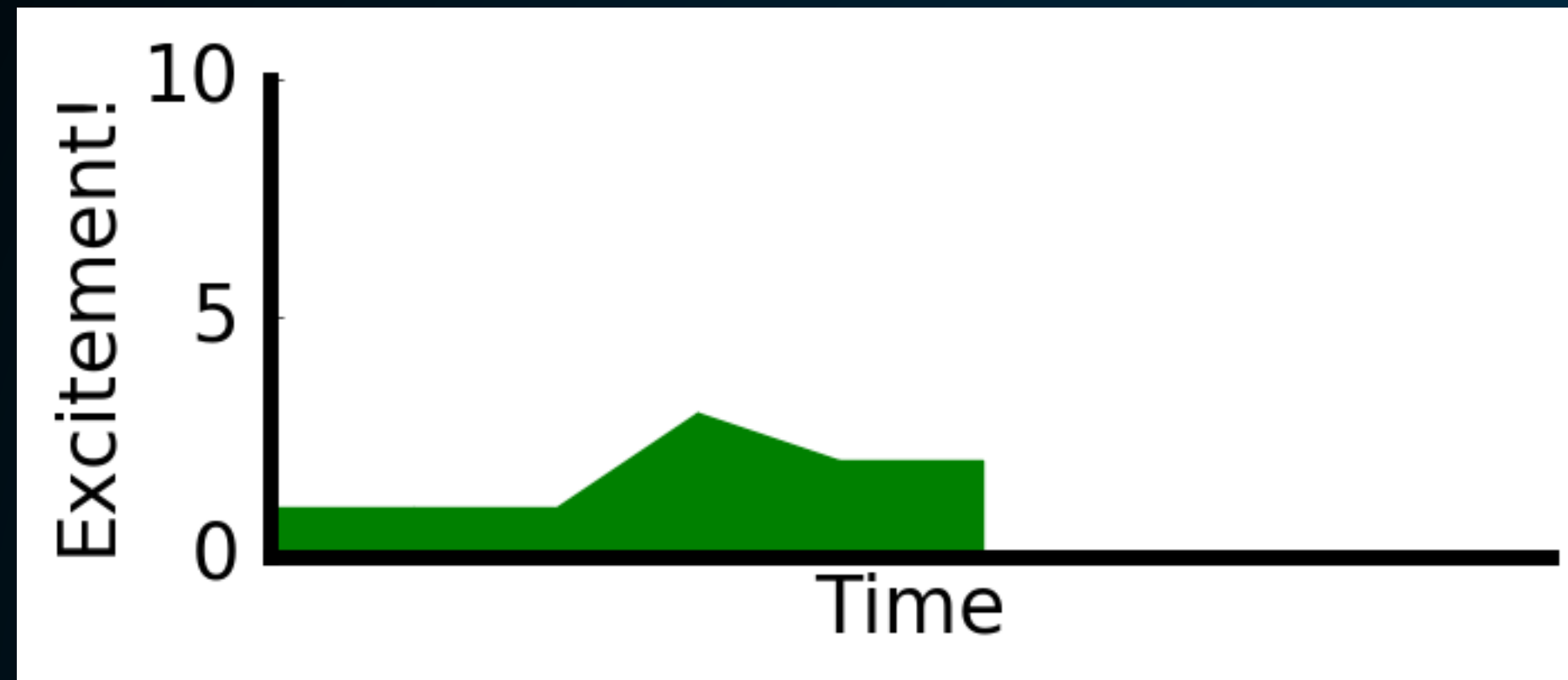
# The Antiproton Excess



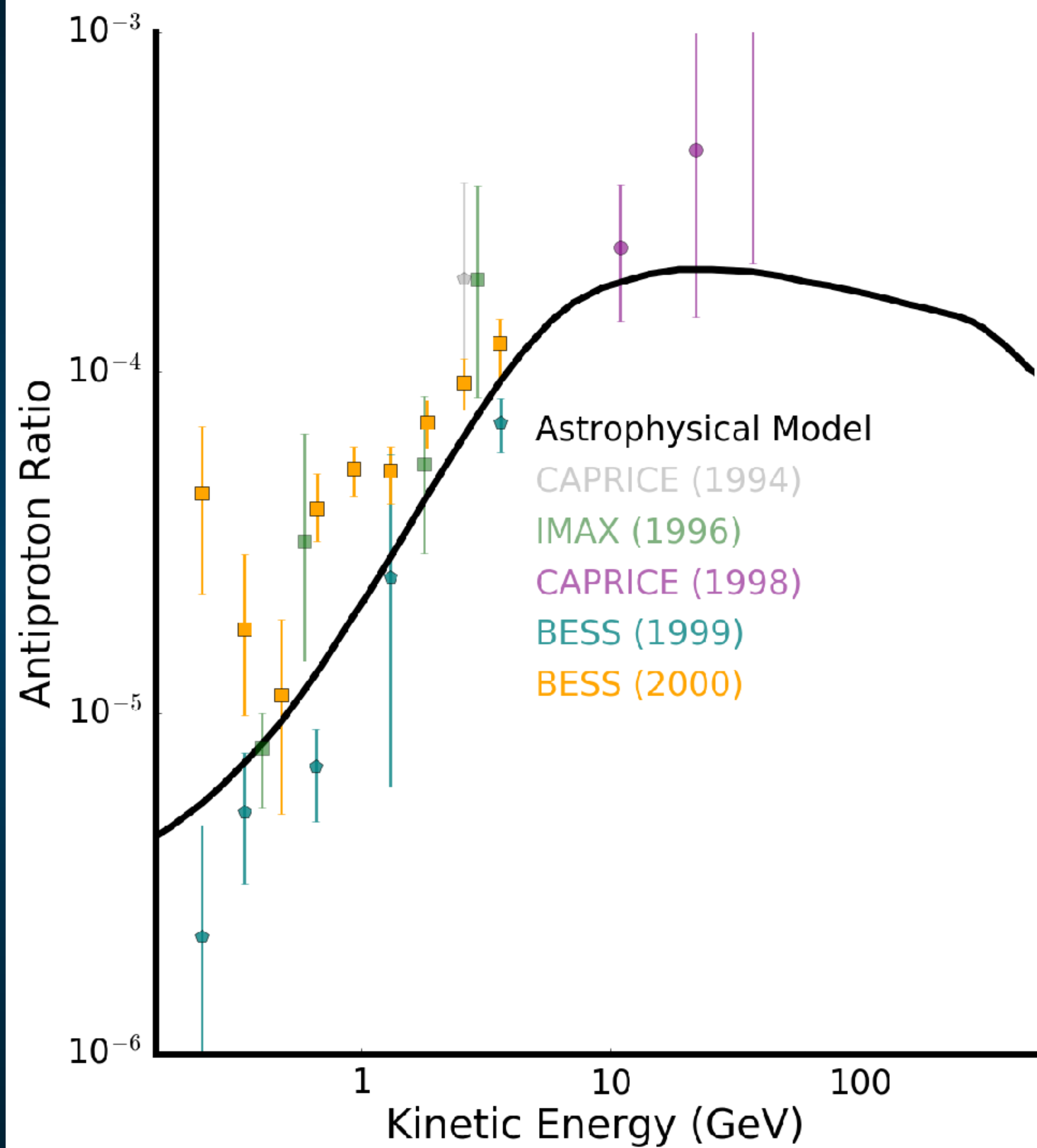
(Not an exhaustive list of observations)



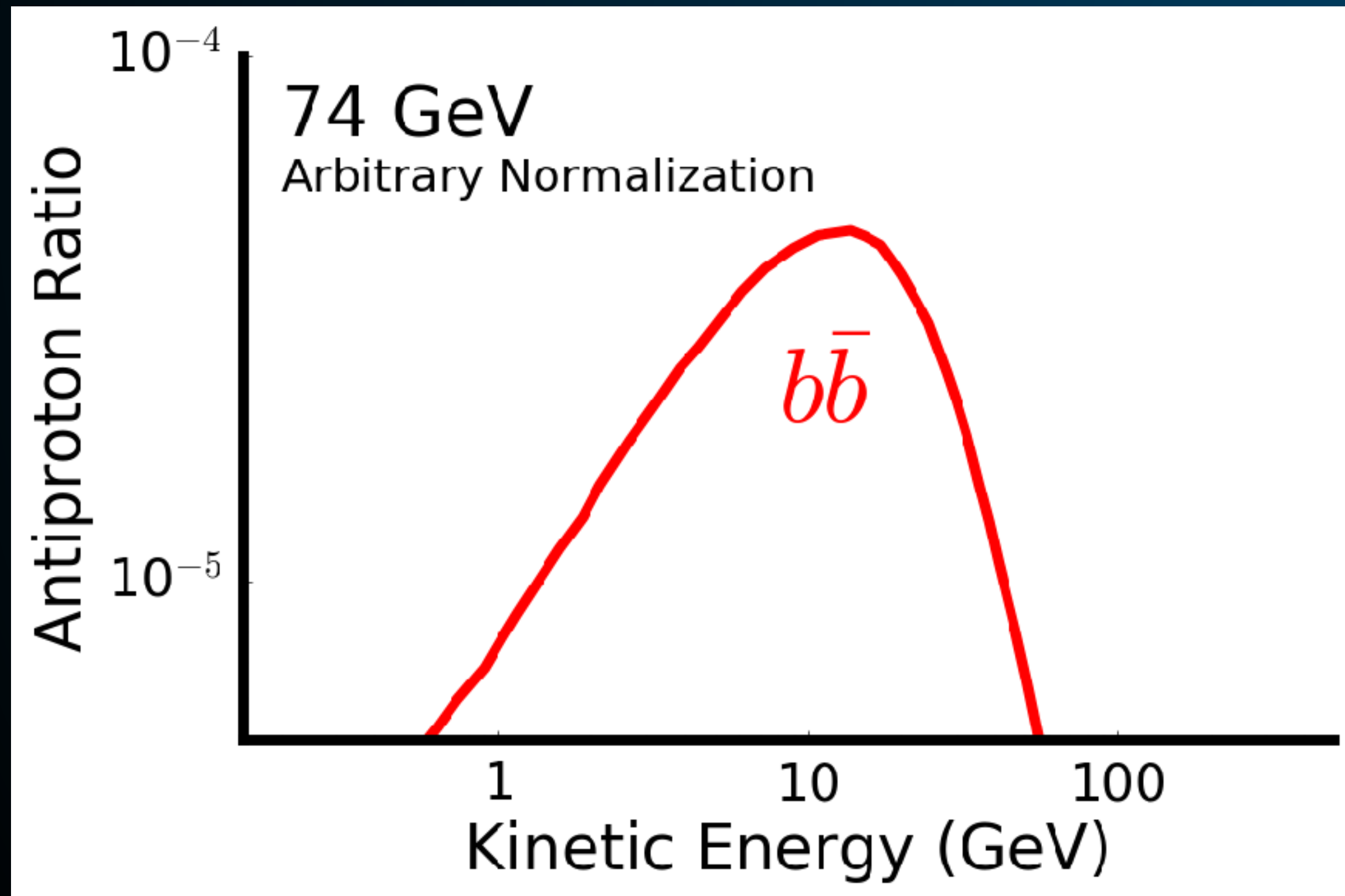
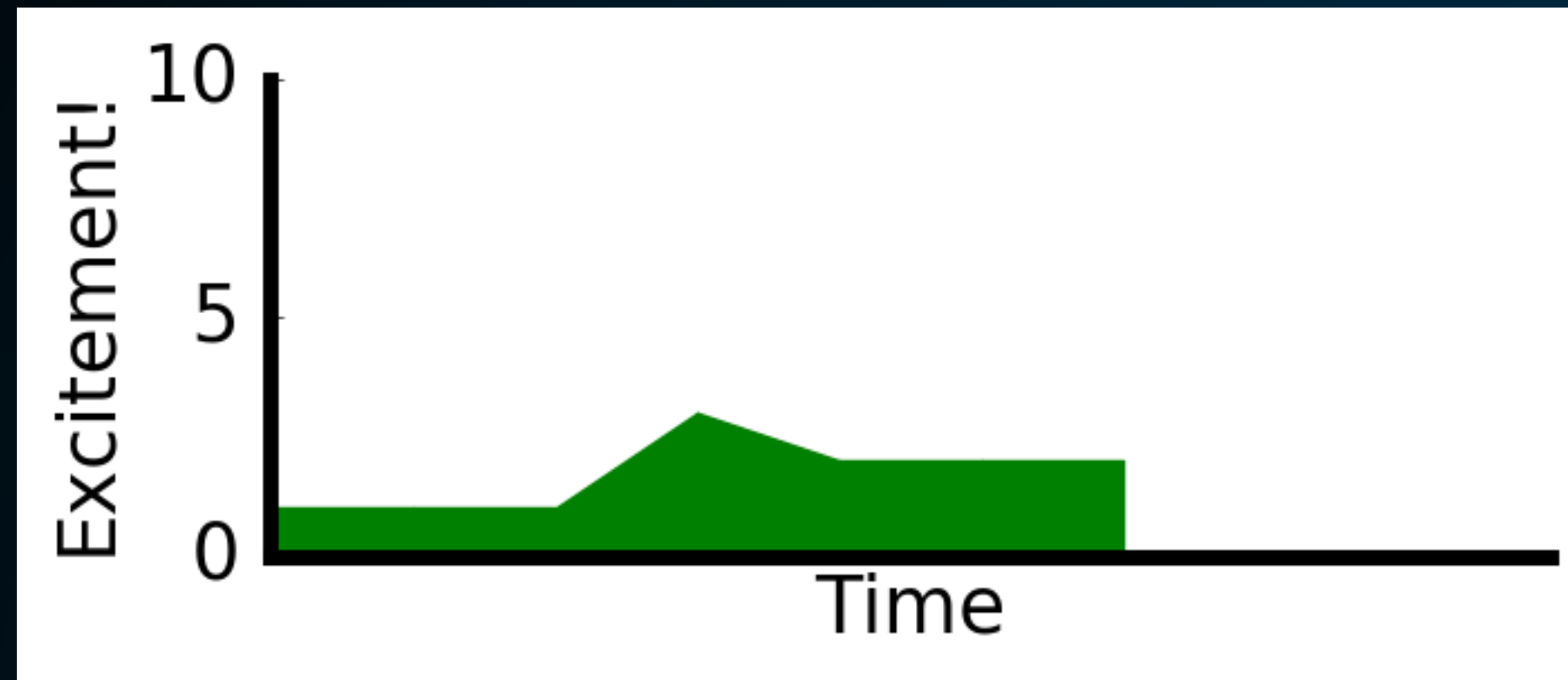
# The Antiproton Excess



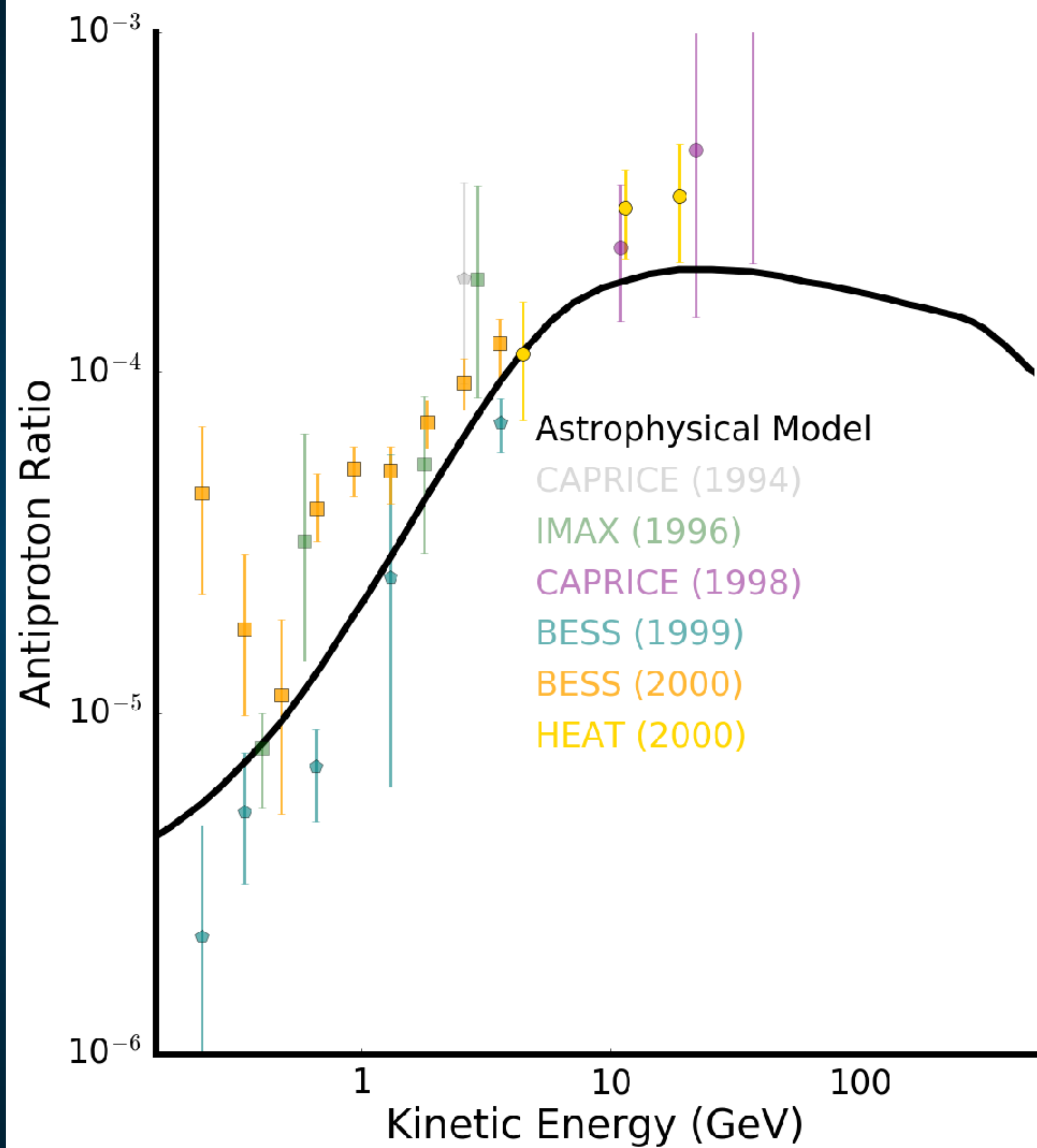
(Not an exhaustive list of observations)



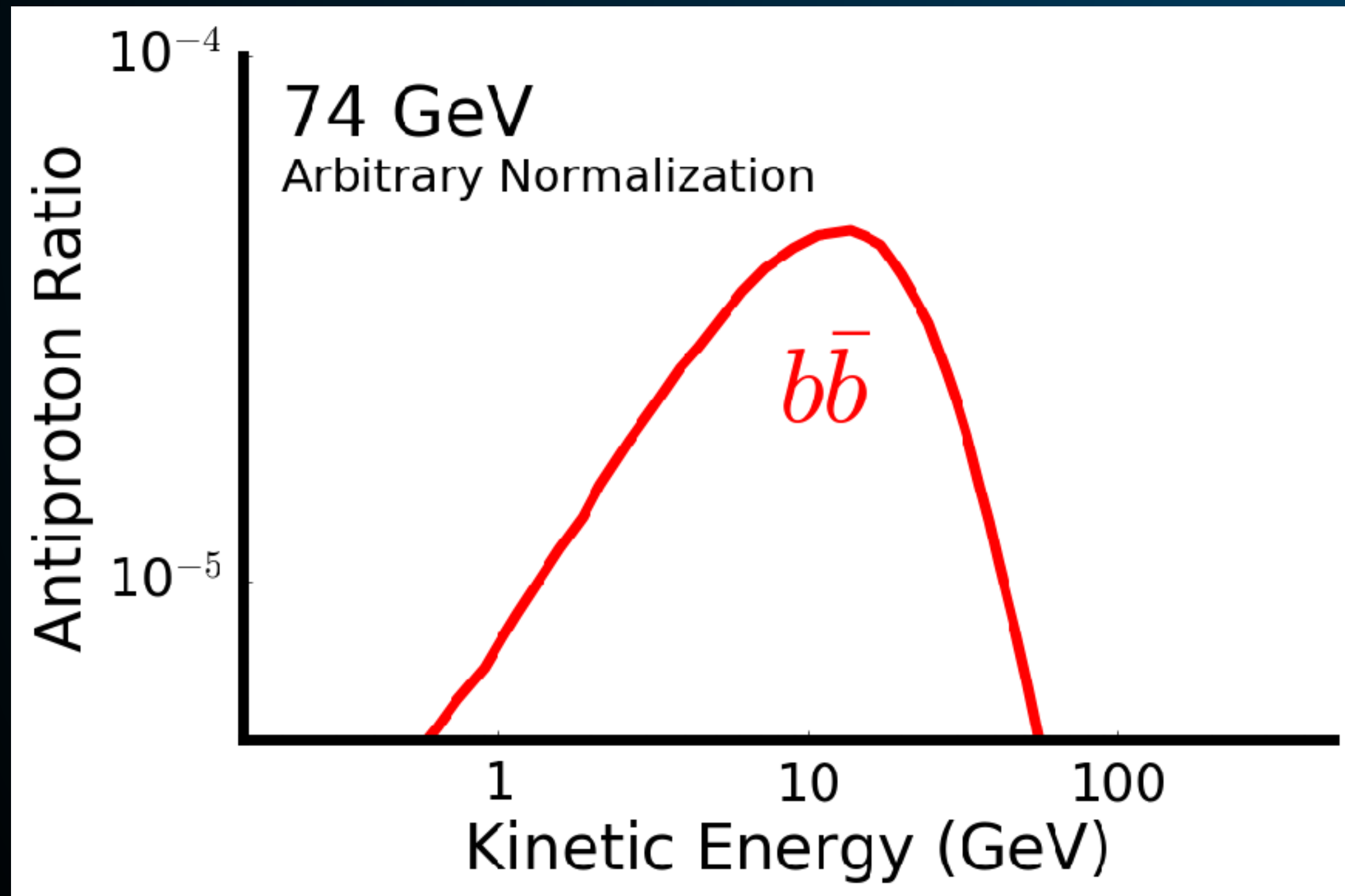
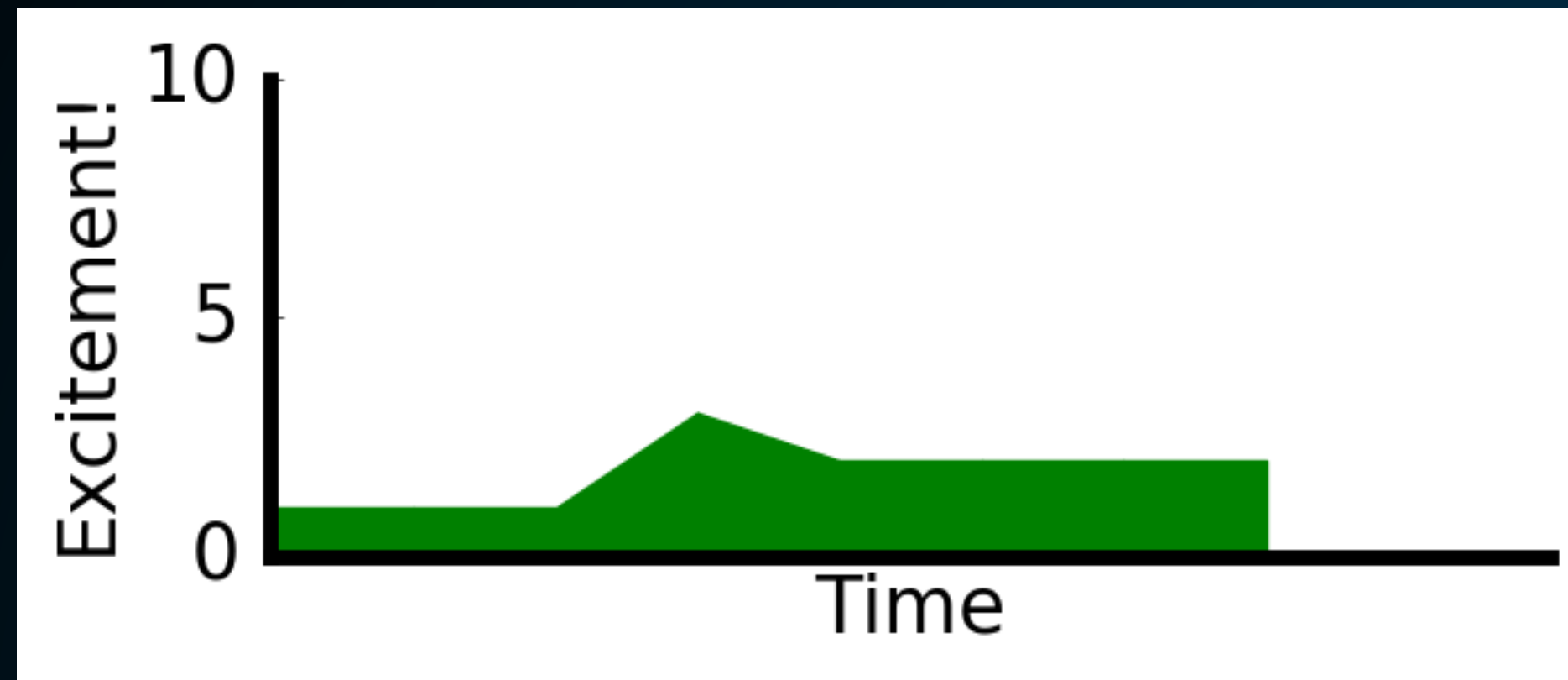
# The Antiproton Excess



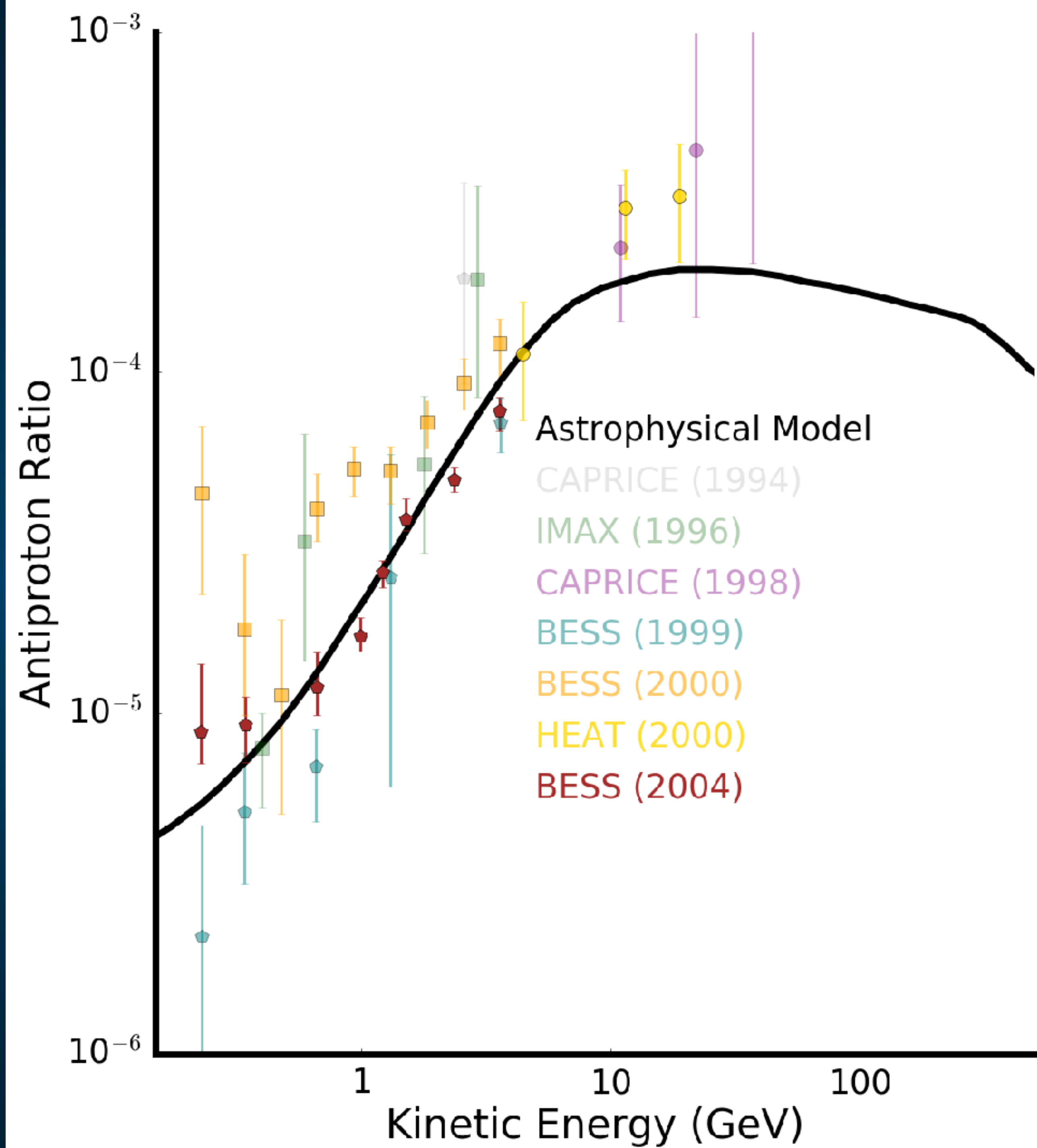
(Not an exhaustive list of observations)



