

# SEARCHING FOR DARK MATTER IN THE MEV SKY

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THE OHIO STATE UNIVERSITY

CENTER FOR COSMOLOGY AND ASTROPARTICLE PHYSICS

# **EVERYTHING WE KNOW ABOUT DARK MATTER**

# CMB

- Galactic Rotation Curves
- Gravitational Lensing
- Baryonic Acoustic Oscillations

Parameter	Planck		Planck+lensing		Planck+WP	
	Best fit	68% limits	Bost fit	68% limits	Best fit	68% limits
Ω.,μ²	0.022068	0.02207 ± 0.00063	0.022242	$0.02217 \pm 0.00033$	0.022032	$0.02205 \pm 0.00028$
Ω <sub>c</sub> h <sup>2</sup>	0.12029	$0.1196 \pm 0.0031$	0.11805	$0.1186 \pm 0.0031$	0.12038	$0.1199 \pm 0.0027$
1000 <sub>MC</sub>	1.04122	$1.04132 \pm 0.00068$	1.04150	$1.04141 \pm 0.00067$	1.04119	$1.04131 \pm 0.00063$
τ	0.0925	$0.097 \pm 0.038$	0.0949	$0.089 \pm 0.032$	0.0925	0.069_0012
N <sub>x</sub>	0.9624	$0.9616 \pm 0.0094$	0.9675	$0.9635 \pm 0.0094$	0.9619	$0.9603 \pm 0.0073$
$\ln(10^{10}A_a)$	3.098	$3.103 \pm 0.072$	3.098	$3.085\pm0.057$	3.0980	3.069_0.027
Ω	0.6825	$0.686 \pm 0.020$	0.6964	$0.693 \pm 0.019$	0.6817	0.685-0.015
Ω	0.3175	$0.314 \pm 0.020$	0.3036	$0.307\pm0.019$	0.3183	0.315-0015
σ1	0.8344	$0.834 \pm 0.027$	0.8285	$0.823 \pm 0.018$	0.8347	$0.829 \pm 0.012$







CMB 13.7 billion years ago

# CMB

- Galactic Rotation Curves
- Gravitational Lensing
- Baryonic Acoustic Oscillations

10<sup>-25</sup> GeV

 $\sigma_X > R_{UFD}$ 

**10<sup>62</sup> GeV** m<sub>X</sub> > M<sub>UFD</sub>

slide concept courtesy of Asher Berlin

## **EVERYTHING WE KNOW ABOUT DARK MATTER**



# So, why look for dark matter in the MeV range?

Because we can.

 Many studies have been sensitive to GeV-scale dark matter, but most observations run out of steam in the MeV range.

Fermi-LAT gamma-ray observations are sensitive to DM heavier than ~5 GeV



A low energy mission could significantly improve this sensitivity.

Space-based cosmic-ray instruments are sensitive to signals from MeV dark matter.





However, the heliospheric potential is ~1 GV.

Significantly uncertain, prevents DM/background differentiation at low energies.

### Cuoco et al. (1610.03071)

Direct detection studies are not currently sensitive to dark matter masses below the proton mass.



May be improved someday with alternative technologies.

Many studies have been sensitive to GeV-scale dark matter, but most observations run out of steam in the MeV range.

Colliders can potentially constrain very light dark matter, but their sensitivity is highly model dependent.



Beam dump experiments may soon significantly increase the sensitivity of collider experiments in the MeV range.

## THE MEV SCALE - A WORLD OF POSSIBILITIES!



We need to transfer this picture of complementarity into the MeV regime.