# The Antiproton Excess

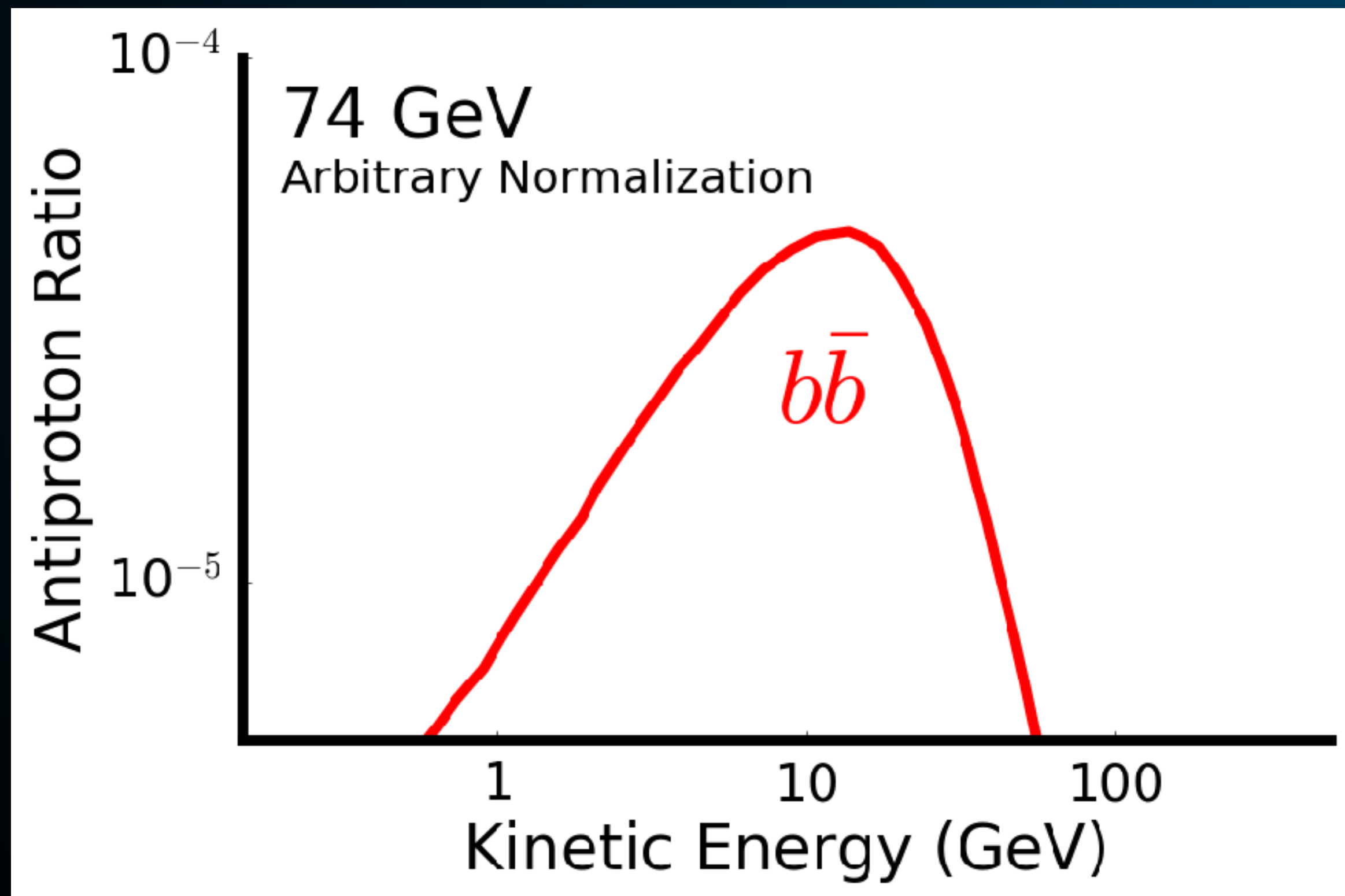
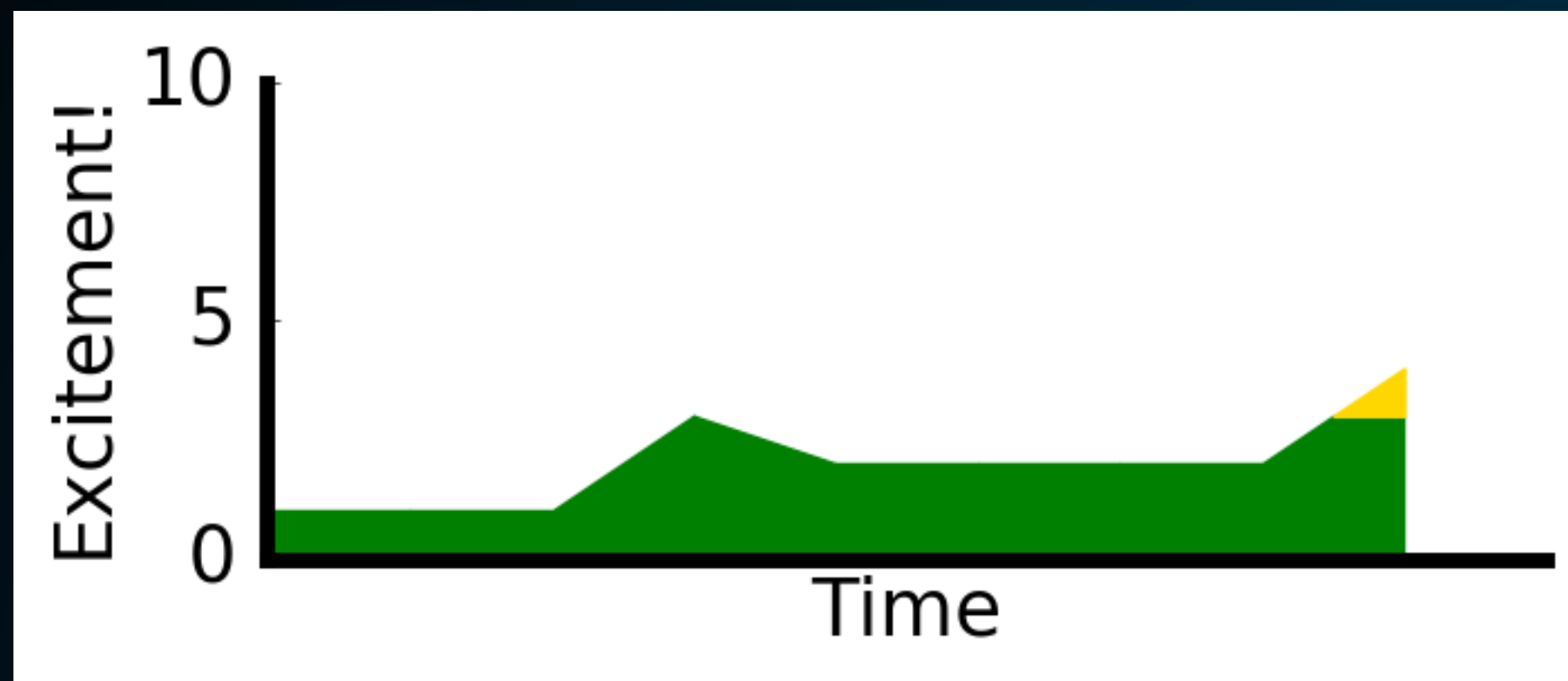


(Not an exhaustive list of observations)

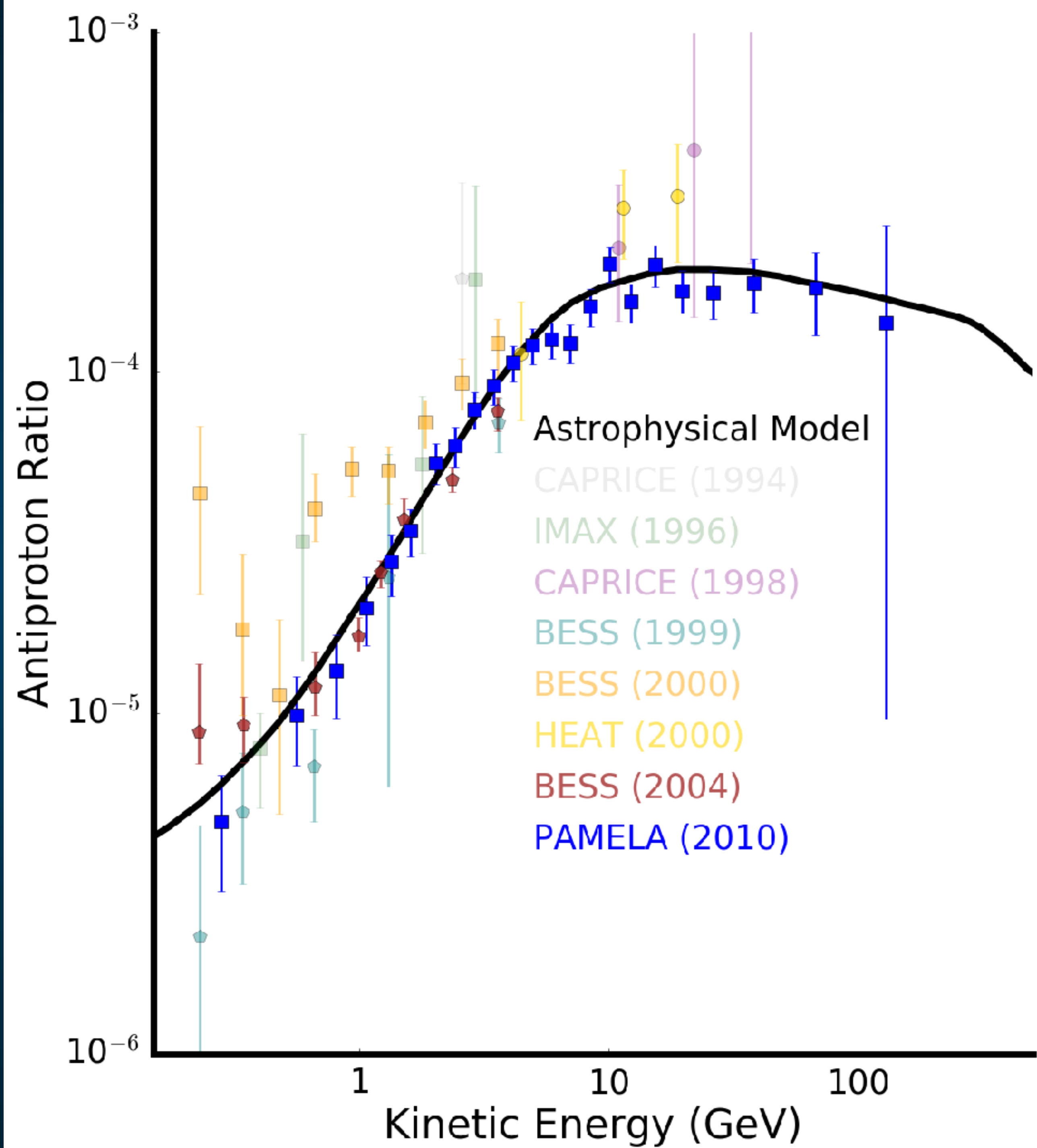




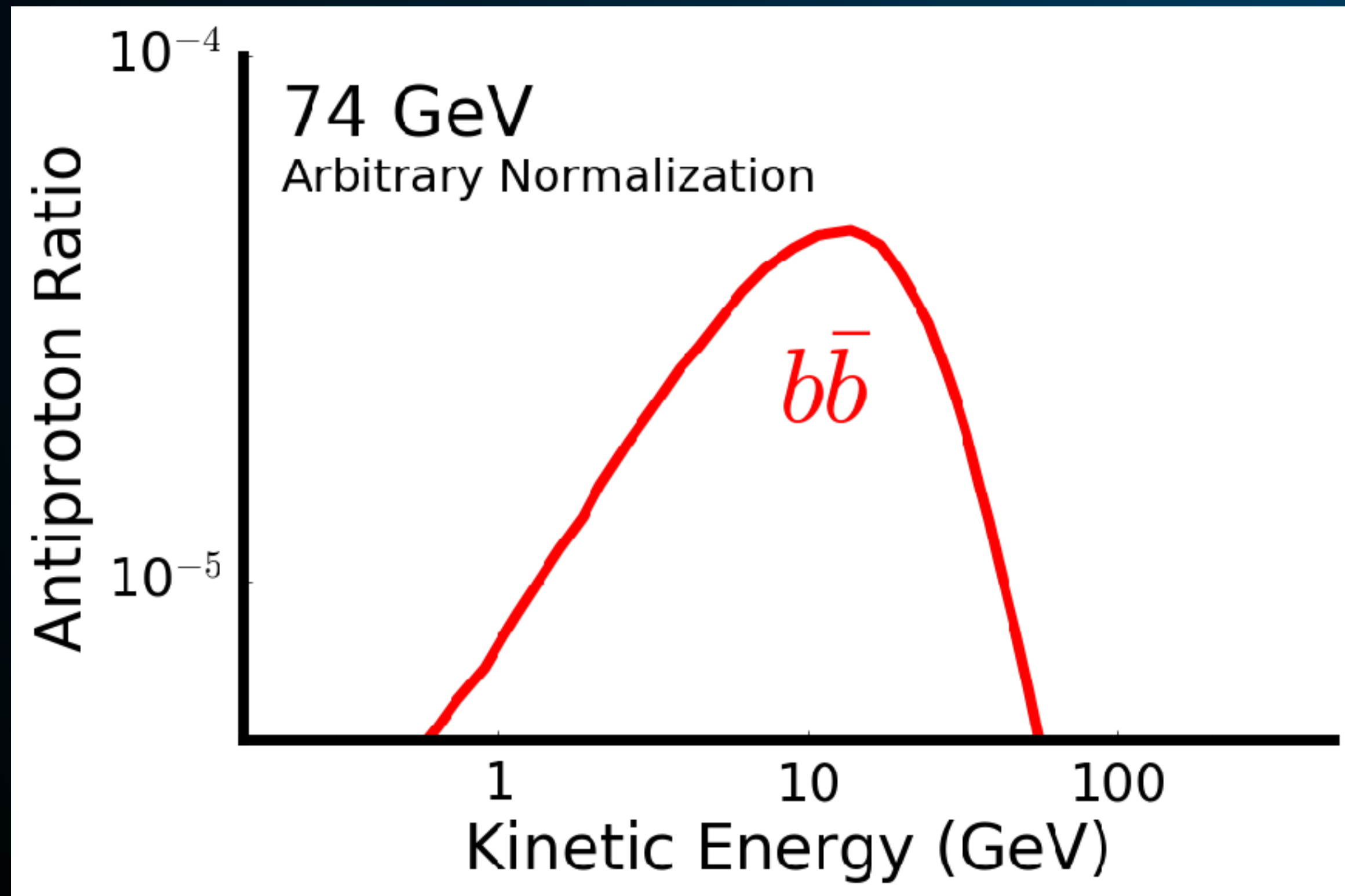
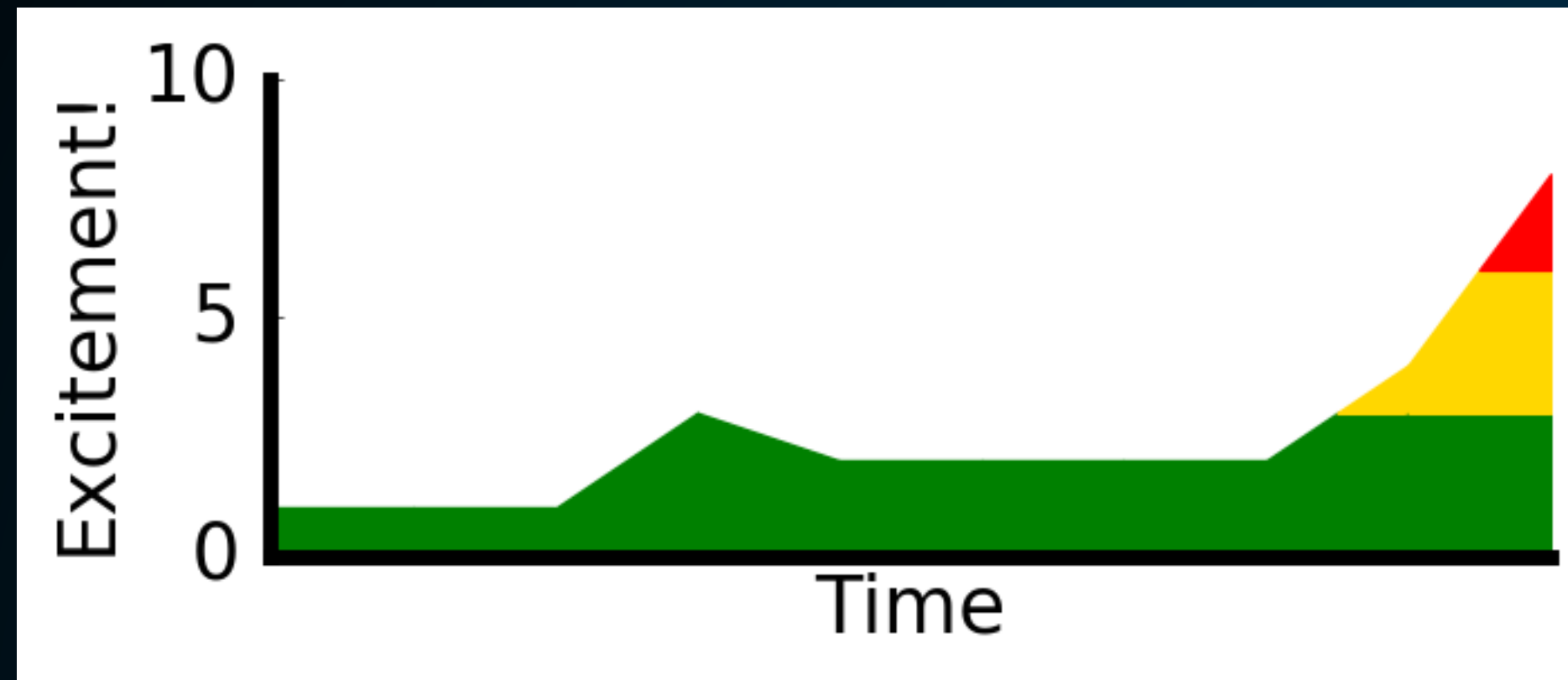
# The Antiproton Excess



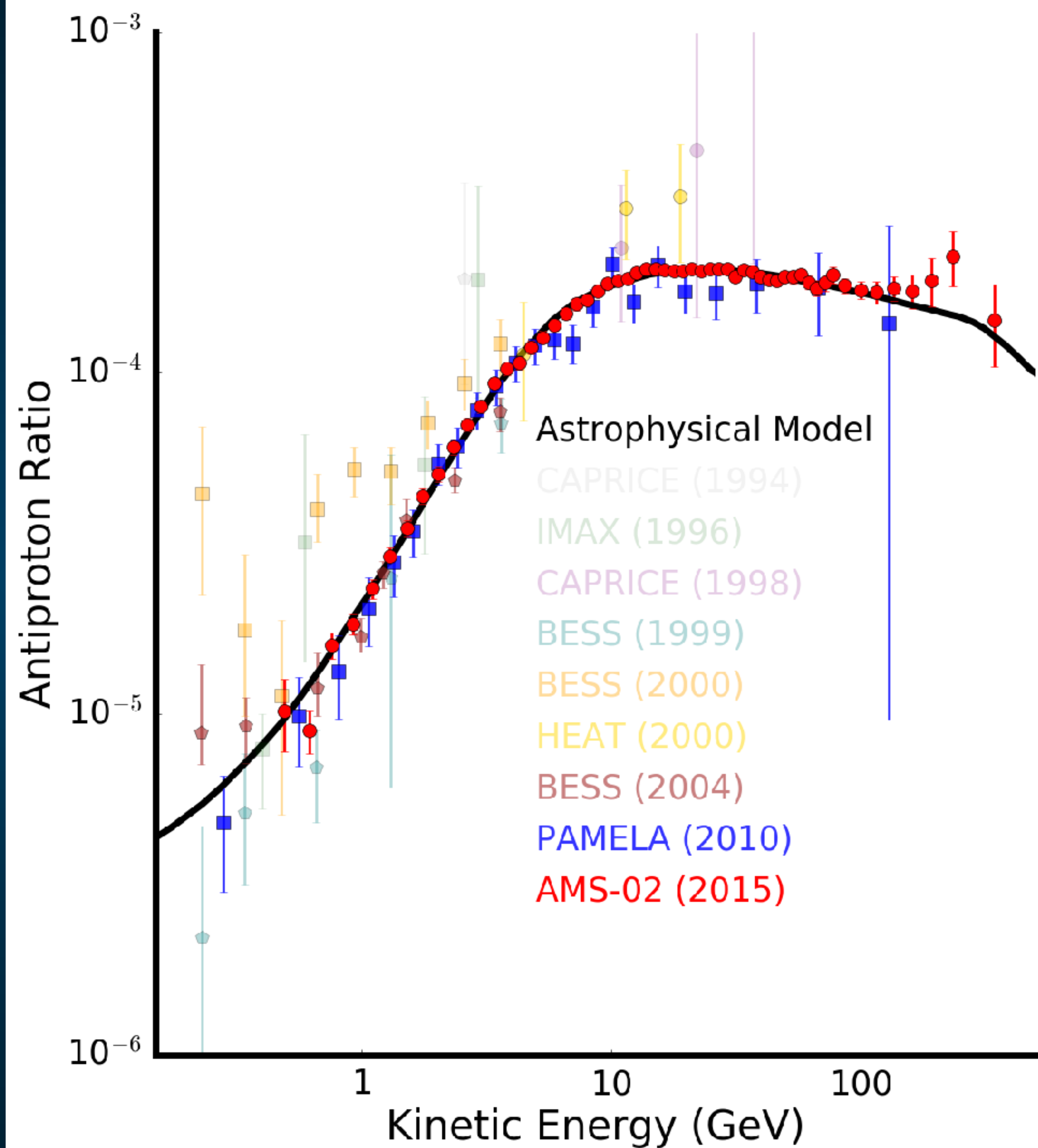
(Not an exhaustive list of observations)



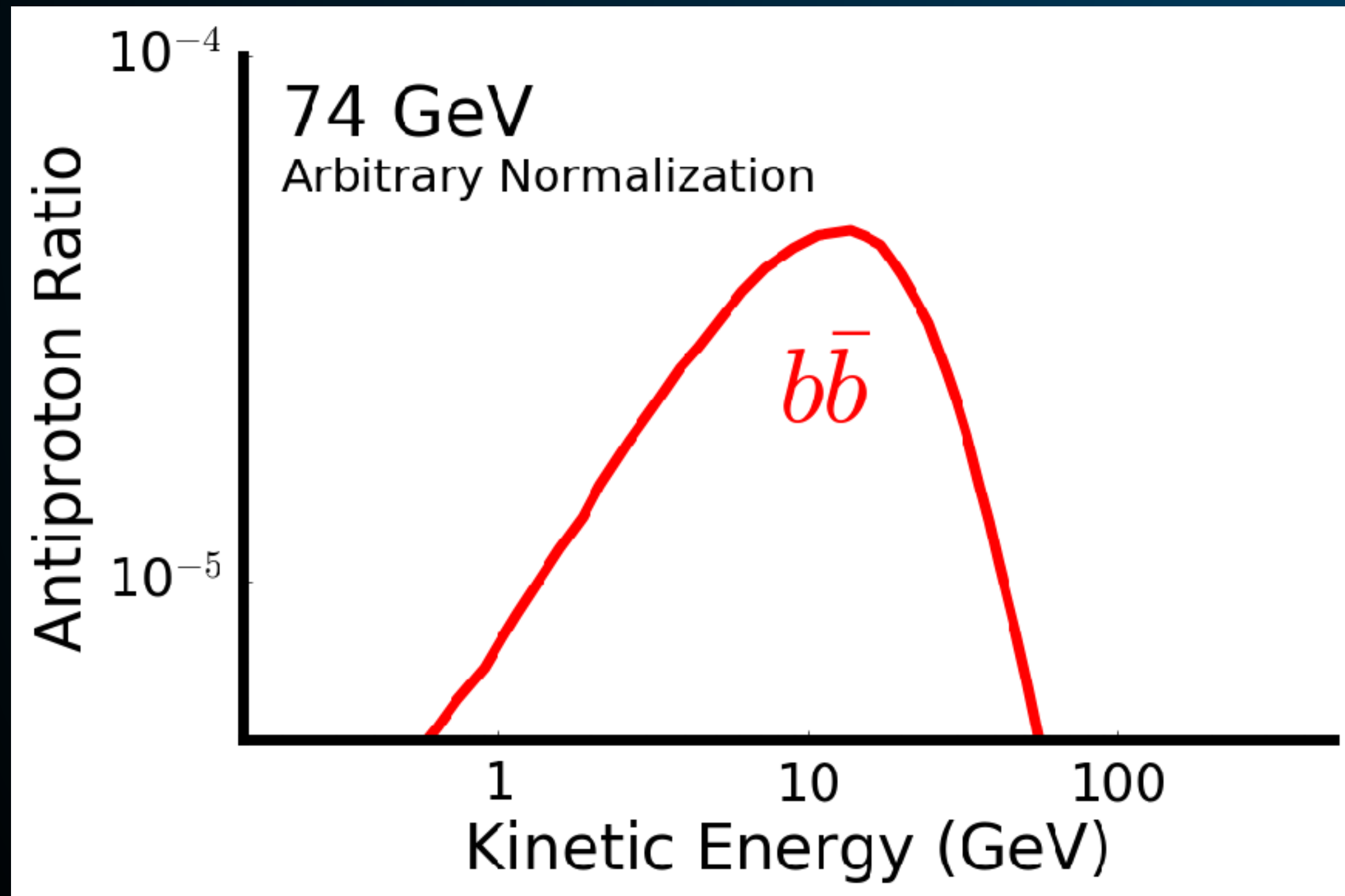
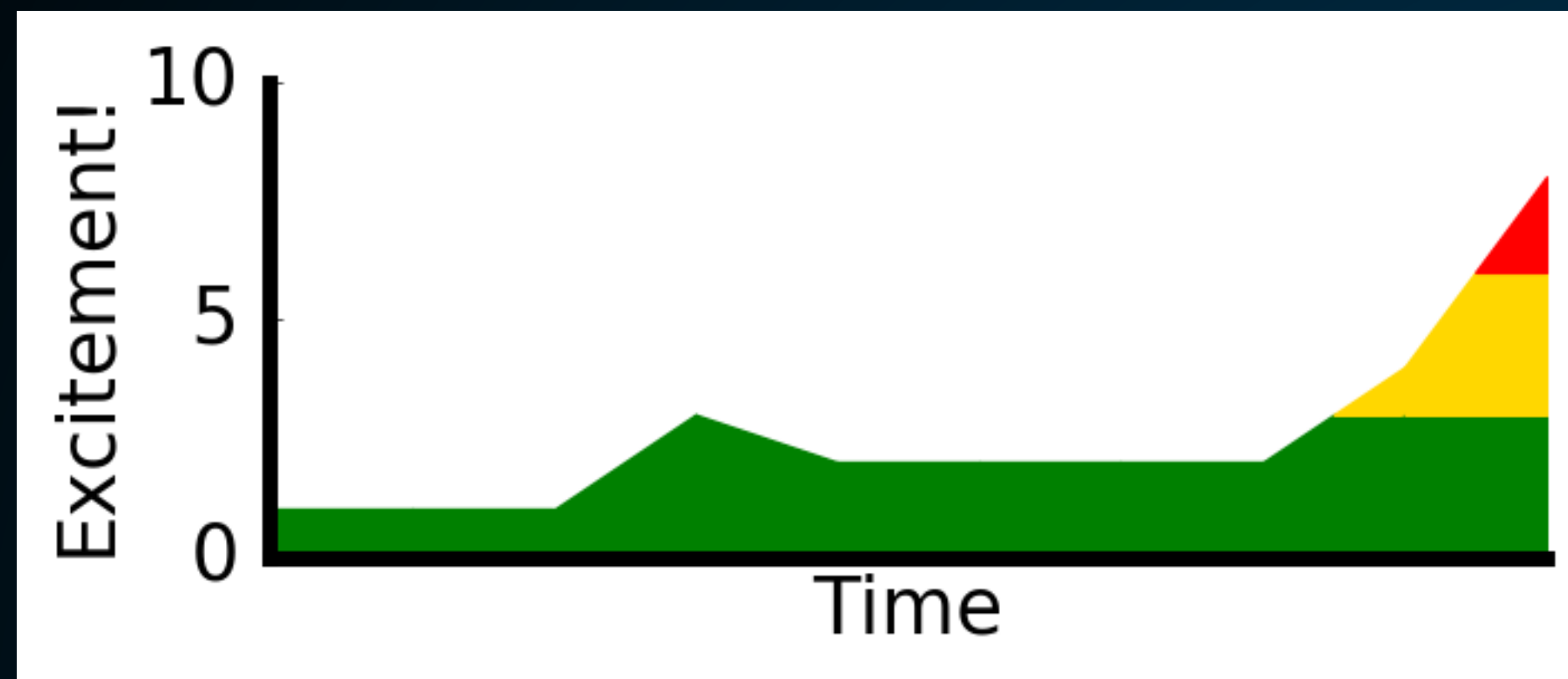
# The Antiproton Excess



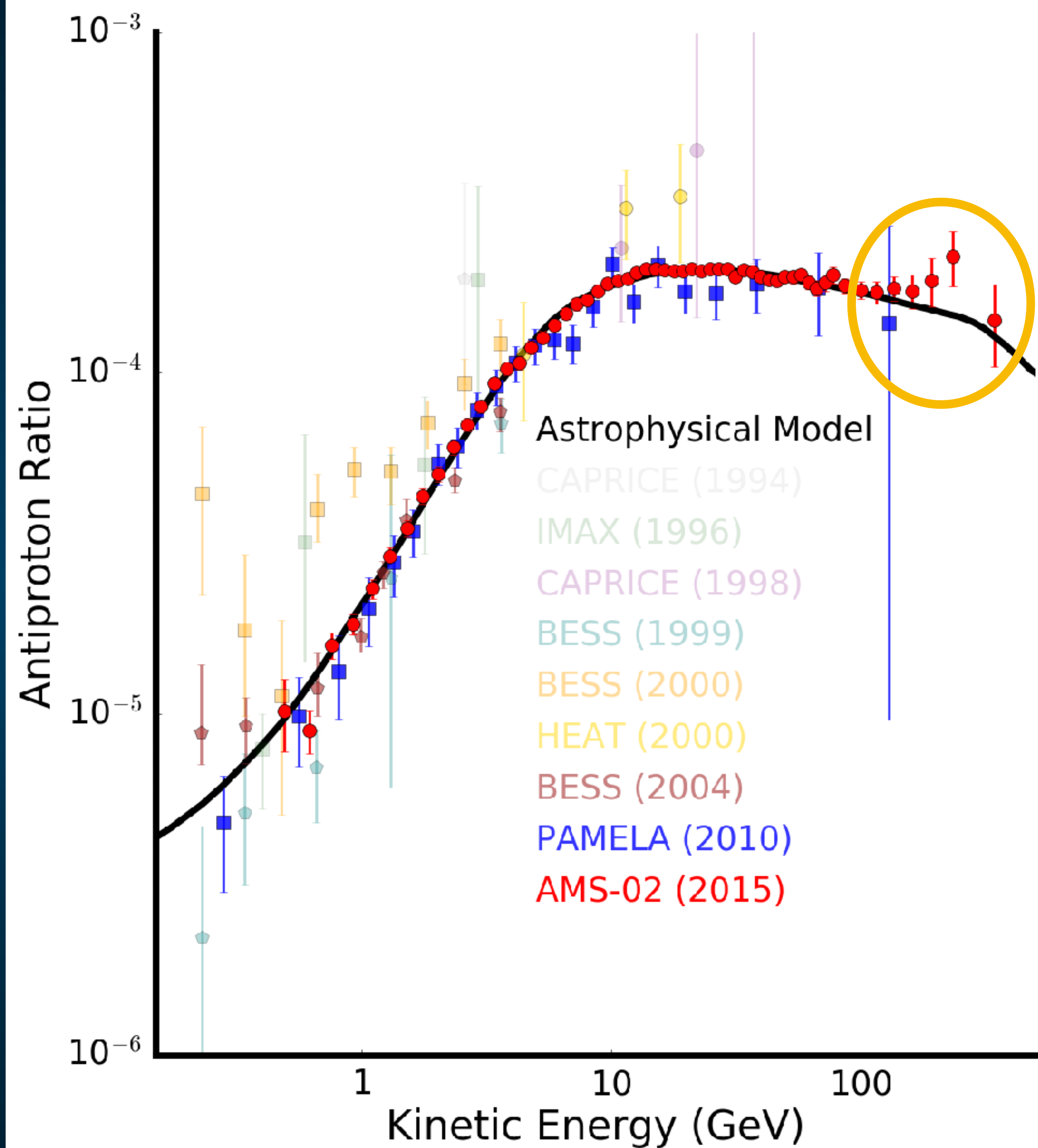
(Not an exhaustive list of observations)



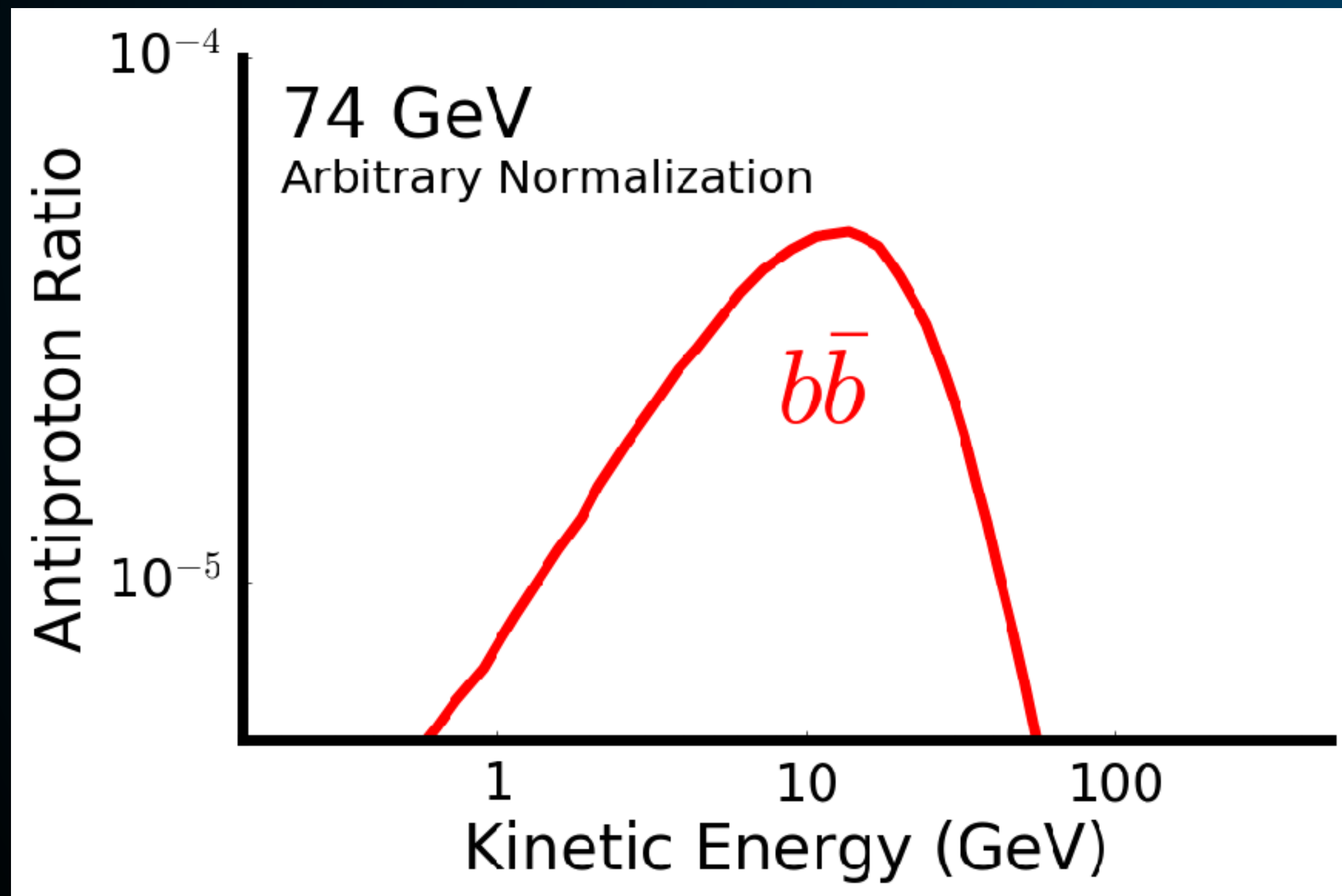
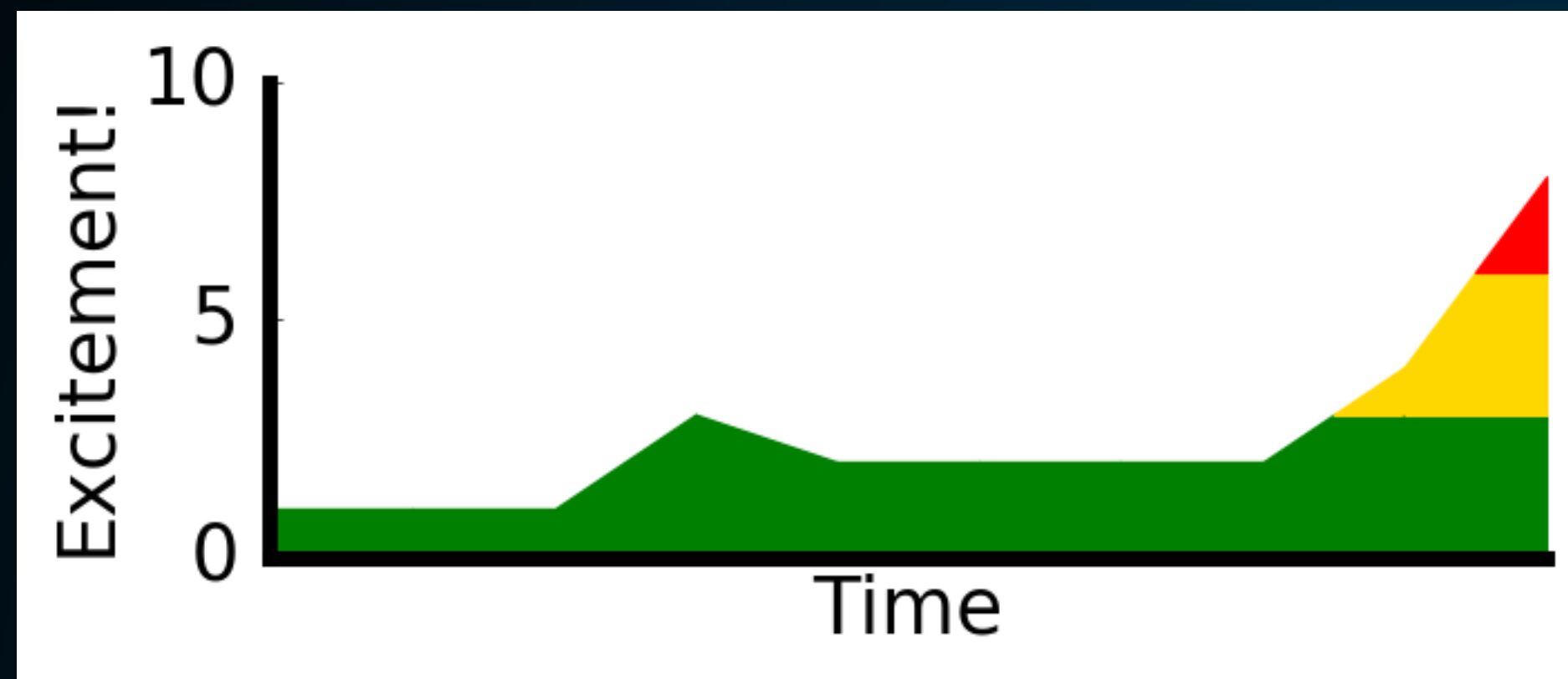
# The Antiproton Excess



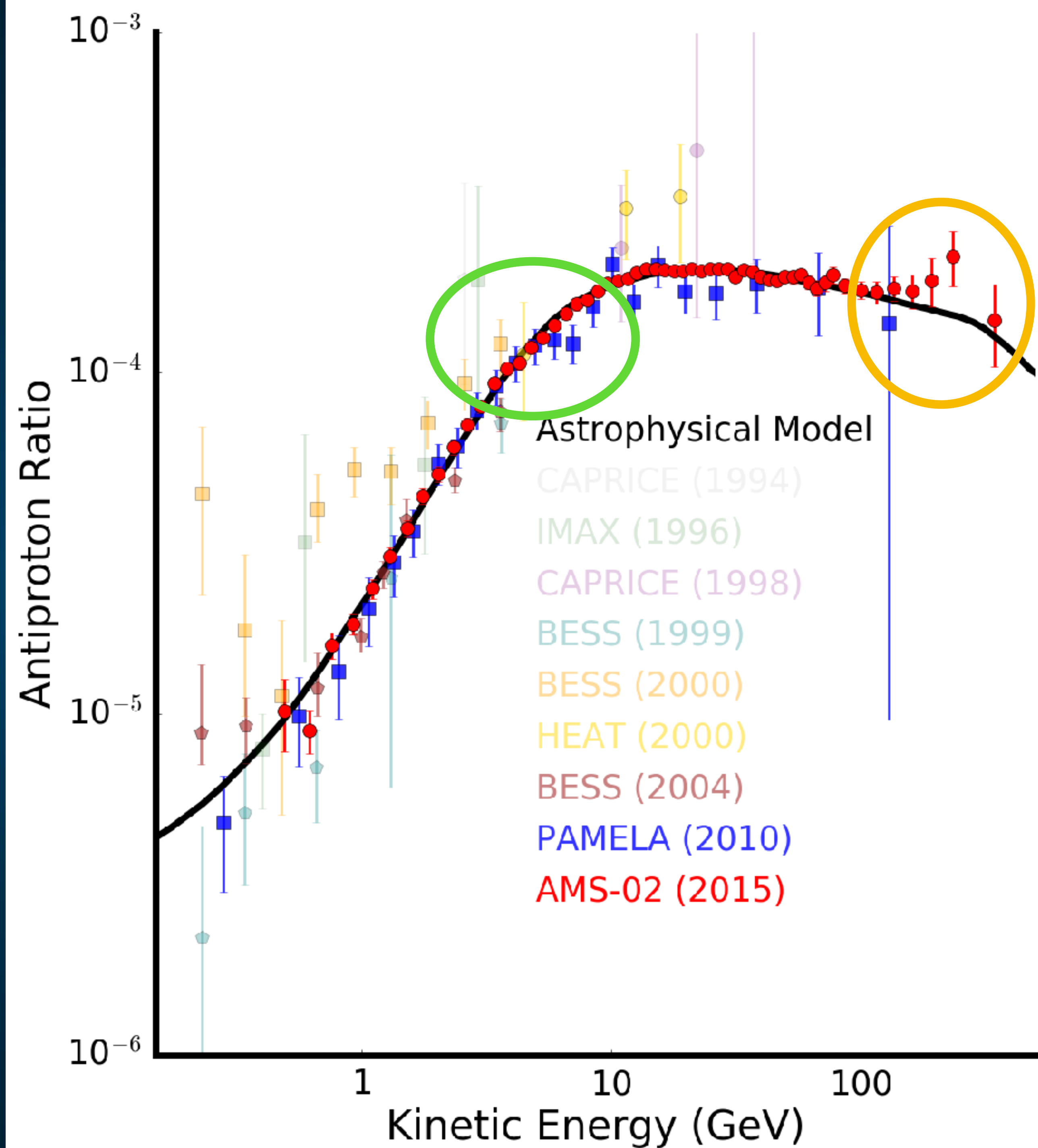
(Not an exhaustive list of observations)



# The Antiproton Excess



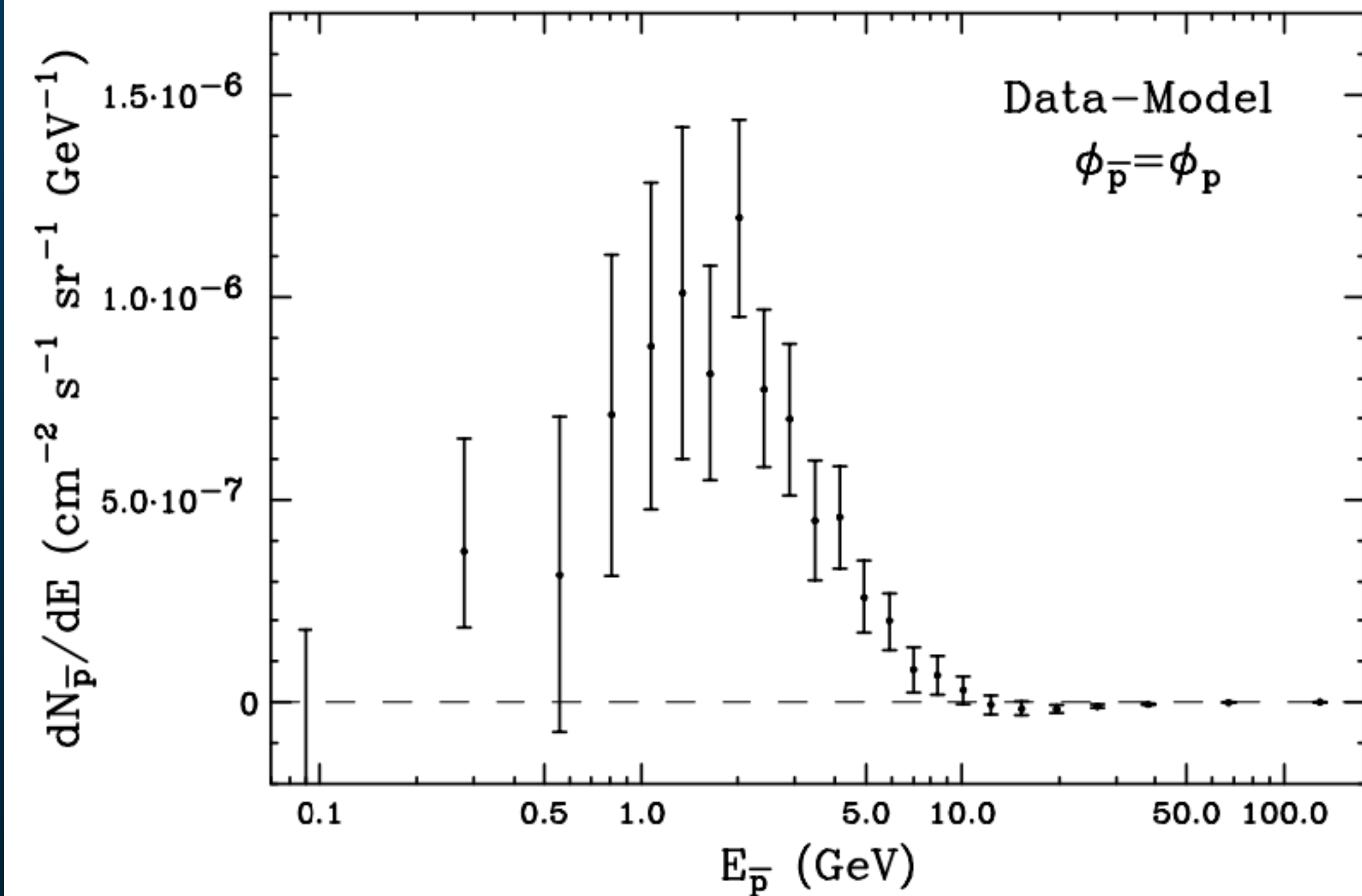
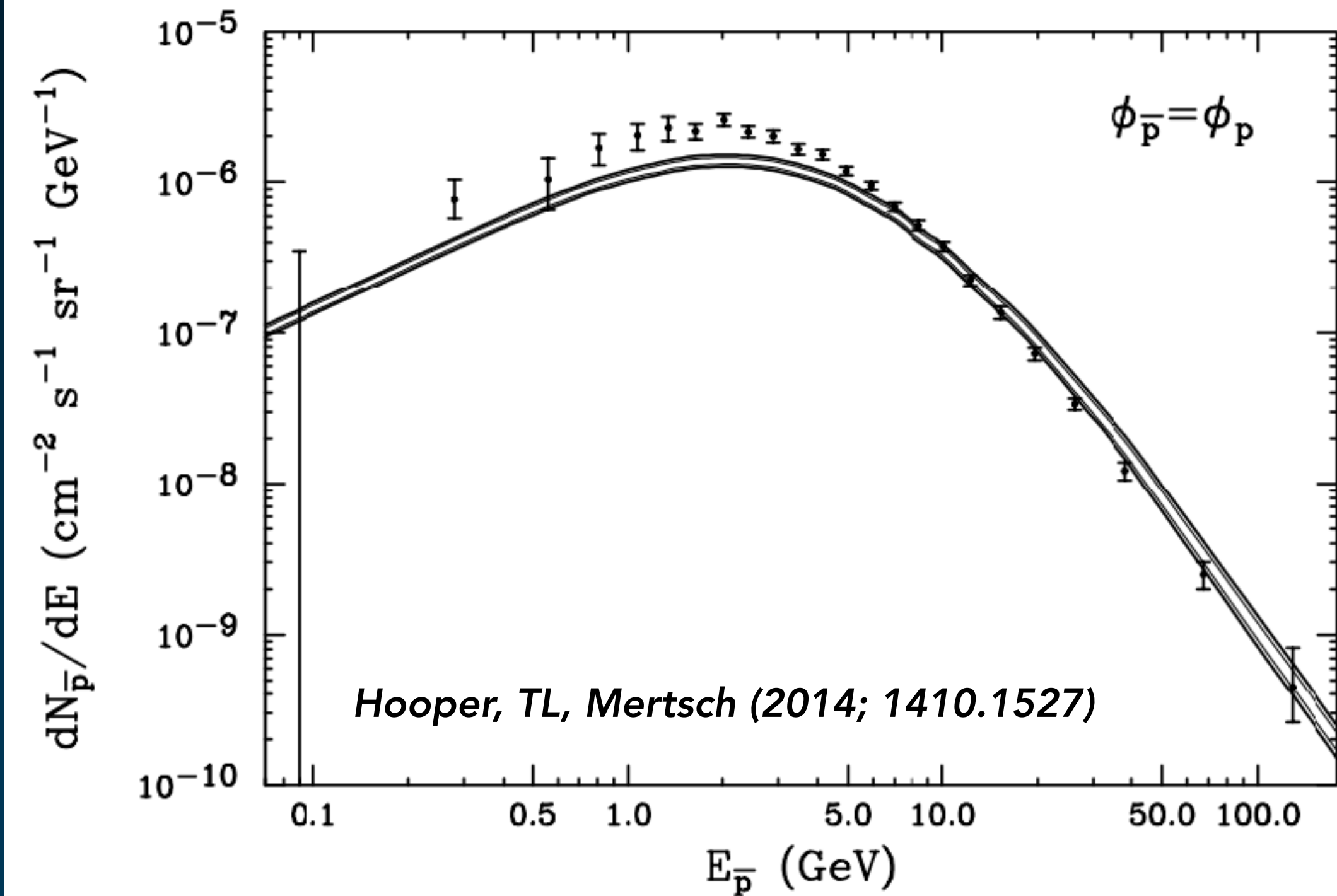
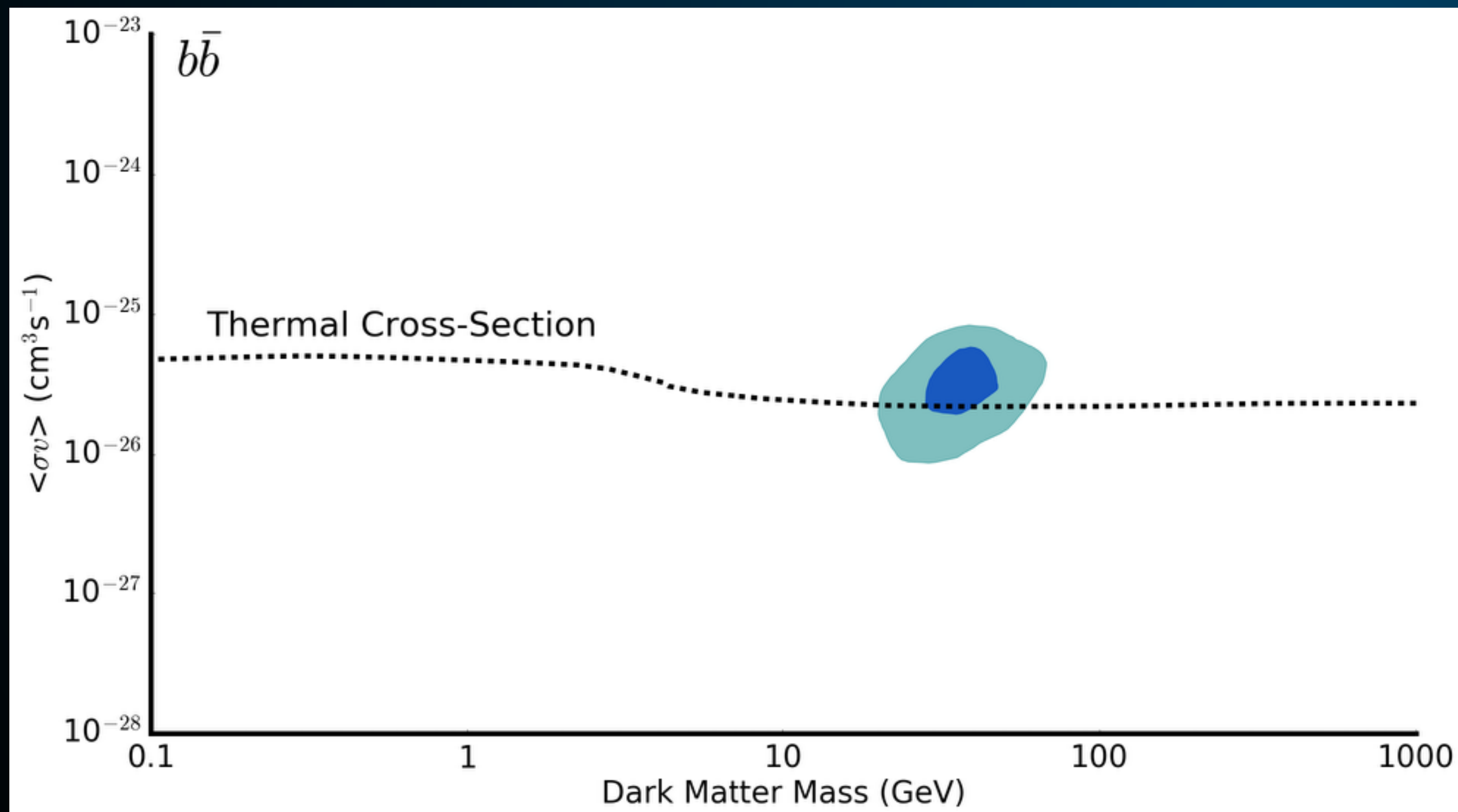
(Not an exhaustive list of observations)



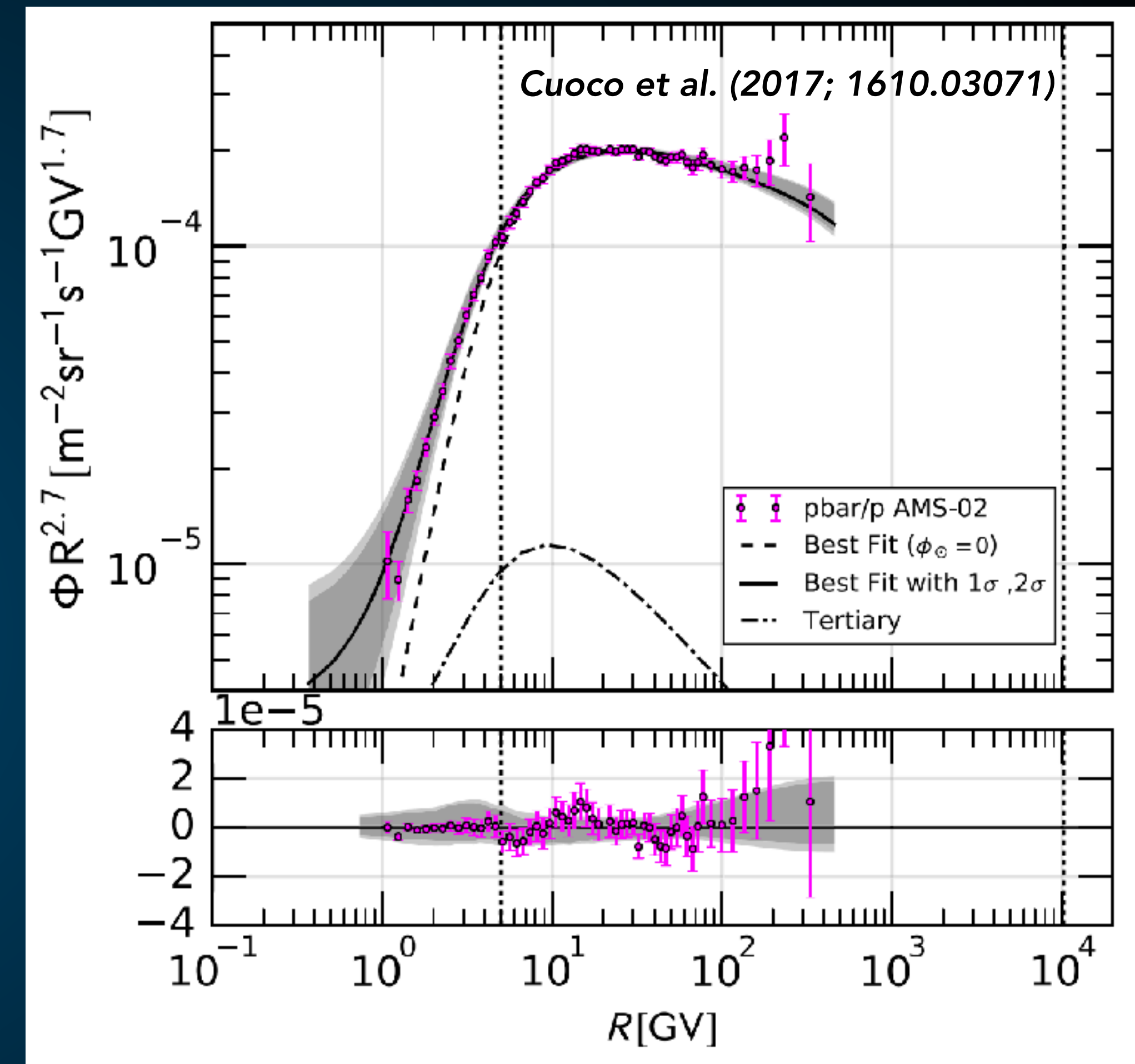
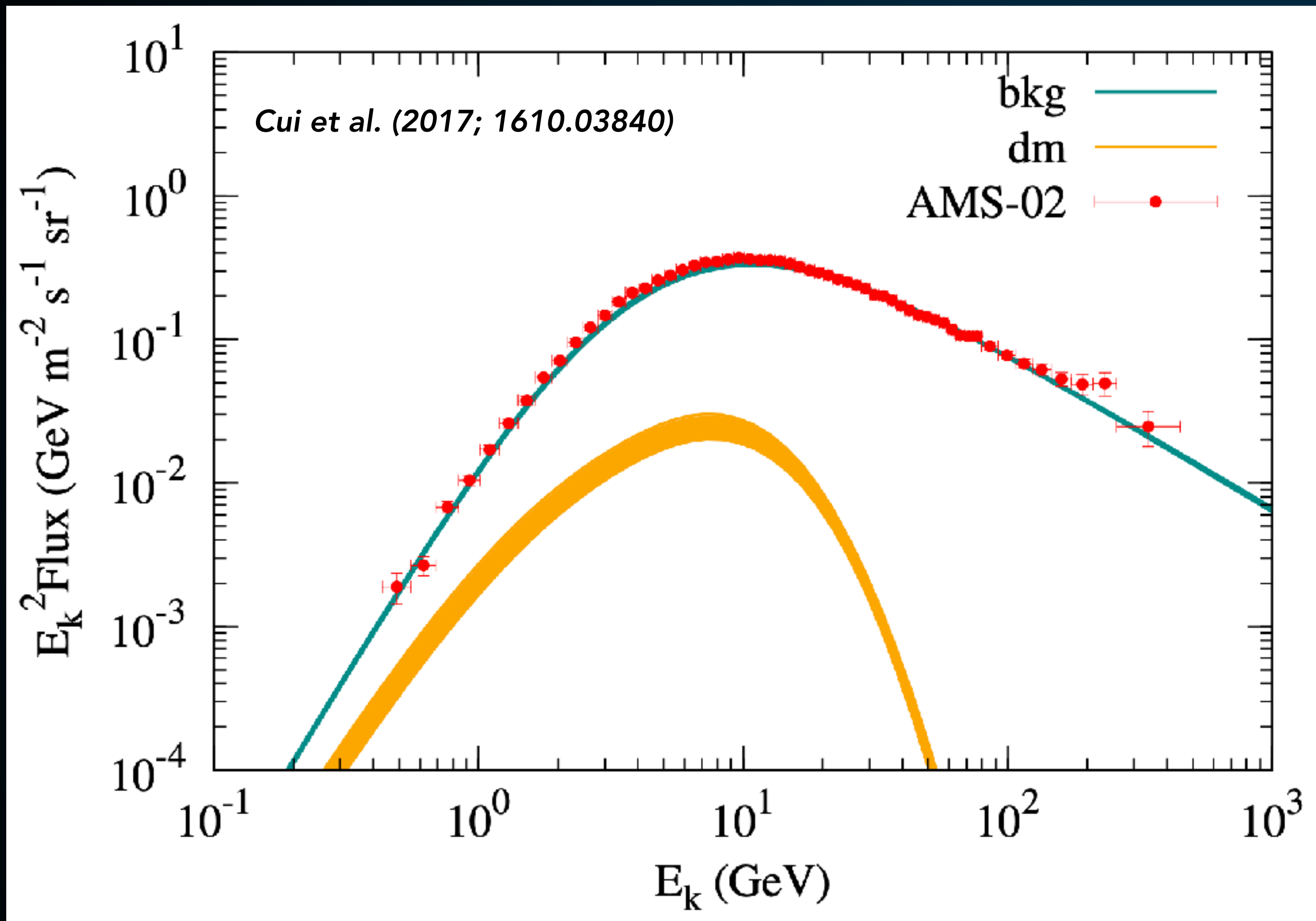
# The Antiproton Excess

Hint of Excess in ~5 GeV antiprotons!