## THE IMPETUS FOR MEV DARK MATTER - LESSONS FROM WIMPS



# So, why look for dark matter in the MeV range?

# What is the landscape of MeV dark matter?

# CMB

- Galactic Rotation Curves
- Gravitational Lensing
- Baryonic Acoustic Oscillations



density increasing number  $< \sigma v >$ comoving  $\mathrm{N}_{\mathrm{EQ}}$ 10-13 10  $10^{1}$  $10^{2}$  $10^{3}$ time- $\rightarrow$ m

A particle with a weak interaction cross-section and a mass on the weak scale is expected to naturally obtain the correct relic abundance through thermal freeze-out in the Earth universe.

$$\left(\frac{\Omega_{\chi}}{0.2}\right) \simeq \frac{x_{\rm f.o.}}{20} \left(\frac{10^{-8} \,\,{\rm GeV^{-2}}}{\sigma}\right)$$

 $\langle \sigma v \rangle \sim 10^{-8} \; {\rm GeV^{-2}} \left( 3 \times 10^{-28} \; {\rm GeV^2} \; {\rm cm^2} \right) \; 10^{10} \; \frac{{\rm cm}}{{\rm s}} = 3 \times 10^{-26} \; \frac{{\rm cm^3}}{{\rm s}}$ 

- At high-masses we have the unitarity bound:
  - Electroweak coupling constant is proportional to m<sub>x</sub><sup>2</sup>.
  - Coupling constant must be 1, or loop contributions diverge.



# What about the lower-limit?



slide concept courtesy of Asher Berlin

# Lee-Weinberg bound:

$$\Omega_{\chi} h^2 \sim 0.1 \frac{10^{-8}}{{\rm GeV}^{-2}} \cdot \frac{1}{G_F^2 \; m_{\chi}^2} \sim 0.1 \left(\frac{10 \; {\rm GeV}}{m_{\chi}}\right)^2$$

# PHYSICAL REVIEW LETTERS

NUMBER 4

Cosmological Lower Bound on Heavy-Neutrino Masses

25 JULY 1977

Eenjamin W. Lee<sup>40</sup> Fermi National Accelerator Laboratory,<sup>(b)</sup> Batovia, Illinois 50510

and

Steven Weinberg<sup>[0]</sup> Stanford University, Physics Department, Stanford, California 94305 (Received 13 May 1977)

The present cosmic mass density of possible stable neutral heavy leptons is calculated in a standard cosmological model. In order for this density not to exceed the upper limit of  $2 \times 10^{-10}$  g/cm<sup>3</sup>, the lepton mass would have to be greater than a lower bound of the order of 2 GeV.

- Under the assumption that the interaction is electroweak (G<sub>F</sub><sup>2</sup>), the dark matter mass must be larger than 10 GeV.
- Two caveats:
  - Dark matter may still be thermal, if the interaction that thermalizes DM is not electroweak.

VOLUME 39

Electroweak models may still be reasonable, but only if a light scalar exists to boost the annihilation rate.

# N<sub>eff</sub> bounds:

If dark matter falls out of thermal equilibrium during or after neutrino decoupling, then dark matter contributes to the effective number of degrees of freedom.

# Bounds dark matter to be >~ 3-10 MeV

#### Limits on MeV Dark Matter from the Effective Number of Neutrinos

Chiu Man Ho and Robert J. Scherrer Department of Physics and Astronomy, Vanderbilt University, Nashville, TN 37235

Thermal dark matter that couples more strongly to electrons and photons than to neutrinos will heat the electron-photon plasma relative to the neutrino background if it becomes nonrelativistic after the neutrinos decouple from the thermal background. This results in a reduction in  $N_{eff}$ below the standard-model value, a result strongly disfavored by current CMB observations. Taking conservative lower bounds on  $N_{eff}$  and on the decoupling temperature of the neutrinos, we derive a bound on the dark matter particle mass of  $m_{\chi} > 3-9$  MeV, depending on the spin and statistics of the particle. For *p*-wave annihilation, our limit on the dark matter particle mass is stronger than the limit derived from distortions to the CMB fluctuation spectrum produced by annihilations near the epoch of recombination.





slide concept courtesy of Asher Berlin

# A NEW PICTURE - NONSTANDARD WIMP MODELS

- Thinking outside the standard thermal weakly interacting particle box:
  - Asymmetric Dark Matter (e.g. Lin et al. 1111.0293)
  - Secluded Dark Matter (e.g. Pospelov et al. 0711.4866)
  - MeV Sterile neutrinos (Huang & Nelson, 1306.6079)
  - Strongly Interacting Massive Particles (e.g. Hochberg et al. 1402.5143)
  - Hidden Dark Sectors (e.g. Hufnagel et al. 1712.03972)
  - Late Decay to DM (e.g. Choquette et al. 1604.01039)
  - Also: Mirror Dark Matter, Atomic Dark Matter, Magnetic Dark Matter, WIMPless dark matter, etc.