Astrophysical Uncertainties can significantly affect the signal.

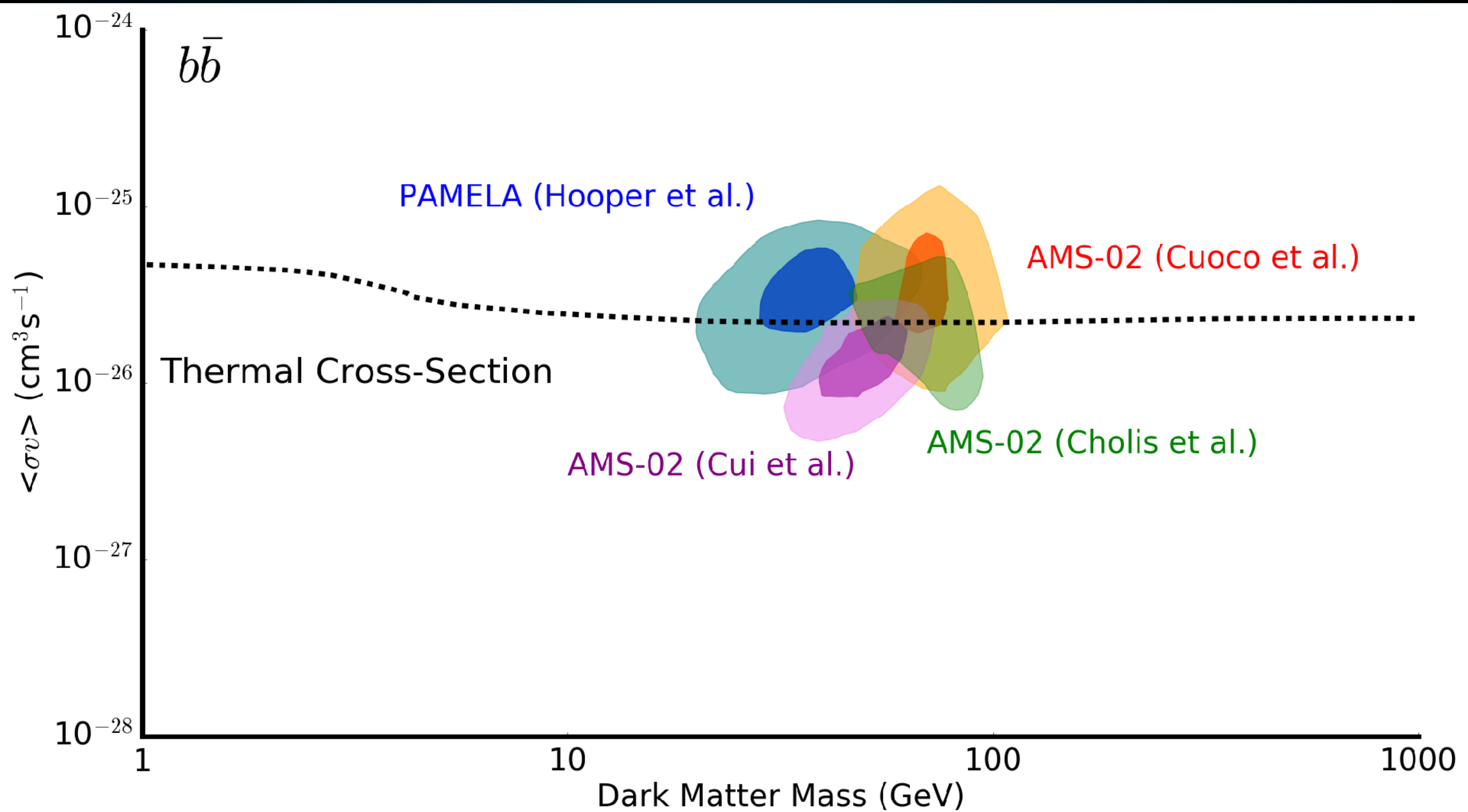


# The Antiproton Excess



**Two papers simultaneously find an excess in the AMS-02 Antiproton Data!**

**Significance approaching (or past)  $5\sigma$  !**



# The Antiproton Excess

With great precision comes great responsibility:

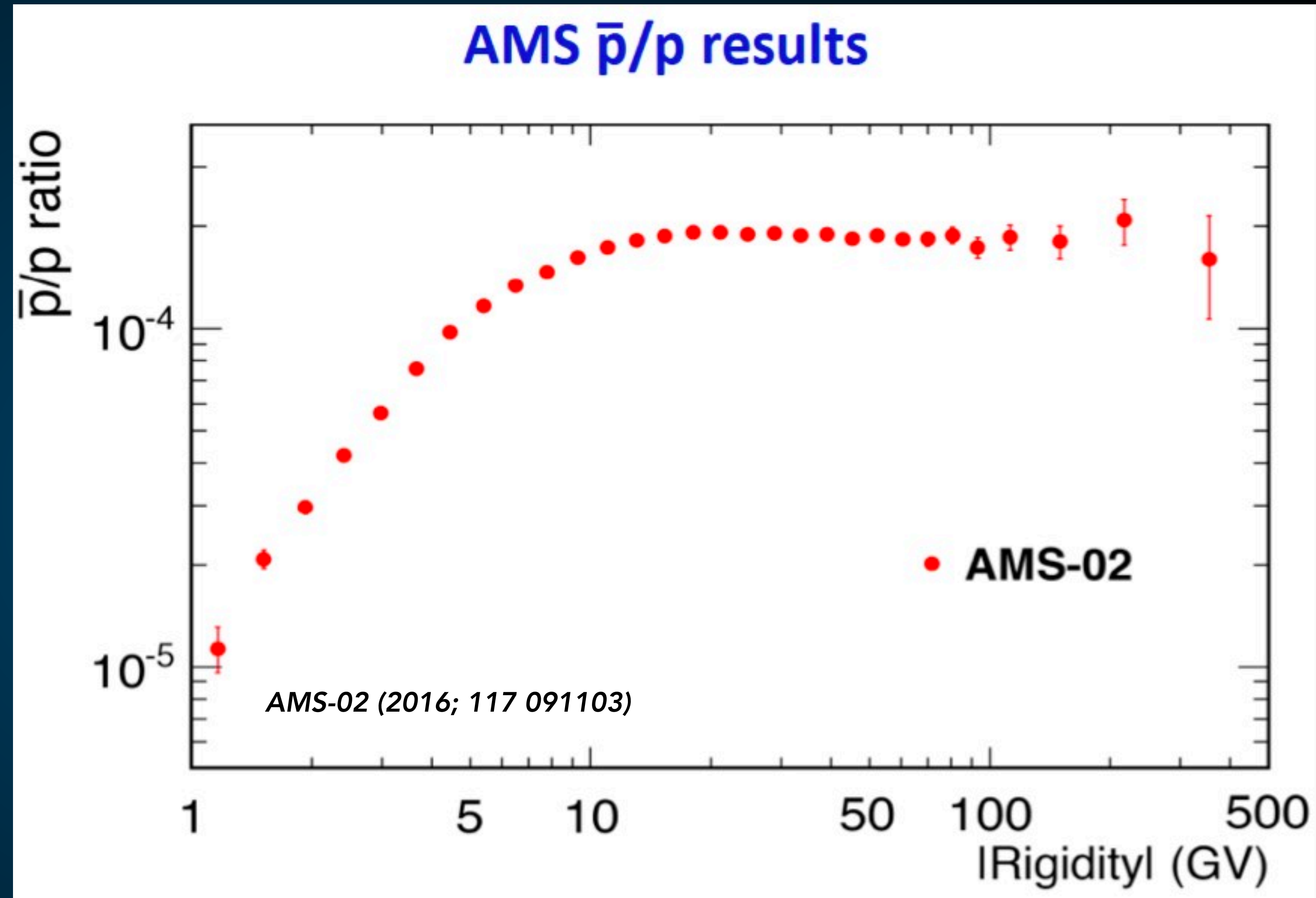
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

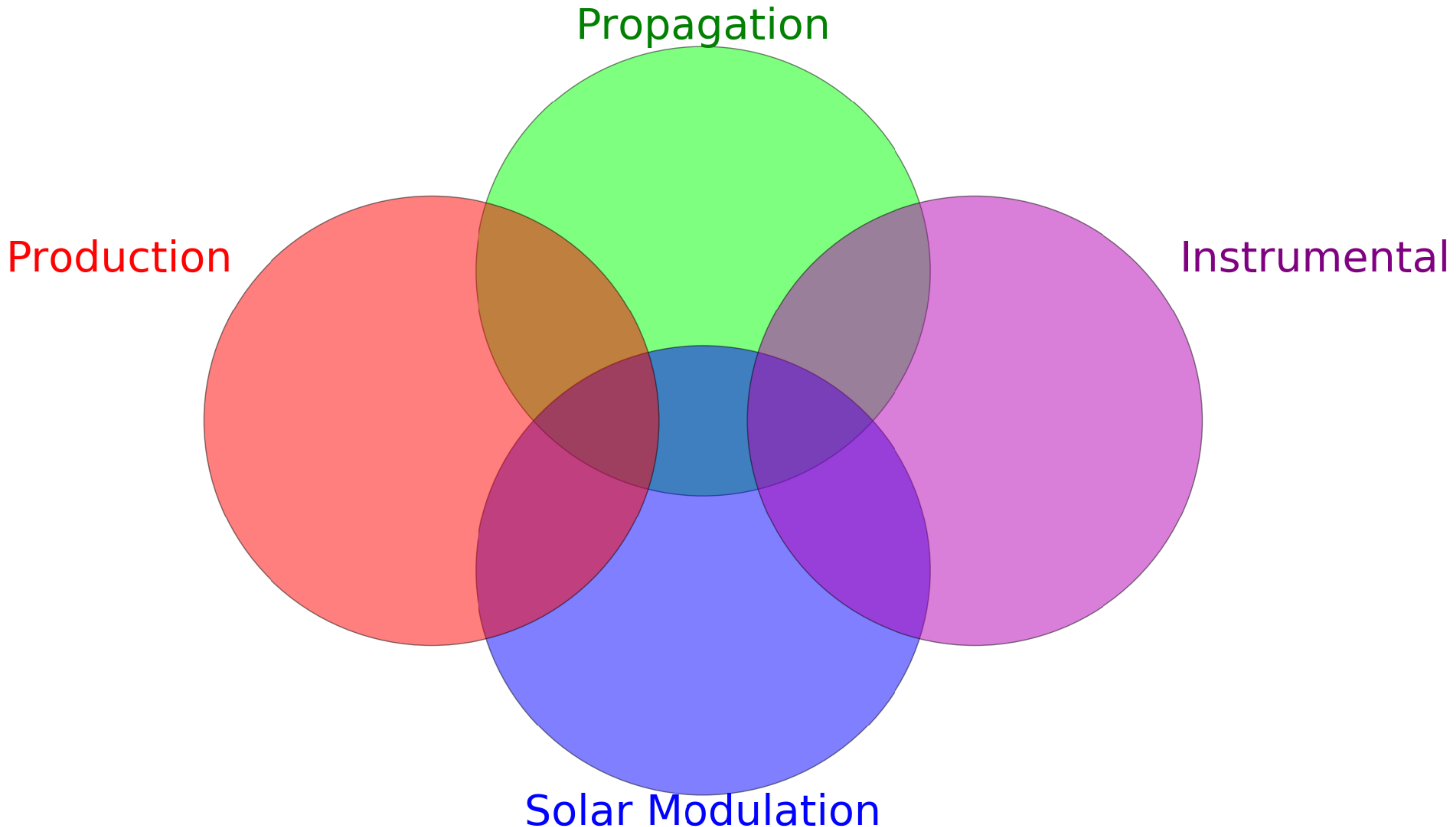
Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties







# The Antiproton Excess

Winkler (2017; 1701.04866)

Reinert, Winkler (2018; 1712.00002)

With great precision comes great responsibility:

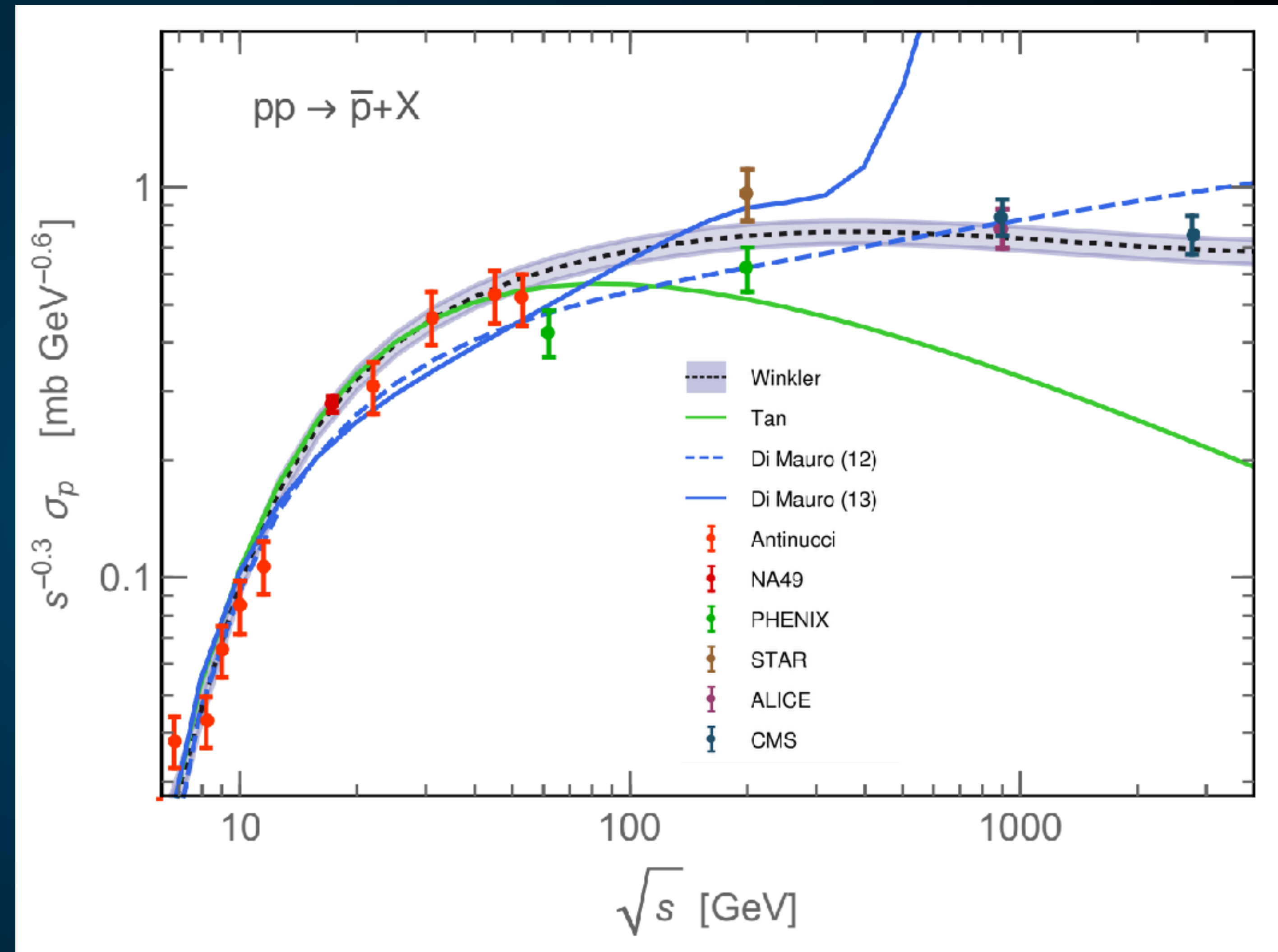
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties



# The Antiproton Excess

With great precision comes great responsibility:

Antiproton Production Cross-Section

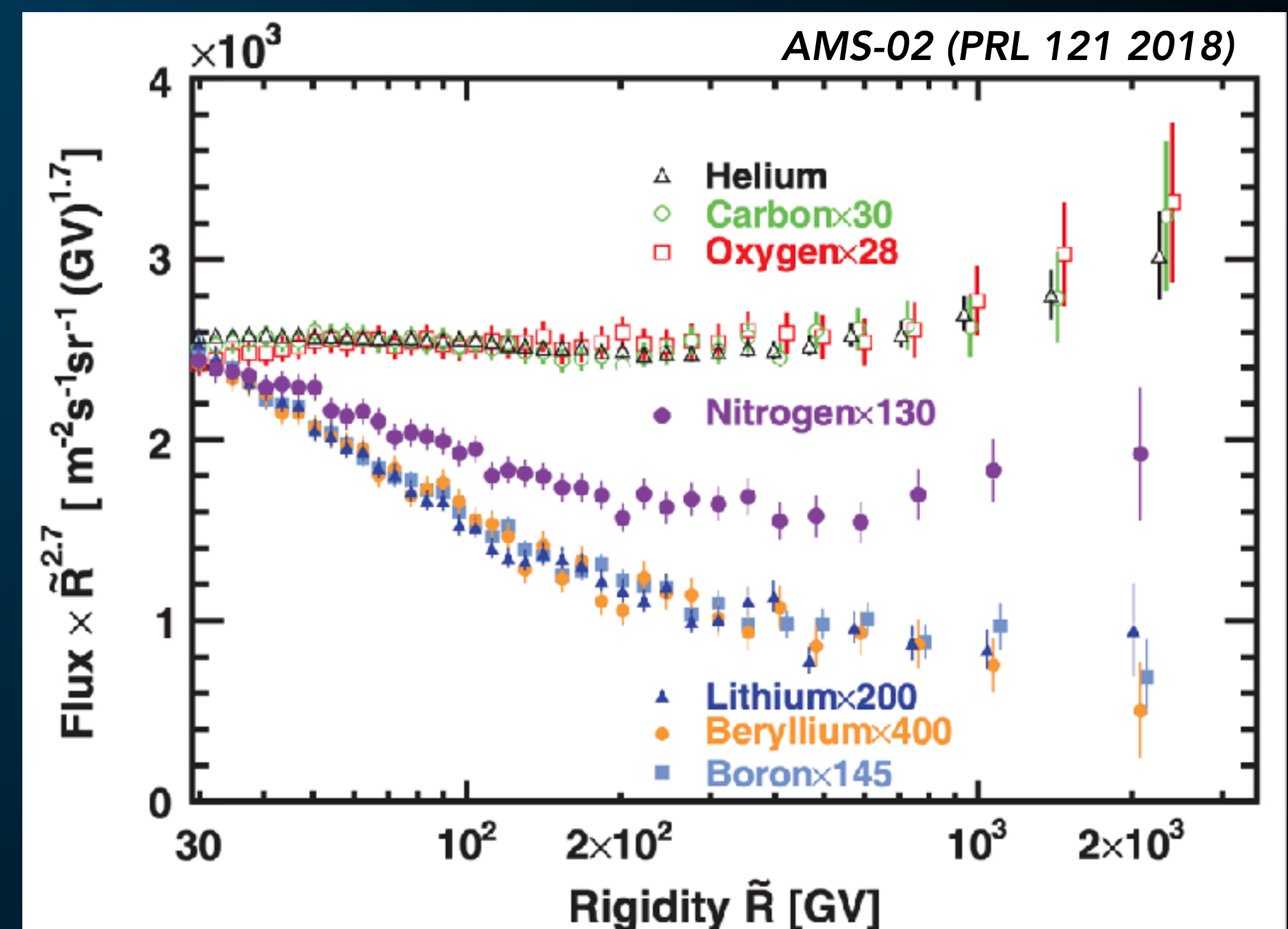
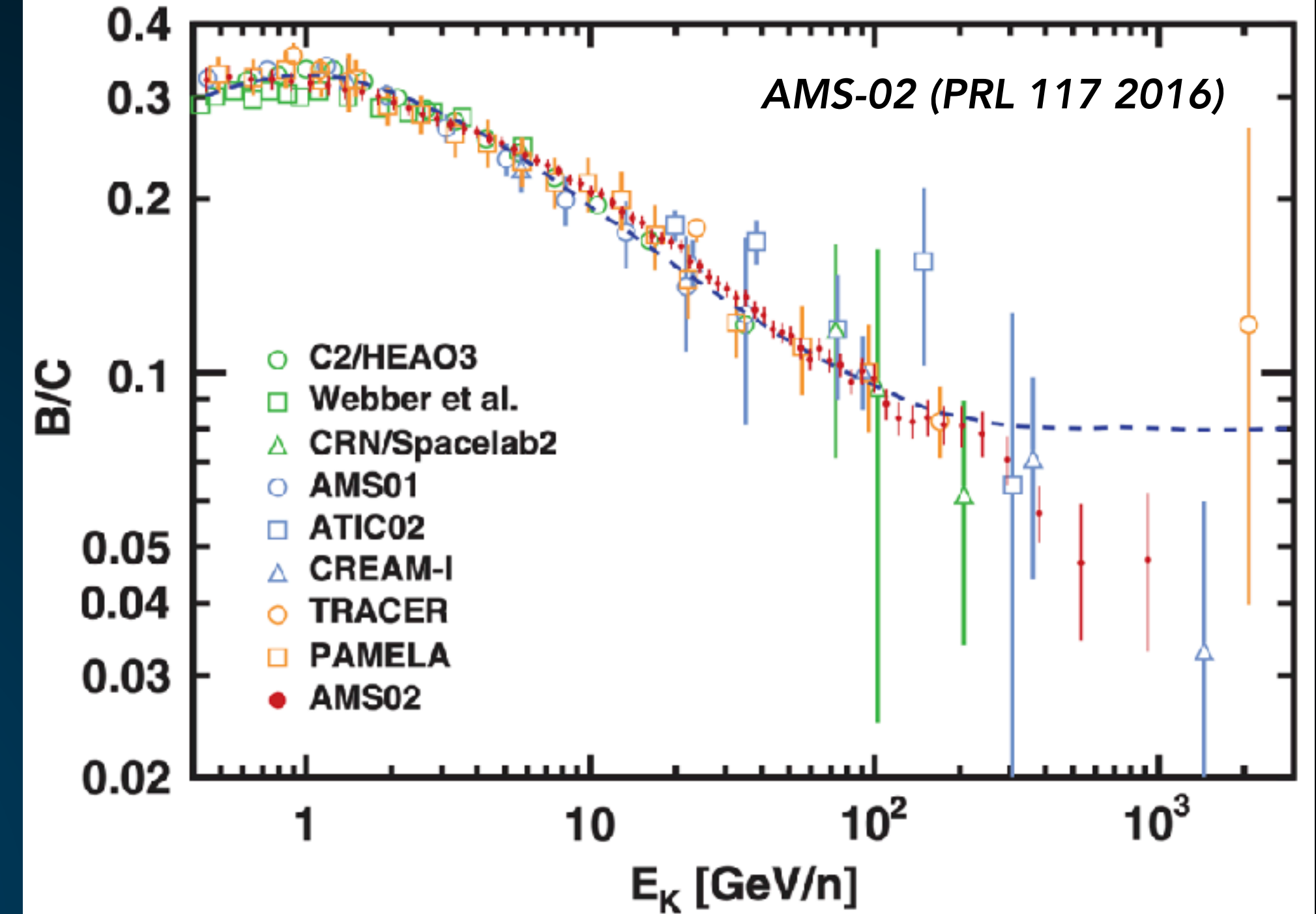
Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties

See e.g., Weinrich et al. (2002; 2002.11406)



# The Antiproton Excess

With great precision comes great responsibility:

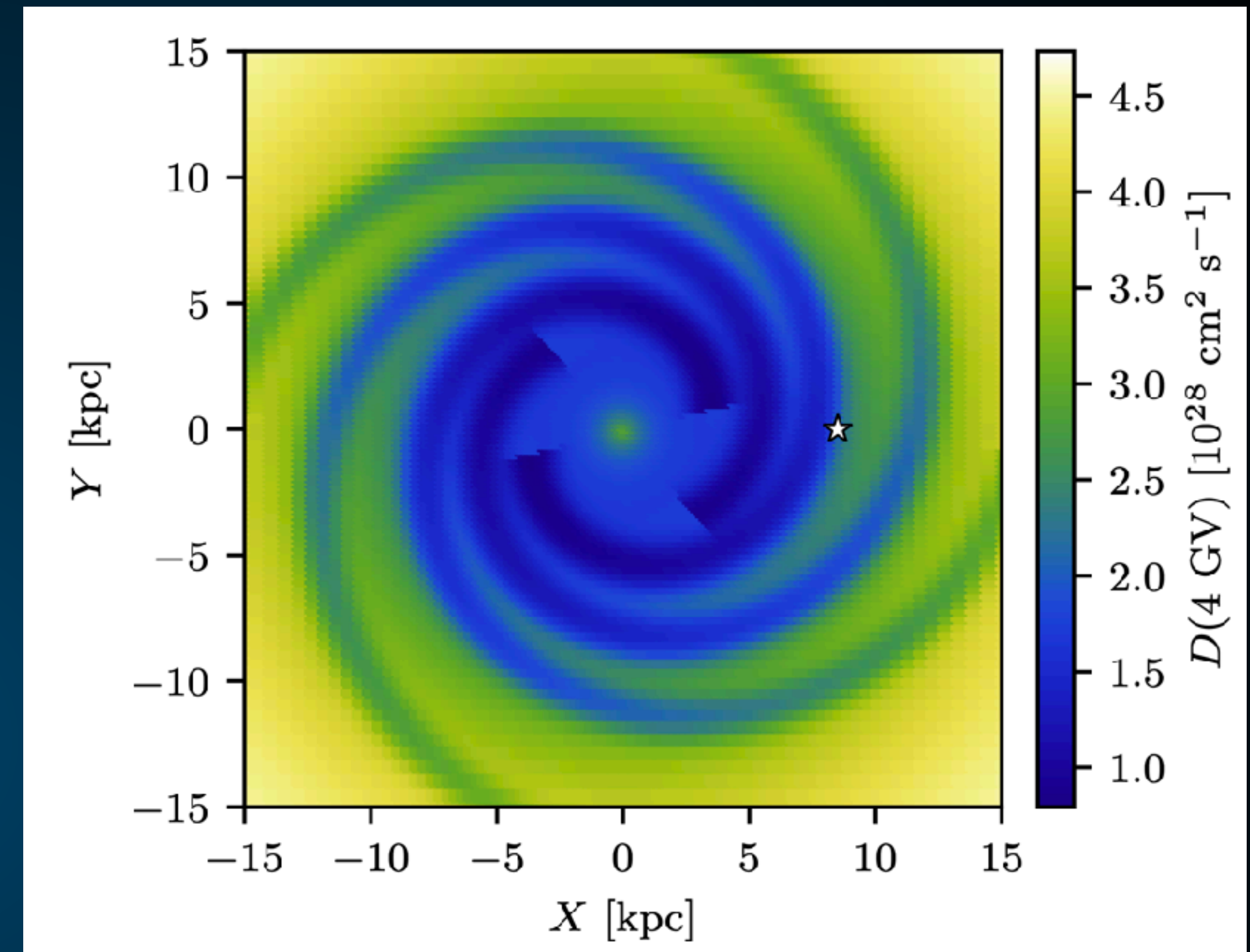
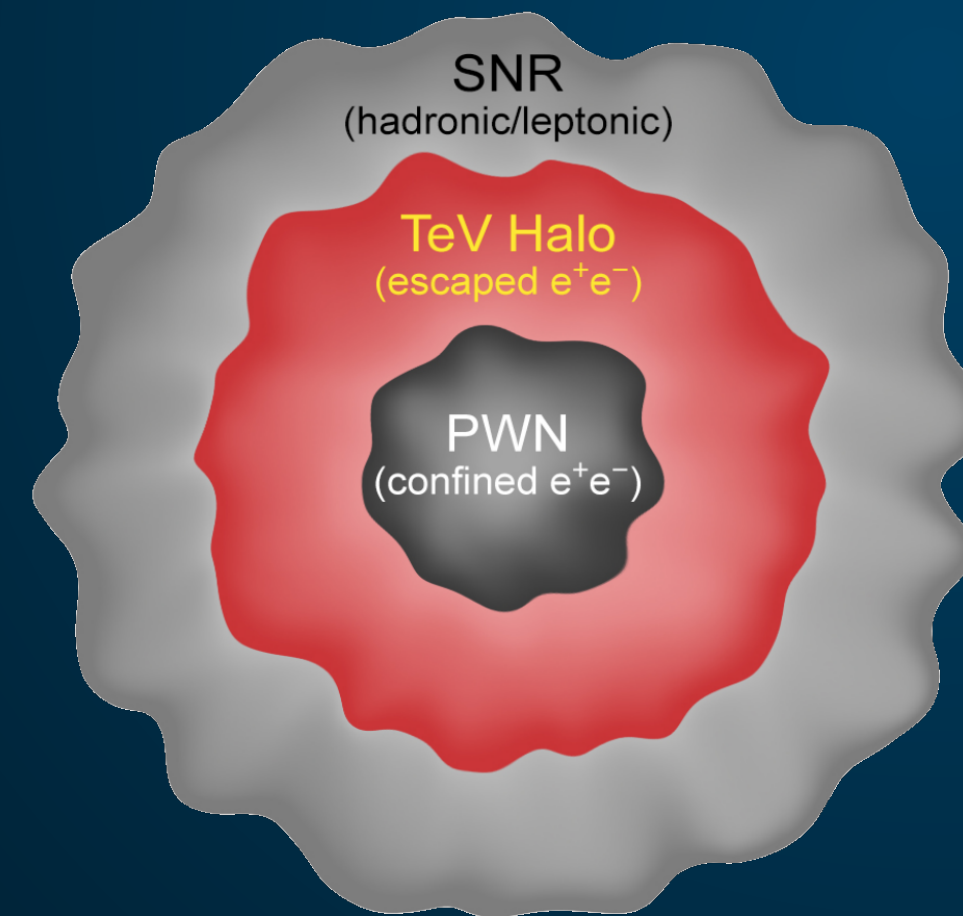
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

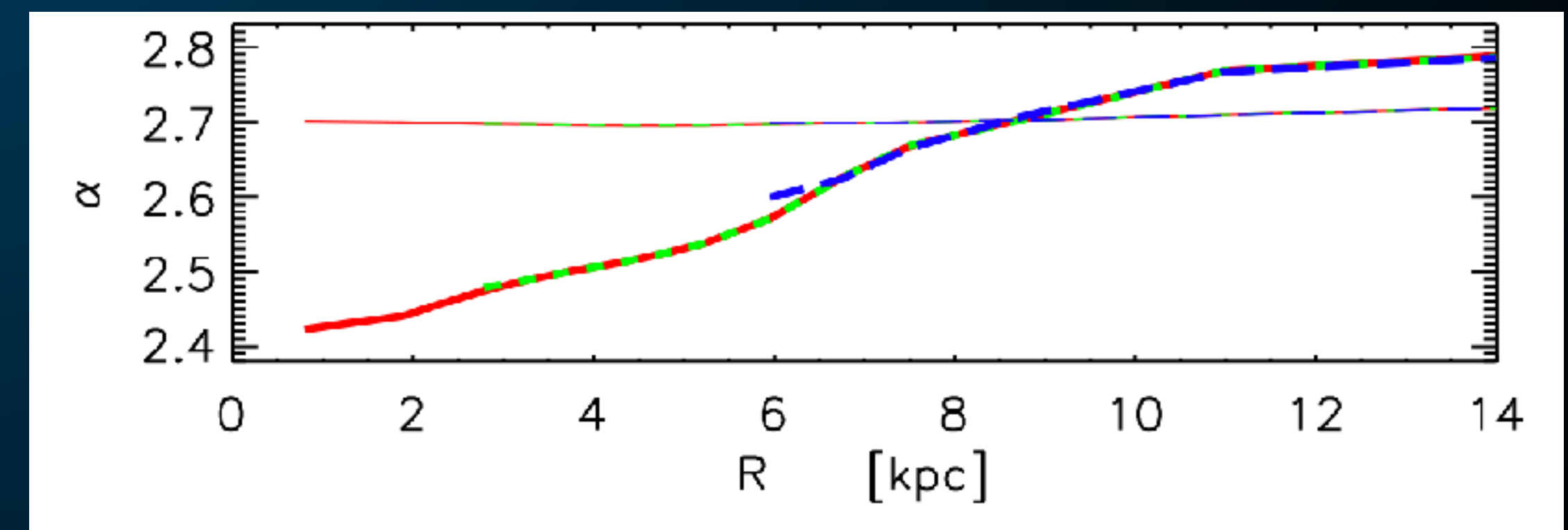
Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties



Evoli et al. (2014; 1411.7623)



# The Antiproton Excess

With great precision comes great responsibility:

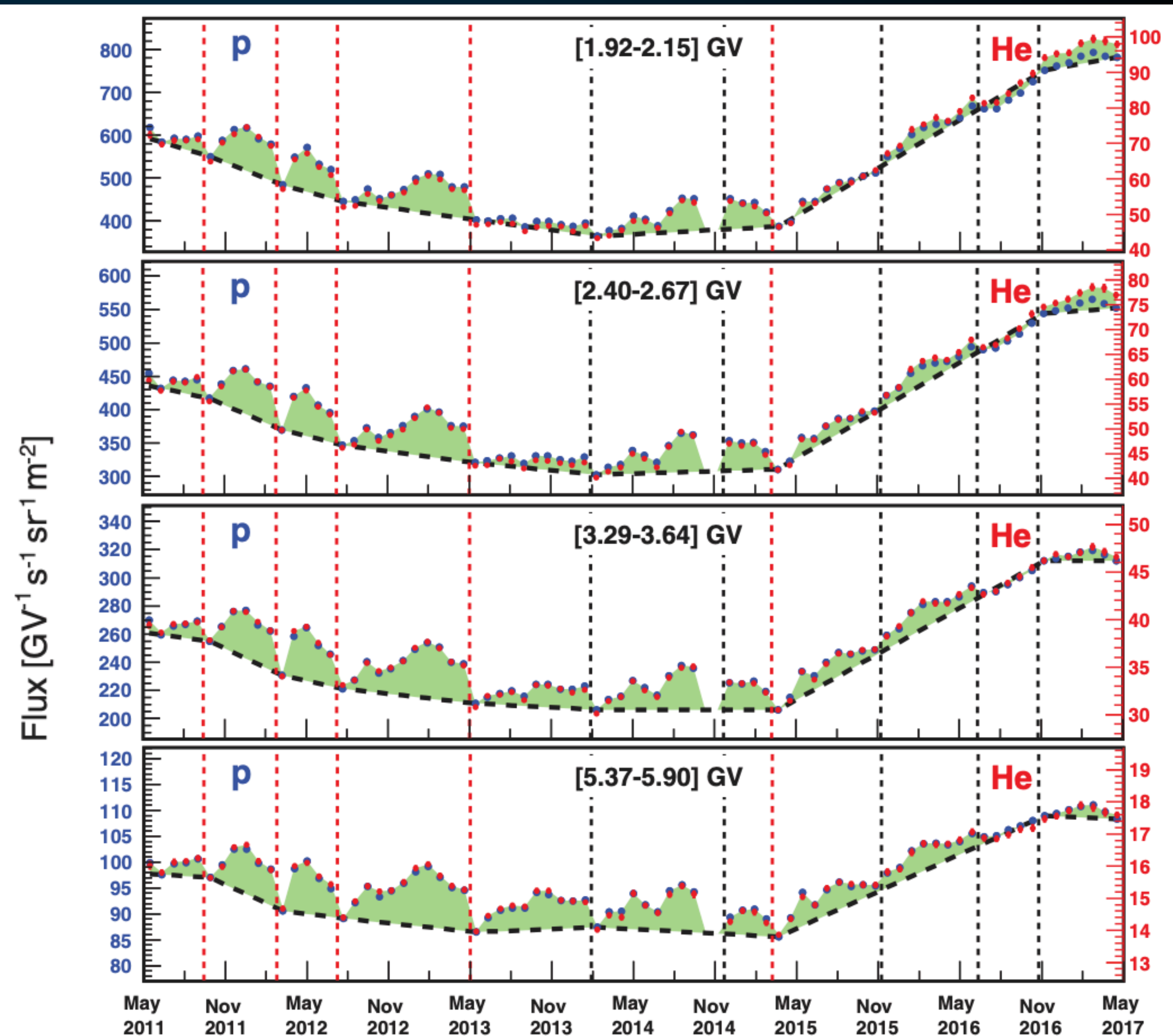
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties



# The Antiproton Excess

With great precision comes great responsibility:

Antiproton Production Cross-Section

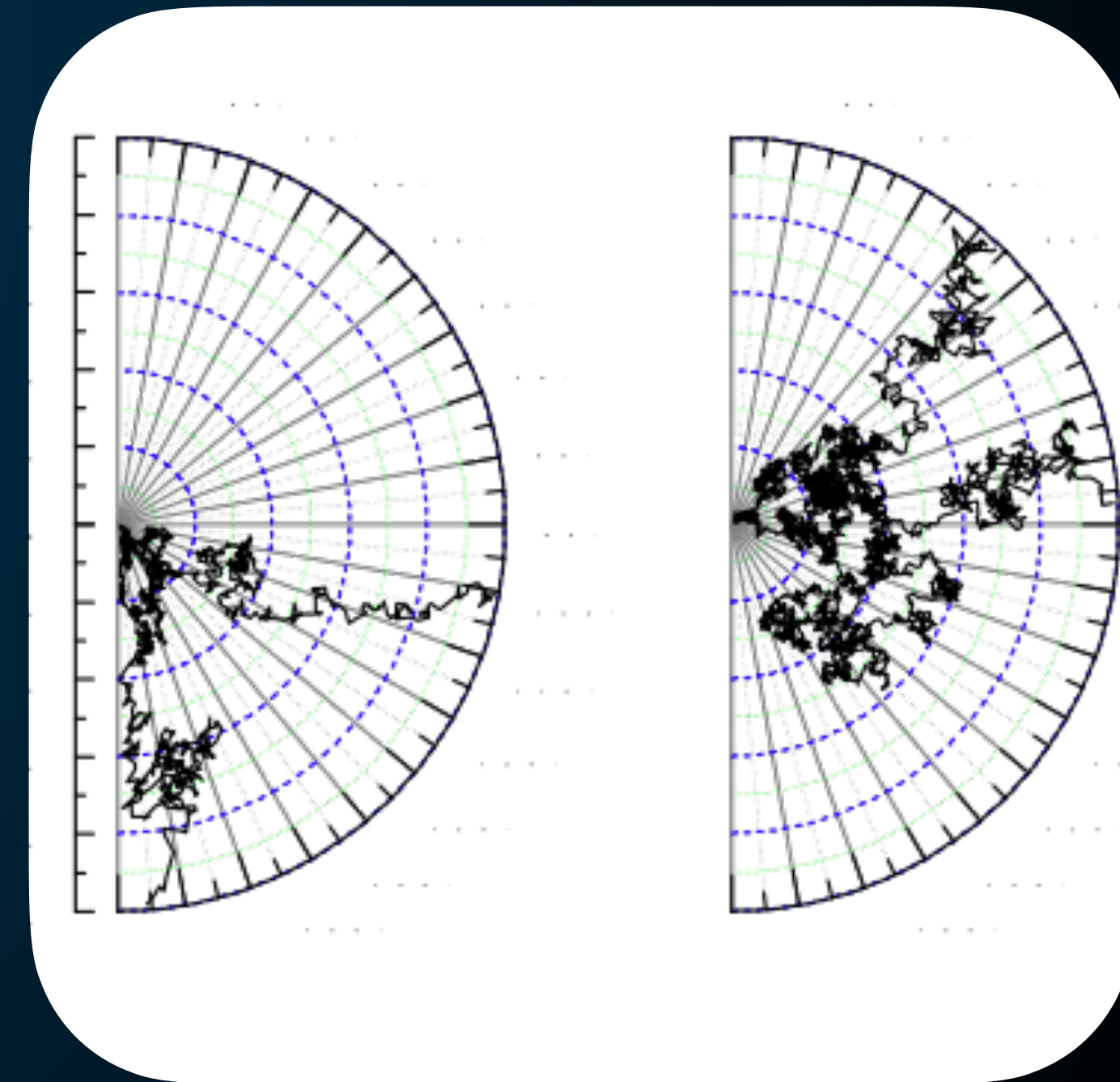
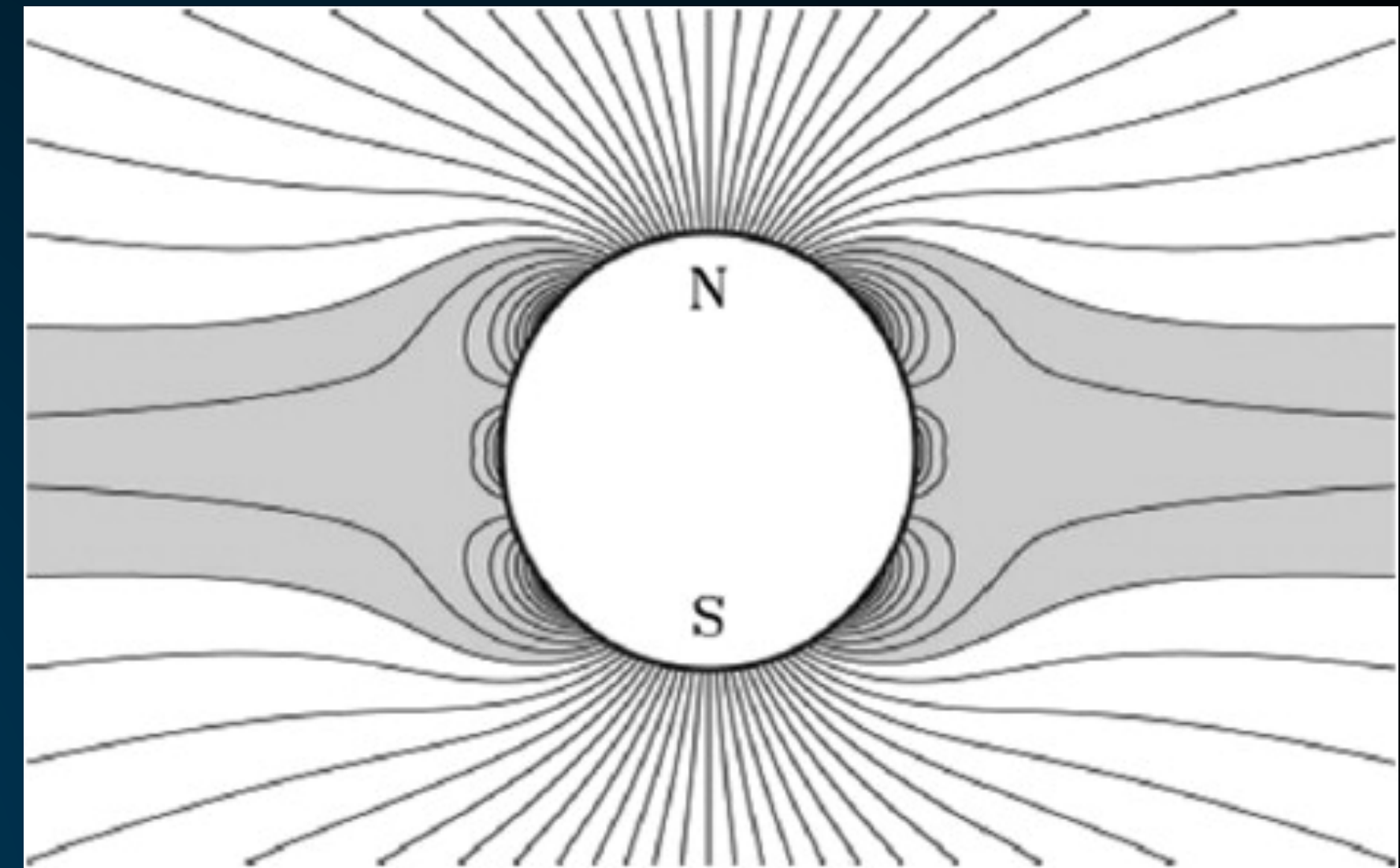
Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties

*Fisk Potential*



*HELMOD Collaboration (2011, 1110.4315)*

# The Antiproton Excess

With great precision comes great responsibility:

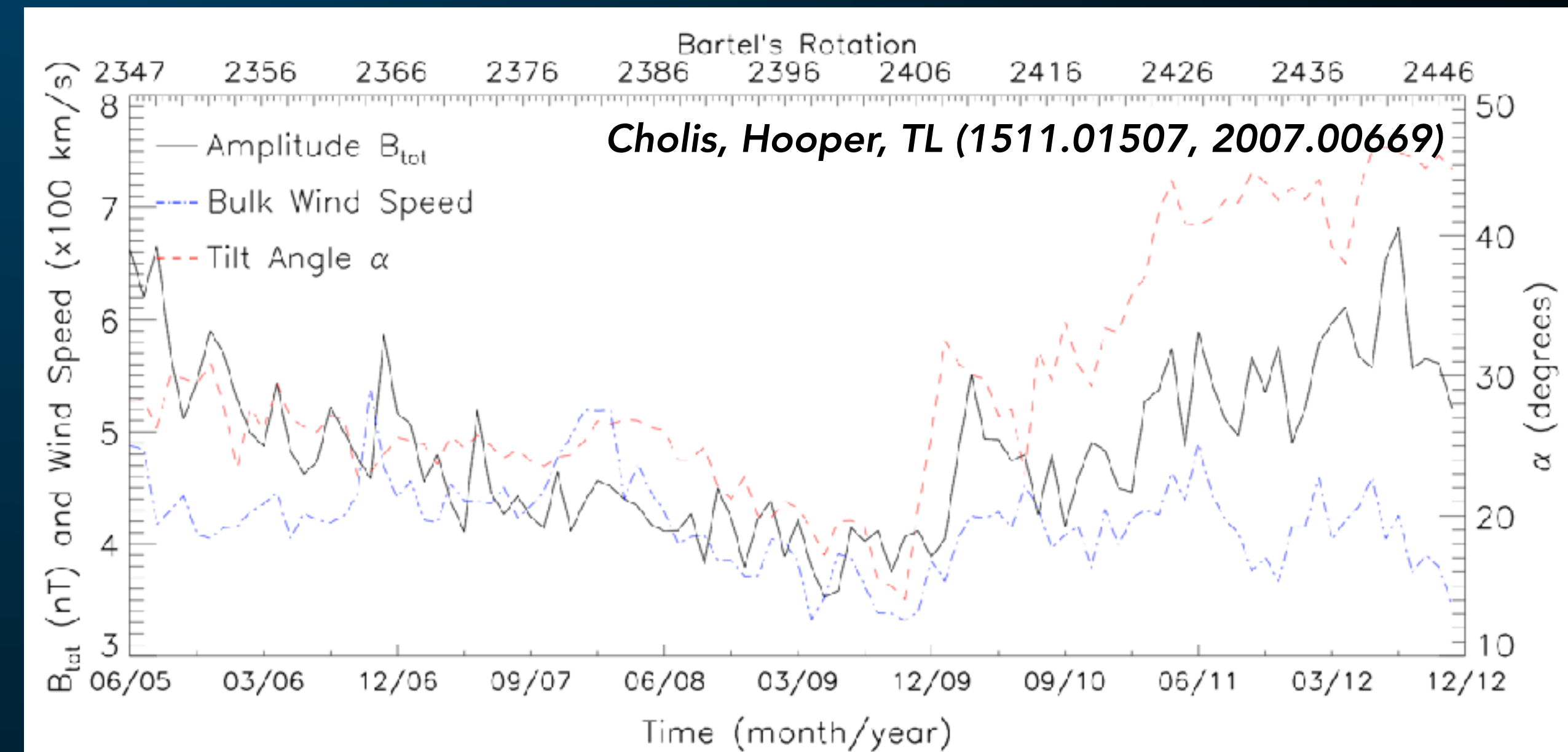
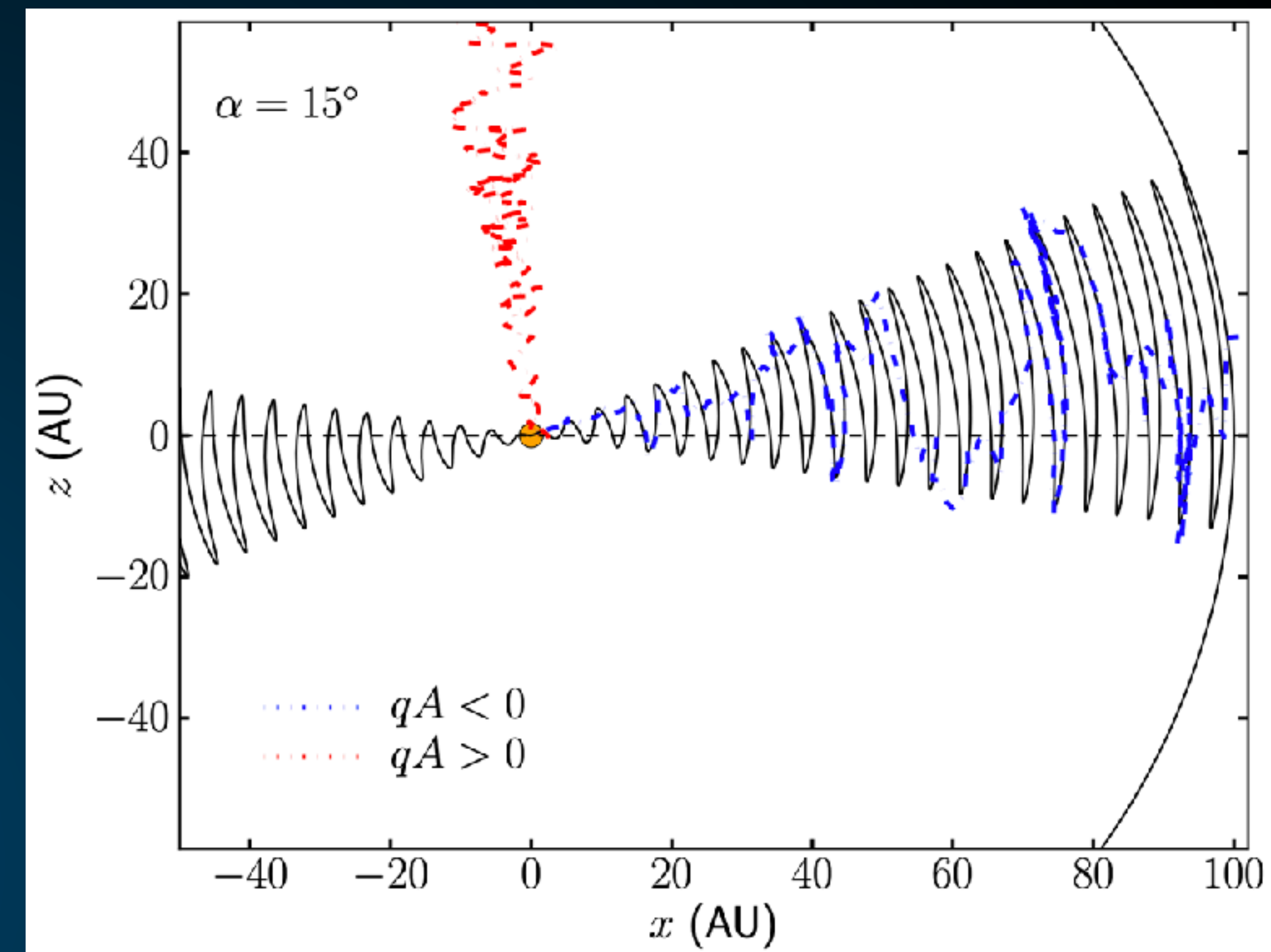
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties



# The Antiproton Excess

Cholis, Hooper, TL (2007.00669)

With great precision comes great responsibility:

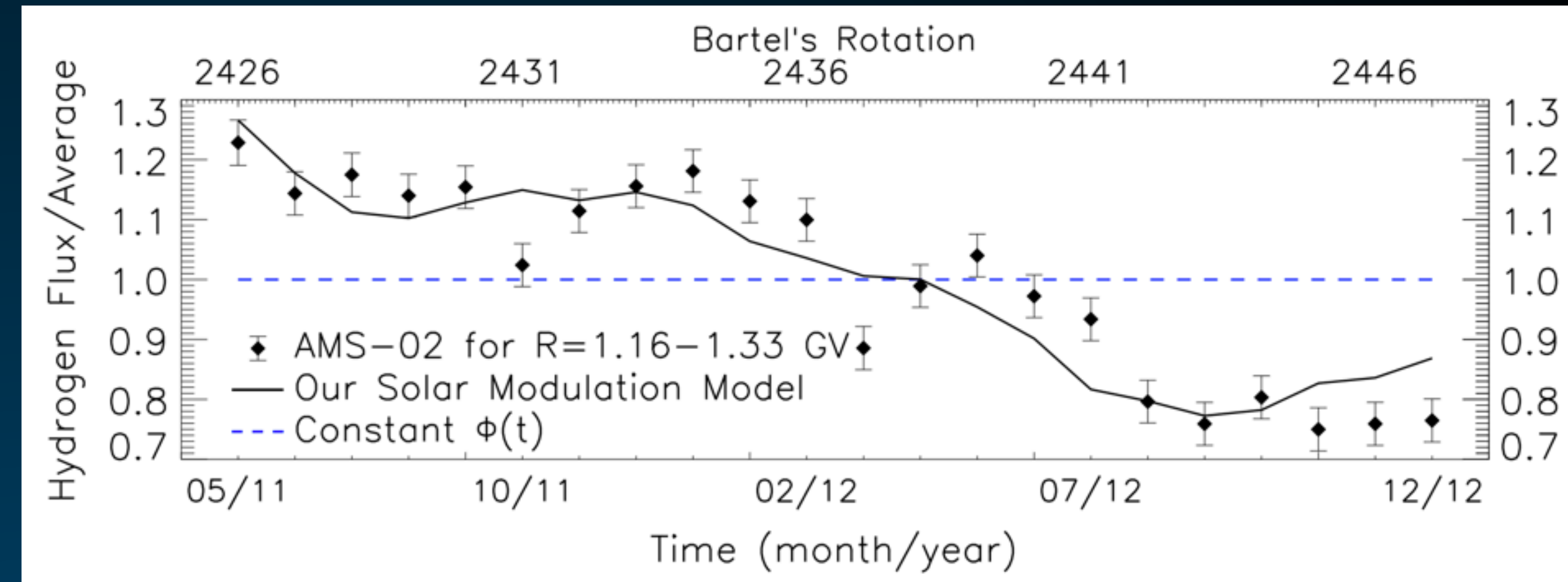
Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