## **INVISIBLE DARK MATTER MODELS**



Need to be careful - many of these models produce no dark matter annihilation signal.

Our dark matter models need to produce <u>a little bit of light.</u>

# A NEW PICTURE

# What can light dark matter annihilate into?

- (i)  $\gamma\gamma$ : Accessible at all energies. The final state is C-even.
- (ii)  $\gamma \pi^0$ : Accessible for  $\sqrt{s} > m_{\pi^0}$ . The final state is C-odd.
- (iii)  $\pi^0 \pi^0$ : Accessible for  $\sqrt{s} \ge 2m_{\pi^0}$ . The final state is *C*-even.
- (iv)  $\pi^+\pi^-$ : Accessible for  $\sqrt{s} \ge 2m_{\pi^{\pm}}$ . The final state is C-even or C-odd.

(v)  $\bar{\ell}\ell$  ( $\ell = e, \mu, \nu$ ): Accessible for  $\sqrt{s} \geq 2m_{\ell}$ . The final state is either C-odd or is weak suppressed.

This significantly limits the possible annihilation channels, especially for masses below 135 MeV.

# THE CMB DOUBLE-BIND

# CMB Bounds:

- Any energy deposition during recombination affects the CMB anisotropy
- The constraint trends as
   <σv> ~ 1/m
- Small changes due to the fraction of total power which is transferred to electrons.





# Slatyer (1506.03811)

# THE CMB DOUBLE-BIND

- Moreover the annihilation rate at low masses should exceed 3 x 10<sup>-26</sup> cm<sup>3</sup> s<sup>-1</sup>.
- CMB bounds appear to remove most of the parameter space for the thermal DM.



- We can try to model build around this, for example invoking p-wave suppressed annihilations:
  - But these typically suppress gamma-ray emission today even more than the CMB!

## **COMPARISON OF CMB AND DWARF GAMMA-RAY CONSTRAINTS**

#### Gonzalez-Morales et al. (1705.00777)



- CMB constraints fall between 1x10<sup>-28</sup> and 1x10<sup>-27</sup> in the 0.1-1 GeV range.
- These constraints are likely to outperform future MeV gamma-ray experiments for annihilation to charged final states.

## **COMPARISON OF CMB AND GALACTIC CENTER GAMMA-RAY CONSTRAINTS**

#### Gonzalez-Morales et al. (1705.00777)



- CMB constraints are typically even more sensitive than annihilation in the Galactic center.
- There is allowable parameter space if the dark matter density near the galactic center is very peaked.

# So the combination of

- (i)  $\gamma\gamma$ : Accessible at all energies. The final state is C-even.
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## **COMPARISON OF CMB AND DWARF GAMMA-RAY CONSTRAINTS**

#### Gonzalez-Morales et al. (1705.00777)



In the case of annihilation to neutral particles:

- The CMB constraints stay the same
- The gamma-ray sensitivity improves drastically.
- Gamma-Rays are the most sensitive channel!

## **COMPARISON OF CMB AND GALACTIC CENTER GAMMA-RAY CONSTRAINTS**

#### Gonzalez-Morales et al. (1705.00777)



- In the Galactic center, nearly every density profile outperforms CMB constraints.
- Can potentially observe the same dark matter signal in two ROIs!

# **DECAYING DARK MATTER IS, AS ALWAYS, REASONABLE**

Dark Matter decay rate is not enhanced in the early universe.

In general, the MeV scale of these models is not highly motivated (similar to GeV range).

Analyzing these models will proceed similarly to GeV searches.