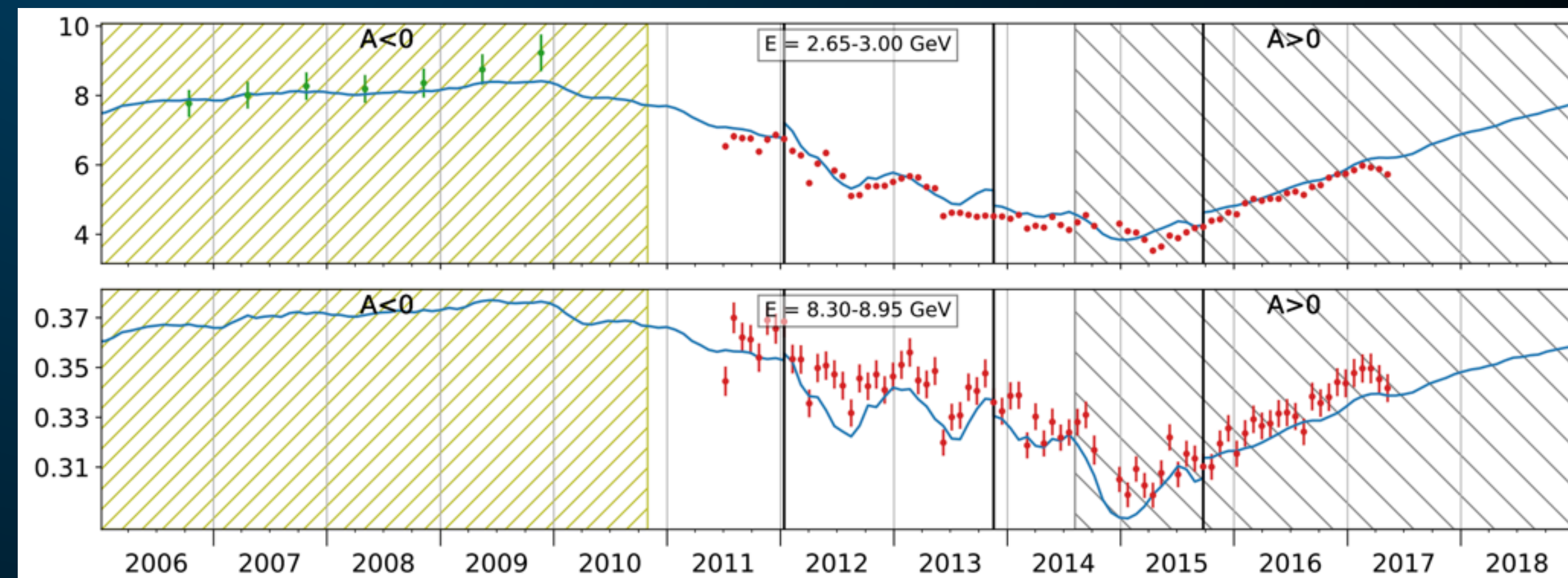
Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties



Kuhlen, Mertsch (1909.01154)





# The Antiproton Excess

With great precision comes great responsibility:

Antiproton Production Cross-Section

Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

Instrumental Uncertainties

| Rigidity [GV] | $\tilde{N}^{\bar{p}}$ | $\Phi^{\bar{p}}$ | $\sigma_{\text{stat}}$ | $\sigma_{\text{syst}}$ | $\Phi^{\bar{p}}/\Phi^p$ | $\sigma_{\text{stat}}$ | $\sigma_{\text{syst}}$ |       |                  |
|---------------|-----------------------|------------------|------------------------|------------------------|-------------------------|------------------------|------------------------|-------|------------------|
| 1.00 – 1.16   | 21                    | (5.94            | 1.31                   | 0.58)                  | $\times 10^{-3}$        | (1.02                  | 0.23                   | 0.08) | $\times 10^{-5}$ |
| 1.16 – 1.33   | 74                    | (5.57            | 0.68                   | 0.51)                  | $\times 10^{-3}$        | (8.93                  | 1.09                   | 0.66) | $\times 10^{-6}$ |
| 1.33 – 1.51   | 233                   | (9.75            | 0.68                   | 0.68)                  | $\times 10^{-3}$        | (1.59                  | 0.11                   | 0.09) | $\times 10^{-5}$ |
| 1.51 – 1.71   | 502                   | (1.06            | 0.05                   | 0.07)                  | $\times 10^{-2}$        | (1.83                  | 0.09                   | 0.09) | $\times 10^{-5}$ |
| 1.71 – 1.92   | 888                   | (1.25            | 0.05                   | 0.08)                  | $\times 10^{-2}$        | (2.33                  | 0.10                   | 0.12) | $\times 10^{-5}$ |
| 1.92 – 2.15   | 1449                  | (1.40            | 0.05                   | 0.08)                  | $\times 10^{-2}$        | (2.90                  | 0.10                   | 0.14) | $\times 10^{-5}$ |
| 2.15 – 2.40   | 2192                  | (1.50            | 0.05                   | 0.09)                  | $\times 10^{-2}$        | (3.50                  | 0.11                   | 0.17) | $\times 10^{-5}$ |
| 2.40 – 2.67   | 3366                  | (1.64            | 0.04                   | 0.09)                  | $\times 10^{-2}$        | (4.36                  | 0.11                   | 0.20) | $\times 10^{-5}$ |
| 2.67 – 2.97   | 4474                  | (1.64            | 0.04                   | 0.09)                  | $\times 10^{-2}$        | (5.05                  | 0.12                   | 0.23) | $\times 10^{-5}$ |
| 2.97 – 3.29   | 6028                  | (1.69            | 0.04                   | 0.09)                  | $\times 10^{-2}$        | (6.07                  | 0.13                   | 0.27) | $\times 10^{-5}$ |
| 3.29 – 3.64   | 7321                  | (1.67            | 0.03                   | 0.09)                  | $\times 10^{-2}$        | (7.05                  | 0.14                   | 0.30) | $\times 10^{-5}$ |
| 3.64 – 4.02   | 8592                  | (1.59            | 0.03                   | 0.08)                  | $\times 10^{-2}$        | (7.96                  | 0.15                   | 0.32) | $\times 10^{-5}$ |
| 4.02 – 4.43   | 1932                  | (1.56            | 0.04                   | 0.08)                  | $\times 10^{-2}$        | (9.31                  | 0.21                   | 0.37) | $\times 10^{-5}$ |
| 4.43 – 4.88   | 3083                  | (1.43            | 0.03                   | 0.07)                  | $\times 10^{-2}$        | (1.03                  | 0.02                   | 0.04) | $\times 10^{-4}$ |
| 4.88 – 5.37   | 3880                  | (1.23            | 0.02                   | 0.06)                  | $\times 10^{-2}$        | (1.07                  | 0.02                   | 0.04) | $\times 10^{-4}$ |
| 5.37 – 5.90   | 4780                  | (1.12            | 0.02                   | 0.05)                  | $\times 10^{-2}$        | (1.19                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 5.90 – 6.47   | 5472                  | (9.80            | 0.13                   | 0.45)                  | $\times 10^{-3}$        | (1.27                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 6.47 – 7.09   | 6538                  | (8.69            | 0.11                   | 0.39)                  | $\times 10^{-3}$        | (1.38                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 7.09 – 7.76   | 7369                  | (7.59            | 0.09                   | 0.34)                  | $\times 10^{-3}$        | (1.49                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 7.76 – 8.48   | 7818                  | (6.54            | 0.08                   | 0.29)                  | $\times 10^{-3}$        | (1.59                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 8.48 – 9.26   | 7821                  | (5.46            | 0.06                   | 0.24)                  | $\times 10^{-3}$        | (1.64                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 9.26 – 10.1   | 20382                 | (4.67            | 0.03                   | 0.20)                  | $\times 10^{-3}$        | (1.74                  | 0.01                   | 0.06) | $\times 10^{-4}$ |
| 10.1 – 11.0   | 19445                 | (3.96            | 0.03                   | 0.17)                  | $\times 10^{-3}$        | (1.83                  | 0.01                   | 0.07) | $\times 10^{-4}$ |
| 11.0 – 12.0   | 18769                 | (3.23            | 0.02                   | 0.14)                  | $\times 10^{-3}$        | (1.86                  | 0.01                   | 0.07) | $\times 10^{-4}$ |
| 12.0 – 13.0   | 16372                 | (2.65            | 0.02                   | 0.11)                  | $\times 10^{-3}$        | (1.89                  | 0.02                   | 0.07) | $\times 10^{-4}$ |
| 13.0 – 14.1   | 16076                 | (2.23            | 0.02                   | 0.09)                  | $\times 10^{-3}$        | (1.96                  | 0.02                   | 0.07) | $\times 10^{-4}$ |
| 14.1 – 15.3   | 15578                 | (1.85            | 0.02                   | 0.08)                  | $\times 10^{-3}$        | (2.02                  | 0.02                   | 0.07) | $\times 10^{-4}$ |
| 15.3 – 16.6   | 14734                 | (1.49            | 0.01                   | 0.06)                  | $\times 10^{-3}$        | (2.02                  | 0.02                   | 0.07) | $\times 10^{-4}$ |
| 16.6 – 18.0   | 15816                 | (1.19            | 0.01                   | 0.05)                  | $\times 10^{-3}$        | (2.00                  | 0.02                   | 0.07) | $\times 10^{-4}$ |
| 18.0 – 19.5   | 15049                 | (9.53            | 0.08                   | 0.37)                  | $\times 10^{-4}$        | (1.99                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 19.5 – 21.1   | 14426                 | (7.72            | 0.07                   | 0.29)                  | $\times 10^{-4}$        | (1.99                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 21.1 – 22.8   | 13511                 | (6.33            | 0.06                   | 0.23)                  | $\times 10^{-4}$        | (2.02                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 22.8 – 24.7   | 12943                 | (5.02            | 0.05                   | 0.18)                  | $\times 10^{-4}$        | (1.99                  | 0.02                   | 0.06) | $\times 10^{-4}$ |
| 24.7 – 26.7   | 11723                 | (4.11            | 0.04                   | 0.14)                  | $\times 10^{-4}$        | (2.02                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 26.7 – 28.8   | 10411                 | (3.32            | 0.04                   | 0.11)                  | $\times 10^{-4}$        | (2.02                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 28.8 – 31.1   | 9508                  | (2.68            | 0.03                   | 0.08)                  | $\times 10^{-4}$        | (2.02                  | 0.02                   | 0.05) | $\times 10^{-4}$ |
| 31.1 – 33.5   | 7876                  | (2.07            | 0.03                   | 0.06)                  | $\times 10^{-4}$        | (1.92                  | 0.02                   | 0.04) | $\times 10^{-4}$ |
| 33.5 – 36.1   | 7212                  | (1.75            | 0.02                   | 0.05)                  | $\times 10^{-4}$        | (2.00                  | 0.03                   | 0.05) | $\times 10^{-4}$ |

(Table continued)

# The Antiproton Excess

With great precision comes great responsibility:

Antiproton Production Cross-Section

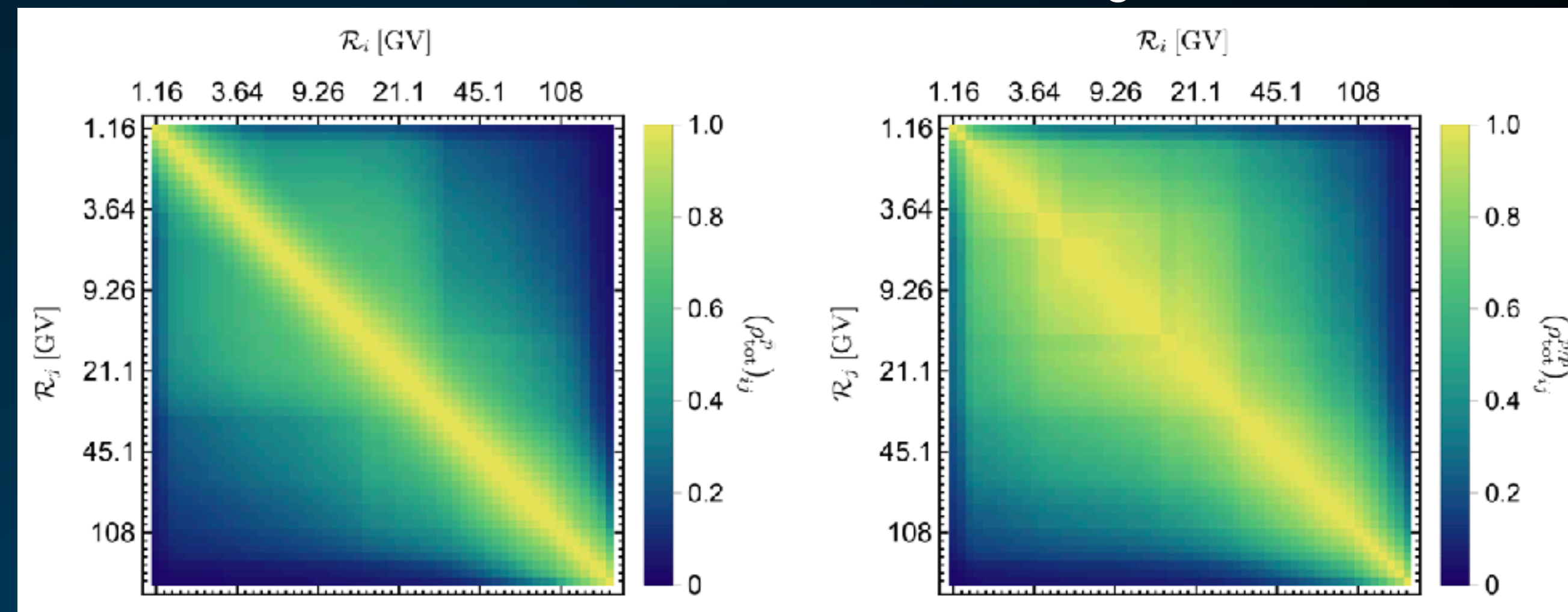
Galactic Primary to Secondary Ratios

Inhomogeneous Diffusion

Solar Modulation

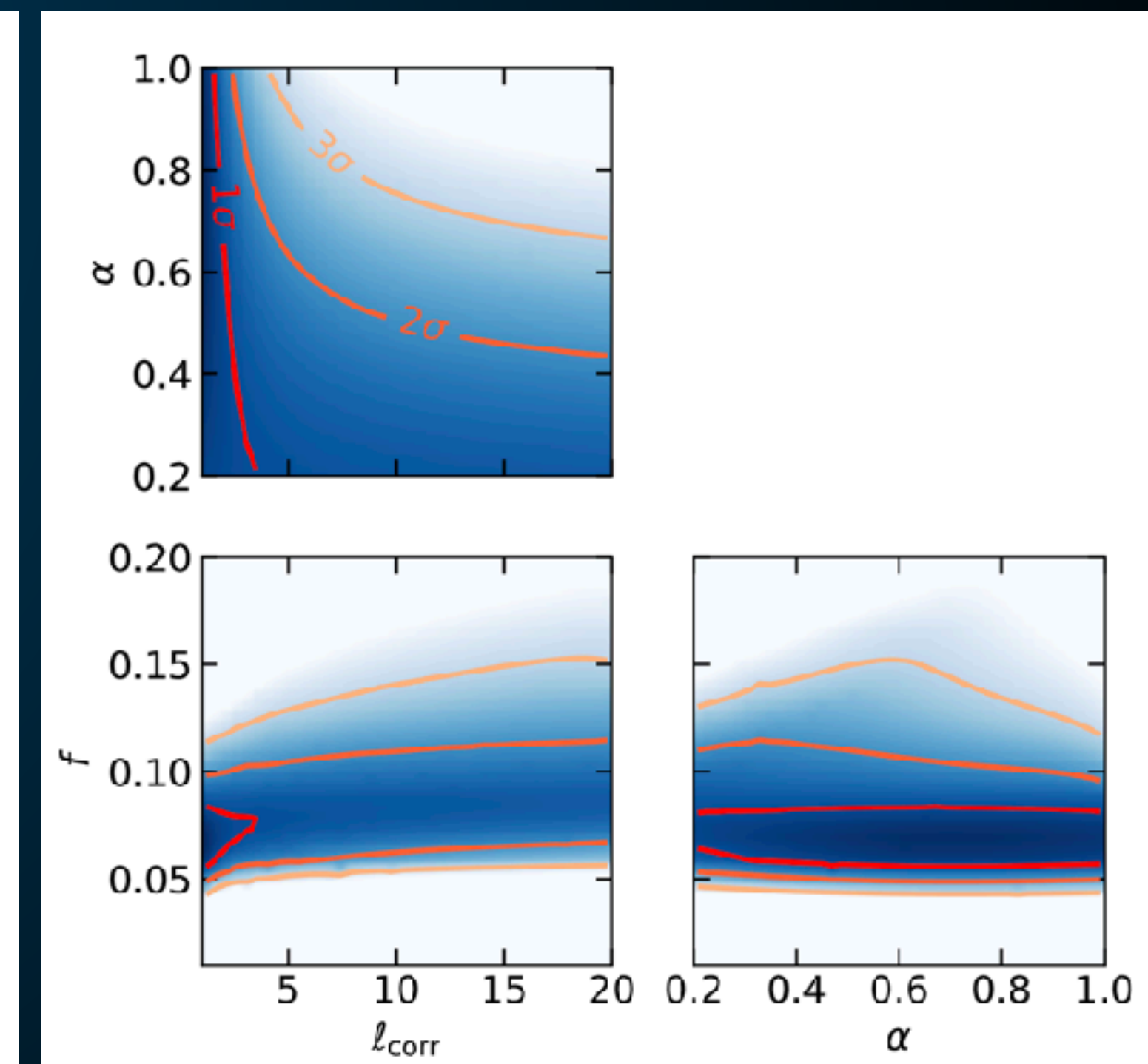
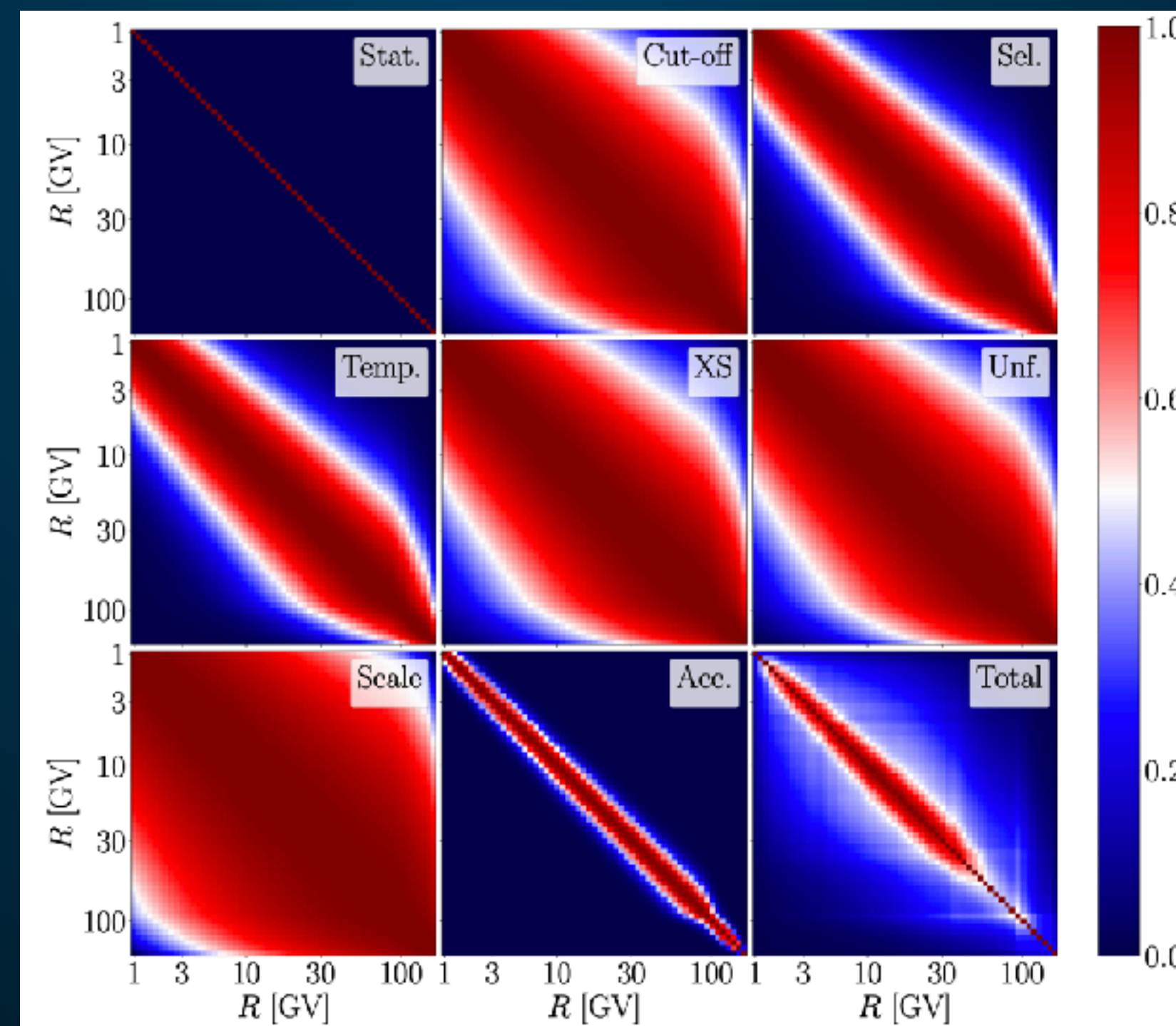
Instrumental Uncertainties

Heisig et al. (2020; 2005.04237)



Boudaud et al. (2019; 1906.07119)

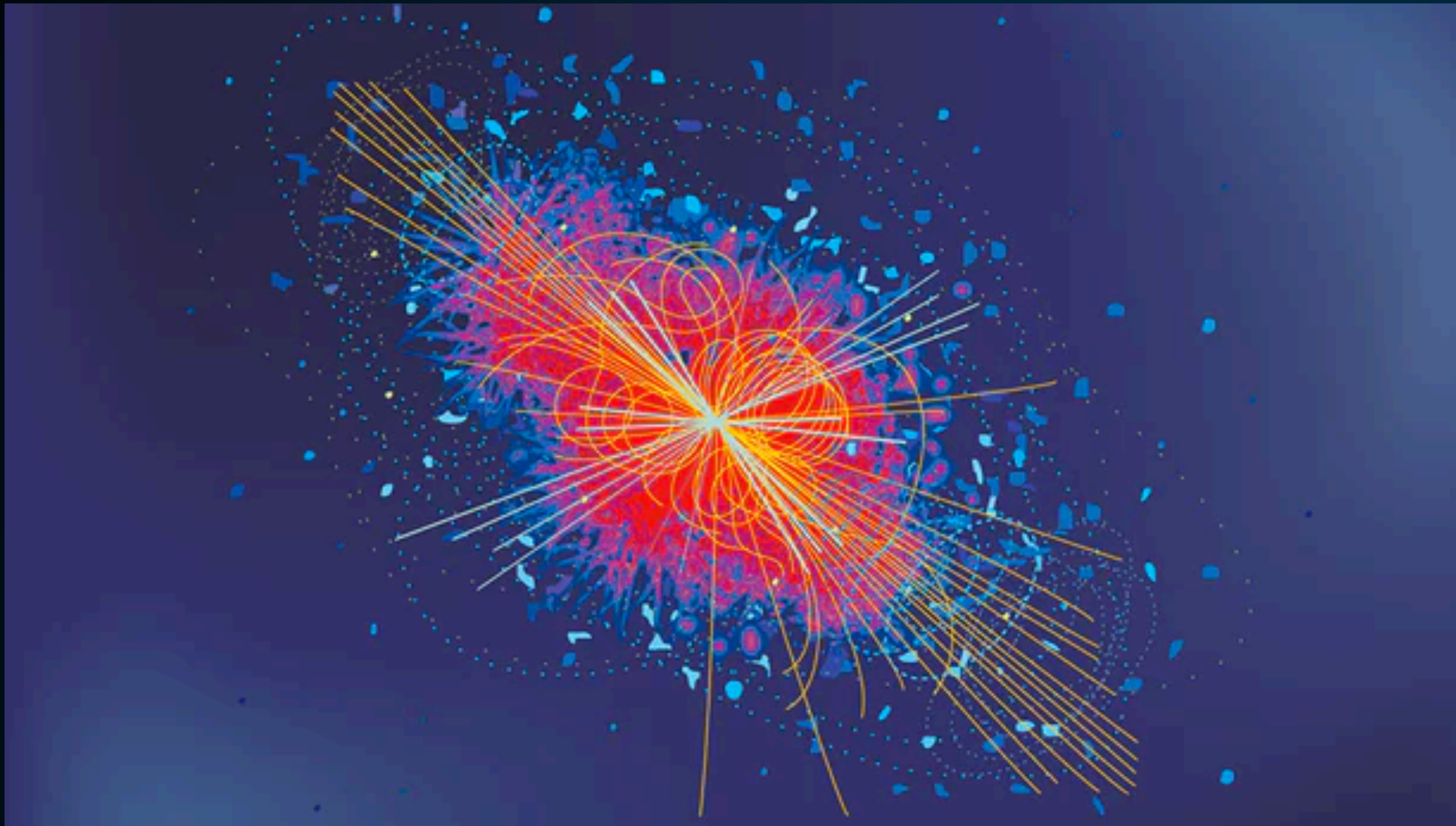
Cuoco et al. (2019; 1903.01472)



Antinuclei!?



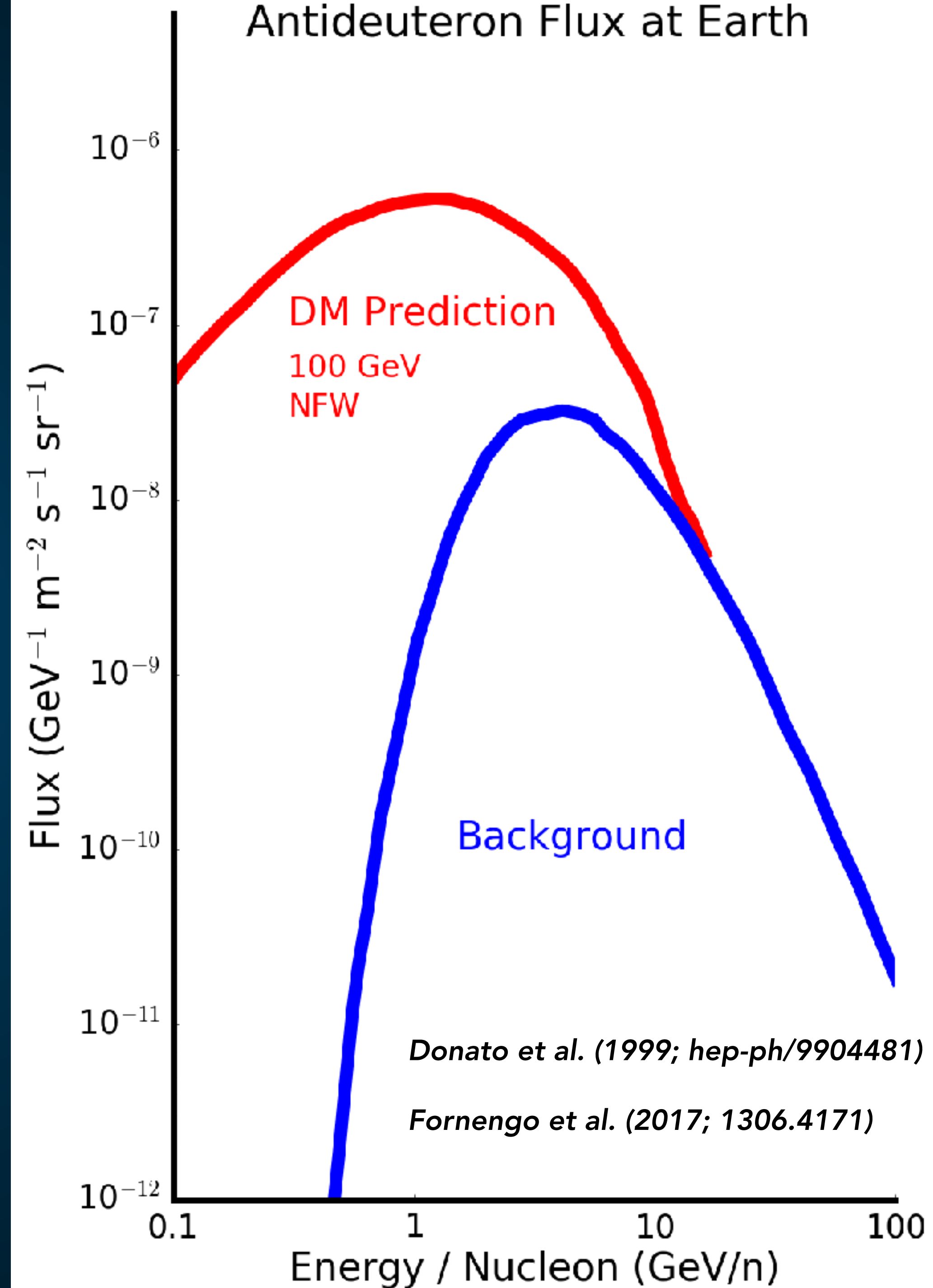
# AntiNuclei - A Clean Search Strategy ?



Antinuclei carry away a significant fraction of the total momentum in a particle collision.

**Astrophysical Antinuclei - Most be moving relativistically!**

**Dark Matter Antinuclei - Can be slow!**



To date, we have observed eight events in the mass region from 0 to 10 GeV with  $Z = -2$ . All eight events are in the helium mass region.