Boddy et al. (1606.07440)



# MOREOVER, CAN DECAY AT LATE TIMES INTO ANNIHILATING COMPONENT

Some well motivated models generate the dark matter asymmetrically, and then decay to the annihilating component.



These models can naturally produce MeV emission.



Hardy et al. (1402.4500)

# A NEW PICTURE

# What can light dark matter annihilate into?

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- (iv)  $\pi^+\pi^-$ : Accessible for  $\sqrt{s} \ge 2m_{\pi^{\pm}}$ . The final state is C-even or C-odd.

(v)  $\bar{\ell}\ell$  ( $\ell = e, \mu, \nu$ ): Accessible for  $\sqrt{s} \geq 2m_{\ell}$ . The final state is either *C*-odd or is weak suppressed.

In general - annihilation into uncharged states is better for gamma-ray observation.

# ANOTHER BENEFIT OF ANNIHILATION INTO UNCHARGED PARTICLES

 If annihilation proceeds to charged particles, secondary emission becomes important.





This signal will diffuse, complicating analyses.

Must be taken into account.

1.) Standard Thermal Electroweak WIMP largely excluded from gamma-ray detection in the MeV range

Can still be thermal if interaction is not electroweak

- Some exceptions, like p-wave suppressed interactions which may be observable in the Galactic center.

2.) Never fear!

- There are many models for producing MeV dark matter!
- Many models will have unique signatures:
  - \* Harder to use a one-size fits all approach
  - \* Easier to distinguish between models if you see something.

## EVERYTHING WE KNOW ABOUT DARK MATTER



# So, why look for dark matter in the MeV range?

# New Techniques for looking at MeV dark matter!

## **EVERYTHING WE KNOW ABOUT DARK MATTER**

Galactic center: Satellites: Good statistics but source Low background confusion/diffuse background and good source ID, Milky Way halo: but low statistics Large statistics but diffuse background

Spectral lines:

No astrophysical uncertainties, good source ID, but low statistics

Galaxy clusters: Low background but low statistics Extragalactic: Large statistics, but astrophysics, Galactic diffuse background

# THE ANNIHILATION SPECTRA WILL UNIVERSALLY BE HARD

 Annihilation of 80 MeV DM -> e+e- produces gammarays primarily through final state radiation and bremsstrahlung.

 These spectra are brightest at the dark matter mass.



Comparison of multiple targets becomes significantly easier at MeV energies - look for identical energy cuts!

# WHAT ABOUT THE BACKGROUNDS?



# ADDITIONAL COMPLEXITIES IN BACKGROUND MODELING

$$D_{xx} = eta D_0 igg( rac{
ho}{4\,{
m GV}} igg)^{\delta}$$

$$egin{array}{rll} t^p_{
m pp} &\simeq& 3.1 imes 10^5 \ {
m yr} \ \left( rac{n_H}{120 \ {
m cm}^{-3}} 
ight)^{-1} \ , \ t^e_{
m inztn} &\simeq& 6.7 imes 10^5 \ {
m yr} \ \left( rac{E}{{
m GeV}} 
ight) \ \left( rac{n_H}{120 \ {
m cm}^{-3}} 
ight)^{-1} \ , \ t^e_{
m brems} &\simeq& 2.4 imes 10^5 \ {
m yr} \ \left( rac{n_H}{120 \ {
m cm}^{-3}} 
ight)^{-1} \ , \end{array}$$

Assuming Kolmogorov Diffusion in the MeV regime:

- A ~10 GeV proton which produces a 1 GeV gamma-ray travels approximately: 2.8 kpc.
- A 100 MeV electron which produces a ~1-10 MeV gamma-ray travels approximately: 0.5 kpc.

MeV diffuse emission will be more sensitive to local injection around 100 MeV.

# ADDITIONAL COMPLEXITIES IN BACKGROUND MODELING



# In general models are not yet equipped for this detail.

# WHAT ABOUT THE BACKGROUNDS?



# WHAT ABOUT THE BACKGROUNDS?



Porter et al. (1708.00816)

# The ISRF energy density is extremely smooth

Most of the substructure washes away, and you get very smooth background gamma-ray emission.