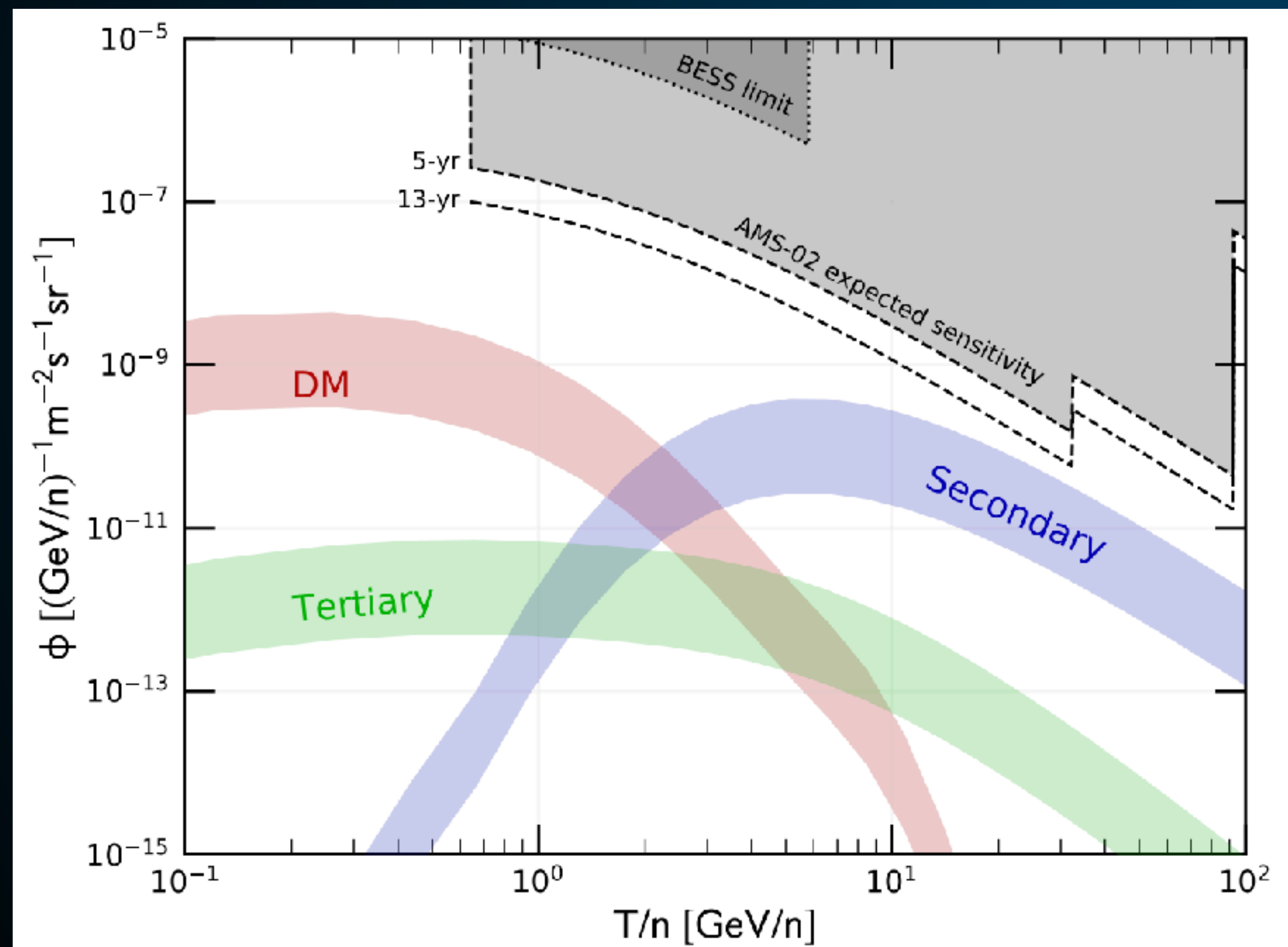
Currently (having used 50 million core hours to generate 7 times more simulated events than measured events and having found no background events from the simulation), our best evaluation of the probability of the background origin for the eight  $\bar{\text{He}}$  events is **less than  $3 \times 10^{-8}$** . For the two  ${}^4\bar{\text{He}}$  events our best evaluation of the probability (upon completion of the current 100 million core hours of simulation) will be less than  $3 \times 10^{-3}$ .

Note that for  ${}^4\bar{\text{He}}$ , projecting based on the statistics we have today, by using an additional 400 million core hours for simulation the background probability would be  $10^{-4}$ . Simultaneously, continuing to run until 2023, which doubles the data sample, the background probability for  ${}^4\bar{\text{He}}$  would be  **$2 \times 10^{-7}$** , i.e., greater than 5-sigma significance.

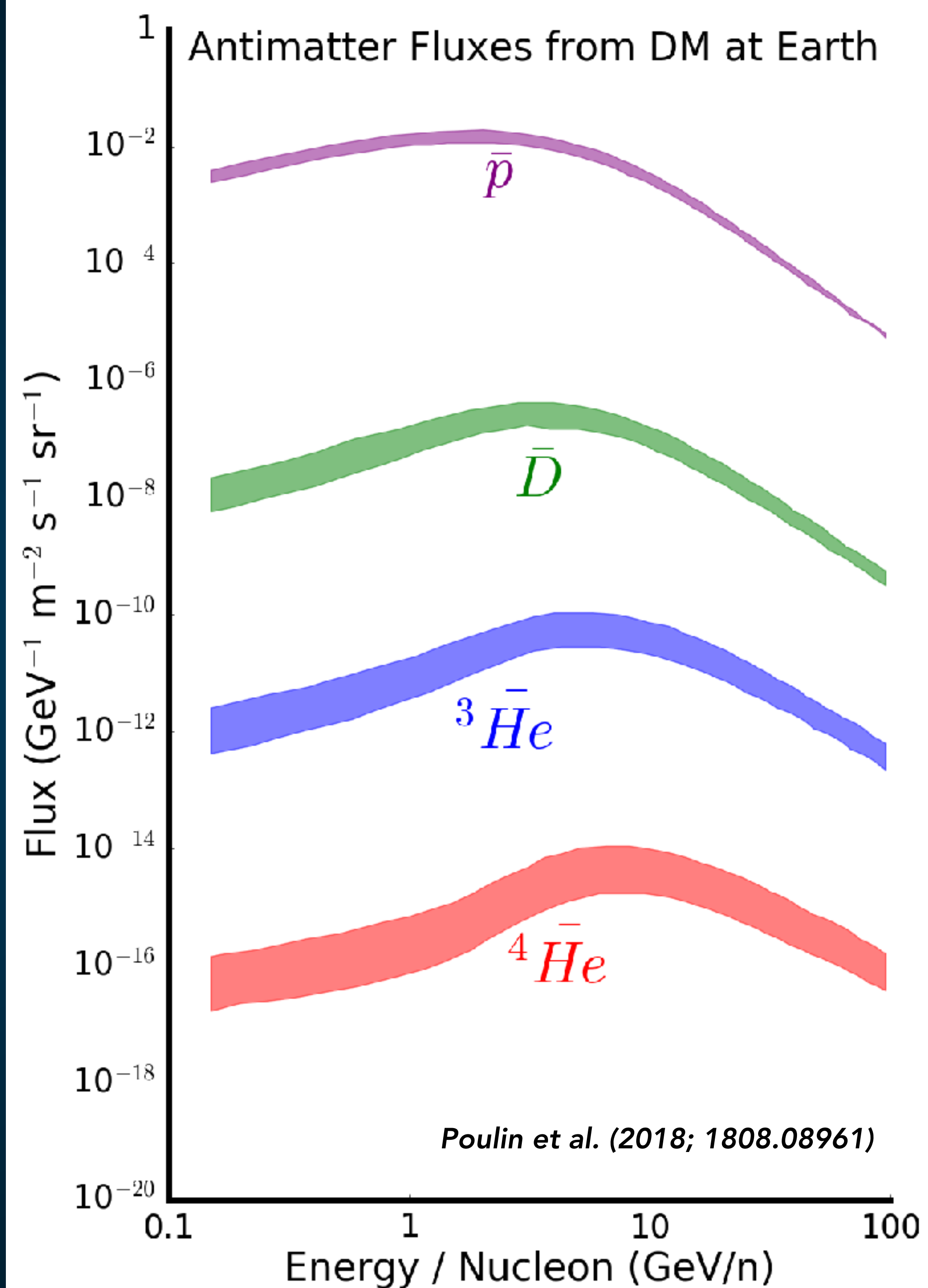
# AntiNuclei - A Clean Search Strategy ?

Antihelium background even cleaner than antideuterons

But the flux is supposed to be much smaller.



Korsmeier (2017; 1711.08465)

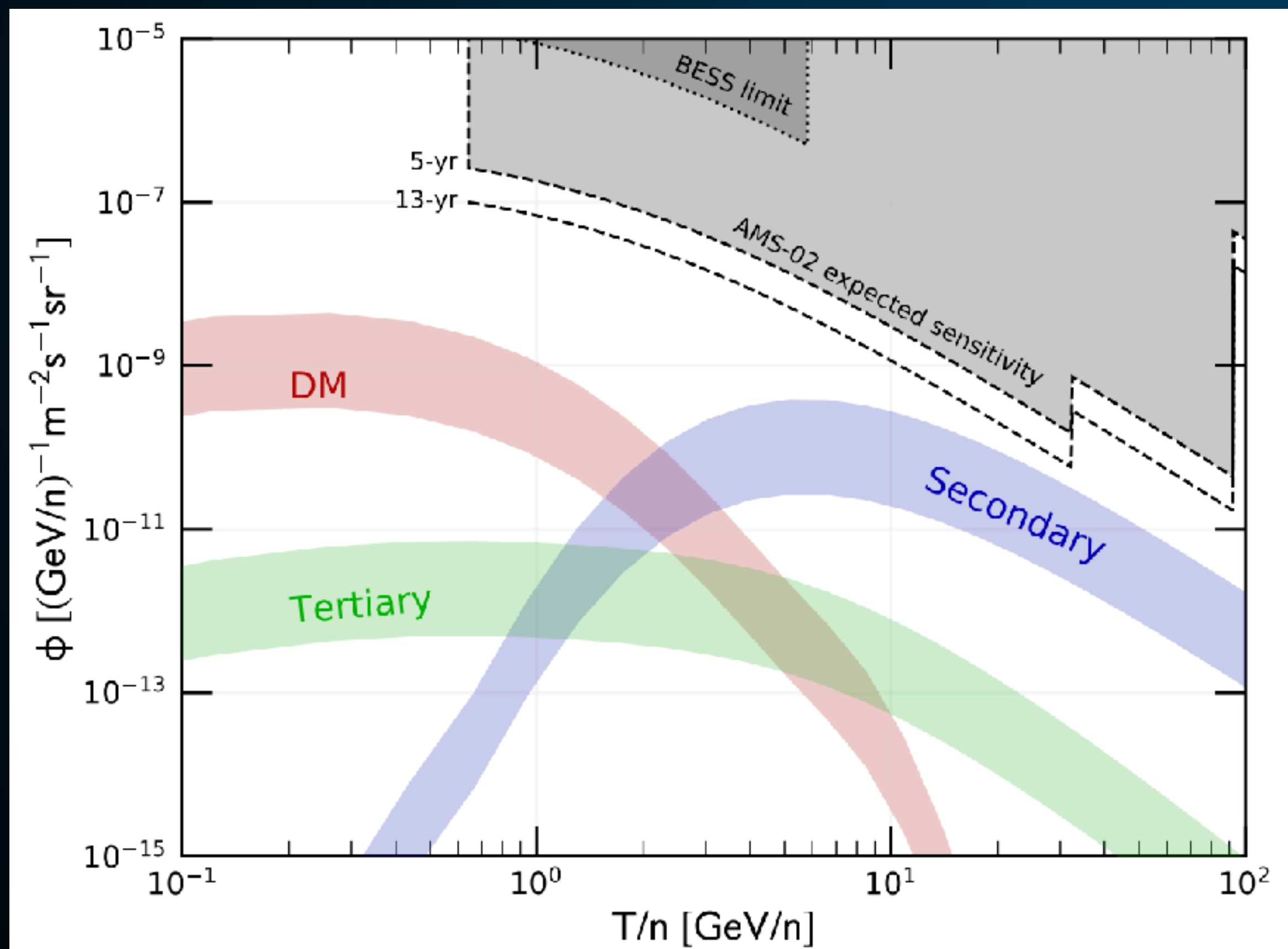


Poulin et al. (2018; 1808.08961)

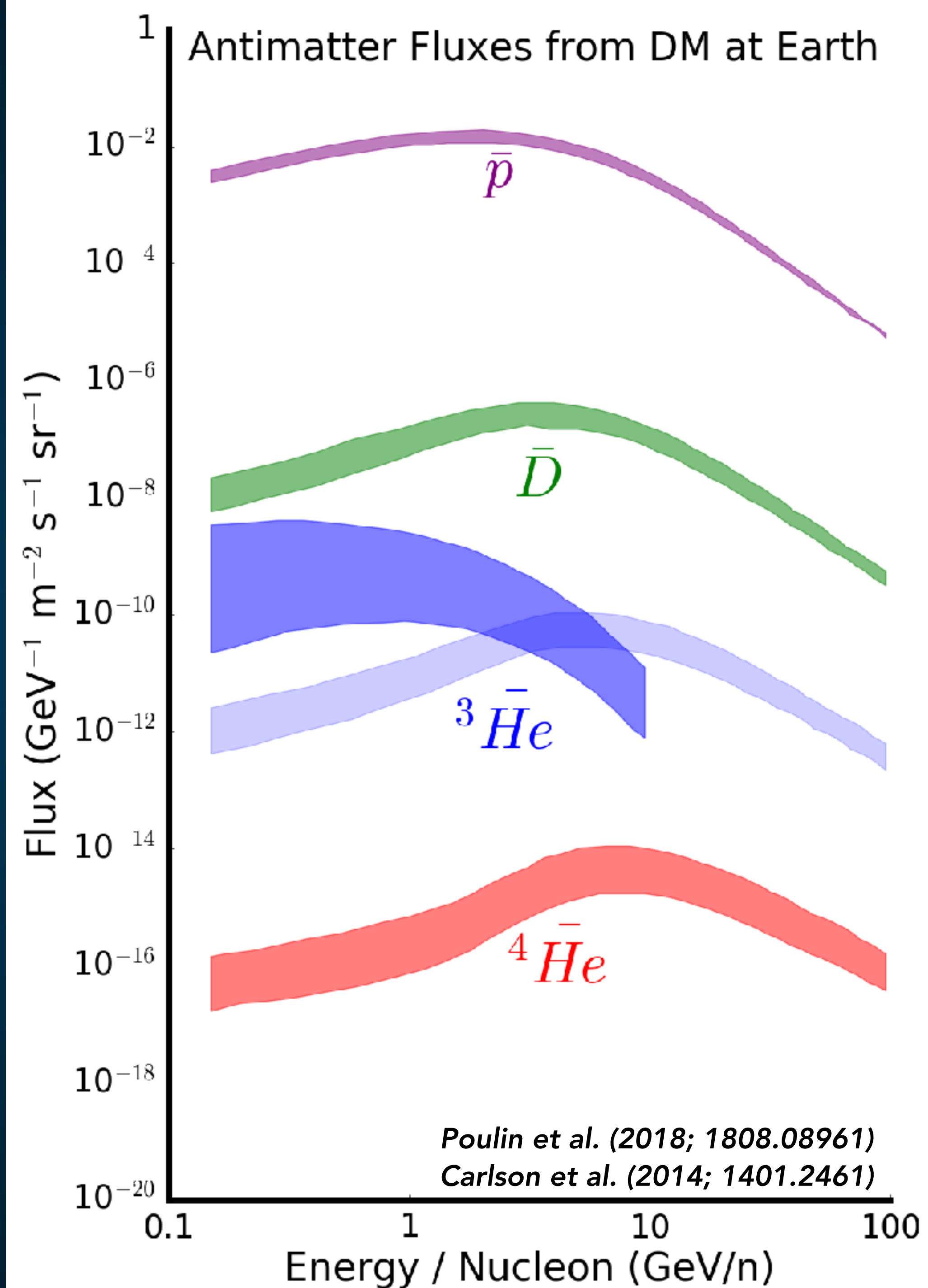
# AntiNuclei - A Clean Search Strategy ?

Antihelium background even cleaner than antideuterons

But the flux is supposed to be much smaller.



Korsmeier (2017; 1711.08465)



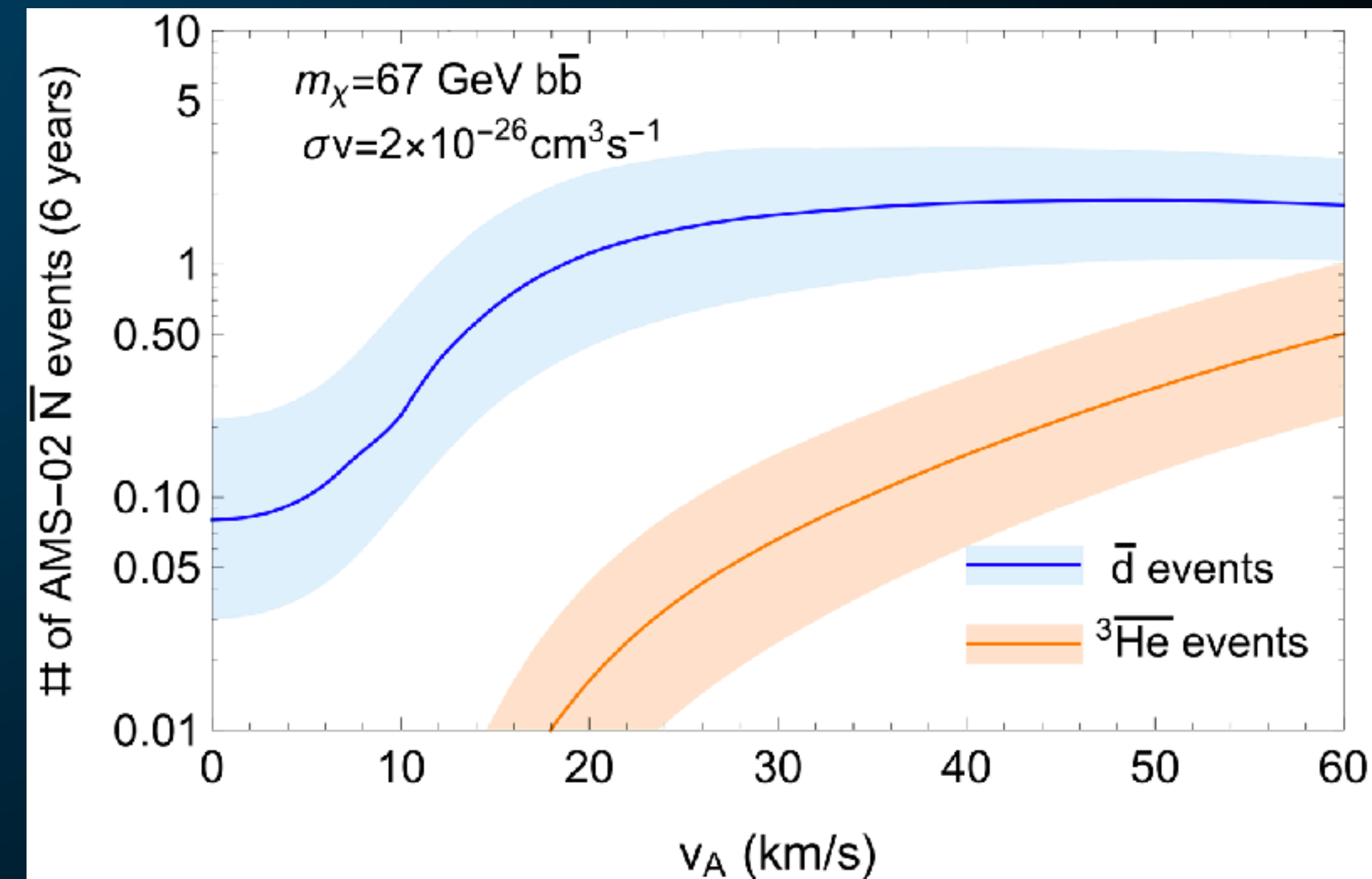
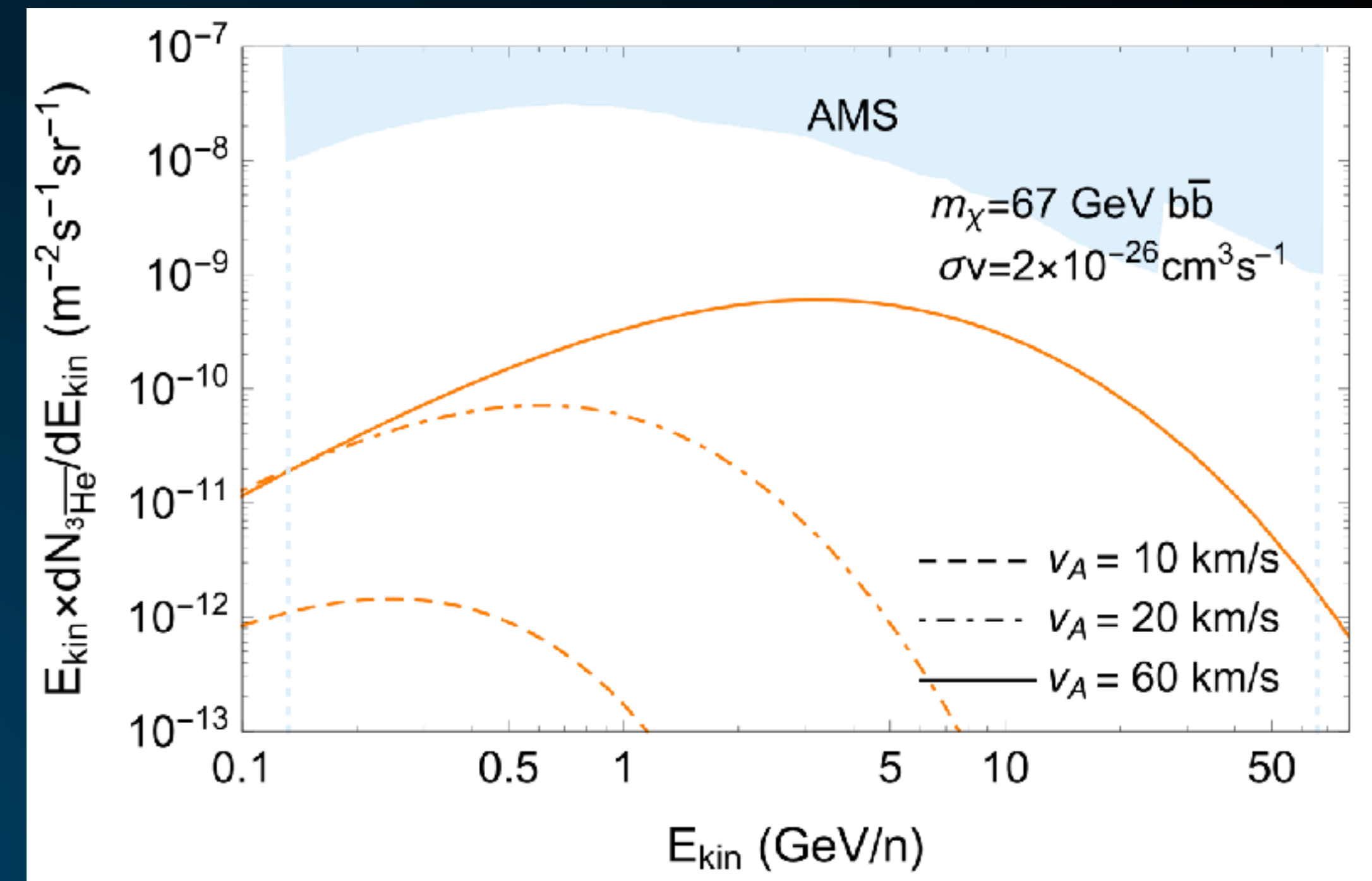
Poulin et al. (2018; 1808.08961)  
Carlson et al. (2014; 1401.2461)

# Enhancing the Dark Matter Flux

## Method I: Astrophysics

The dark matter induced-antihelium bump is at too low of an energy to be detected.

Use reaccelerating to boost the antihelium events into a detectable range.





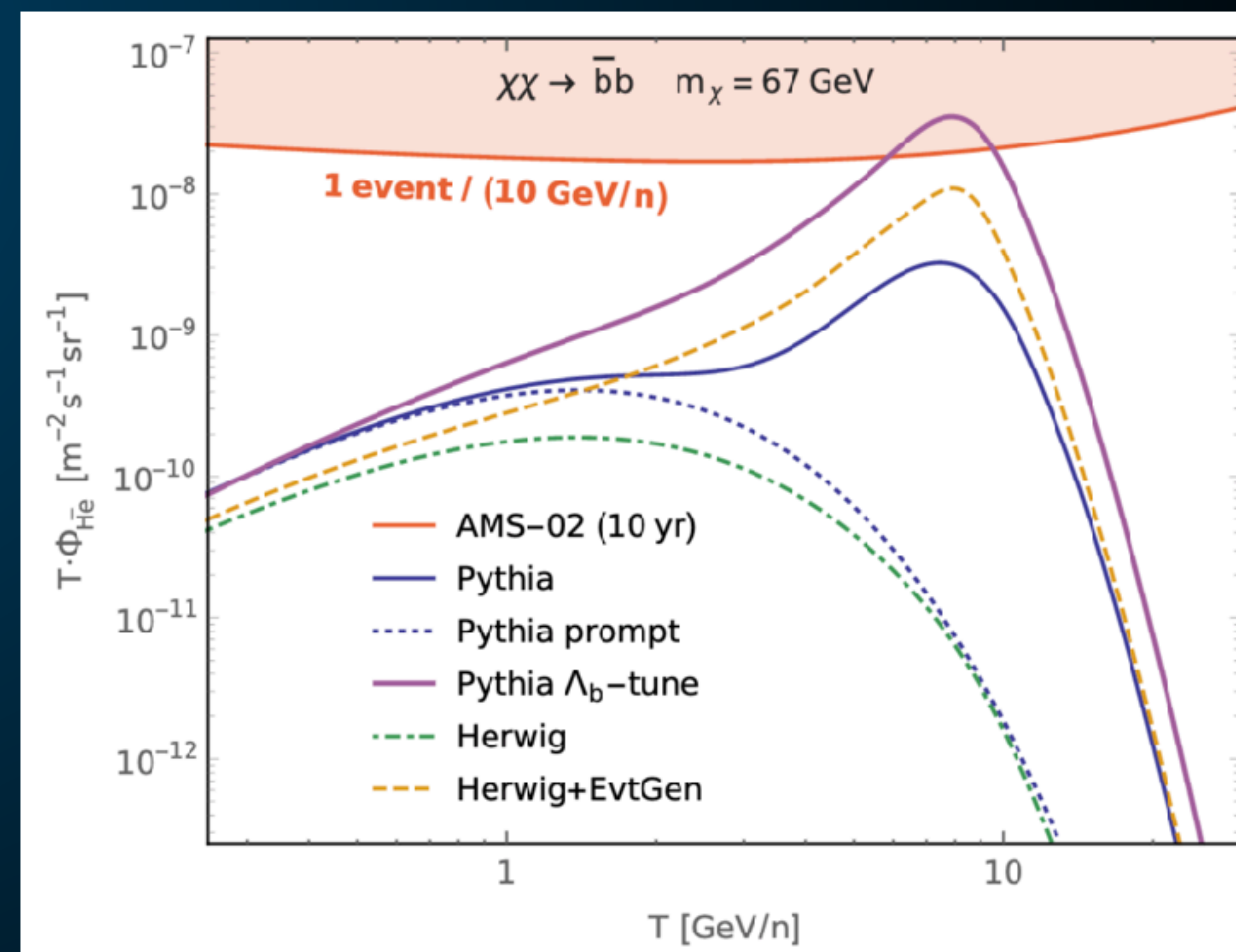
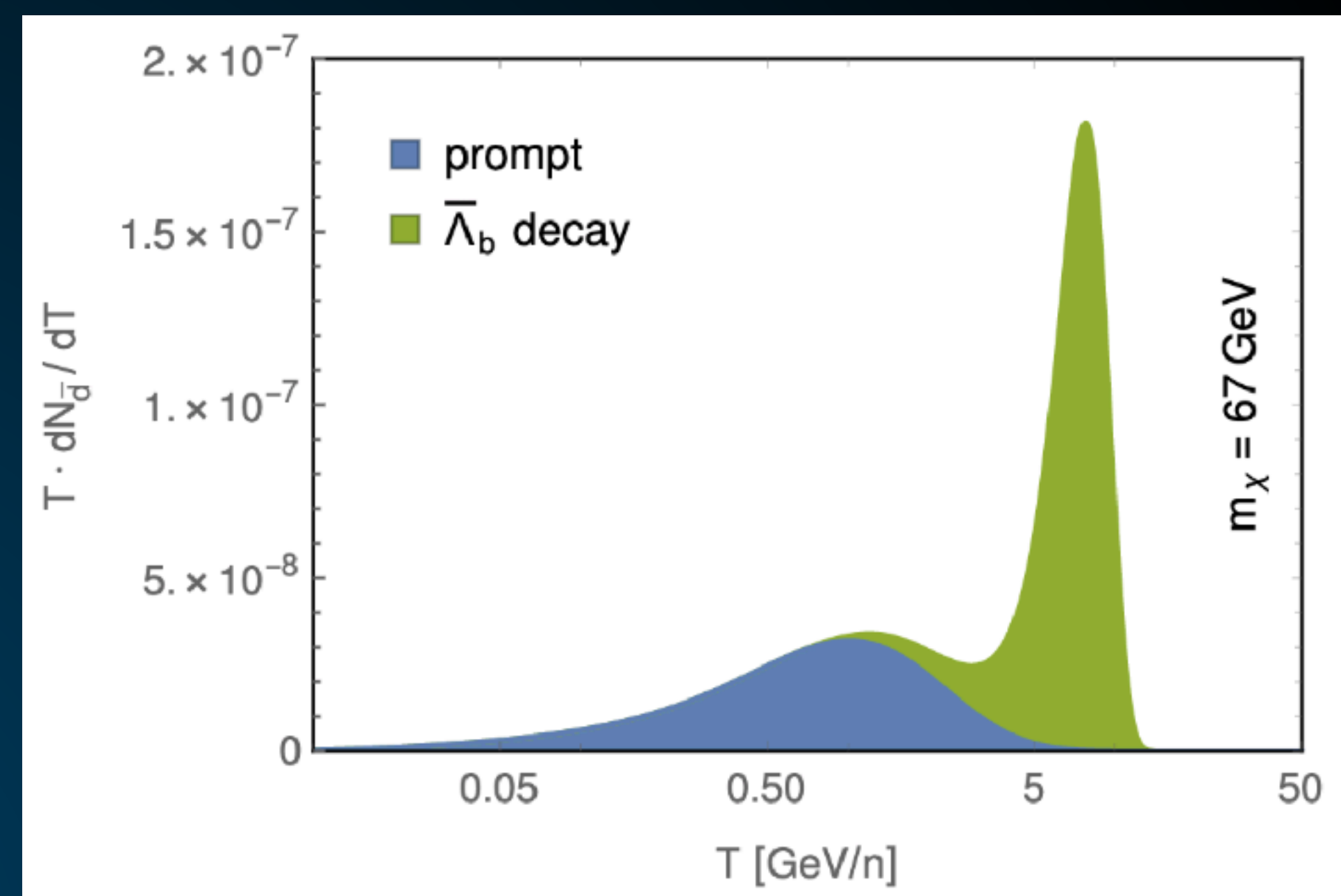
# Enhancing the Dark Matter Flux

## Method II: Particle Physics

Previous analyses may have missed the dominant antihelium production pathway from dark matter.