- Propagation in 10 Myr at 100 GeV:
  - Diffusion: ~2.1 kpc
  - Convection (20 km/s) ~100 pc

- Propagation in 10 Myr at 10 MeV:
  - Diffusion: ~466 pc
  - Convection (20 km/s) ~200 pc

# **DIFFICULTIES IN DIFFUSE BACKGROUND MODELING**



# **OPPORTUNITIES IN DIFFUSE BACKGROUND MODELING!**

The MeV sky might shed significant new light on galactic cosmic-ray diffusion!

We will soon have a diffuse sky from ~1 MeV to ~100 TeV (HAWC). Exciting potential for significant advancement.



Cosmic-Ray Propagation Codes are rising to the challenge (e.g. PICARD; 1701.07285).

- Direct gamma-rays from dark matter annihilation will have a profile that is relatively energy independent.
- In the GeV range, excesses at 10 GeV are also likely to be excesses at 500 MeV
  - same production Mechanism

Multi-wavelength models may provide more systematic cross-checks in the MeV regime.

## **EVERYTHING WE KNOW ABOUT DARK MATTER**



# So, why look for dark matter in the MeV range?

# **Does the MeV lamppost connect to the GeV range?**

- Fermi-LAT Constraints on 100 GeV dark matter
  - 1-5 GeV most sensitive

- For 5-10 GeV dark matter?
  - Fermi already shows signs of lost sensitivity.
  - AMEGO can compete in this range of "standard" dark matter.
  - Interesting range due to asymmetries!





Should not forget the capability of MeV instrumentation to solve current uncertainties from Fermi-LAT observations.

The separation of ~10-100 GeV DM from background (including π⁰-decay and blazars) depends sensitively on the strength of ~100 MeV gamma-ray limits.





Hooper & TL (1503.06209)

- Another example is the observation of spatial extension in unassociated Fermi-LAT sources
  - In this scenario, improved angular resolution at ~100 MeV is particularly important.

Because many of these sources are bright, instruments like AMEGO should have plenty of sensitivity.





synergy with radio telescopes, dependence on leptonic final states

Calore et al. (1409.0042)

## **COMPLEMENTARITY WITH TEV DARK MATTER**

Exciting dark matter is another natural model for MeV instruments to investigate.

#### An X-Ray Line from eXciting Dark Matter

Douglas P. Finkbeiner<sup>1</sup> and Neal Weiner<sup>2</sup> <sup>1</sup>Center for Particle Astrophysics, Hurvard University, Cambridge, MA <sup>3</sup>Center for Cosmology and Particle Physics, Department of Physics, New York University, New York, NY 10003 (Dated: March 4, 2014) Abstract

The eXciting Dark Matter (XDM) model was proposed as a mechanism to efficiently convert

- For many TeV -> PeV dark matter sectors, the momentum transfer in a dark matter scattering collision is in the keV to MeV range.
  - Gamma-Ray Morphology will track dark matter mass and rotation velocity.
- MeV sensitivity is complementary for known excesses (e.g. PAMELA).

# **1.) Current dark matter limits in the MeV range are significantly lacking.**

# **2.)** Dark Matter model building is not as straightforward in the MeV range, compared to GeV energies.

a.) More diversity in models -> more diversity in signals.
b.) More diversity in models -> more diversity in techniques.

**3.) MeV sensitivity is key to unlock the full power of Fermi in the GeV range.** 

## THE MEV SCALE - A RANGE OF POSSIBILITIES!

# The next decade is likely to include a significant push to MeV energies.



**Cosmic Visions Whitepaper** 

Direct Detection will utilize electron recoils and superconductors.

# THE MEV SCALE - A WORLD OF POSSIBILITIES!

# The next decade is likely to include a significant push to MeV energies.



Izaguirre et al. (1411.1404)

Fixed-target experiments can make colliders significantly more sensitive to MeV-scale dark matter

# THE MEV SCALE - A RANGE OF POSSIBILITIES!



This vision is obtainable at MeV energies!