Including this term boosts the antihelium production rate by a factor of 100!

| Generator                    | P           | P [ $\Lambda_b$ -tune] | H     | H+EvtGen |
|------------------------------|-------------|------------------------|-------|----------|
| ${}^3\bar{\text{He}}$ events | 0.1 (0.007) | 0.9                    | 0.003 | 0.3      |
| $\bar{\text{d}}$ events      | 3.7 (3.5)   | 4.2                    | 1.7   | 2.1      |



# Enhancing the Dark Matter Flux

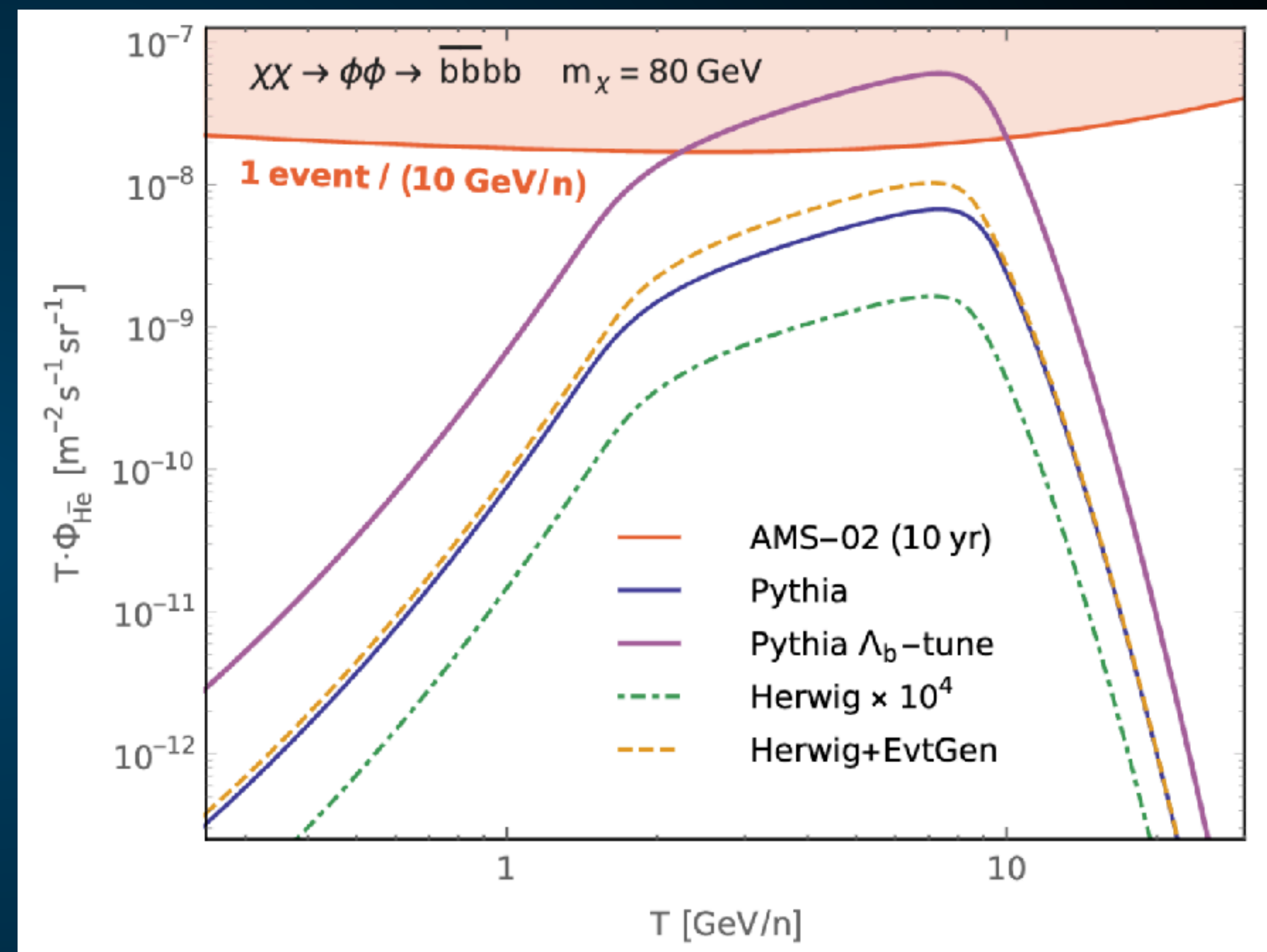
## Method III: Tuned Models

If the goal is to fit the antihelium data.

Can combine astrophysical and particle physics mechanisms.

Can develop non-standard particle physics models.

*Winkler & Linden (2020; 2020.16251)*



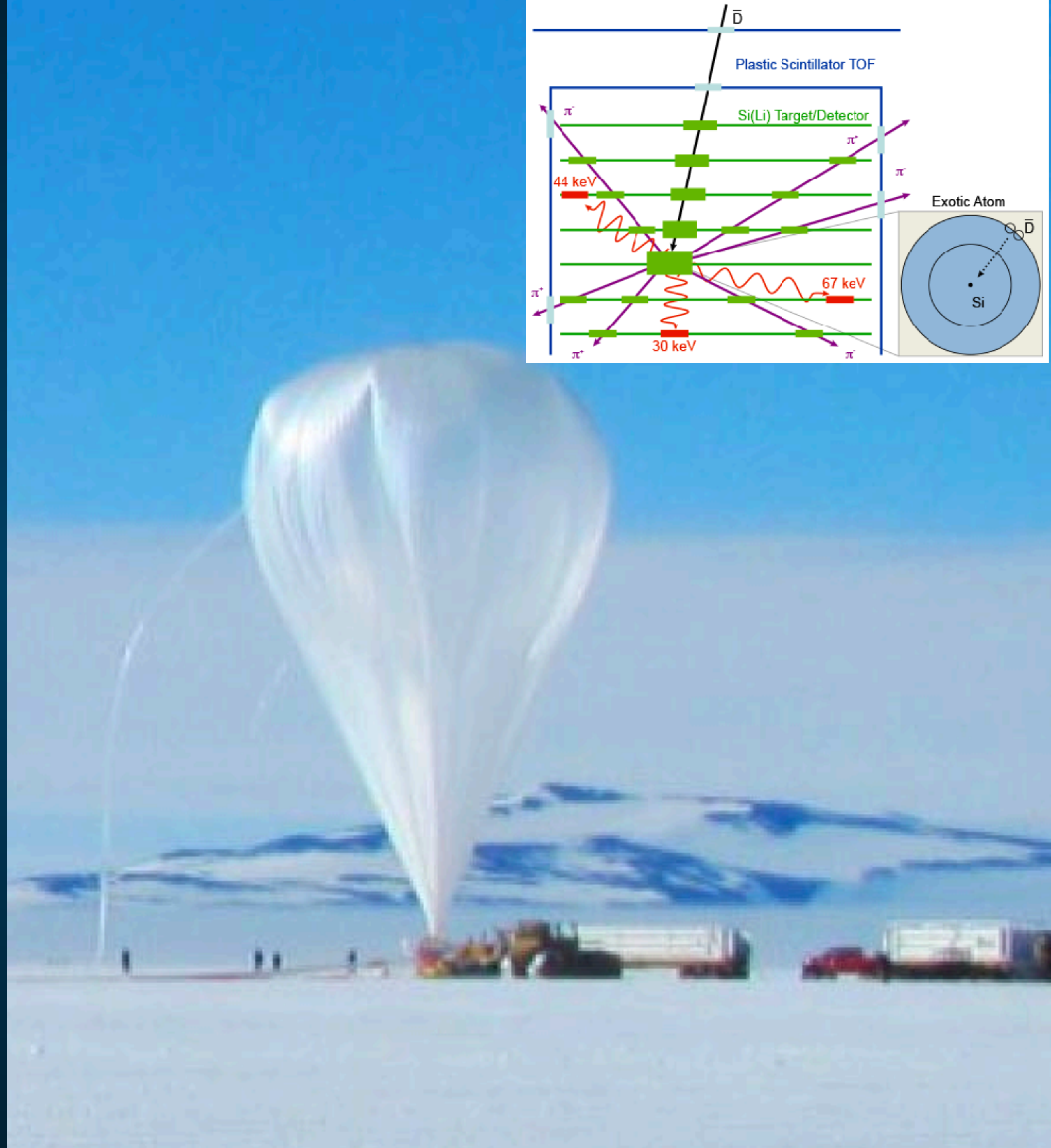
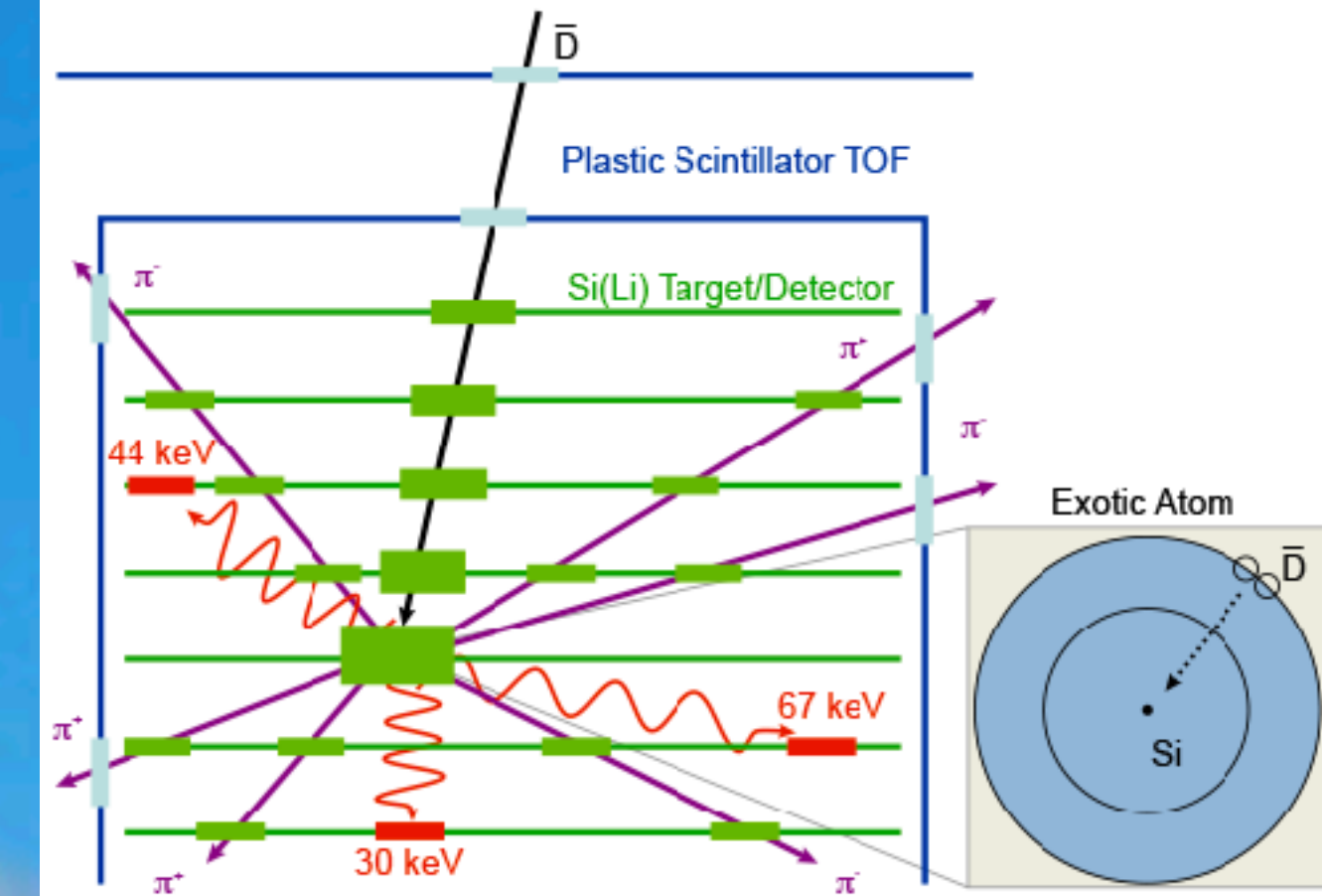
# Difficulties

## Method III: Tuned Models

Can we verify the observation of antihelium against a much larger background?

Can we produce enough  ${}^3\text{He}$  without violating antideuteron and antiproton constraints?

What about  ${}^4\text{He}$ ?



# Difficulties

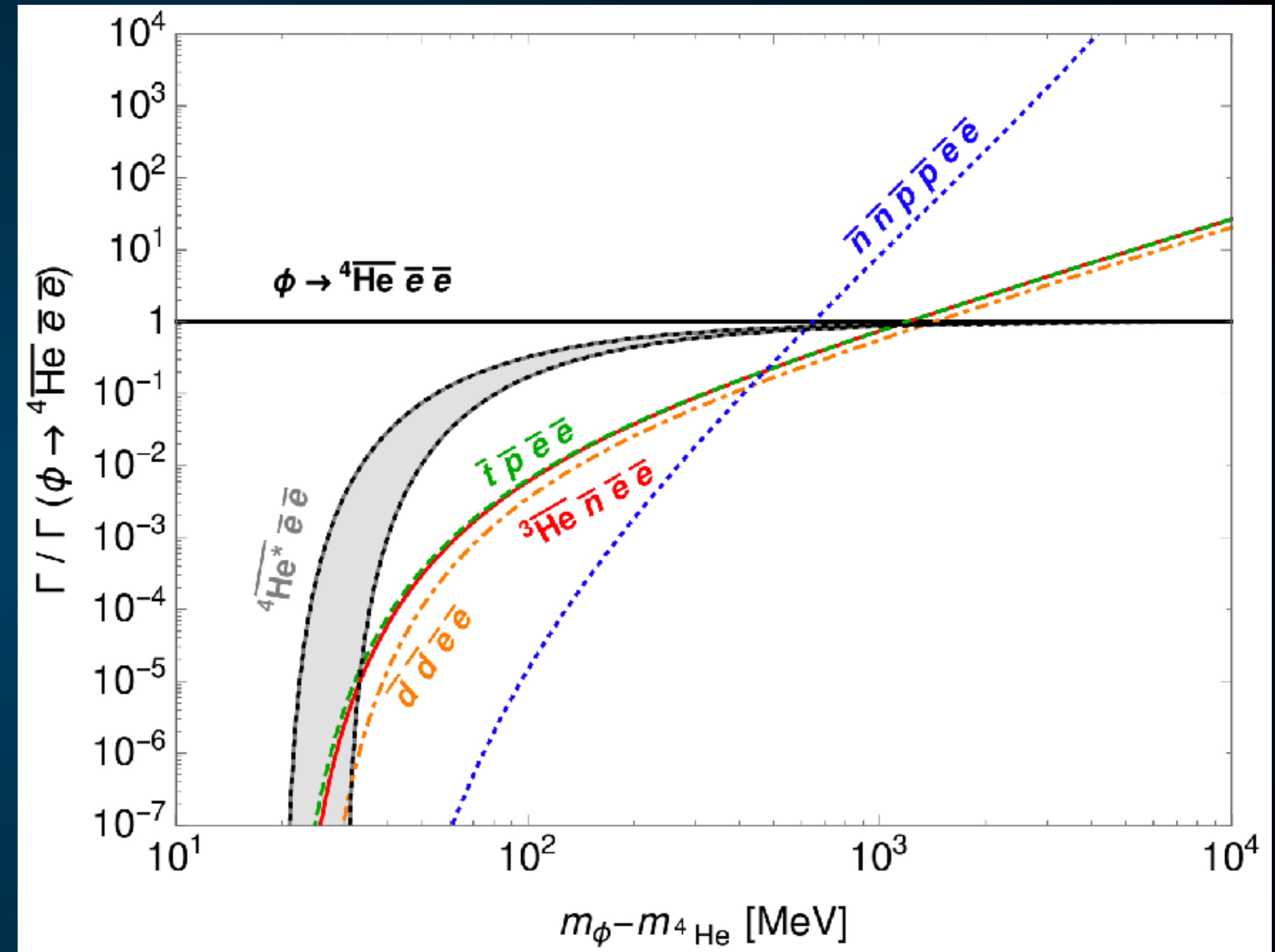
## Method III: Tuned Models

Heeck & Rajaraman (2019; 1906.01667)

Can we verify the observation of antihelium against a much larger background?

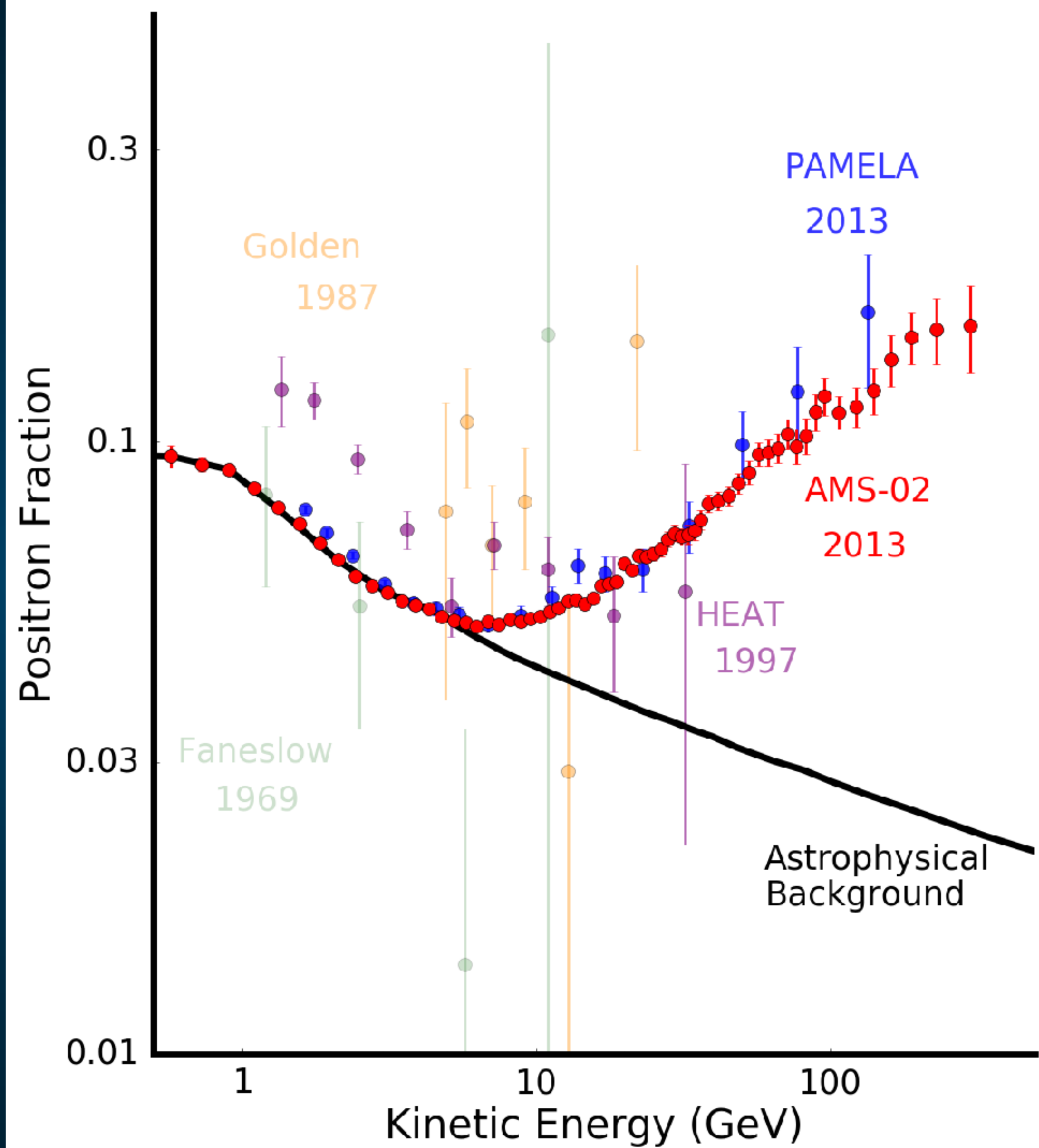
Can we produce enough  ${}^3\text{He}$  without violating antideuteron and antiproton constraints?

What about  ${}^4\text{He}$ ?



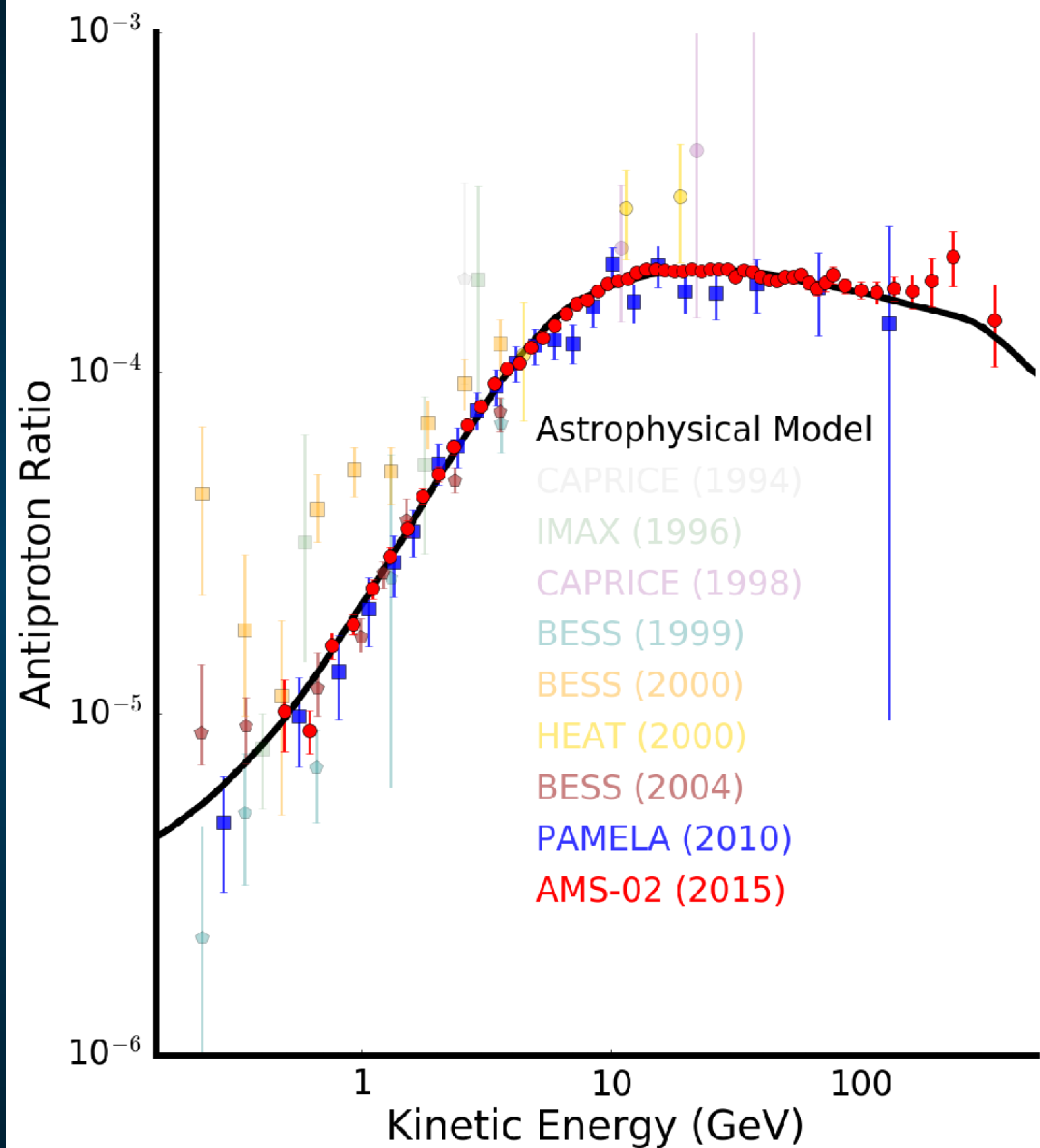
# Provocative Questions

Can we produce a robust calculation of  $e^+e^-$  from pulsars that allows us to search for dark matter?



# Provocative Questions

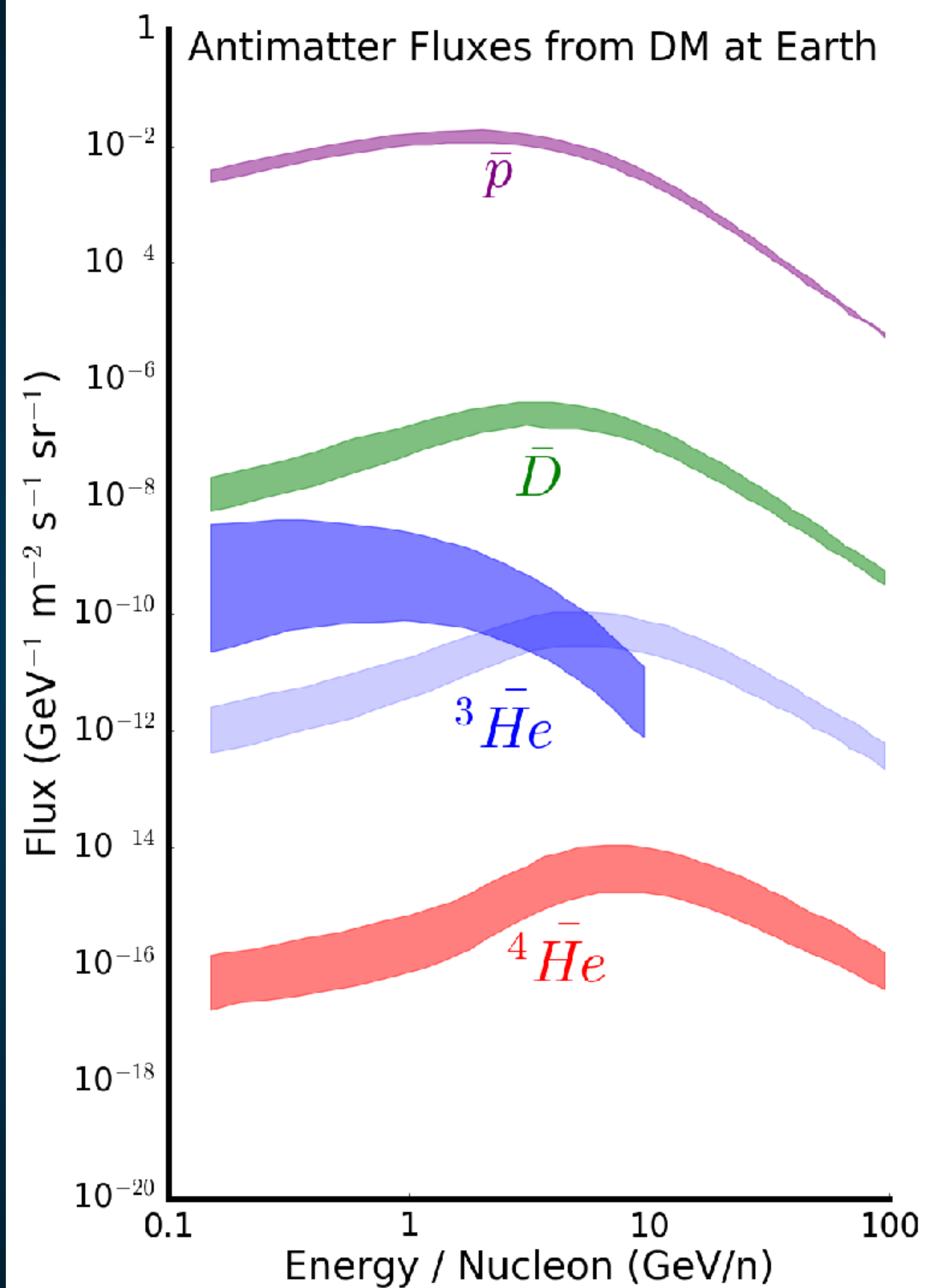
Is it possible to find dark matter as a  $O(1\%)$  effect?



# Provocative Questions

Can we understand nucleon coalescence sufficiently to predict the dark matter induced flux?

${}^4\text{He}$ ?



# Dark Matter Searches with Cosmic-Rays

Yesterday, Today, and Tomorrow

**Need to produce a complete model of antiprotons/antideuterons/antihelium from a dark matter annihilation model.**

**Need to constrain systematic uncertainties: instrumental, astrophysical, solar modulation.**

**More surprises may be in store!**





